

**Symbology,
Cluster Algebras
and
Amplitudes+**

**Anastasia Volovich
Brown University**

Corfu 2026

Outline

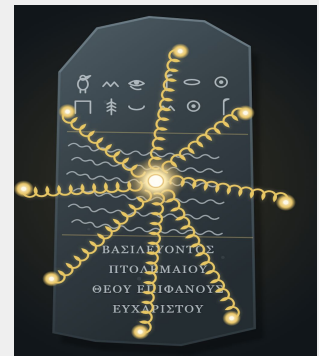
- Introduction
- N=4 Yang-Mills Amplitudes
- Tools: Symbolology and Cluster Algebras
- Applications:
 - Feynman integrals
 - Energy correlator integrals
 - Cosmological correlator integrals

(Mark's talk)
- Outlook

N=4 SYM: Quantum Rosetta Stone

Planar N=4 Yang-Mills (SYM) amplitudes have the most amazing mathematical structures:

- Dual Conformal/Yangian Symmetry
- Amplituhedron
- Cluster Algebras



It is a simple model that (likely) can be solved exactly, and provides general tools and intuition.

Arkani-Hamed, Dixon, McLeod, Spradlin, Trnka, AV [2207.10636](#)

n-point L-loop amplitudes in N=4 SYM: status

- all n, all k L=0
- $n < 6$ all L
- $n=6$ L=8
- $n=7$ L=5
- all n MHV L=2
- $n=8$ MHV L=3
- $n=8, 9$ NMHV L=2

Britto, Cachazo, Feng, Witten (BCFW) [0501052](#)

Bern, Dixon, Smirnov (BDS) [0505205](#)

Caron-Huot, Dixon, Drummond, Dulat,
Foster, Gurdogan, von Hippel, McLeod,
Papathanasiou, review: [2005.06735](#)

Dixon, Liu [2308.08199](#) He et al [2511.09669](#)

Caron-Huot [1105.5606](#)

Li, Zhang [2110.00350](#)

He, Li, Zhang [1911.01290](#)
[2009.11471](#)

As of today, Dixon and Liu 8-loop results stand at the forefront of perturbative calculations in N=4 SYM (and quantum field theory in general) and, in my opinion, serve as a benchmark - or a “challenge” - against which any “AI takeover” should be measured.

Lance's Bootstrap Program
antipodal duality • 8-loop 6-point amplitude • Steinmann relations • cluster adjacency

3-point form factor kinematic space
 $\hat{u}$
 $\hat{v}$
 $\hat{w}$
Simplex $\Delta \geq 0$

ANTIPODAL DUALITY

6-point amplitude kinematic space
1 2 3 4 5 6
Parity-preserving slice $\Delta = 0$

8-LOOP 6-POINT AMPLITUDE
WEIGHT 16
HEXAGON
FUNCTIONS
 $\{11, 1, 1, 1, 1, 1\}$
COPRODUCTS

$\Delta = 0$ SURFACE IN KINEMATIC SPACE

BOOTSTRAP
SYMBOLIC METHODS
HEXAGON FUNCTIONS
ITERATED INTEGRALS

LANCE DIXON
Stanford University

PHYSICAL SINGULARITIES
STEINMANN RELATIONS
CLUSTER ADJACENCY
PARITY PRESERVING
HEXAGON & HEPTAGON BOOTSTRAP

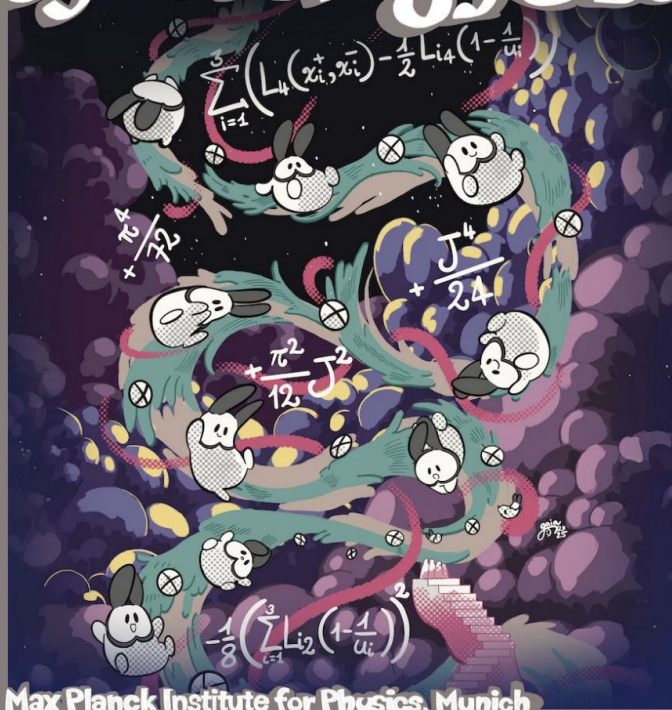
How are there computations done?

There are two important ingredients:

- 6- and 7-point amplitudes are MPLs, i.e. they can be studied via **SYMBOLLOGY**
- symbol alphabet for n-point amplitudes is described by **$\text{Gr}(4,n)$ CLUSTER ALGEBRA**

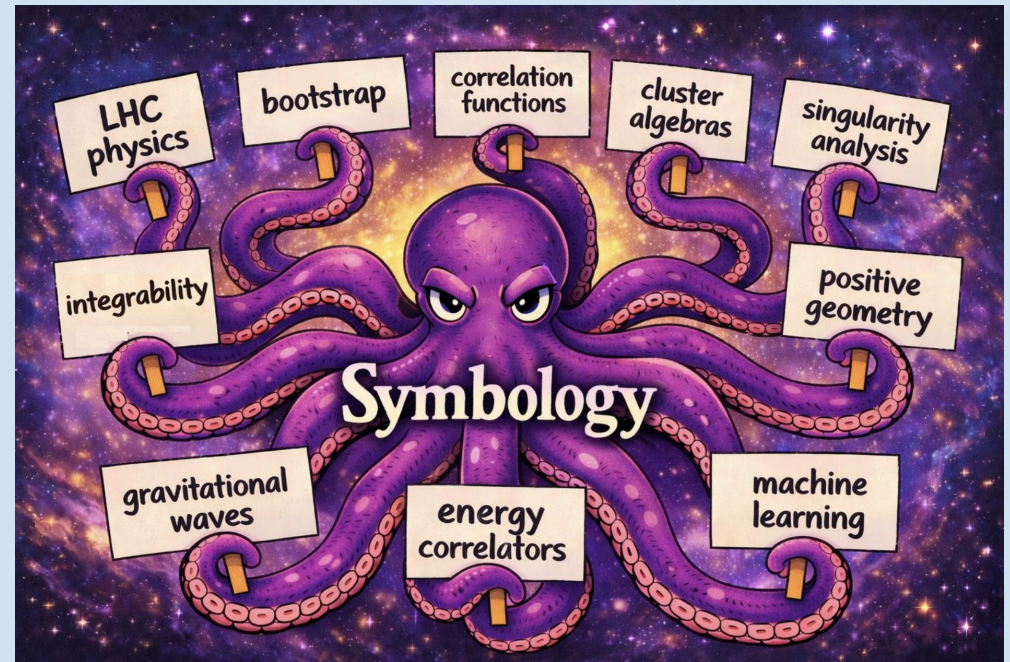
15-18 December 2025

Symbology@15



Max Planck Institute for Physics, Munich

Organizers: Johannes Henn, Marcus Spradlin, Anastasia Volovich
Scientific Coordinator: Diana López-Falcón
Design: Gaë Fontana



SYMBOLLOGY

- Uniform Transcendentality **m=2L @ L-loops**
- Polylogarithms

$$Li_m(x) = \int_0^x \frac{dt}{t} Li_{m-1}(t) \quad Li_1(x) = -\log(1-x)$$

- Generalized Polylogarithms (MPLs)

$$G(a_1, \dots, a_m; x) = \int_0^x \frac{dt}{t - a_1} G(a_2, \dots, a_m; t) \quad G(; x) = 1$$

- Elliptic Polylogarithms (eMPLs) **n=10 L=2 k=3**

Bourjaily, Broedel, Chaubey, Duhr, Frellesvig, Hidding, Marzucca, McLeod, Spradlin, Tancredi, Vergu, Volk, AV, Hippel, Weinzierl, Wilhelm, Zhang [2203.07088](#)

SYMBOLGY

Goncharov, Spradlin, Vergu, AV 1006.5703

$$dF^{(m)} = \sum_i F_i^{(m-1)} d \log x_i$$

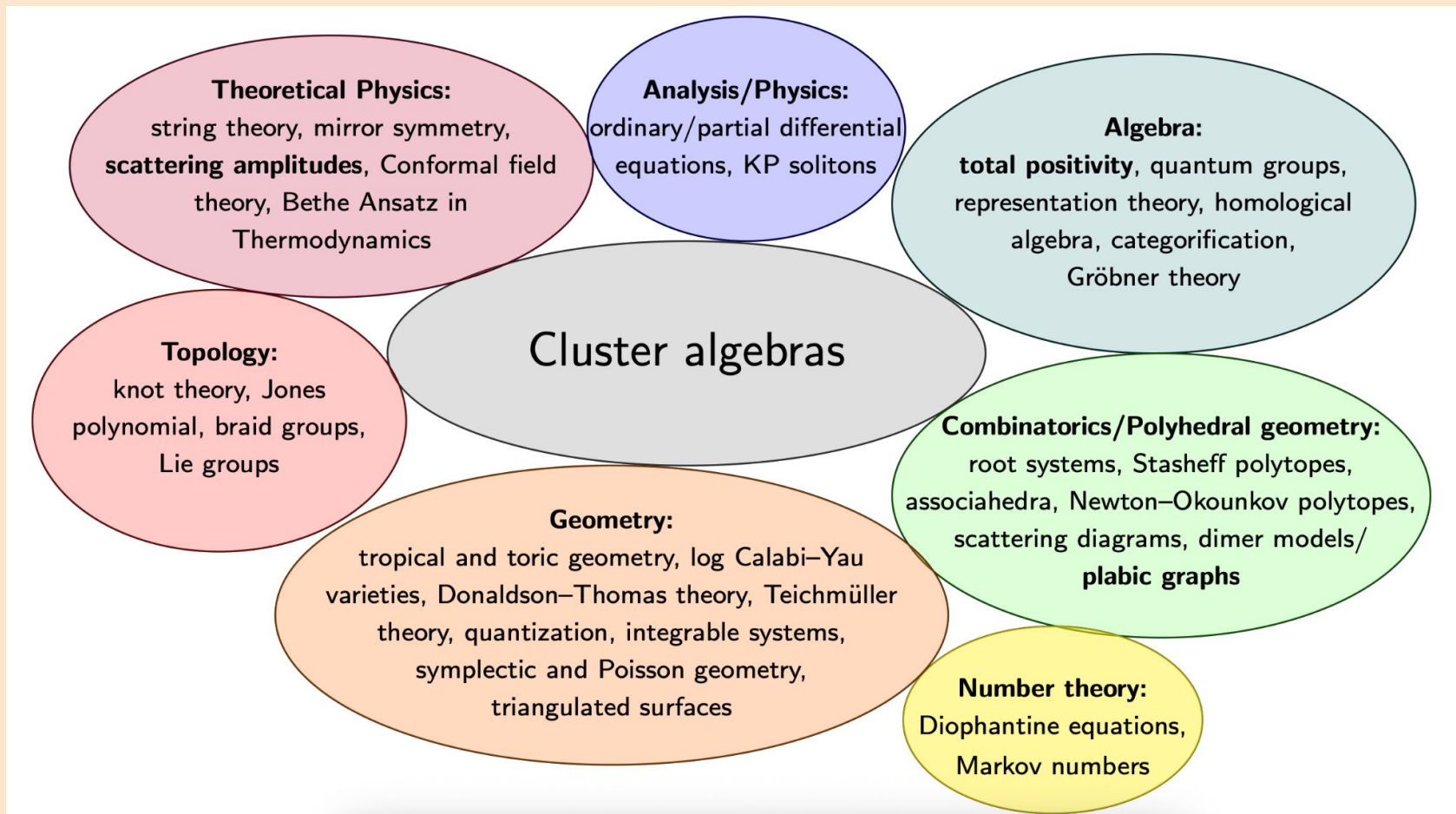
$$S(F^{(m)}) := \sum_i S(F_i^{(m-1)}) \otimes x_i$$

$$S(\log x) = x \quad S[Li_2(x)] = -(1 - x) \otimes x$$

The arguments that appear in a symbol are called **symbol letters** and collectively they are called **symbol alphabet**.

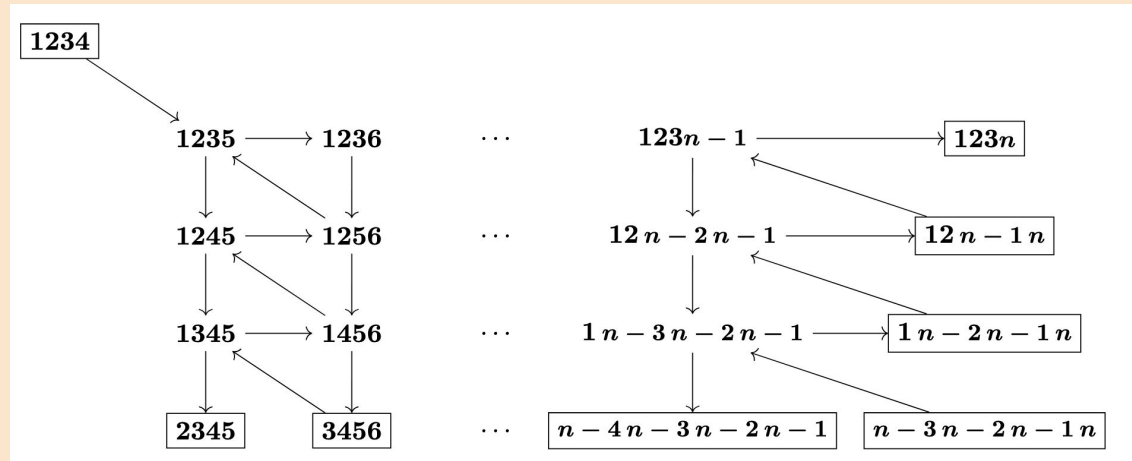
CLUSTER ALGEBRAS

MSC13F60



Grassmannian Cluster Algebra $\text{Gr}(4,n)$

Initial Seed



Mutation Rule

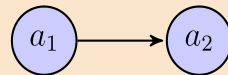
$$a_k \rightarrow a'_k = \frac{1}{a_k} \left(\prod_{\text{arrows } i \rightarrow k} a_i + \prod_{\text{arrows } k \rightarrow j} a_j \right)$$

Cluster Variables

$$\{a_k\}$$

Fomin, Zelevinsky [0104151](#), Scott [0311148](#), Gekhtman, Shapiro, Vainshtein [0208033](#), Williams [1212.6262](#), Fomin, Williams, Zelevinsky [“Introduction to Cluster Algebras”](#)

A₂ Cluster Algebra



$$a_k \rightarrow a'_k = \frac{1}{a_k} \left(\prod_{\text{arrows } i \rightarrow k} a_i + \prod_{\text{arrows } k \rightarrow j} a_j \right)$$

$$a_1, a_2, a_3 = \frac{1 + a_2}{a_1}, a_4 = \frac{1 + a_1 + a_2}{a_1 a_2}, a_5 = \frac{1 + a_1}{a_2}$$

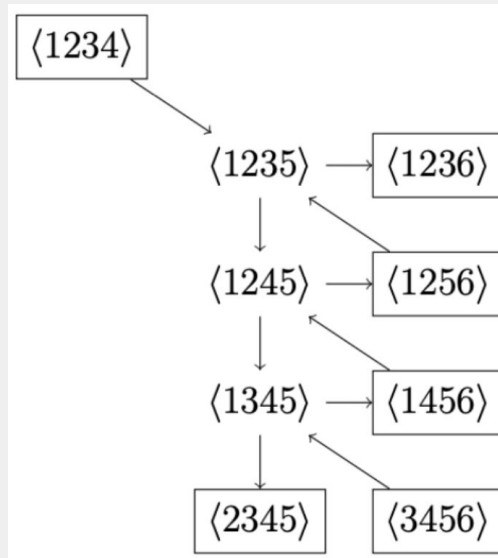
Fomin, Williams, Zelevinsky “Introduction to Cluster Algebras”

So far,
I have explained to you two main ingredients:

- 6-, 7-point amplitudes are MPLs, i.e. they can be studied via **SYMBOLGY**
- symbol alphabet for n-point amplitudes is described by **Gr(4,n) CLUSTER ALGEBRA**

Next,
let us look how these appear in
N=4 SYM amplitudes.

Gr(4,6) Cluster Algebra for 6-point SYM



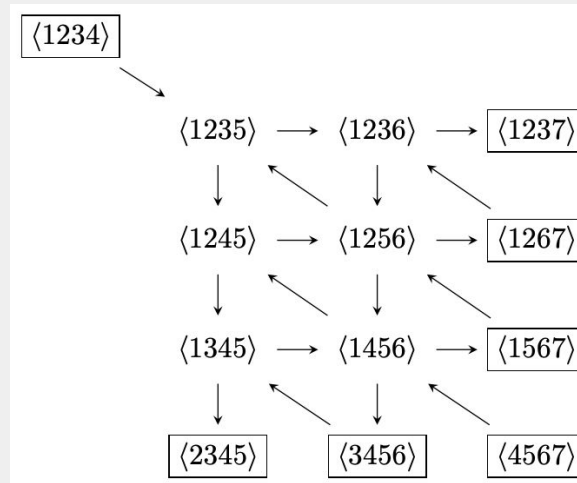
Golden, Goncharov
Spradlin, Vergu, AV
[1305.1617](#)

Start with the initial seed
Perform all mutations
Obtain 15 cluster variables

$$\{\langle a \ a+1 \ b \ c \rangle\}$$

Match 2-loop 6-point symbol alphabet!

Gr(4,7) Cluster Algebra for 7-point SYM



Golden, Goncharov
Spradlin, Vergu, AV
[1305.1617](#)

Start with the initial seed

Perform all mutations

Obtain 35+7+7 cluster variables

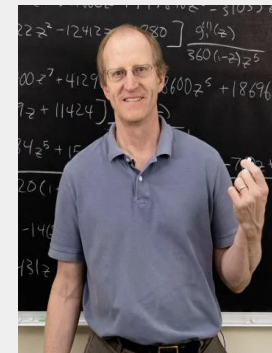
$\{\langle \mathbf{a} \mathbf{a} + 1 \mathbf{b} \mathbf{c} \rangle, \langle 1(23)(45)(67) \rangle, \langle 1(27)(34)(56) \rangle\}$

Match 2-loop 7-point symbol alphabet!

Cluster Bootstrap for N=4 SYM

Dixon et al have computed 6- and 7-point amplitudes to higher number of loops (**8 and 5**, respectively) by starting with the symbol alphabet presented on the last two slides and using a bootstrap approach, based on mathematical and physical inputs.

Constraint	$L = 1$	$L = 2$	$L = 3$	$L = 4$	$L = 5$	$L = 6$
1. $\mathcal{H}_6$	6	27	105	372	1214	3692?
2. Symmetry	(2,4)	(7,16)	(22,56)	(66,190)	(197,602)	(567,1795?)
3. Final-entry	(1,1)	(4,3)	(11,6)	(30,16)	(85,39)	(236,102)
4. Collinear	(0,0)	(0,0)	(0*,0*)	(0*,2*)	(1* ³ ,5* ³)	(6* ² ,17* ²)
5. LL MRK	(0,0)	(0,0)	(0,0)	(0,0)	(0*,0*)	(1* ² ,2* ²)
6. NLL MRK	(0,0)	(0,0)	(0,0)	(0,0)	(0*,0*)	(1*,0* ²)
7. NNLL MRK	(0,0)	(0,0)	(0,0)	(0,0)	(0,0)	(1,0*)
8. N ³ LL MRK	(0,0)	(0,0)	(0,0)	(0,0)	(0,0)	(1,0)
9. Full MRK	(0,0)	(0,0)	(0,0)	(0,0)	(0,0)	(1,0)
10. T^1 OPE	(0,0)	(0,0)	(0,0)	(0,0)	(0,0)	(1,0)
11. T^2 OPE	(0,0)	(0,0)	(0,0)	(0,0)	(0,0)	(0,0)



**Caron-Huot, Dixon, Drummond, Dulat, Foster, Gurdogan, von Hippel, McLeod, Papathanasiou, review: [2005.06735](#)
Dixon, Liu [2308.08199](#)
He, Jiang, Li, Liu [2511.09669](#)**

New Features at $n > 7$

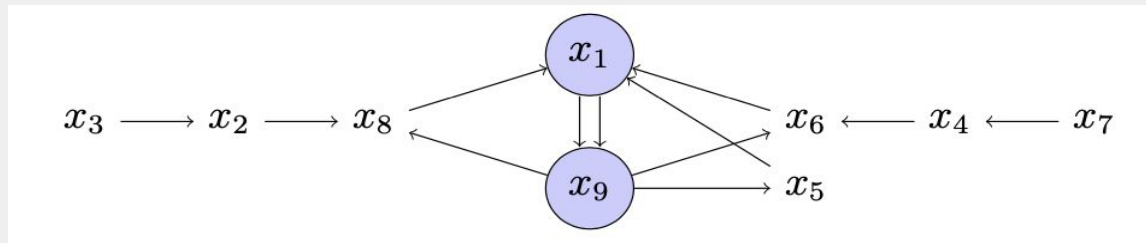
- $\text{Gr}(4, n)$ cluster algebra is infinite for $n > 7$
- Some symbol letters are not cluster coordinates, but involve square roots when written in terms of Pluckers

Li, Zhang [2110.00350](#)

He, Li, Zhang [1911.01290](#)
[2009.11471](#)

Algebraic Letters from Infinite Paths

All algebraic symbol letters known to appear in 8 and 9-point amplitudes are known to be associated to infinite mutation sequences inside $\text{Gr}(4,8)$ and $\text{Gr}(4,9)$.



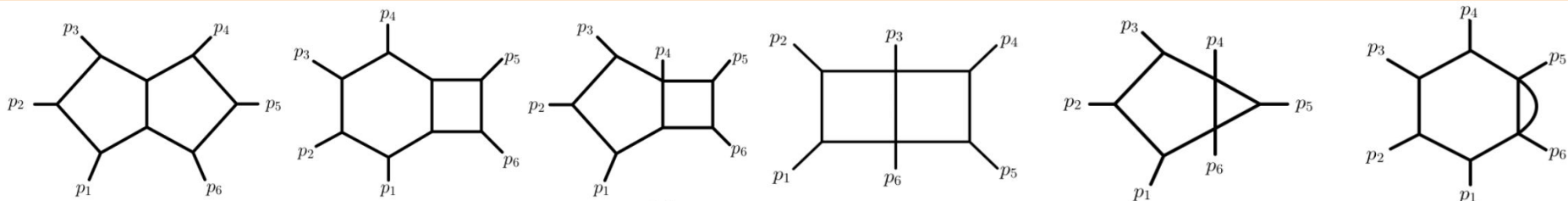
Drummond, Foster, Gürdoğan, Kalousios [1912.08217](#)

Applications:

1. Feynman integrals

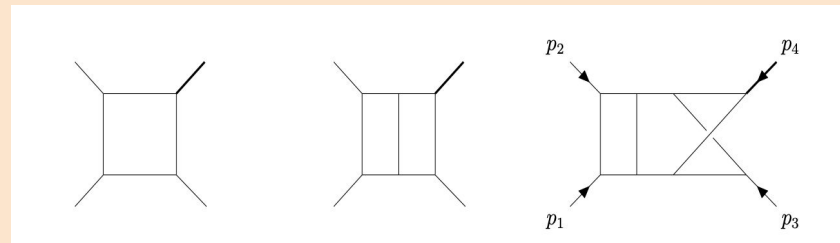
Feynman integrals: status

- **$n=4$ $L=3$ one mass** Henn, Lim, Torres Bobadilla [2302.12776](#); Gehrmann, Henn, Jakubcik, Lim, Mella, Syrrakos, Torres Bobadilla [2410.19088](#)
- **$n=5$ $L=2$ massless** Gehrmann, Henn, Lo Presti [1511.05409](#); Chicherin, Henn, Mitev [1712.09610](#); Chicherin, Gehrmann, Henn, Wasser, Zhang, Zoia [1812.11160](#); Abreu, Dixon, Herrmann, Page, Zeng [1812.08941](#), [1902.08563](#)
- **$n=5$ $L=3$ massless** Liu, Matijasic, Miczajka, Xu, Xu, Zhang [2411.18697](#)
- **$n=5$ $L=2$ one mass** Abreu, Chicherin, Ita, Page, Sotnikov, Tschernow, Zoia [2306.15431](#)
- **$n=5$ $L=2$ two masses** Abreu, Chicherin, Sotnikov, Zoia [2406.05201](#)
- **$n=6$ $L=2$ massless** Henn, Matijasic, Miczajka, Peraro, Xu, Zhang [2403.19742](#) [2501.01847](#); Abreu, Monni, Page, Usovitsch [2412.19884](#)



Cluster algebras for 4-point integrals

Chicherin, Henn, Papathanasiou [2012.12285](#),
 Aliaj, Papathanasiou [2408.14544](#) found cluster algebras
 structure in certain Feynman integrals relevant to QCD.



A_2, C_2, G_2
 for
 $L=1,2,3$

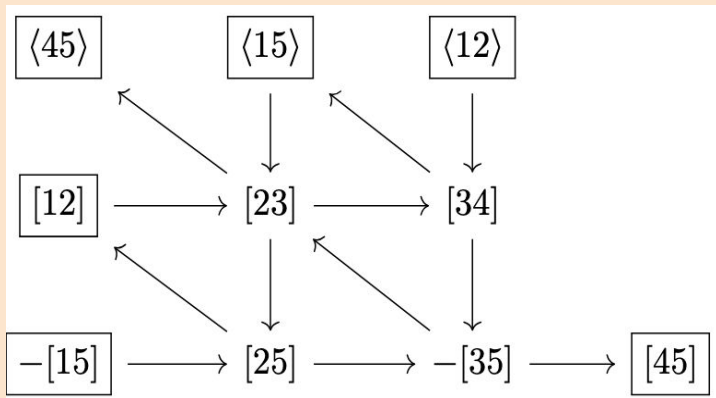
$$\{a_1, a_2\}$$

$$a_{m+1} = \begin{cases} \frac{1+a_m}{a_{m-1}} & m \text{ odd} \\ \frac{1+a_m^L}{a_{m-1}} & m \text{ even} \end{cases}$$

5,6,8 cluster variables match
 symbol alphabet after the
 change of variables

$$a_1 = \frac{s - p_4^2}{t} \quad a_2 = \frac{s}{t} \frac{p_4^2 - s - t}{p_4^2}$$

F(2,4;5) Flag Cluster Algebra for 5-point integrals



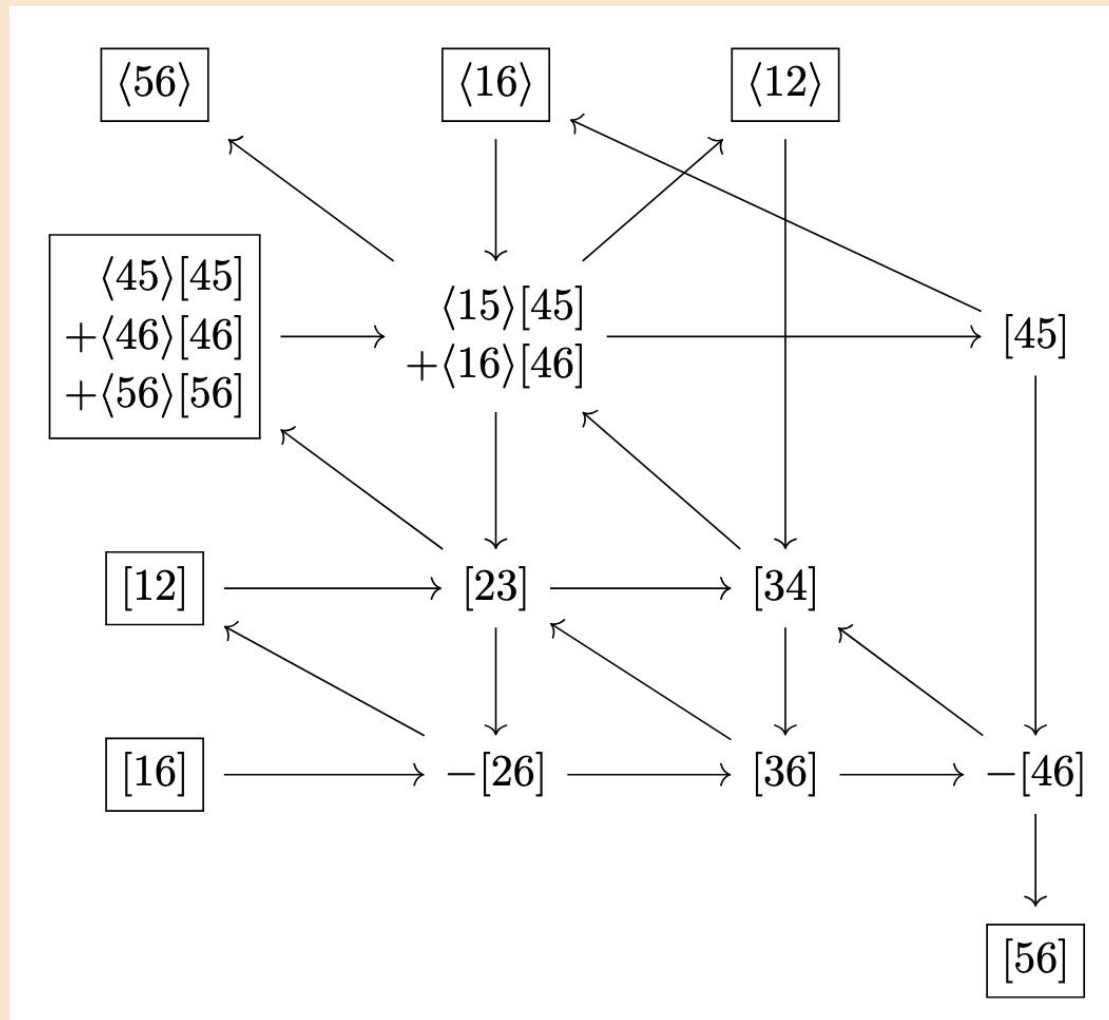
22 cluster variables

$$\{\langle ij \rangle, [ij], \langle 23 \rangle [23] - \langle 45 \rangle [45], \langle 12 \rangle [12] - \langle 34 \rangle [34]\}$$

If we take the union under cyclic permutation (or under all permutations) it produces exactly 25 planar (and 5 non-planar) letters (except for one letter $\sqrt{\Delta}$ that drops out of all known amplitudes in four dimensions).

Pokraka, Spradlin, AV, Weng 2506.11895

F(2,4;6) Flag Cluster Algebra for 6-point integrals



F(2,4;6) Cluster Algebra & 2-loop 6-point Symbol Alphabet

245 symbol letters through all orders in dim reg

- 7 are analogs of $\sqrt{\Delta}$ in the 5-point case
- 135 are F(2,4;6) cluster variables (including permutations)
- 40 are algebraic and we reproduce all of them from infinite mutation sequences
- 27 of them are rational, are NOT cluster variables, but only appear at $O(\epsilon)$ so are not relevant in four dimensions
- **36 remain somewhat mysterious...**

(24 of 36 appear in Carrolo, Chicherin, Henn, Yang, Zhang [2505.01245](#))

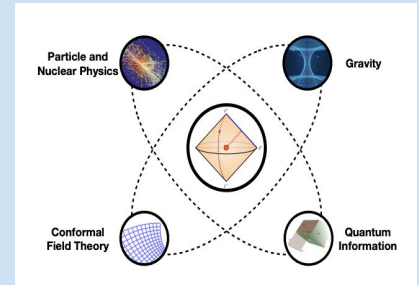
Pokraka, Spradlin, AV, Weng [2506.11895](#)

Applications:

2. Energy correlators integrals

Energy Correlators

- Energy Correlators are of central importance in both collider phenomenology and in the formal high energy theory.
- They are defined as energy-weighted cross sections and measure energy deposited in detectors as a function of the angular separation between detector pairs.
- The study of energy correlators has a very long history... see Moul-Zhu 114 pages review [2506.09119](#)



Energy Correlators: A Journey From Theory to Experiment

Ian Mould^{1,*} and Hua Xing Zhu^{2,3,†}

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Peking University, Beijing 100871,
China*

2506.09119

Collider experiments offer a unique opportunity to explore the Standard Model (SM), and to search for new physics, new interactions, and new principles of nature. The theoretical abstraction of a collider, namely the study of correlations in asymptotic fluxes, provides one of the most basic examples of an observable in quantum field theory (QFT) and quantum gravity. Energy flux is described in QFT by energy flow operators, a particular example of light-ray operators. In addition to their central role in the theoretical description of collider physics, energy flow operators play an important role in diverse areas of formal QFT and gravity, providing a connection between real world collider phenomenology, and the deep underlying principles of QFT.

Recently it has become possible to measure correlation functions of energy flow operators in a wide variety of collider experiments, providing an exciting new connection between collider physics and formal theory. In this review, we provide a survey of recent progress in our understanding of energy operators and their correlators, highlighting their importance in both formal theory and collider phenomenology, and in particular, their great potential for bridging these areas to provide new ways to understand the real world. We intend this article as a resource for both formal theorists interested in understanding how light-ray operators are being applied in particle and nuclear physics, as well as for experimentalists interested in the theoretical motivation for these observables. Most importantly, we aim to stimulate further interaction between the formal, phenomenological and experimental communities through the common lens of energy correlators.

[hep-ph] 10 Jun 2025

Energy Correlators: Four-Point

The Collinear Limit of the Four-Point Energy Correlator in $\mathcal{N} = 4$ Super Yang-Mills Theory

Dmitry Chicherin,^a Ian Mould,^b Emery Sokatchev,^a Kai Yan,^{c,d} Yunyue Zhu^c

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^b *Department of Physics, Yale University, New Haven, CT 06511, USA*

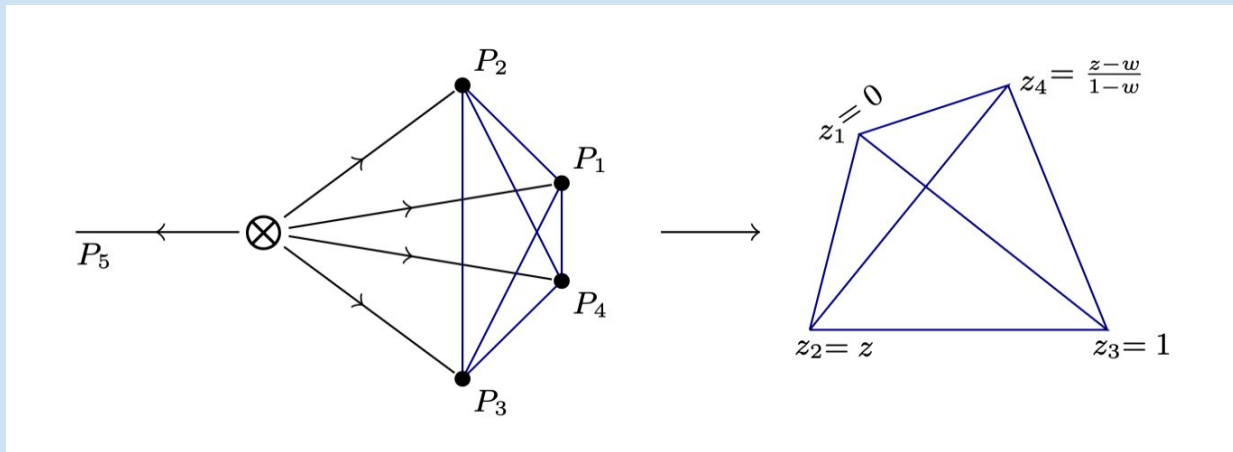
^c *School of Physics and Astronomy, Shanghai Jiao Tong University, Shanghai 200240, China*

^d *Key Laboratory for Particle Astrophysics and Cosmology (MOE), Shanghai 200240, China*

We present a compact formula, expressed in terms of classical polylogarithms up to weight three, for the leading order four-point energy correlator in maximally supersymmetric Yang-Mills theory, in the limit where the four detectors are collinear. This formula is derived by combining a simplified, manifestly dual conformal invariant form of the $1 \rightarrow 4$ splitting function obtained from the square of the tree-level five-particle form factor of stress-tensor multiplet operators, with a novel integration-by-parts algorithm operating directly on Feynman parameter integrals. Our results provide valuable data for exploring the structure of physical observables in perturbation theory, and for calculations of jet substructure observables in quantum chromodynamics.

2401.06463

Energy Correlators: Four-Point



$$E^4 C_{\mathcal{N}=4}^{\text{LO, coll}}(z_{ij}) = \frac{\sum_{I=1}^{51} R_I(z, \omega) F_I(z, \omega)}{|z_{12}|^2 |z_{23}|^2 |z_{34}|^2} + \text{perm}(1, 2, 3, 4)$$

```
{F[1] -> (2*froot[a])/3 + (2*froot[b])/3 + (2*froot[c])/3 +
  (2*Pi^2*G[0, v/(-1 + v)])/3 - 2*G[0, 1/(u*ub)]*G[0, u/(u - v)]*
  G[0, v/(-1 + v)] - G[0, u/(u - v)]^2*G[0, v/(-1 + v)] +
  (G[0, 1/(u*ub)]*G[0, v/(-1 + v)]^2)/2 + (2*G[0, v/(-1 + v)]^3)/3 +
  G[0, 1/(u*ub)]*G[0, v/(-1 + v)]*G[0, 1/(v*vb)] +
  2*G[0, u/(u - v)]*G[0, v/(-1 + v)]*G[0, 1/(v*vb)] +
  G[0, v/(-1 + v)]^2*G[0, 1/(v*vb)] + (2*Pi^2*G[0, vb/(-1 + vb)])/3 +
  G[0, 1/(u*ub)]*G[0, 1/(v*vb)]*G[0, vb/(-1 + vb)] +
  (G[0, 1/(u*ub)]*G[0, vb/(-1 + vb)]^2)/2 +
```

```
{R[1] -> -2 + u + ub + v + vb + (-1 + v)^2*(-1 + vb)*q[1] - v^2*vb*q[2] +
  (-1 + u)^2*(-1 + ub)*q[3] - u^2*ub*q[4] + (-1 + v)*(-1 + vb)^2*q[7] -
  v*vb^2*q[8] + (-1 + u)*(-1 + ub)^2*q[13] - u*ub^2*q[14],
```

```
{q[1] -> (-1 + u)^(-1), q[2] -> u^(-1), q[3] -> (-1 + v)^(-1),
  q[4] -> v^(-1), q[5] -> (u - ub)^(-1), q[6] -> (-ub + v)^(-1),
  q[7] -> (-1 + ub)^(-1), q[8] -> ub^(-1), q[9] -> (-1 + u*ub)^(-1),
```

Energy Correlators: Four-Point

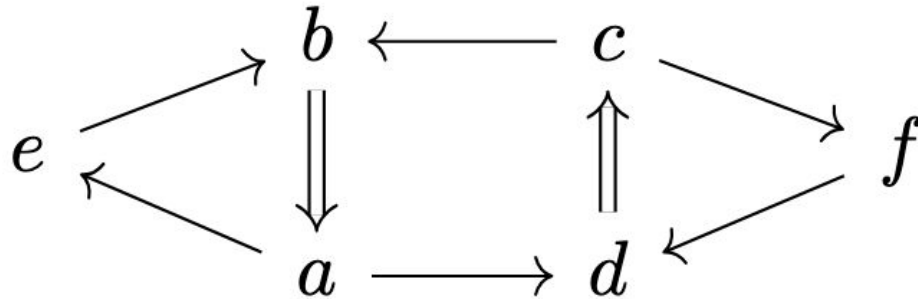
MPLs up to weight 3 with Symbol Alphabet that involves **CUBIC** roots

$$z, \bar{z}, 1 - z, 1 - \bar{z}, 1 - |z|^2, z - \bar{z}, w\bar{z} - \bar{w}z, \\ w\bar{z} - \bar{w}, z\bar{w} - w, 1 - w - \bar{w} + z\bar{w}, 1 - w - \bar{w} + w\bar{z}$$

$$|z|^2 - |w|^2, z - |w|^2, \bar{z} - |w|^2, \frac{a}{b}, \frac{a + |w|^2}{b + |w|^2}, \frac{a + w\bar{b} + \bar{w}}{b + w\bar{a} + \bar{w}}$$

$$abc = -|z|^2|w|^2, \quad a + b + c = 1 - w - \bar{w} - z - \bar{z}, \\ a^{-1} + b^{-1} + c^{-1} = 1 - z^{-1} - \bar{z}^{-1} - w^{-1} - \bar{w}^{-1}.$$

Cube Roots and Cluster Algebra



Cluster Algebras, Cube Roots, and Energy Correlators in $\mathcal{N} = 4$ SYM

Dani Kaufman,¹ Marcus Spradlin,^{2,3} Anastasia Volovich,² and Kai Yan^{4,5}

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²*Department of Physics, Brown University, Providence, RI 02912, USA*

³*Brown Center for Theoretical Physics and Innovation, Brown University, Providence, RI 02912, USA*

⁴*School of Physics and Astronomy, Shanghai Jiao Tong University, Shanghai 200240, China*

⁵*Key Laboratory for Particle Astrophysics and Cosmology (MOE), Shanghai 200240, China*

We provide a cluster algebraic construction of the cube root symbol letters that appear at leading order in the near-collinear expansion of the four-point energy correlator in $\mathcal{N} = 4$ SYM theory.

2608.03717

Applications:

3. Cosmological correlators integrals

Mark's talk

Conclusions

Cluster Algebras appear—for reasons that are still mysterious—in a growing number of amplitude related contexts:

**Feynman integrals,
Cosmological correlator integrals,
Energy correlator integrals.**

They enable us to make progress in previously untrackable computations.

They point towards richer structure yet to be understood and explored.

Thank you!