

A new solvable holographic dual for sine-Liouville gravity

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based on

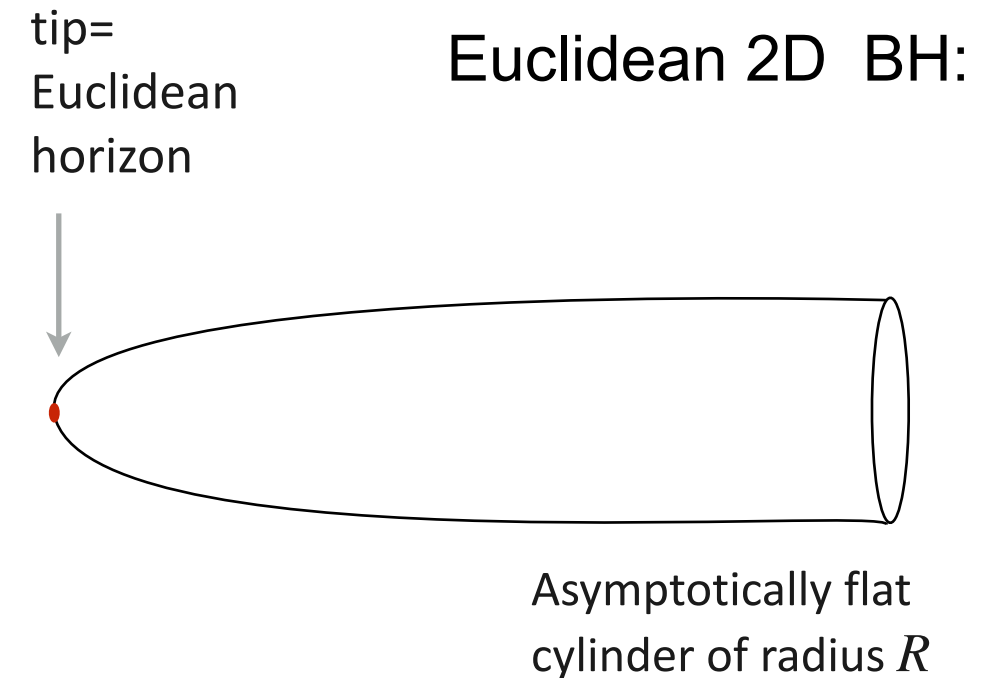
arXiv:2512.18916 — JHEP 06 (2026) 080

arXiv:2606.22651 (with Andre Alves, Mateus Da Silva Junca)

Corfu 2026: Quantum Gravity & Strings Programme

Intro

- Sine-Liouville gravity = sine-Gordon matter coupled to 2D gravity: studied for 30+ years (Moore '92; Hsu–Kutasov '93)
- Related (for $R = 3/2$ or $R_{\text{dual}} = 2/3$) by the FZZ duality to the Euclidean black hole (cigar) CFT, a rare solvable window on black-hole backgrounds in string theory.
- Holographic dual for Sine-Liouville gravity: Matrix Quantum Mechanics (MQM) perturbed by winding or momentum condensates (Toda integrability). Claim [KKK, 2001]: Euclidean black hole emerges at the extremity of the flow generated by the perturbation.
- **This work: an alternative holographic dual with a neat statistical interpretation — the 7-vertex model on dynamical planar graphs. Main observation: the compactification radius renormalises along the flow. E.g. $R = 2/3 \rightarrow R_{\text{end-of-flow}} = 1/3$. ==> The KKK proposal must be reconsidered.**



Plan:

1. sine-Gordon & sine-Liouville phases
2. the 7-vertex model & its matrix dual
3. spectral curve, EOS, phase diagram
4. disk amplitudes: new kind of branes
5. comparison with MQM / 2D strings

The 7-vertex model on flat/random lattice = solvable discretisation of the Sine-Gordon model without/with coupling to gravity.

The 7-vertex model unveils a connection we have missed in 2001:

The flow from Liouville gravity to Sine-Liouville theory is the gravitational analogue of the **massless RG flow in the Sine-Gordon theory for purely imaginary mass coupling constant**

[Fendley–Saleur–Al. Zamolodchikov '93].

This analogy leads to new understanding with important consequences

RG flows in Sine-Gordon QFT

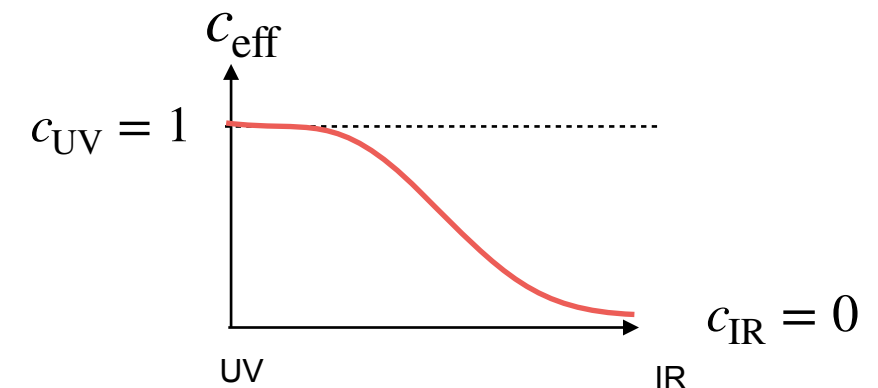
SG model= free compactified boson ($\varphi = \varphi + \pi$) perturbed by a pair of exponentials:

$$\mathcal{A}_{SG}[\varphi] = \int \frac{d^2x}{4\pi} \left[\left(\partial_\mu \varphi \right)^2 + \sqrt{t} \cos \left(\frac{2}{R} \varphi \right) \right]$$

$$\begin{aligned} 1 &\leq R \leq \sqrt{2} \\ R_{s.d} &= 1 \\ \beta_{SG} &= \sqrt{8\pi}/R \end{aligned}$$

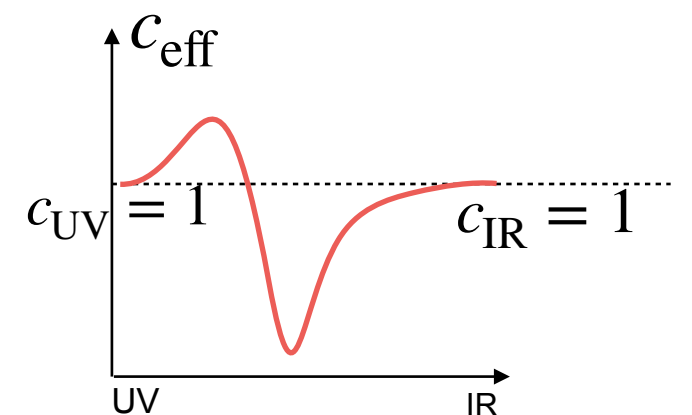
■ Coulomb gas of positive and negative charges,
perturbation is relevant: dimension $x_{SG} = 2 \times 2/R^2 - 2 < 2$.

■ $t > 0$: **massive** RG flow to a trivial plasma phase with
finite correlation length (zero central charge).



■ $t < 0$: **massless** flow between two $c = 1$ fixed points
[Nienhuis'84] [Fendley–Saleur–Al. Zamolodchikov '93]:
free bosons at two different radii,

$$R = R_{UV}^{SG} \rightarrow R_{IR}^{SG}; \quad (R_{IR}^{SG})^2 = 2 - (R_{UV}^{SG})^2$$



— Non-unitary flow: the c-theorem does not apply

Sine-Liouville gravity= gravitational SG model

Couple the sine-Gordon matter to 2D gravity; in conformal gauge (Liouville field ϕ):

$$\mathcal{A}_{\text{SL}}[\phi, \varphi] = \frac{1}{4\pi} \int_{\mathcal{M}} d^2z \sqrt{\hat{g}} \left[(\hat{\nabla} \phi)^2 + (\hat{\nabla} \varphi)^2 + 2\phi \hat{R} \right] + \text{ghosts} \\ + \mu \int d^2z \sqrt{\hat{g}} e^{2\phi} + \sqrt{t} \int d^2z \sqrt{\hat{g}} e^{(2-\frac{1}{R})\phi} \cos(\varphi/R)$$

t = hermal coupling; μ = cosmological constant; scaling: $t \sim \mu^{2-1/R}$.

Gravitational massless flow: connects two Liouville + compactified-boson theories.

$$t = 0: R_{\text{UV}}^{\text{SL}} = R; \quad t \rightarrow \infty: R_{\text{IR}}^{\text{SL}} < R_{\text{UV}}^{\text{SL}}$$

Our main result:

Exact renormalisation of the compactification radius along the flow:

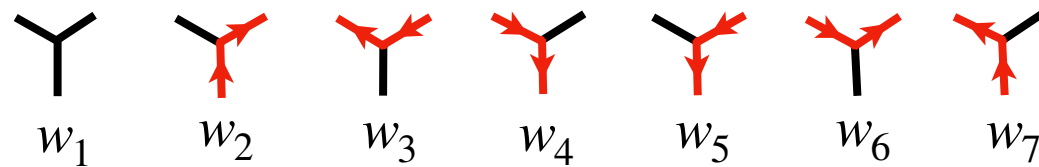
$$R_{\text{IR}}^{\text{SL}} = 1 - R_{\text{UV}}^{\text{SL}}$$

(Compare with $(R_{\text{IR}}^{\text{SG}})^2 = 2 - (R_{\text{UV}}^{\text{SG}})^2$ for the SG)

(the relation will be derived from the matrix model)

KPZ-like scaling: $2R^{\text{SL}} = (2R^{\text{SG}})^2$

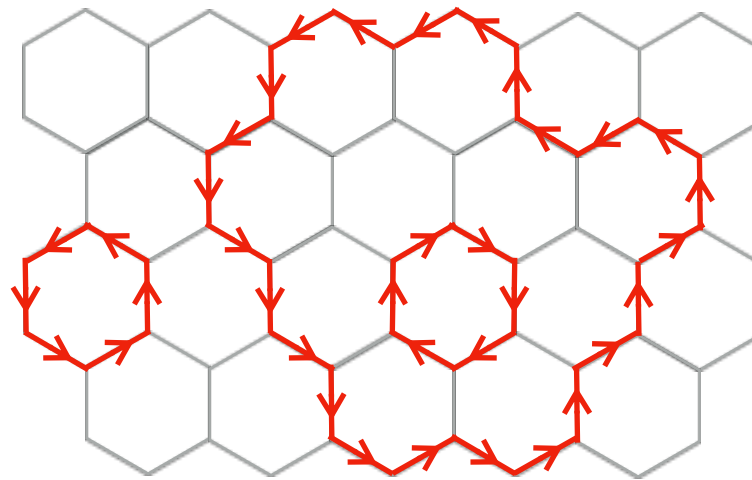
The 7-vertex model on a honeycomb lattice



B. Nienhuis'84, R. Baxter'87,
M. Batchelor & H. Blöte'89,
H. Blöte'89 & B. Nienhuis,
N. Reshetikhin'91 (for $T=0$)

$$w_1 = T, \quad w_2 = w_3 = w_4 = w, \quad w_5 = w_6 = w_7 = w^{-1}$$

The partition function of the vertex model is a sum over all configurations of self and mutually avoidant oriented polymers (loops) on the lattice

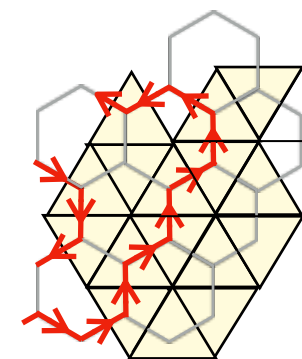


$\#[\text{left turns}] - \#[\text{right turns}] = \pm 6. \Rightarrow \text{weight} [\text{loop}] = w^6 + w^{-6} = n.$ With $w = e^{i\pi/6p}$, $n = 2 \cos(\pi/p)$.

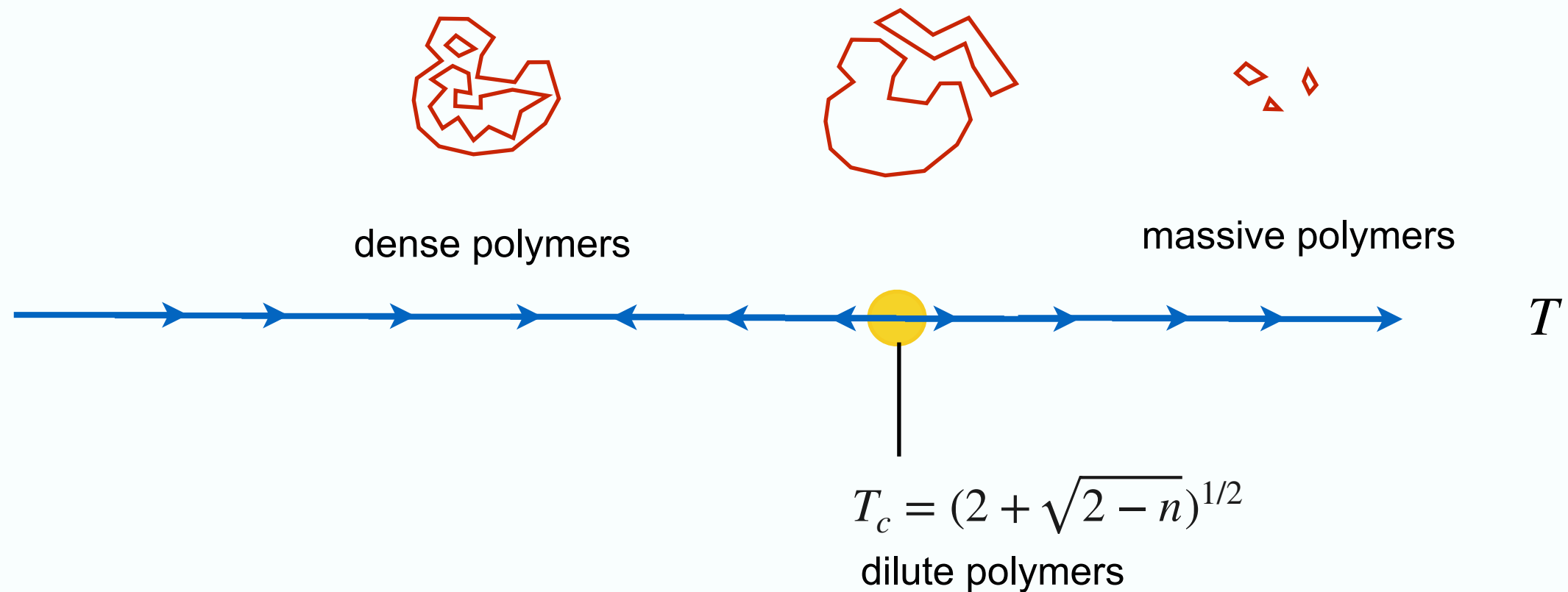
Partition function: $\mathcal{Z} = \sum T^{V_{\text{vac}}} n^{\#\text{loops}}$ (V_{vac} = volume non-occupied by loops)

Dual view:

SOS height model on the dual triangular lattice; loops = domain walls, heights jump by ± 1 .



Phase diagram of the 7-vertex model

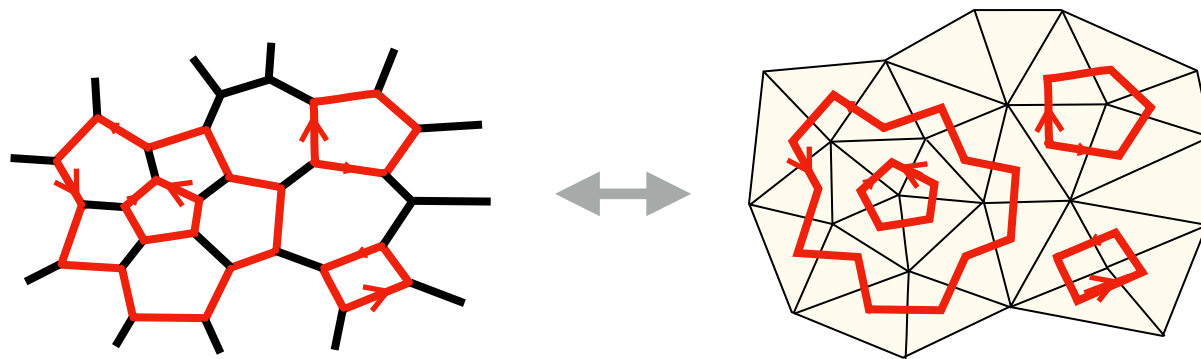


- The 7-vertex model = solvable discretisation of the Sine-Gordon model [B. Nienhuis'84]

$$t \sim T - T_c$$

7v model on dynamical planar graphs

The 7v model can be defined on an arbitrary trivalent planar graph $\mathcal{G}$.



Dual view:
SOS height model on random
triangulations; loops = domain walls,
heights jump by ± 1 .

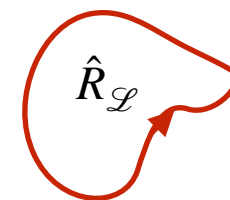
- On a lattice curvature defects $\hat{R}_i = \frac{2\pi}{3}(6 - C_i)$ the loop weights are NOT topological — the phase of a loop is a holonomy of a lattice spin connection

- Discrete geometrical phase = geodesic curvature: with $w = e^{i\pi\lambda/2}$,

- Phase = $\frac{1}{2}\pi\lambda$ [(left turns) – (right turns)] = $\frac{3}{2}\lambda \hat{K}_{\mathcal{L}}$;

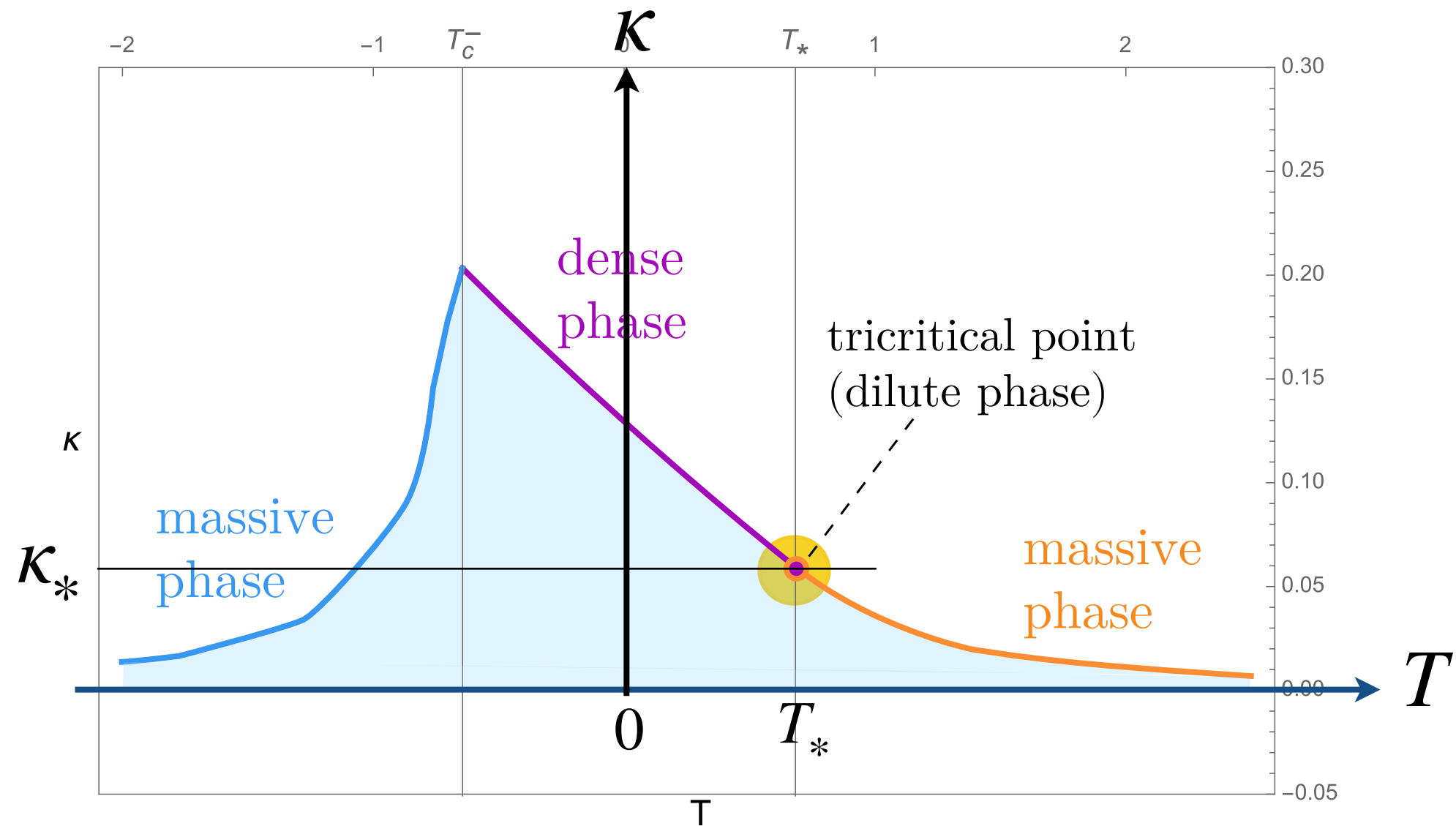
- Gauss–Bonnet ties it to the curvature inside the loop:

$$\hat{R}_{\mathcal{L}} + 2\hat{K}_{\mathcal{L}} = 4\pi, \quad \hat{R}_{\mathcal{L}} \equiv \sum_{i \in \mathcal{I}_{\mathcal{L}}} \hat{R}_i$$



- Partition function:

$$\mathcal{F}_0(\kappa, T) = \sum_{\text{graphs } \mathcal{G}} \kappa^{V_{\mathcal{G}}} Z_{\mathcal{G}}$$



● The scaling limit = Sine-Liouville gravity

Stands for the scale
parameter along the RG flow

$$\mu \sim K_* - K$$

$$t \sim T - T_*$$

Solution of the in the scaling limit

Grand canonical ensemble: $\mu \rightarrow \tilde{\mu}$; Legendre $\tilde{F}_0(\tilde{\mu}) = -\tilde{\mu}\mu + F_0(\mu)$

$$\tilde{\mu} = \partial_{\mu} F_0, \quad \partial_{\tilde{\mu}} \tilde{F}_0 = -\mu$$

Equation of state:
$$\tilde{\mu} = \frac{1+\lambda}{2} e^{\frac{1+\lambda}{2}\tilde{u}} - \frac{1-\lambda}{2} t e^{\frac{1-\lambda}{2}\tilde{u}} \quad \tilde{u} \equiv \frac{\partial^2 \tilde{F}_0}{\partial \tilde{\mu}^2}$$

Same equation of state as MQM with a sine-Liouville condensate (Kazakov–IK–Kutasov '01)

$$t \rightarrow -\infty: \quad \tilde{F}_0(\tilde{\mu}) \equiv \frac{2}{1-\lambda} \frac{(\tilde{\mu}/t)^2}{2} \log |\tilde{\mu}/t| \quad \Rightarrow \quad R_{\text{IR}}^{\text{SL}} \equiv R \Big|_{t \rightarrow \infty} = \frac{1-\lambda}{2}$$

— Follows also from the six-vertex model result for $T = 0$ [IK, 1999]

$$t \rightarrow 0: \quad \tilde{F}_0(\tilde{\mu}) = \frac{2}{1+\lambda} \frac{\tilde{\mu}^2}{2} \log \tilde{\mu} \quad \Rightarrow \quad R_{\text{UV}}^{\text{SL}} \equiv R \Big|_{t=0} = \frac{1+\lambda}{2}$$

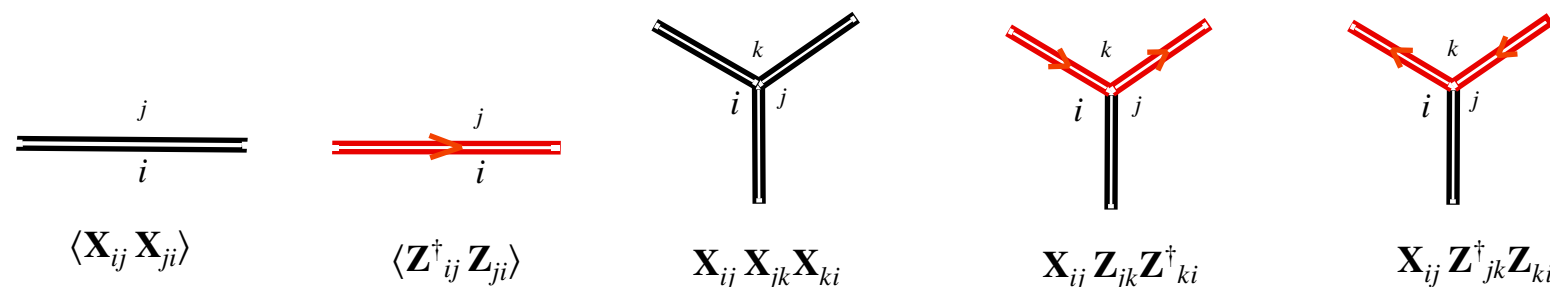
The 7-vertex matrix model

Holographic dual: a one-parameter deformation of Ginsparg's 6-vertex matrix model,
 X Hermitian, Z complex, $N \times N$:

$$\mathcal{Z}_N = \int d\mathbf{X} d\mathbf{Z} d\mathbf{Z}^\dagger \exp \left(-\frac{1}{\hbar} \mathcal{S} \right) \quad \mathcal{S} = \text{Tr} \left(\frac{1}{2} \mathbf{X}^2 + \mathbf{Z} \mathbf{Z}^\dagger - \frac{T}{3} \mathbf{X}^3 - w \mathbf{X} \mathbf{Z} \mathbf{Z}^\dagger - \bar{w} \mathbf{X} \mathbf{Z}^\dagger \mathbf{Z} \right)$$

$$\hbar = \kappa/N$$

Planar Feynman graphs of this action
 = trivalent graphs with the seven
 vertices of fig. 1; $\hbar \sim 1/N$ organises
 the genus expansion.



Integrating out Z : eigenvalue integral with a dipole-like pairwise interaction,

$$\mathcal{Z}_N \sim \int \prod_{i=1}^N dx_i e^{-\frac{1}{\hbar} \left(\frac{1}{2} x_i^2 - \frac{T}{3} x_i^3 \right)} \frac{\prod_{k < j} (x_k - x_j)^2}{\prod_{k \neq j} (1 - w x_k - w^{-1} x_j)}$$

$T = 0$ reproduces the 6v matrix model (Kostov '99; P. Zinn-Justin '99).

Grand ensemble, collective field, Virasoro constraint

Grand canonical ensemble (chemical potential $\tilde{\kappa}$ = lattice cosmological constant)

— a Fredholm determinant of free fermions:

$$\mathcal{Z}(\bar{\mu}) = \sum_{N=1}^{\infty} e^{-N\tilde{\kappa}/\hbar} \mathcal{Z}_N = \text{Det}(1 - K) \quad K(x, x') = \frac{\exp(-\frac{\tilde{\kappa} + V'}{\hbar})}{wx + w^{-1}x'}$$

Dyson–Schwinger equations (Virasoro constraint) on the collective field

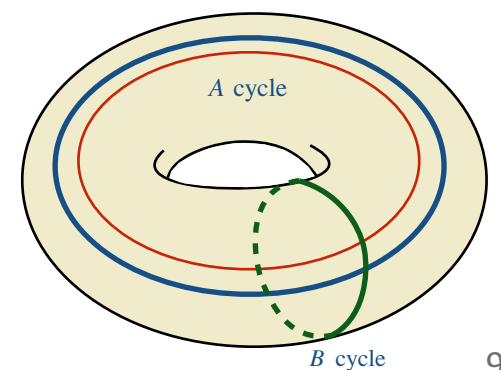
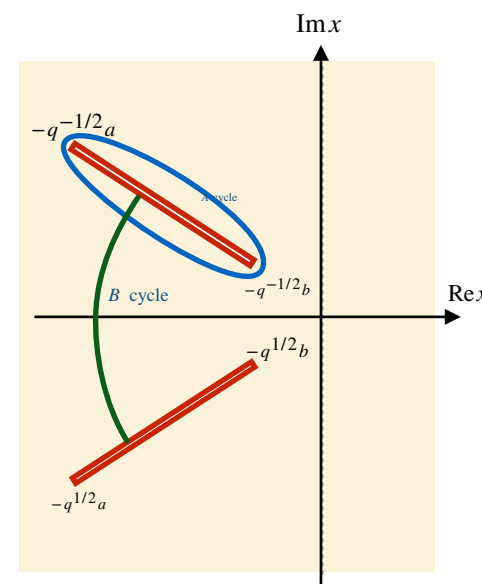
$$\Omega(x) = \log \left\langle \prod_{i=1}^N (x - x_i) \right\rangle$$

Classical Virasoro constraint = sewing condition across the eigenvalue cut:

$$\tilde{\kappa} = V(x) - \Omega(x + i0) - \Omega(x - i0) + \Omega(-iwx) + \Omega(iw^{-1}x) \quad \text{for } x \in [-a, -b]$$

Can be formulated as a quasi-periodicity conditions on the B-cycle of the torus for

$$\Phi(x) = \frac{\Omega(-iwx) - \Omega(iw^{-1}x)}{i} + \text{Pol}$$



Classical spectral curve in the scaling limit

General solution of the quasi-periodicity condition in terms of θ -functions

[M. da Silva Junca, a IK, A. Alves Lima, arXiv: arXiv:2606.22651]

In the scaling limit when the torus degenerates into a cylinder: $\tilde{\mu} \sim \tilde{\kappa}_* - \tilde{\kappa}$, $t \sim T - T_*$

$$x(s) = -2Me^{-\lambda s} \sinh(s)$$

$$y(s) \equiv \partial_x \Phi = e^{(1+\lambda)s} M^{\frac{1+\lambda}{1-\lambda}} - t e^{-(1-\lambda)s} M^{\frac{1-\lambda}{1+\lambda}}$$

$$\tilde{u} \equiv \partial_{\tilde{\mu}}^2 \tilde{F} = \frac{4}{1-\lambda^2} \log(M\epsilon)$$

$$\mu = \frac{1+\lambda}{2} M^{\frac{2}{1-\lambda}} - \frac{1-\lambda}{2} t M^{\frac{2}{1+\lambda}}$$

Same equation of state as MQM with a sine-Liouville condensate (Kazakov–IK–Kutasov '01)

Summary and outlook

- A microscopic, integrable realisation of sine-Liouville gravity: the 7-vertex model on dynamical planar graphs — a new class of models where loop weights couple to curvature.
- Gravitational analogue of the sine-Gordon massless flow, with $R_{IR} = 1 - R_{UV}$
- Disk amplitudes are given by a deformation (Krätzel) of Bessel-K functions- a new type of branes in 2D string theory.

Thank you!