

Supersymmetric scale-separated vacua were meant IIB

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Based on [VH, 2502.04791]

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Hiding extra dimensions

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$$10 \neq 4$$

Hiding extra dimensions

$$10 \neq 4$$

$$11 \neq 4$$

Hiding extra dimensions

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$$11 \neq 4$$

Dimensionality of best
understood string
theories/M-theory

$\neq$

Observed dimensionality
of the universe

Hiding extra dimensions

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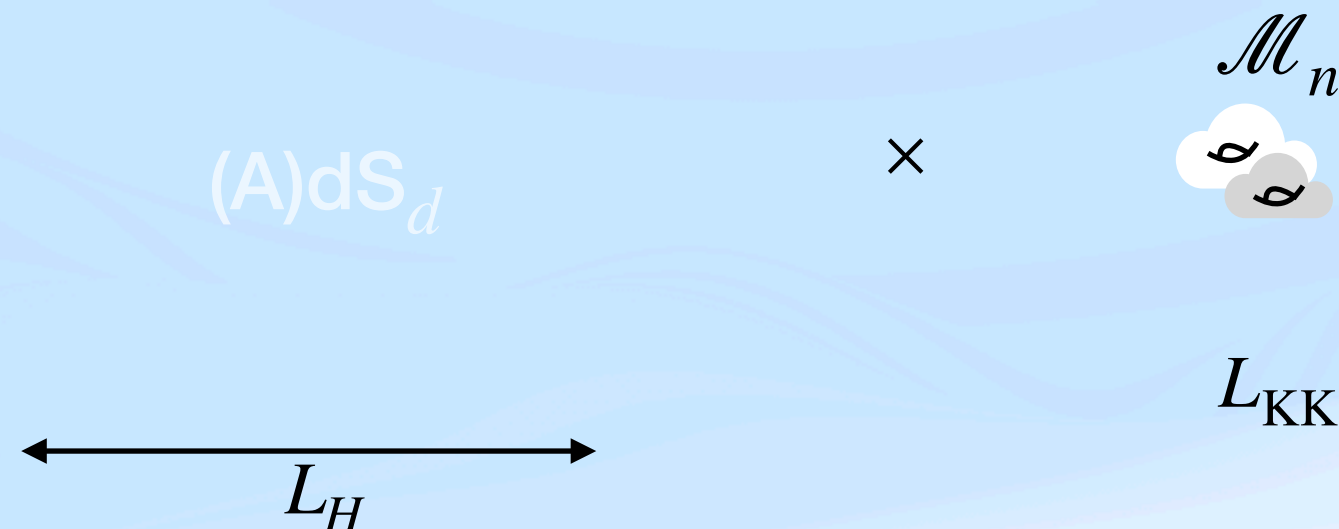
$\neq$

Observed dimensionality
of the universe

Solution



Compactifications & scale separation



The diagram illustrates the compactification of a higher-dimensional space. On the left, the text $(A)dS_d$ is shown. To its right is a multiplication symbol $\times$, followed by a representation of an n -dimensional manifold $\mathcal{M}_n$, depicted as a cluster of three overlapping circles. Below $\mathcal{M}_n$ is the label L_{KK} . A horizontal double-headed arrow is positioned below the $(A)dS_d$ term, with the label L_H centered underneath it.

$$ds_{d+n}^2 = g_{\mu\nu} dx^\mu dx^\nu + g_{mn} dy^m dy^n$$

Lower-dimensional effective field theory: **scale separation:**

$$L_{\text{KK}} \ll L_H$$

Extra dimensions smaller
than cosmological radius

$$\Lambda_d \ll M_{\text{KK}}^2$$

Cosmological constant
smaller than the UV cut-off

Moduli stabilisation

with fluxes, branes and O-planes

See e.g. [Van Riet, Zoccarato, 2023]

$$ds_{10}^2 = e^{\frac{2\Phi_d}{d-2}} g_{\mu\nu} dx^\mu dx^\nu + L^2 g_{mn} dy^m dy^n$$

$$S_{10} = \frac{M_{\text{Pl},10}^8}{2} \int \sqrt{-g_{10}} e^{-2\phi} (R_{10} + \dots) \longrightarrow S_d = \frac{M_{\text{Pl},d}^{d-2}}{2} \int \sqrt{-g_d} \left(R_d - g_{ij} \partial\phi^i \partial\phi^j - V(L, \Phi_d) \right)$$

$$\text{schematically } V = -R_n + \sum_k |F_k|^2 + V_{\text{D-branes}} - V_{\text{O-planes}}$$

Moduli stabilisation

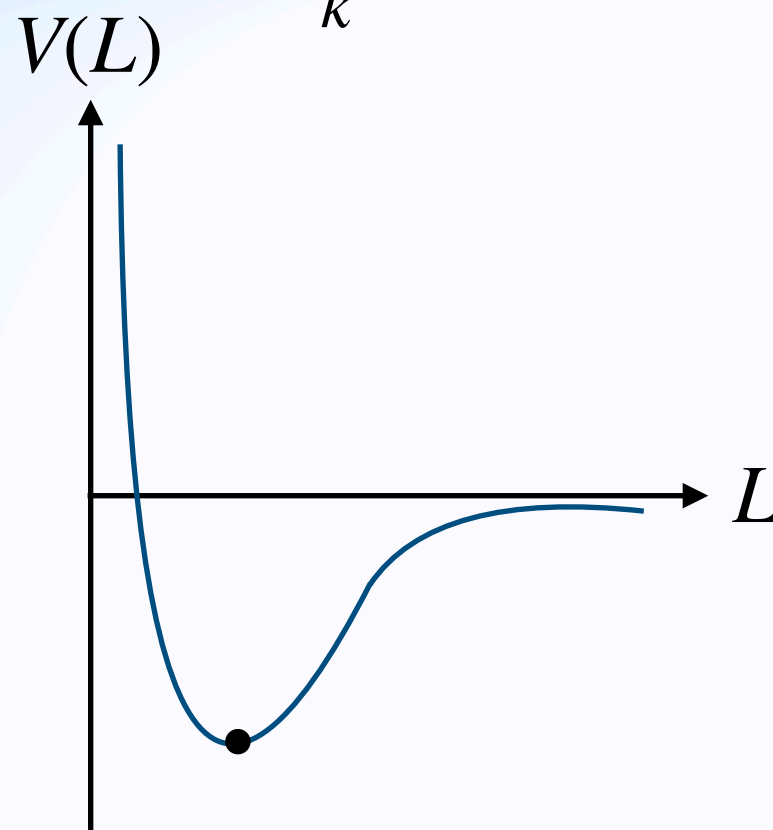
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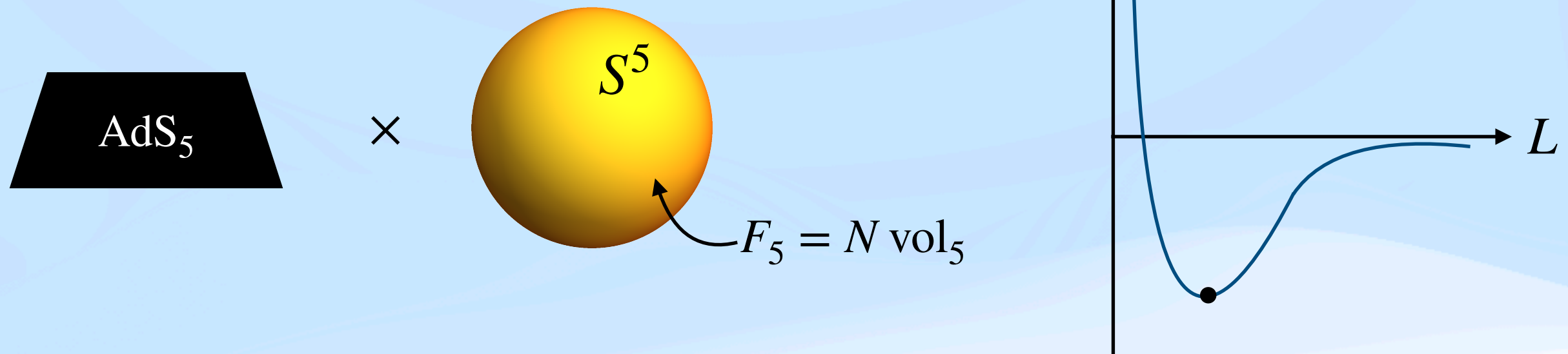
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Moduli stabilisation

e.g. $\text{AdS}_5 \times S^5$



Both scales are completely determined by the five-form flux

$$L = (4\pi g_s N)^{1/4} \ell_s = L_{\text{AdS}}$$

So **no scale separation!** It is possible for compactifications with

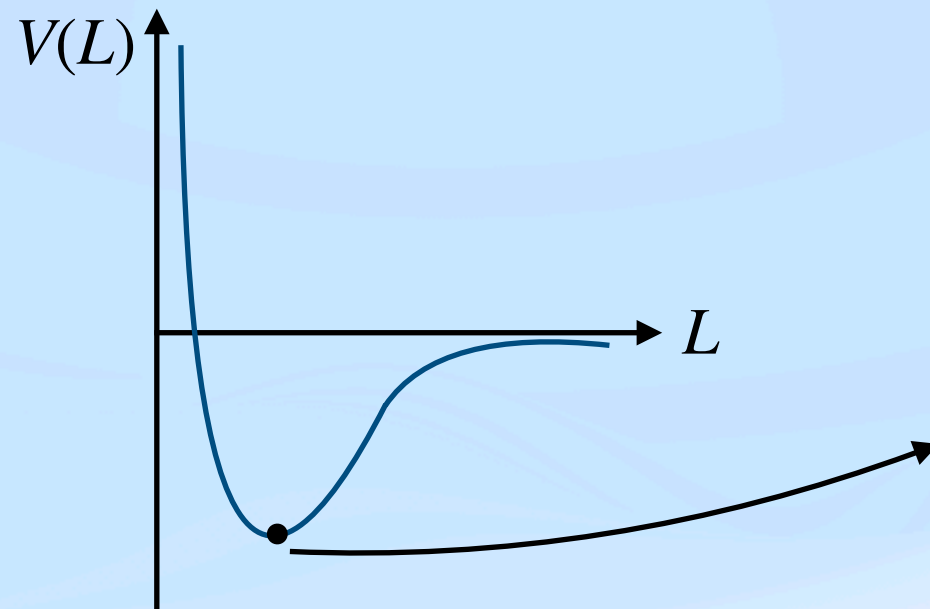
$$L_{\text{KK}} \sim N^\alpha \ell_s \quad \text{and} \quad L_{\text{AdS}} \sim N^\beta \ell_s$$

with $0 < \alpha < \beta$.

We do have such examples in type IIA string compactifications

DGKT-CFI, 4d & 3d variants
by:
Cribiori, Junghans, VVH,
Van Riet, Wrase, Andriot,
Horer, Marconnet, Farakos,
Tringas, Van Riet, Emelin,
Morittu, Arboleya,
Guarino,...]

No scale separation?



$$\text{If } R_n \equiv L_R^{-2} \sim L_{\text{KK}}^{-2}$$

No scale separation

Theorems

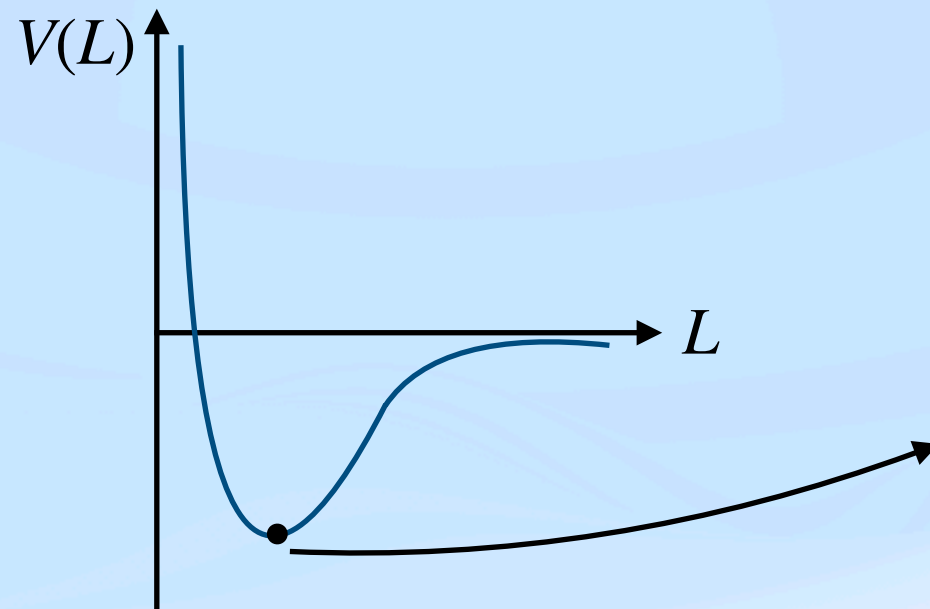
[Gautason, Schillo, Van Riet, Williams, 2016],
[De Luca, Tomasiello, 2021], [Tringas, Wrase, 2025]

Strong AdS distance conjecture: with SUSY, $m \sim \Lambda^{1/2}$ [Lüst, Palti, Vafa, 2019]

Arguments against scale separation with extended
SUSY and $Q \geq 8$

[Cribiori, Dall'agata, 2022]
[Cribiori, Montella, 2023]
[Cribiori, Farakos, Liatsos, 2024]
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Arguments against scale separation with extended
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Not for $\mathcal{N} = 1$ in $d \leq 4$

[Cribiori, Dall'agata, 2022]
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[Perlmutter, 2024],...

Option 1: no curvature

$\mathcal{N} = 1$ state of the art with parametric scale separation

massive Type IIA ($F_0 \neq 0$)

Type IIB

O6-planes, H

$$F_4 \sim N$$

4d: $\mathcal{M}_6 = \text{CY}_3$ [Dewolfe, Giryavets, Kachru, Taylor, 2005],
[Camara, Font, Ibáñez, 2005], ...

3d: $\mathcal{M}_7 = \mathcal{M}_{G_2}$ [Farakos, Tringas, Van Riet, 2020]



massless Type IIA ($F_0 = 0$)



Option 2: curvature

Can the curvature scale decouple?

L_{KK}^{-2} estimated by λ_1 = lowest eigenvalue of scalar Laplacian

$$|R_n| \ll \lambda_1$$

**Positively curved
manifolds**



No examples known

Claimed to be impossible by
[Collins, Jafferis, Vafa, Xu, Yau, 2022]
However, see
[Cribiori, Junghans, VVH, Van Riet,
Wrase, 2021], [VVH, 2024]

**Negatively curved
manifolds**



Nilmanifolds

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**Negatively curved
manifolds**



Nilmanifolds

(Negative energy density necessary, such as by O-planes)

The harmonic oscillator of nilmanifolds

$$ds_3^2 = (L_1 e^1)^2 + (L_2 e^2)^2 + (L_3 e^3)^2 \quad e^1 = dx^2 + \omega x^2 dx^3, \quad e^{2,3} = dx^{2,3}$$

$$de^1 = \omega e^2 \wedge e^3, \quad de^{2,3} = 0$$

Laplacian spectrum

$$m_{p,q}^2 = p^2 \left(\frac{2\pi}{L_2} \right)^2 + q^2 \left(\frac{2\pi}{L_3} \right)^2$$

$$m_{k,l,n}^2 = k^2 \left(\frac{2\pi}{L_1} \right)^2 + (2n+1) |k\omega| \frac{2\pi}{L_2 L_3}$$

[Andriot, Tsimpis, 2018]

Curvature

$$R_3 = - \left(\omega \frac{L_1}{L_2 L_3} \right)^2$$

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Take $L_1 \ll L_3 \leq L_2$
and ω fixed

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$$\lambda_1 = m_{1,0}^2 = \left(\frac{2\pi}{L_2} \right)^2$$

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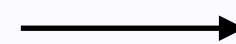
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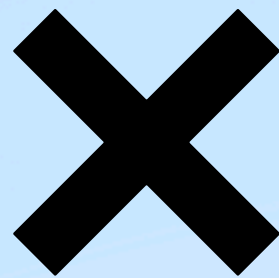


$$|R_3| \ll \left(\frac{\omega}{L_2} \right)^2 \sim \lambda_1$$

$\mathcal{N} = 1$ state of the art with parametric scale separation

massive Type IIA ($F_0 \neq 0$)

Type IIB



massless Type IIA ($F_0 = 0$)

O6-planes,

$$H = 0, F_6 \sim N^a,$$

$$F_{2,2} \sim N^b, F_{2,3} \sim N^c$$

In 4d only

[Cribiori, Junghans, VVH, Van Riet, Wrase, 2021]
see also [Koerber, Lüst, Tsimpis, 2008], [Caviezel,
Koerber, Kors, Lüst, Tsimpis, Zagermann, 2009],...
[Carrasco, Coudarchet, Marchesano, Prieto, 2023]

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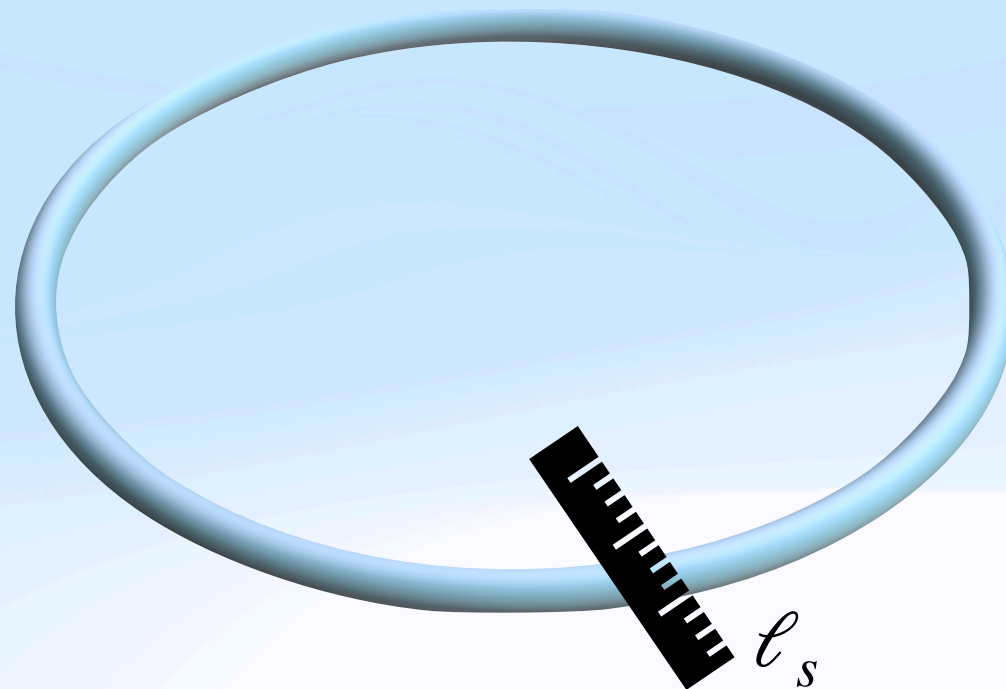
IIB or not IIB?

Supersymmetric history

Efforts only for AdS_4 : $\text{SU}(2)$ -structures

[Cavaziel, Wrase, Zagermann 2009]
[Petrini, Solard, Van Riet, 2013]

Problem:



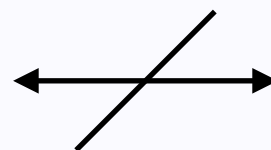
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For AdS_3 : co-closed G_2 -structures?

→ If too many O5-planes:

[Emelin, Farakos, Tringas, 2021]
[Passias, Prins, 2019]
[VVH, 2022]

Geometric flux quantisation



Scale separation

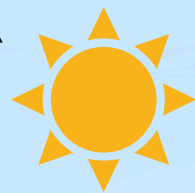
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Problem:

$$\text{Vol}(\Sigma_p)/\ell_s^p$$



[Cribiori, Junghans, VVH, Van Riet, Wrase, 2021]

g_s

For AdS_3 : co-closed G_2 -structures?

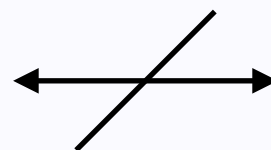
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→ If too many O5-planes:

Geometric flux quantisation



Scale separation

$\mathcal{N} = 1$ scale-separated AdS
solutions of type IIB

with nilmanifolds

$\mathcal{N} = 1$ AdS_3 solutions of type IIB

On G_2 -structure manifolds

[Dibitetto, Lo Monaco, Passias, Tomasiello, 2018]

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[VVH, 2022]

Co-closed G_2 -structure

$$d\Phi = W_1 \star \Phi + W_{27}, \quad d \star \Phi = 0, \quad W_1 = 12/(7L_{\text{AdS}})$$

SUSY variations imply

$$\begin{aligned} g_s F_3 &= - \star d\Phi + 2\mu\Phi & F_1 &= 0 \\ g_s F_7 &= - 2\mu \text{vol}_7 & F_5 &= 0 \\ & & H &= 0 \end{aligned}$$

Non-trivial Bianchi identity

$$dF_3 = \sum_i j_{\text{D5},i} - \sum_i j_{\text{O5},i}$$

$\mathcal{N} = 1$ AdS_3 solutions of type IIB

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$\mathbb{Z}_2^3$ toroidal orbifolds

[Joyce, 1996]²,...
[Dall'Agata, Prezas, 2005,],
[Emelin, Farakos, Tringas, 2021],...

Joyce orbifolds

$\mathbb{Z}_2^3$ orbifold generators

x^i	1	2	3	4	5	6	7
α	-	-	+	+	-	-	+
β	+	+	-	-	-	-	+
γ	-	+	-	+	-	+	-

+ shifts

Orientifold involution

x^i	1	2	3	4	5	6	7
σ_α	-	-	+	+	-	-	+

Invariant cycles/forms

Σ_3	1	2	3	4	5	6	7
$e^{\cdots}$	127	347	567	145	235	136	426

+ dual 4-cycles/forms

$$\Phi = L_1 L_2 L_7 e^{127} + L_3 L_4 L_7 e^{347} + L_5 L_6 L_7 e^{567} \\ + L_1 L_4 L_5 e^{145} + L_2 L_3 L_5 e^{235} + L_1 L_3 L_6 e^{136} - L_2 L_4 L_6 e^{246}$$

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+ shifts

Orientifold involution

x^i	1	2	3	4	5	6	7
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$$de^a = \frac{1}{2} f^a_{bc} e^b \wedge e^c \quad \rightarrow \quad \text{must be invariant}$$

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→ W_1, W_{27}

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+ shifts

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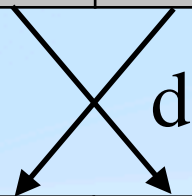
$\rightarrow W_1, W_{27}$

The harmonic oscillator: $\mathfrak{n}_2 = \text{Nil}_3 \times \mathbb{T}^4$

$$de^7 = \omega e^1 \wedge e^2$$

Σ_3	1	2	3	4	5	6	7
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Σ_3	1	2	3	4	5	6	7
F_3							



Σ_4	1	2	3	4	5	6	7
$\perp \text{ O5}$		j_{O5}	j_{O5}				

$\frac{L_{1,2}^2}{L_{\text{AdS}}^2} \sim \frac{1}{\text{orange box}}^2$

$\frac{L_{3,4,5,6}^2}{L_{\text{AdS}}^2} \sim \frac{1}{\text{orange box}}^2$

$\frac{L_7^2}{L_{\text{AdS}}^2} \sim \frac{1}{\text{orange box}}^4$

$g_s \sim \frac{\text{blue box}^{1/2}}{\text{orange box}^{1/2} \text{ orange box}^2}$

	L_1	L_2	L_3	L_4	L_5	L_6	L_7
F_3							
F_3							
F_3							
F_3							
F_3							
F_3							
F_3							
F_7							

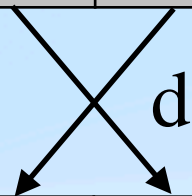
$$L_i = \frac{\prod (\text{indicated fluxes})^{1/2}}{\prod (\text{all fluxes})^{1/4}} \times l_i$$

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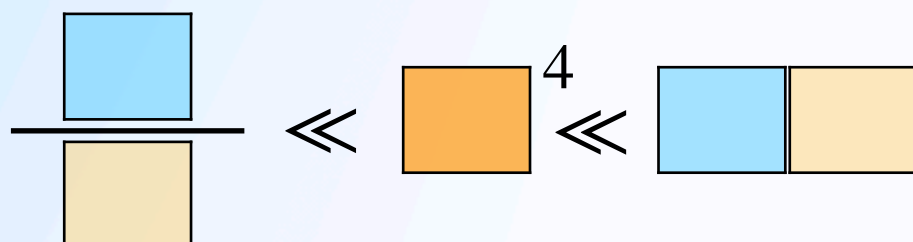


Σ_4	1	2	3	4	5	6	7
$\perp \text{ O5}$		j_{O5}	j_{O5}				

Scale separation

+

Well-controlled regime if 



	L_1	L_2	L_3	L_4	L_5	L_6	L_7
F_3							
F_3							
F_3							
F_3							
F_3							
F_3							
F_3							
F_7							

$$m^2 L_{\text{AdS}}^2 = \{120, 8^7\}$$

$$\Delta = \{12, 4^7\}$$

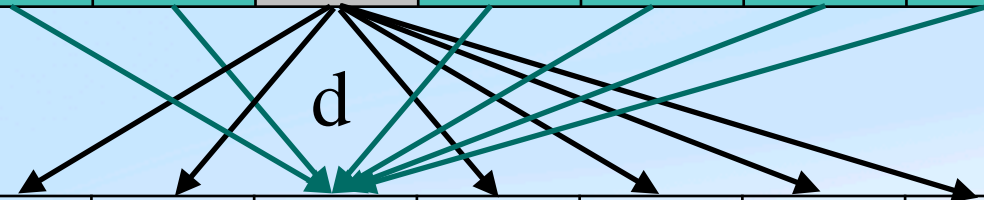
[Conlon, Ning, Revello, 2021],...

37D₁

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Σ_3	1	2	3	4	5	6	7
F_3							

Σ_4	1	2	3	4	5	6	7
$\perp O5$	$\dot{j}_{O5}$	$\dot{j}_{O5}$		$\dot{j}_{O5}$	$\dot{j}_{O5}$	$\dot{j}_{O5}$	$\dot{j}_{O5}$



Non-trivially unbounded fluxes

Similar scale separated limit

	L_1	L_2	L_3	L_4	L_5	L_6	L_7
F_3							
F_3							
F_3							
F_3							
F_3							
F_3							
F_3							
F_7							

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37D₁

Σ_3	1	2	3	4	5	6	7
$e^{\cdots}$	127	347	567	145	235	136	426

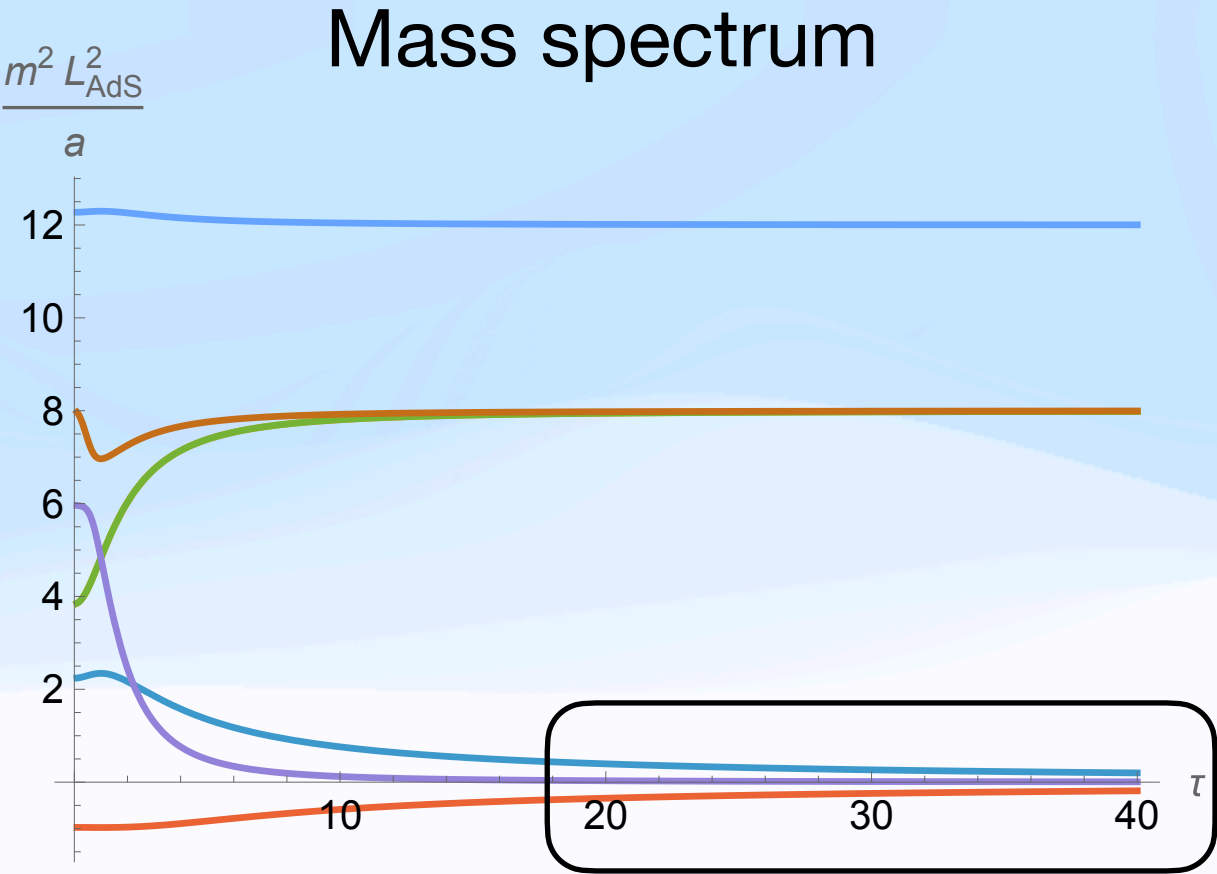
Σ_3	1	2	3	4	5	6	7
F_3							

Σ_4	1	2	3	4	5	6	7
$\perp O5$	j_{O5}	j_{O5}		j_{O5}	j_{O5}	j_{O5}	j_{O5}



Non-trivially unbounded fluxes

Similar scale separated limit



3 scalars
become light

Similar as in [Farakos, Tringas, 2025],
related by T-duality

Remark about solvmanifold constructions

Similar solvmanifold constructions on $\mathbb{Z}_2^3$ orbifolds

Hierarchy between cycle volumes and AdS scale:

$$\frac{\text{Vol}(\Sigma_p)}{L_{\text{AdS}}^p} \ll 1 \text{ for } p = 3, 4$$

But λ_1 decouples from cycle volumes:

$$L_{\text{KK}} \sim \lambda_1^{-1/2} \propto \text{Vol}(\Sigma_p)^{1/p}$$

[VH, 2025]

Also true for non-supersymmetric cousins

[Arbolea, Guarino, Morittu, 2024 & 2025]
[Arbolea, Guarino, Morittu, Sudano, 2025]

Conclusions

First supersymmetric, scale-separated solutions of type IIB

Nilmanifolds with the required property $|R_7| \ll \lambda_1$ (or $|W_1| \ll \lambda_1$)

Co-closed G_2 -structure, with only RR fluxes

Simplest example is $\text{AdS}_3 \times (\text{Nil}_3 \times \mathbb{T}^4)/\mathbb{Z}_2^3$ with only 2 sets of O-planes

Outlook

Go beyond $\mathbb{Z}_2^3$ orbifolds

Investigating T-duals in IIA into $SU(3)$ -structures

[Passias, Prins, 2020]

Going beyond the smearing approximation

[Junghans, 2020]
[Marchesano, Palti, Quirant, Tomasiello, 2020]
[Cribiori, Junghans, VVH, Van Riet, Wrase, 2021],
[Junghans, 2023], [VVH, 2024]
[Emelin, Farakos, Tringas, 2022], [Emelin, 2024],...

Formulating the holographic dual field theory?

[Apers, Montero, Valenzuela, 2025]



Thank you!