

Resolving G_2 Manifolds in Scale separated models in Type IIA

Based on work with Fotis Farakos and Fien Apers. (To appear soon)

Aviral Aggarwal
aviral@mail.muni.cz

Faculty of Science, Masaryk University

September 11, 2026

Motivation

String theory is consistent in 10 D.

Motivation

String theory is consistent in 10 D.
But our world is 4d!

Motivation

String theory is consistent in 10 D.

But our world is 4d!

String theory has to give a good low energy EFT. (That describes our world)

What do we demand from our string models?

- Trustable effective theory in 4d
- Minimal supersymmetry $\mathcal{N} = 1$
- The standard model of particle physics.

The Landscape

As a first step we look for consistent scale separated solutions.

We demand the following features -

- $m_{kk}L_{AdS} \gg 1$ or $L_{kk} \ll L_{ext}$
- N=1 SUSY.
- Parametric Scale separation.

The Landscape

As a first step we look for consistent scale separated solutions.

We demand the following features -

- $m_{kk} L_{AdS} \gg 1$ or $L_{kk} \ll L_{ext}$
- N=1 SUSY.
- Parametric Scale separation.

We have considerable amount of models in literature- DGKT, KKLT, LVS

The Landscape

As a first step we look for consistent scale separated solutions.

We demand the following features -

- $m_{kk} \gg m_{AdS}$ or $L_{kk} \ll L_{ext}$
- $N=1$ SUSY.
- Parametric Scale separation.

We have considerable amount of models in literature- **DGKT**, KKLT, LVS

Arguments against them -

- Evidence from Swampland conjectures
- Lack of explicit dual CFTs
- Inconsistency of Uplifting the models/ Smearing approximations
- Construction requires singular manifolds with intersecting O-Planes

Review of String compactification

- We begin by assuming a product structure of the 10d spacetime - $X_d \times M$
- Start with a Ricci flat manifold, allow for the metric to fluctuate
- Fluctuation of metric can be written in terms of fluctuation of internal cycles
- These show up as moduli (scalars) in Low energy EFT.
- We dimensionally reduce the theory to a low dimensional one.
- We stabilize these moduli by turning on fluxes and generating a potential.

Scale separated solution in AdS_3

We have in particular an orbifold of the form

$$T^7/(Z_2)^3, \quad (1)$$

where the $(Z_2)^3$ group is given by the following three involutions -

$$\begin{aligned} \Theta_\alpha : y^j &\rightarrow (-y^1, -y^2, -y^3, -y^4, y^5, y^6, y^7), \\ \Theta_\beta : y^j &\rightarrow (-y^1, 1/2 - y^2, y^3, y^4, -y^5, -y^6, y^7), \\ \Theta_\gamma : y^j &\rightarrow (1/2 - y^1, y^2, -y^3, y^4, -y^5, y^6, -y^7), \end{aligned} \quad (2)$$

This space is Equipped with G_2 structure $(\Phi) \rightarrow$ metric on manifold.

We also have a dual closed 4-form $(\Psi) \rightarrow$ Ricci Flat Manifold.

Scale separated solution in AdS_3 -Overview

Important relation is -

$$\text{vol}(X) = \frac{1}{7} \int \Phi \wedge \star \Phi, \quad (3)$$

Scale separated solution in AdS_3 -Overview

Important relation is -

$$\text{vol}(X) = \frac{1}{7} \int \Phi \wedge \star \Phi, \quad (4)$$

Ricci flat deformation can be written as deformation of Φ . (No. of moduli = b^3)

Scale separated solution in AdS₃-Overview

Important relation is -

$$\text{vol}(X) = \frac{1}{7} \int \Phi \wedge \star \Phi, \quad (5)$$

Ricci flat deformation can be written as deformation of Φ . (No. of moduli = b^3)

- These moduli are given by periods of the invariant three forms on associated cycles
 $s_i = \int_{C^i} \Phi$
- These G_2 manifold preserves $N = 2$ susy.
- We induce further involution by introducing O2 planes and O6 planes. (N=1) susy
- Ingredients- moduli(s_i, ϕ).
- All other moduli dont satisfy the parity wrt orientifolds or simply the cycles dont exist

Scale separated solution in AdS₃-Overview

Important relation is -

$$\text{vol}(X) = \frac{1}{7} \int \Phi \wedge \star \Phi, \quad (6)$$

Ricci flat deformation can be written as deformation of Φ . (No. of moduli = b^3)

- These moduli are given by periods of the invariant three forms on associated cycles
 $s_i = \int_{C^i} \Phi$
- These G_2 manifold preserves $N = 2$ susy.
- We induce further involution by introducing O2 planes and O6 planes. (N=1) susy
- Ingredients- moduli(s_i, ϕ).
- All other moduli don't satisfy the parity wrt orientifolds or simply the cycles don't exist
- Singularities can be resolved and the manifold remains G_2 [Joyce]

AdS3 type IIA fully moduli stabilized models

We also add more ingredients like D branes to cancel the following tadpoles

$$dF_p = H_3 \wedge F_{p-2} + J_{p+1} \quad (7)$$

AdS3 type IIA fully moduli stabilized models

We also add more ingredients like D branes to cancel the following tadpoles

$$dF_p = H_3 \wedge F_{p-2} + J_{p+1} \quad (8)$$

Insert fluxes $F_0(m)$, $F_4(f)$, $H_3(h)$.

- Since we have N=1 Susy, we find the supersymmetric minima which is naturally gives vacua for full theory.
- This superpotential is obtained by dimensional reduction of 10d to 3d.

AdS3 type IIA fully moduli stabilized models

We also add more ingredients like D branes to cancel the following tadpoles

$$dF_p = H_3 \wedge F_{p-2} + J_{p+1} \quad (9)$$

Insert fluxes $F_0(m), F_4(f), H_3(h)$.

Superpotential Extremization gives-

$$e^\phi \simeq \frac{h}{m^{1/4} f^{3/4}}, \quad s_{1,2,3,4,5,6,7} \simeq \frac{h}{m} e^{-7\phi/4}, \quad V_0 \simeq \frac{f^{49/16}}{m^{21/16} h^{7/4}}. \quad (10)$$

We also see from the minima of the potential -

$$\frac{L_{\text{KK}}^2}{L_{\text{AdS}}^2} \sim f^{-1} \quad (11)$$

Resolving the Orbifolds

The general procedure is the following-

- Locally each orbifold singularity has a resolution. (Depending upon structure of the fixed point)
- We can locally cut out the ball of radius a and replace the ball by a smooth euclidean Einstein solution.

Resolving the Orbifolds

The general procedure is the following

- Locally each orbifold singularity has a resolution. (Depending upon structure of the fixed point)
- We can locally cut out the ball of radius a and replace the ball by a smooth euclidean Einstein solution.

Heuristic argument of bosonic analysis of blow up moduli in DGKT-

$$S \subset S_{\text{kinetic}} \simeq \frac{1}{2\kappa_{10}^2} \int d^{10}x \sqrt{-g} e^{-2\phi} (R + e^{2\phi} m_0^2) \quad (12)$$

Reducing them to 4d and converting to Einstein frame -

$$S_{4d} = \int \sqrt{-g^E} (\tilde{R}_4 + \frac{e^{4\phi}}{Vol_{int}} m_0^2) \quad (13)$$

$$Vol_{int} = Vol_0 - Ba^4.$$

We see that extremization gives us $a = 0$..(Singular again)

Resolving the Orbifolds

The general procedure is the following

- Locally each orbifold singularity has a resolution. (Detailed structure of the singularity)
- We can locally cut out the ball of radius a and replace the ball by a smooth euclidean Einstein solution.

Heuristic argument of bosonic analysis of blow up moduli

$$S \subset S_{\text{kinetic}} \simeq \frac{1}{2\kappa_{10}^2} \int d^{10}x \sqrt{-g} e^{-2\phi} (R + e^{2\phi} m_0^2) + \frac{1}{2\kappa_{10}^2} \int d^{10}x \sqrt{-g_{BL}} \sqrt{-g_{Ext}} \left| \tilde{F}_4 \right|^2 \quad (14)$$

Reducing them to 4d and converting to Einstein frame -

$$S_{4d} = \int \sqrt{-g^E} \left(\tilde{R}_4 + \frac{e^{4\phi}}{\text{Vol}_{int}} m_0^2 \right) \quad (15)$$

Resolving the Orbifolds

The general procedure is the following

- Locally each orbifold singularity has a resolution. (Detailed structure of the singularity)
- We can locally cut out the ball of radius a and replace the ball by a smooth euclidean Einstein solution.

$$S \subset S_{\text{kinetic}} \simeq \frac{1}{2\kappa_{10}^2} \int d^{10}x \sqrt{-g} e^{-2\phi} (R + e^{2\phi} m_0^2) + \frac{1}{2\kappa_{10}^2} \int d^{10}x \sqrt{-g_{BL}} \sqrt{-g_{Ext}} \left| \tilde{F}_4 \right|^2 \quad (16)$$

Reducing them to 4d and converting to Einstein frame -

$$S_{4d} \sim \int \sqrt{-g^E} (\tilde{R}_4 + \frac{e^{4\phi}}{Vol_{int}} m_0^2) + \int \sqrt{-g^E} \frac{e^{4\phi}}{Vol_{int}^2} \frac{f^2}{a^4} \quad (17)$$

Non Zero vev for a^*

Resolution of G_2 manifold for AdS_3

- For these Joyce Orbifolds we have a singularity structure as- $\mathbb{T}^4/Z_2 \times T^3$
- Proven that this can be resolved by Eguchi Hansen space (Similar to the well known $K3$ manifolds)

$$ds^2 = \Delta^{-1} dr^2 + \frac{1}{4} r^2 \Delta (d\psi + \cos \theta d\phi)^2 + \frac{1}{4} r^2 d\Omega_2^2 \quad (18)$$

with $\Delta = 1 - (\rho/r)^4$ and $d\Omega_2^2 = d\theta^2 + \sin^2 \theta d\phi^2$.

Resolution of G_2 manifold for AdS_3

- For these Joyce Orbifolds we have a singularity structure as- $\mathbb{T}^4/Z_2 \times T^3$
- Proven that this can be resolved by Eguchi Hansen space (Similar to the well known $K3$ manifolds)

$$ds^2 = \Delta^{-1} dr^2 + \frac{1}{4} r^2 \Delta (d\psi + \cos \theta d\phi)^2 + \frac{1}{4} r^2 d\Omega_2^2 \quad (19)$$

with $\Delta = 1 - (\rho/r)^4$ and $d\Omega_2^2 = d\theta^2 + \sin^2 \theta d\phi^2$.

The above procedure can not be followed in our case since the Flux terms dont contribute to the scaling of potential.

Resolution of G_2 manifold for AdS_3

- For these Joyce Orbifolds we have a singularity structure as- $\mathbb{T}^4/Z_2 \times T^3$
- Proven that this can be resolved by Eguchi Hansen space (Similar to the well known $K3$ manifolds)

$$ds^2 = \Delta^{-1} dr^2 + \frac{1}{4} r^2 \Delta (d\psi + \cos \theta d\phi)^2 + \frac{1}{4} r^2 d\Omega_2^2 \quad (20)$$

with $\Delta = 1 - (\rho/r)^4$ and $d\Omega_2^2 = d\theta^2 + \sin^2 \theta d\phi^2$.

The above procedure can not be followed in our case since the Flux terms dont contribute to the scaling of potential.

So what next...?

Resolution of G_2 manifold for AdS_3

- For these Joyce Orbifolds we have a singularity structure as- $\mathbb{T}^4/Z_2 \times T^3$
- Proven that this can be resolved by Eguchi Hansen space (Similar to the well known $K3$ manifolds)

The above procedure can not be followed in our case since the Flux terms dont contribute to the scaling of potential.

Instead we directly stabilize using the superpotential-

$$P = \frac{1}{8\mathcal{V}^2} \int \left(e^{-\phi/2} H_3 \wedge \star \Phi + e^{\phi/4} F_4 \wedge \Phi + e^{5\phi/4} \star F_0 \right), \quad (21)$$

Volume correction, [LM,2003]

$$\mathcal{V} = V_0 \Omega(u), \quad \Omega(u) = 1 - \frac{8}{3} \sum_{\tau,n} \sum_{a=1}^3 \frac{(u^{(\tau,n,a)})^2}{s^{A(\tau,a)} s^{B(\tau,a)}} + O(u^4). \quad (22)$$

Resolution of G_2 manifold for AdS_3

Once again the fluxes are expanded in the associated form basis-

$$\eta_I = (\Phi_i, \chi_i) \in H^3(X_7), \quad \Psi_I = (\psi^i, \tilde{\chi}^i) \in H^4(X_7),$$

$$\Phi = s^i \phi_i + u_a \chi_a$$

$$H_3 = h_i \phi_i$$

$$F_4 = e_i \psi_i + q_a \tilde{\chi}^a$$

(23)

Resolution of G_2 manifold for AdS_3

Once again the fluxes are expanded in the associated form basis-

$$\begin{aligned}\eta_I &= (\Phi_i, \chi_i) \in H^3(X_7), \quad \Psi_I = (\psi^i, \tilde{\chi}^i) \in H^4(X_7), \\ \Phi &= s^i \phi_i + u_a \chi_a \\ H_3 &= h_i \phi_i \\ F_4 &= e_i \psi_i + q_a \tilde{\chi}^a\end{aligned}\tag{24}$$

A look into the superpotential to leading order

$$P = P_0 + \sum_{a=1}^3 L_a u_a + \frac{1}{2} \sum_{a=1}^3 B_a u_a^2 + \text{higher order terms}\tag{25}$$

Resolution of G_2 manifold for AdS_3

Extremize it for bulk values at minima !

$$\frac{u_a^*}{\sqrt{s_a s_b}} \sim \frac{q_a}{f} \quad (26)$$

With their correction to bulk superpotential suppressed $O((q_a/f)^2)$.

We also check that the collapsing cycle still is parametrically large in l_s .

0 planes on resolved spaces

0 planes on resolved spaces

From covering torus perspective we have a problem-

Singular locus	Covering-space orientifold loci	Local sectors
F_α	$06_1/06_2$	one $06/06$ pair
F_β	$06_1/06_3$	one $06/06$ pair
F_γ	$02/06_2, 06_3/06_4$	one $02/06$ sector and one $06/06$ sector

Table: Incidences before resolving the local $\mathbb{C}^2/\mathbb{Z}_2$ singularities.

Let us track where does each plane sit from the local patch.
Considering one particular fixed point generated by F_α

0 planes intersection

On the local patch we define complex local coordinates - (z_1, z_2) .

Invariant coordinates

$$U = z_1^2, \quad V = z_2^2, \quad W = z_1 z_2, \quad (27)$$

which obey

$$UV - W^2 = 0. \quad (28)$$

The space of directions approaching the singularity is $\mathbb{CP}^1$ parameterized by w .

And we have t orthogonal to the $\mathbb{CP}^1$

$$t = z_1^2, \quad w = \frac{z_2}{z_1}. \quad (29)$$

Now we have a map from a singular point (U, V, W) to an Exceptional sphere $(t = 0, w)$

O Plane Intersection*

Tracking down our O plane locus from global coordinates y_i we check.
Both meet the exceptional divisor at

$$E_{\mathbb{R}} = \{t = 0, w \in \mathbb{R} \cup \{\infty\}\} \cong \mathbb{RP}^1 \cong S^1. \quad (30)$$

Which is a smooth non-self intersecting surface.

O Plane Intersection*

Tracking down our O plane locus from global coordinates y_i we check.
Both meet the exceptional divisor at

$$E_{\mathbb{R}} = \{t = 0, w \in \mathbb{R} \cup \{\infty\}\} \cong \mathbb{RP}^1 \cong S^1. \quad (31)$$

Which is a smooth non-self intersecting surface.

So we have shown that there are no intersection of the O planes on resolved orbifold but they are the infact the same planes sitting on top of each other.

Caveat

All the above story is valid of O_6 planes in the setup.

Looking at the O_2 planes sitting in our internal space. We see that they are precisely sitting on the fixed points.

Is the Tadpole ruined and is our theory inconsistent?*

Caveat

All the above story is valid of O_6 planes in the setup.

Looking at the O_2 planes sitting in our internal space. We see that they are precisely sitting on the fixed points.

Is the Tadpole ruined and is our theory inconsistent?*

NO!!

Thanks to the higher order curvature coupling of the WZ terms of O_6 planes.

Conclusion and Outlook

- We have shown that for these models with such singularity structure the Singular orbifold can be resolved
- It resolves at the right regime with keeping the scale separation intact
- O planes do not intersect

Thank You for Your Attention!