

# The Double Expansion in gravitational EFTs

Alvaro Herraez

Based on:

[arXiv:2310.07708 ] with A. Castellano, L.E. Ibáñez

**[arXiv:2501.14880] with J. Calderón-Infante, A. Castellano**

**MAX-PLANCK-INSTITUT**  
FÜR PHYSIK



Workshop on Quantum Gravity and Strings, Corfu Summer Institute 2025

Sept 10, 2025

# **What is the Structure of Gravitational EFTs?**

**Can we exploit the multi-scale structure to bound  
gravitational Wilson Coefficients?**

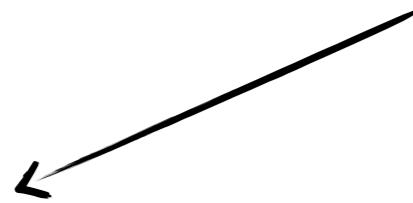
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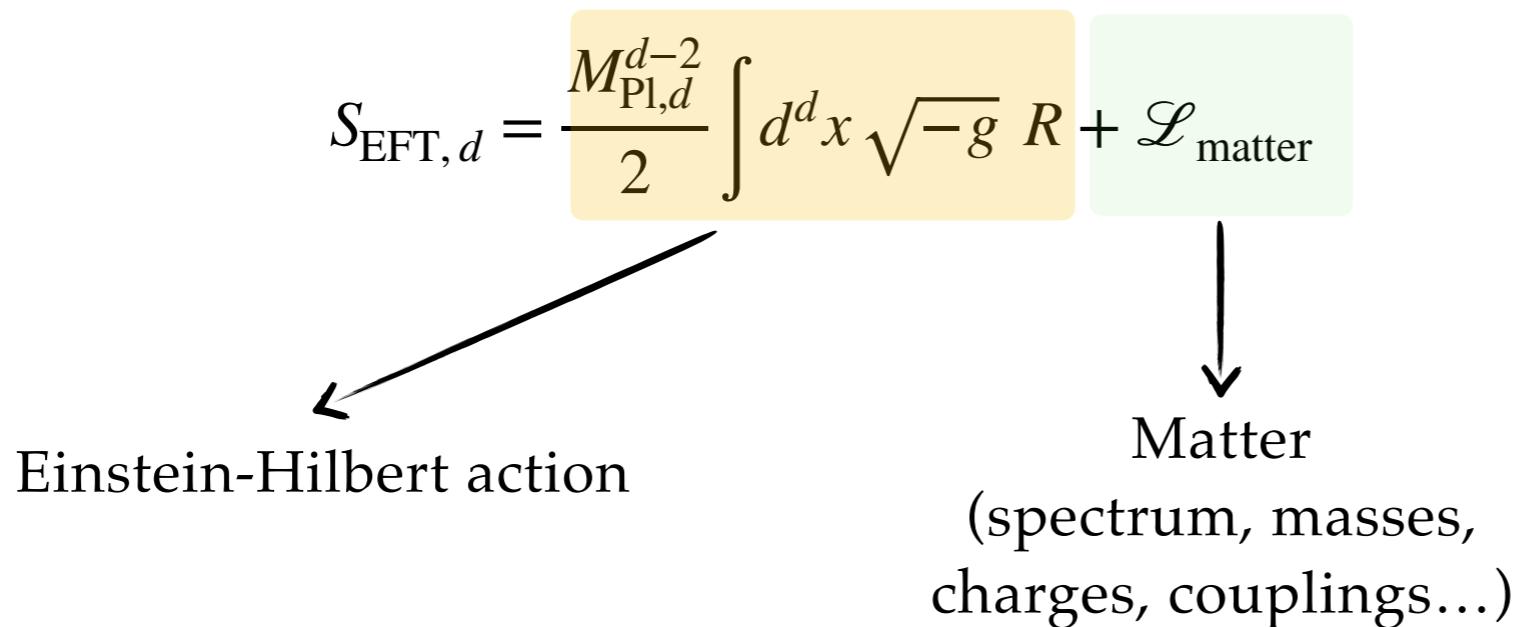
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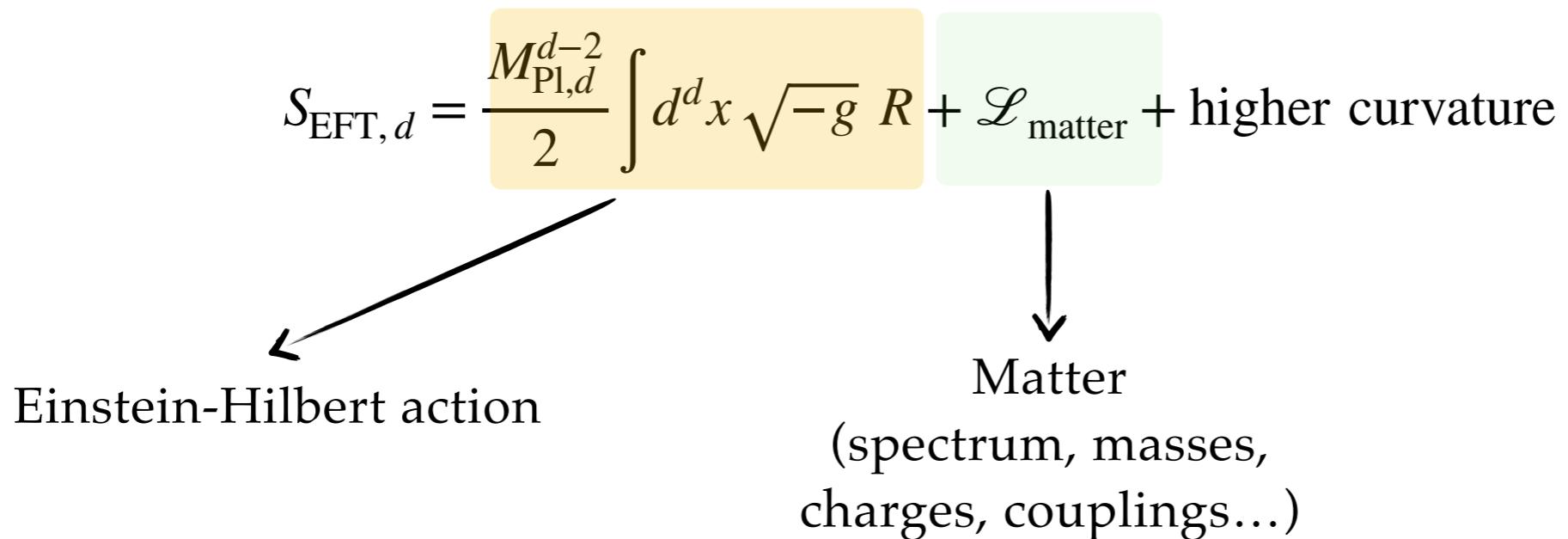


Einstein-Hilbert action

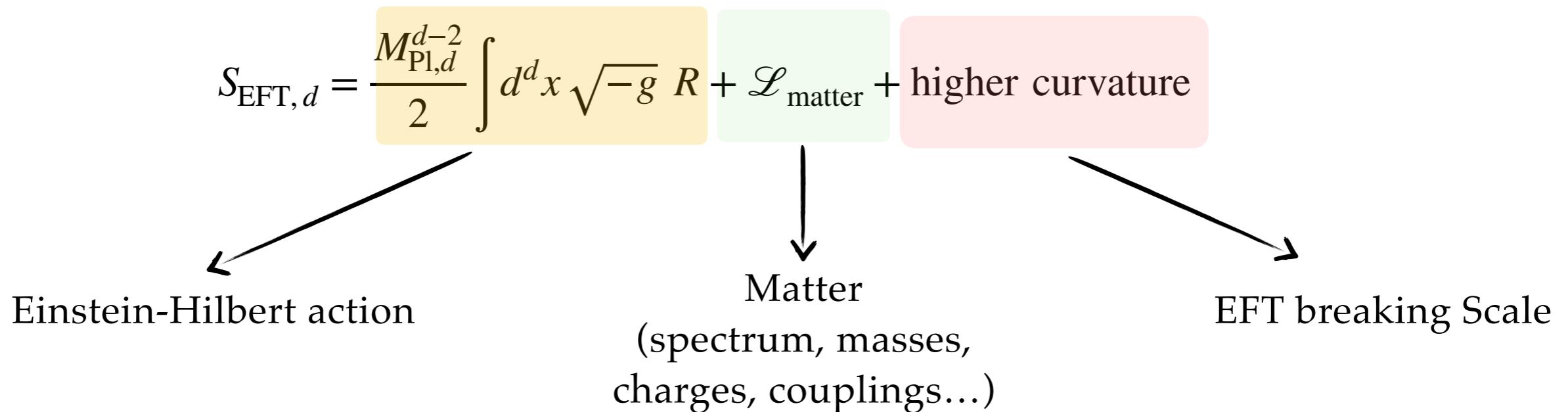
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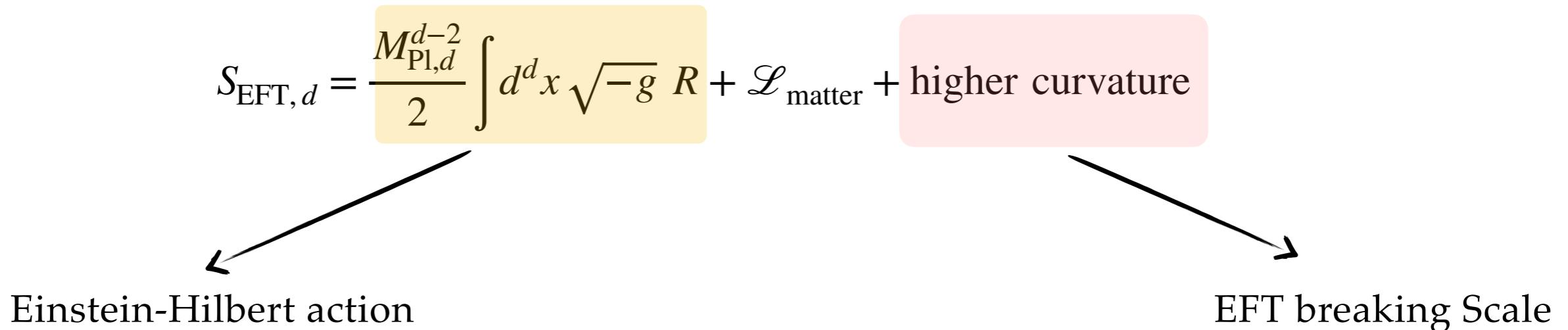
$$S_{\text{EFT},d} = \frac{M_{\text{Pl},d}^{d-2}}{2} \int d^d x \sqrt{-g} R + \mathcal{L}_{\text{matter}} + \text{higher curvature}$$

Einstein-Hilbert action

EFT breaking Scale

- What is the structure of the gravitational part of the action?

# The structure of gravitational EFTs

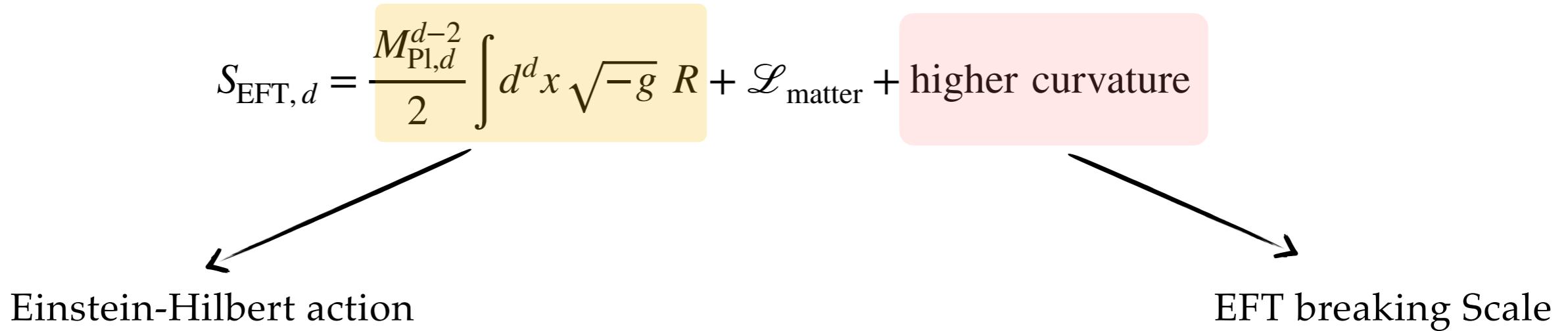


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[Calderón-Infante, Castellano, AH '25]

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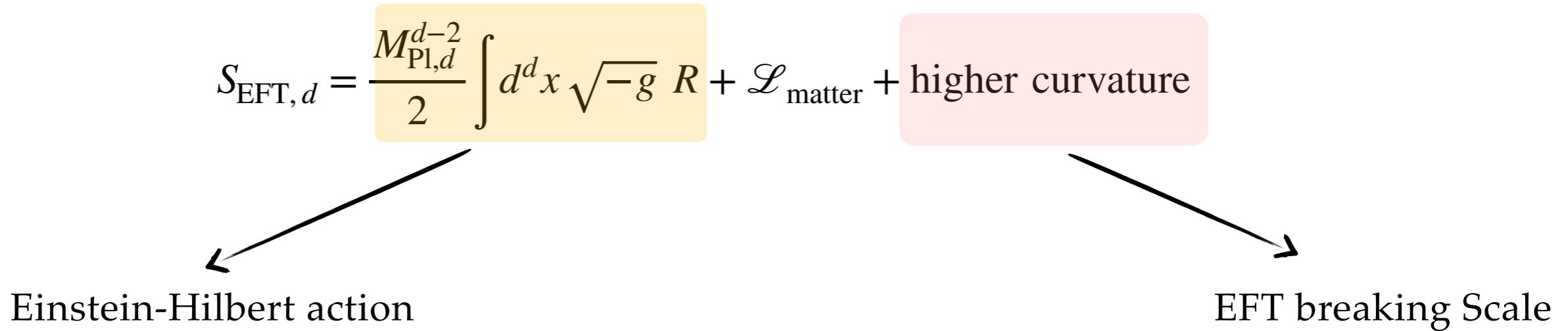
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Quantum Gravity cutoff

[Calderón-Infante, Castellano, AH '25]

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Quantum Gravity cutoff      Mass of a (tower of) state(s) outside original EFT

[Calderón-Infante, Castellano, AH '25]

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- Reaching the scale where genuine Quantum-Gravitational effects kick-in  $\longrightarrow$  Intrinsic non-localities, no local EFT framework

[van de Heisteeg, Vafa, Wiesner,(Wu) '22-'24]  
 [Castellano, AH, Ibáñez '23]  
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$\longrightarrow$  Particularly interesting in the *Asymptotic Regime*:  $M_{\text{Pl},d} \gg \Lambda_{\text{QG}} \gtrsim M_t$

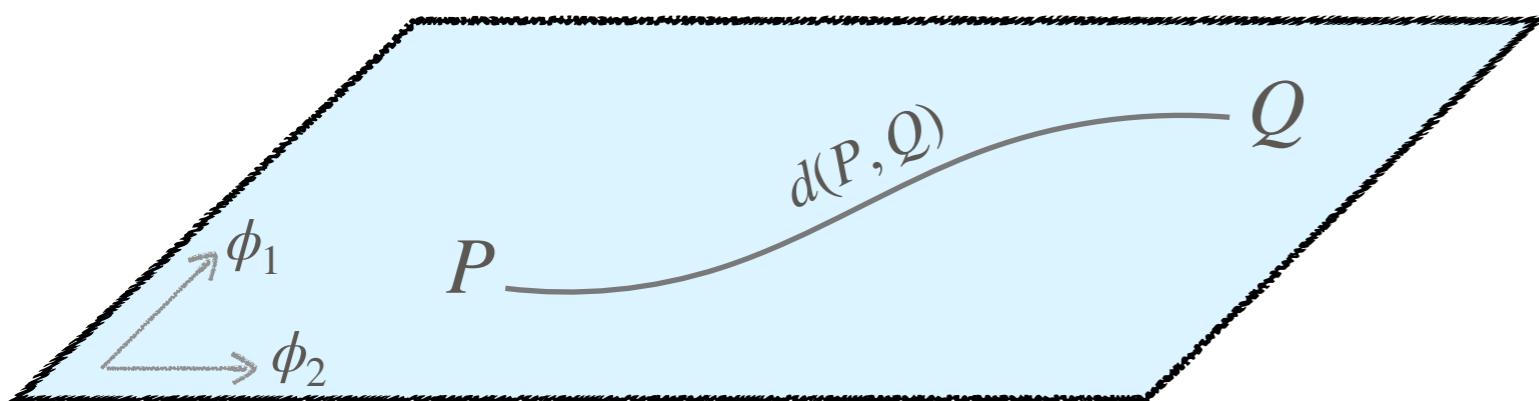
# The Asymptotic regime I: The Distance Conjecture

In a theory of QG, moving in moduli space from a point P to a point Q an infinite distance away, an infinite tower of states become light **exponentially in the geodesic distance**

$$M(Q) \sim M(P) e^{-\alpha \Delta_\phi(P,Q)}$$

[Ooguri, Vafa '06]

Scalar manifold with metric  $g_{ij}(\phi_i)$  from kinetic terms



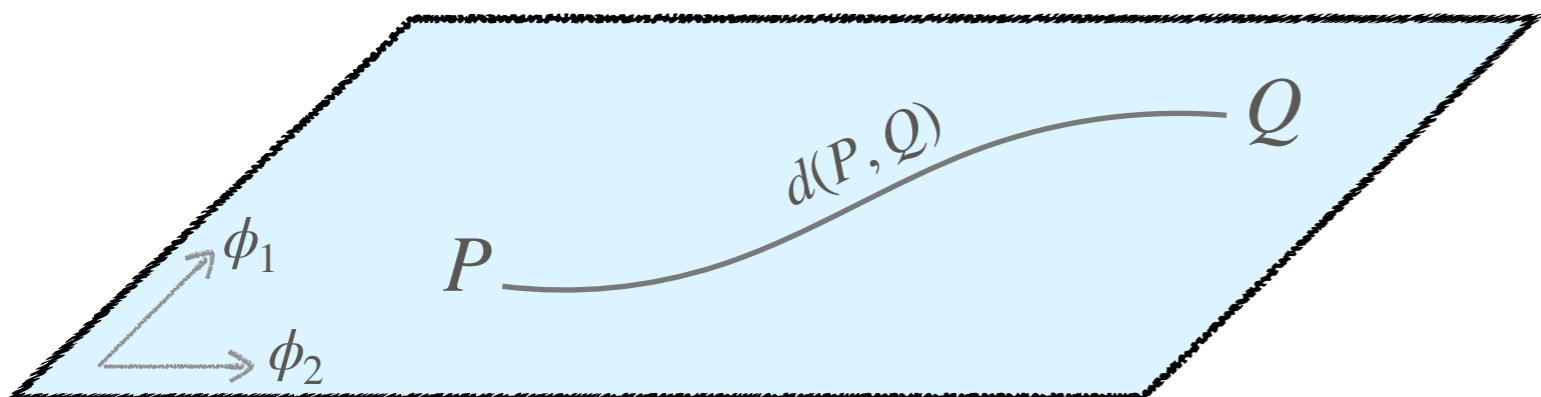
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Emergent String Conjecture: The tower becoming light is either:

- (Dual to) Kaluza-Klein tower
- Excitation modes of weakly coupled string

[Lee, Lerche, Weigand '19]

## The Asymptotic regime II: The Species Scale

- Maximum UV cut-off in QG in the presence of  $N$  light species [Dvali '07] [Dvali, Reedi '08]  
[Dvali, Lüst '10] [Dvali, Gómez '10]

$$M_{\text{Pl,d}} \longrightarrow \Lambda_{\text{QG}} = \frac{M_{\text{Pl,d}}}{N_{\text{sp}}^{\frac{1}{d-2}}} = N_{\text{sp}}^{1/p} M_t$$

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$$\Lambda_{\text{QG}} \simeq M_t \left( \frac{M_{\text{Pl},d}}{M_t} \right)^{\frac{d-2}{d+p-2}} \longrightarrow M_s \quad N_{\text{sp}} \simeq \left( \frac{M_{\text{Pl},d}}{M_t} \right)^{\frac{p(d-2)}{d+p-2}}$$

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[Castellano, AH, Ibáñez '21]

**Species Scale of the d-dim EFT**  $\begin{cases} M_{\text{Pl},d+p} \\ M_{\text{str}} \end{cases}$  (decompactification of  $p$  dimensions)  
(weakly coupled string limits)

# The double EFT expansion: Leading Contributions

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What is this scale in the gravitational sector? → Which term dominates in general?

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EXAMPLE: Maximal 10d supergravity (type IIA) → 4 graviton amplitude

- \* Decompactification limit ( $g_s \rightarrow \infty$ )
- \* Weak string coupling limit ( $g_s \rightarrow 0$ )

$$\mathcal{A}_4 = \mathcal{A}_4^{\text{sugra}} f(s, t), \quad \mathcal{A}_4^{\text{sugra}} = -\hat{K} \ell_{\text{Pl},10}^8 g_s^{-4} \frac{64}{\alpha'^3 stu},$$

$$f(s, t) = 1 + \ell_{\text{Pl},10}^6 f_{R^4}(s, t) + \ell_{\text{Pl},10}^{10} f_{D^4 R^4}(s, t) + \ell_{\text{Pl},10}^{12} f_{D^6 R^4}(s, t) + \dots,$$

[Green, Gutperle, Kwan, Russo, Tseytlin, Vanhove '97-'07]

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$$\Lambda_{\text{QG}} \simeq M_t^{\frac{p}{d+p-2}}$$

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$$f_{D^4 R^4} = \frac{(4\pi)^{5/4}}{(2\pi)^{10}} \frac{stu(s^2 + t^2 + u^2)}{2^{10}} \left( \zeta(5)g_s^{-5/2} + \frac{2\pi^4}{135}g_s^{3/2} \right) \longrightarrow \ell_{\text{Pl},10}^{10} g_s^{3/2} = \ell_{\text{Pl},10}^8 M_{KK}^{-2} \neq \Lambda_{\text{QG}}^{-10}$$

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# The double EFT expansion: Leading Contributions

## (Duals of) Decompactification limits

$$\Lambda_{\text{QG}} \simeq M_t^{\frac{p}{d+p-2}}$$

- \*  $\frac{M_{\text{Pl},d}^{d-2}}{\Lambda_{\text{QG}}^{n-2}}$  if  $n < d + p \longrightarrow$  UV divergent (Relevant in the decomp. th)
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$\longrightarrow \ell_{\text{Pl},10}^{12} g_s^3 = \ell_{\text{Pl},10}^8 M_{KK}^{-4} \neq \Lambda_{\text{QG}}^{-12}$

# The double EFT expansion: Leading Contributions

## Weakly coupled (emergent) string limits

$$\Lambda_{\text{QG}} \simeq M_{\text{tower}} = M_s$$

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# The double EFT expansion: General Lessons

$$S_{\text{EFT},d} = \int d^d x \sqrt{-g} \frac{M_{\text{Pl},d}^{d-2}}{2} \left( R + \sum_{n>2} \frac{\mathcal{O}_n(R)}{\Lambda_{\text{QG}}^{n-2}} \right) + \int d^d x \sqrt{-g} \sum_{n \geq d} \frac{\mathcal{O}_n(R)}{M_t^{n-d}}$$

- Decompactification limits: Non-vanishing contribution to an infinite tower of UV convergent operators of the form  $\frac{D^{2\ell} R^4}{M_t^{8+2\ell-d}} \longrightarrow$  Field-theoretic series diverges as the scale  $M_t$  is approached  $\longrightarrow$  Resummed into higher dim. 2-derivative with more dof

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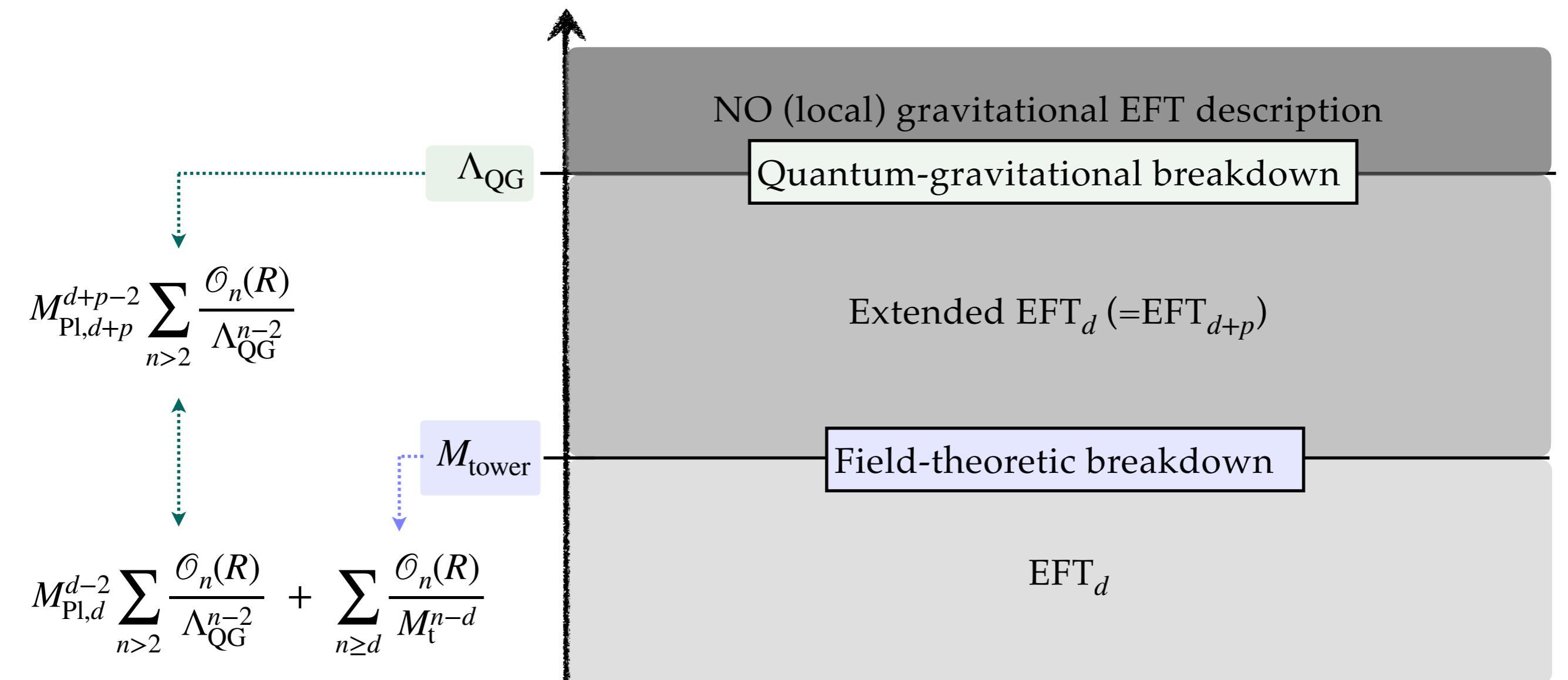
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In all observed examples, the double EFT expansion captures the leading contribution for every (non-vanishing) Wilson Coefficient

# The double EFT expansion: General Lessons



[Calderón-Infante, Castellano, AH '25]

**Can we exploit this multi-scale structure to bound  
gravitational Wilson Coefficients?**

# The Bottom-up Perspective: Constraints on Wilson Coefficients

- Assuming the Double EFT expansion, with the leading piece for a (non-vanishing) gravitational Wilson coefficient being given by the FT or QG pieces

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Expansion

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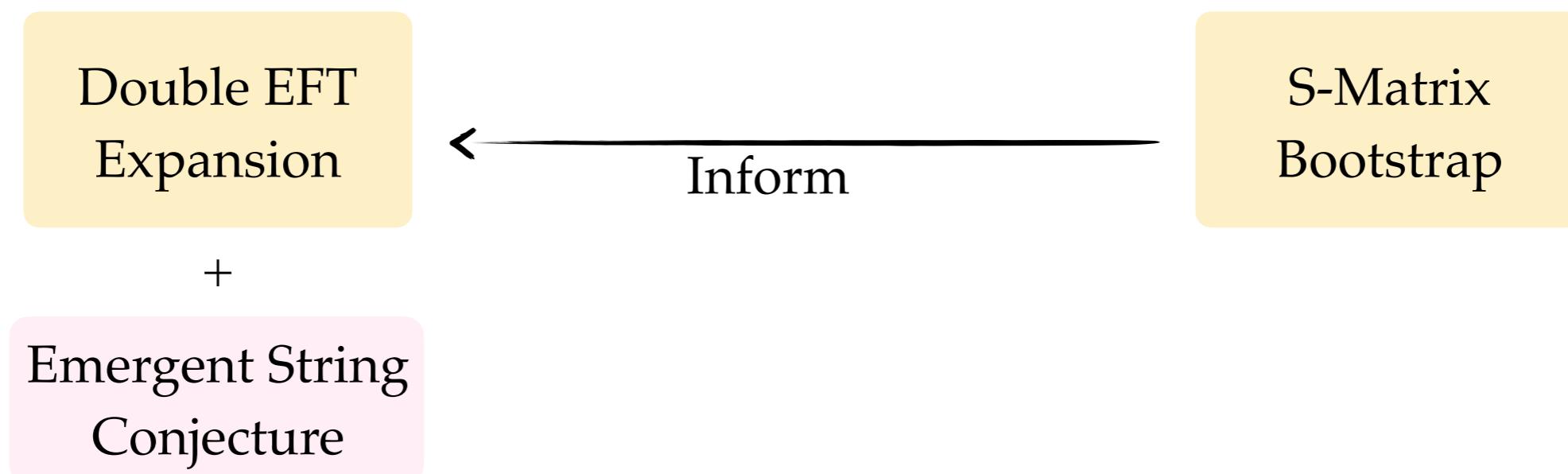
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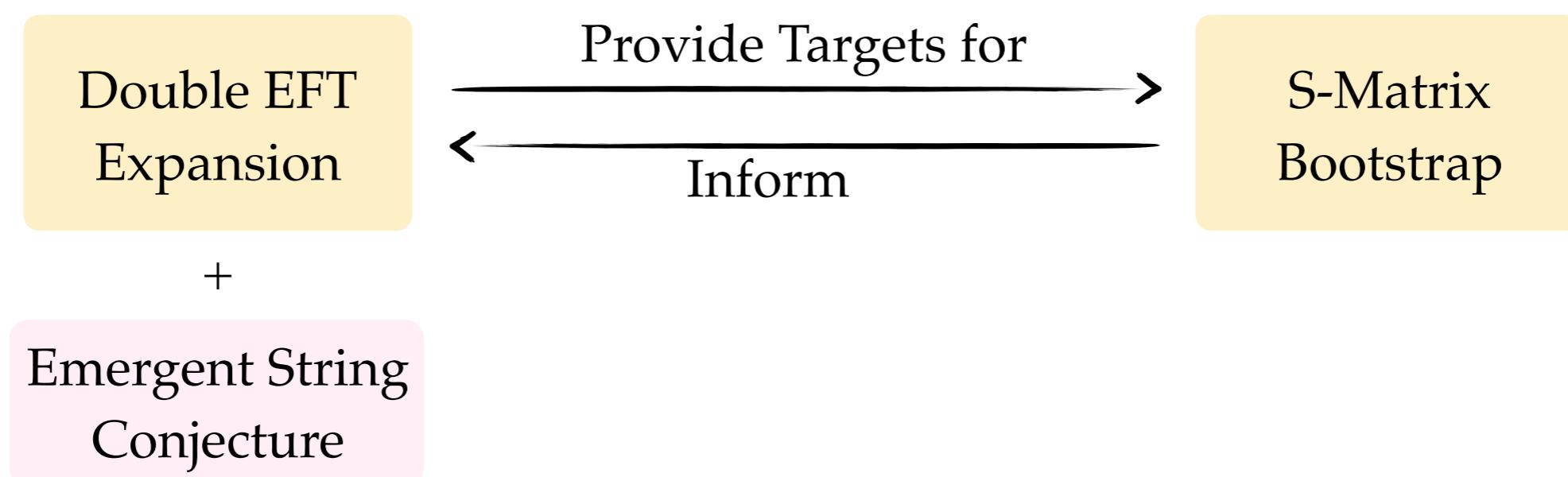


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- We have two scales,  $\Lambda_{\text{QG}}$  and  $M$  and an infinite number of Wilson Coefficients  
—————> Solve for the scales in terms of Wilson coefficients and substitute back

$$\Lambda_{\text{QG}} \simeq M^\gamma \quad 0 \leq \gamma \leq 1$$

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- Consider **10d maximal SUGRA**:

$$\alpha_{R^4} = a_{R^4} M^{-6\gamma} + b_{R_4} M^2 + \dots \quad \text{Assume } a_{R^4} \neq 0 \quad \alpha_{R^4} \sim M^{-6\gamma}$$

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[ Guerrieri, (Murali,) Penedones, Vieira '21 '23]

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Bounds on  $\alpha_{D^4 R^4} = a_{D^4 R^4} M^{-10\gamma} + b_{D^4 R^4} M^{-2}$  as  $\alpha_{R^4} \rightarrow \infty$

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$$\alpha_n = \frac{a_n}{\Lambda_{\text{QG}}^{n-2}} + \frac{b_n}{M^{n-d}} + \dots \quad (\text{In Planck Units})$$

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 —→ Solve for the scales in terms of Wilson coefficients and substitute back

$$\Lambda_{\text{QG}} \simeq M^\gamma \quad 0 \leq \gamma \leq 1$$

$$\text{ESC: } \gamma = \frac{p}{d+p-2} \longrightarrow \frac{1}{9} \leq \gamma \leq 1$$

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Stronger than  
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# Summary

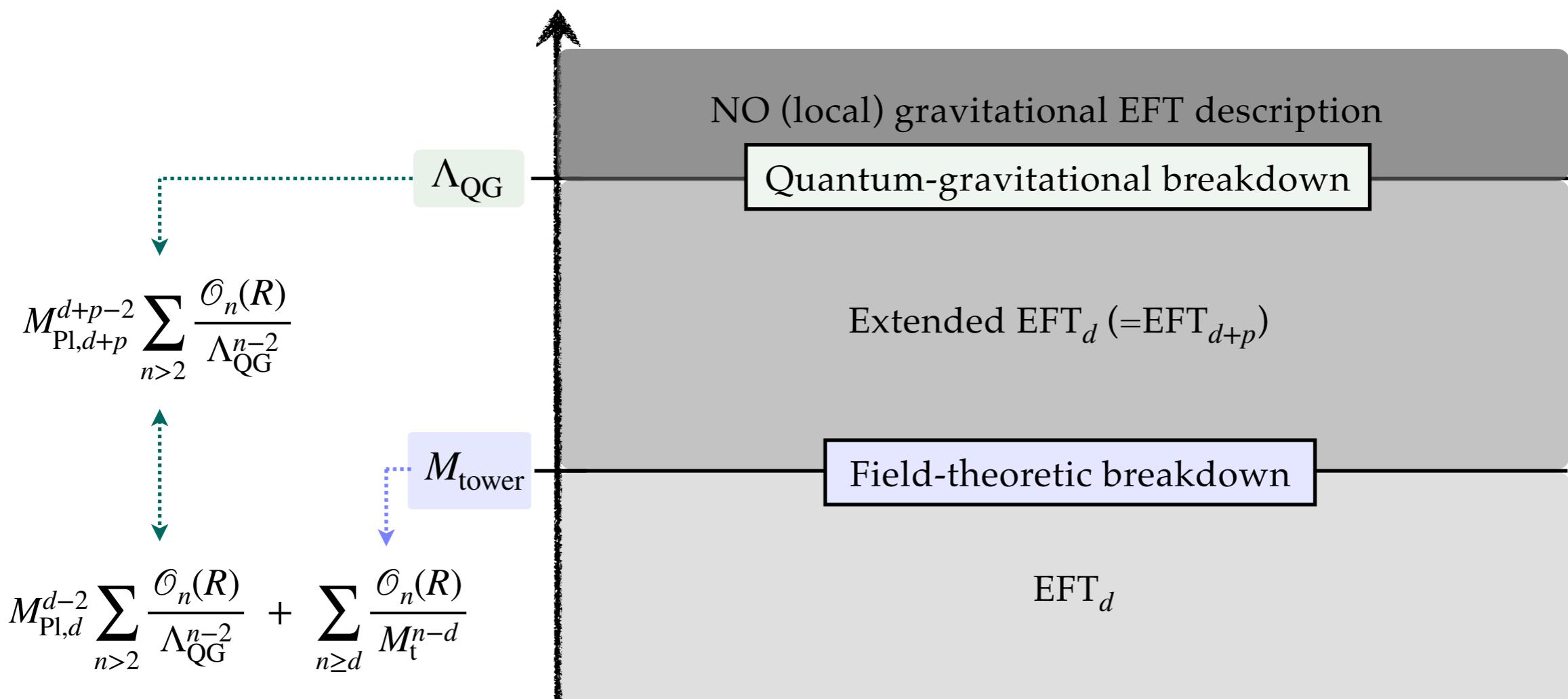
**What is the Structure of Gravitational EFTs?**

**Can we exploit this multi-scale structure to bound  
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# Summary

## What is the Structure of Gravitational EFTs?

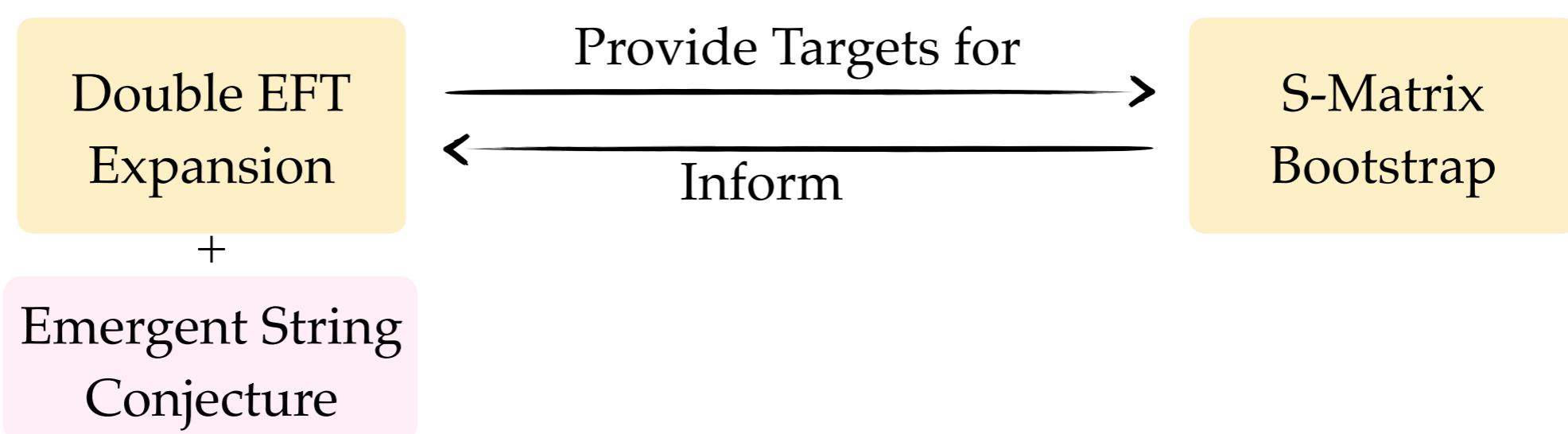
$$S_{\text{EFT},d} = \int d^d x \sqrt{-g} \frac{M_{\text{Pl},d}^{d-2}}{2} \left( R + \sum_{n>2} \frac{\mathcal{O}_n(R)}{\Lambda_{\text{QG}}^{n-2}} \right) + \int d^d x \sqrt{-g} \sum_{n \geq d} \frac{\mathcal{O}_n(R)}{M_t^{n-d}}$$



# Summary

Can we exploit this multi-scale structure to bound gravitational Wilson Coefficients?

$$\alpha_n = \frac{a_n}{\Lambda_{\text{QG}}^{n-2}} + \frac{b_n}{M^{n-d}} + \dots$$



Explore upper/lower bounds (in combination with emergent string conjecture) and relative bounds in the asymptotic regime

# Outlook

- Study subleading terms and multi-scale structure
- Investigate further top-down constructions
- Explore connections with S-Matrix Bootstrap constraints
- Generalize understanding of corrections to *minimal* Black Holes in neutral case
- ...

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# Thank you!

# Back-up slides

# The double EFT expansion: Decompactification Limits

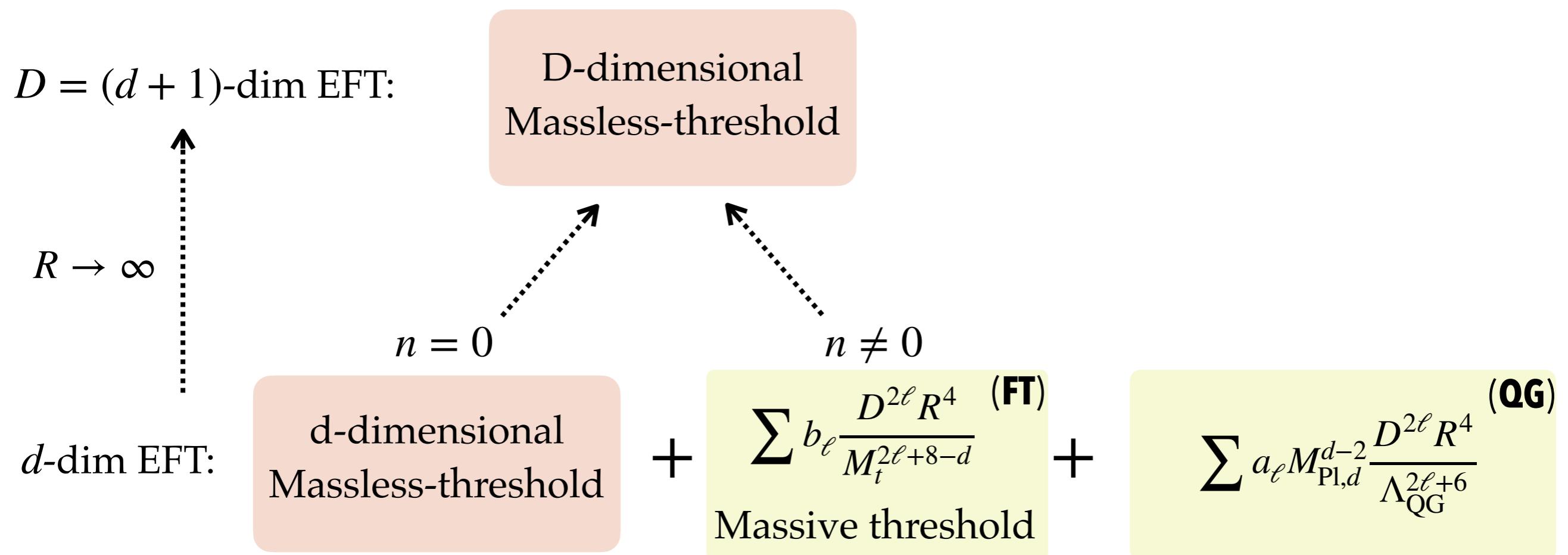
- 4-graviton amplitude: Non-vanishing contribution to an infinite tower of UV convergent operators of the form  $\frac{D^{2\ell}R^4}{M_t^{8+2\ell-d}}$   $\longrightarrow$  Change EFT at scales around  $M_t \leq \Lambda_{\text{QG}}$

- Massive vs. massless thresholds and different EFTs:  $\mathcal{A}_4(s, t) = \mathcal{A}_4^{\text{non-an}}(s, t) + \mathcal{A}_4^{\text{an}}(s, t)$   
 [Green, Gutperle, Kwan, Russo, Tseytlin, Vanhove '97-'07]  
 [Calderón-Infante, Castellano, AH '25]

$$\begin{aligned}
 & n = 0 & n \neq 0 \\
 d\text{-dim EFT: } & \text{d-dimensional Massless-threshold} & + \sum b_\ell \frac{D^{2\ell}R^4}{M_t^{2\ell+8-d}} \text{ (FT)} \\
 & + \sum a_\ell M_{\text{Pl},d}^{d-2} \frac{D^{2\ell}R^4}{\Lambda_{\text{QG}}^{2\ell+6}} \text{ (QG)}
 \end{aligned}$$

# The double EFT expansion: Decompactification Limits

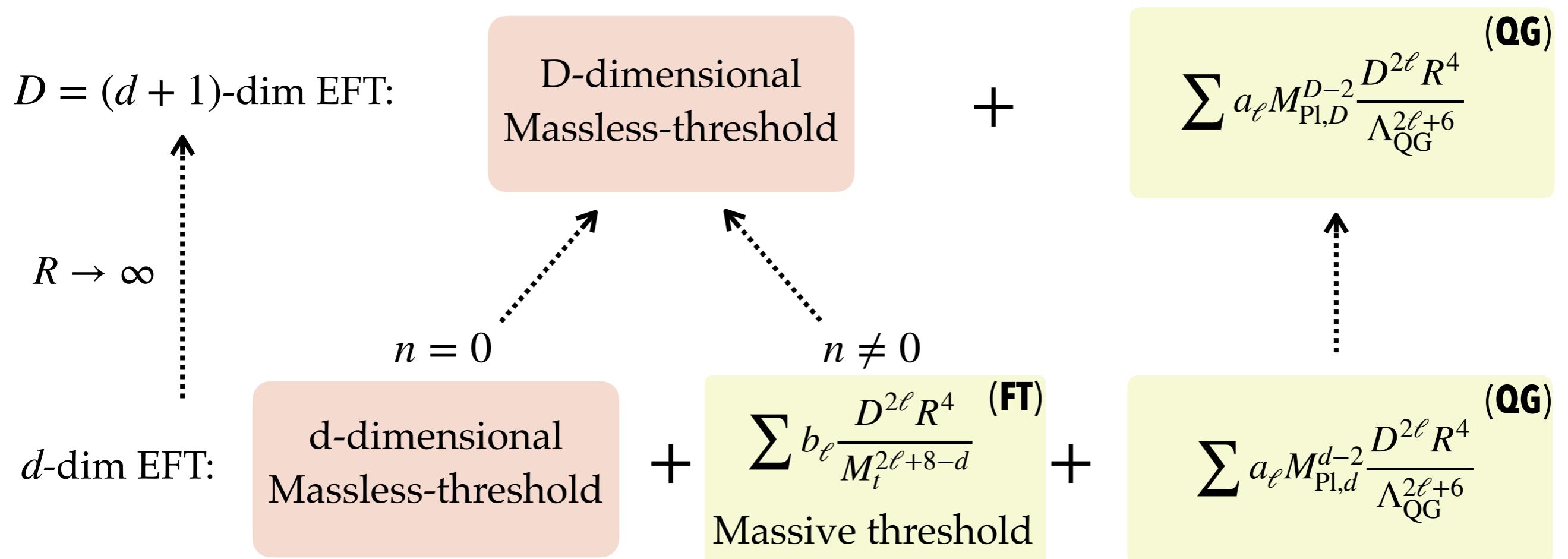
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# The Bottom-up Perspective: Absolute bounds in terms of EFT cutoff

$$\alpha_n = \frac{a_n}{\Lambda_{\text{QG}}^{n-2}} + \frac{b_n}{M^{n-d}} + \dots \quad (\text{In Planck Units})$$

*Asymptotic Regime:  $M_{\text{Pl},d} \gg \Lambda_{\text{QG}} \gtrsim M$*

- Upper bounds in terms of  $M$  as  $\alpha_n \rightarrow \infty$ . Both terms bounded by:

$$|\alpha_n| \lesssim \frac{1}{M^{n-2}}$$

- Lower bounds in terms of  $M$  as  $\alpha_n \rightarrow \infty$  (assuming non-vanishing coefficient):

For  $n \leq d$

**(FT)**  $|\alpha_n| \gtrsim \frac{1}{M^{n-d}}$

Emergent String  
Conjecture ( $p = 1$ ):

$$M_t^{\frac{p}{d+p-2}} \gtrsim \Lambda_{\text{QG}}$$

For  $n > d$

**(QG)**  $|\alpha_n| \gtrsim \frac{1}{\Lambda_{\text{QG}}^{n-d}} \gg \mathcal{O}(1)$

For  $n > d + 1$

**(QG)**  $|\alpha_n| \gtrsim \frac{1}{M_t^{\frac{n-2}{d-2}}}$

For  $n \leq d + 1$

**(FT)**  $|\alpha_n| \gtrsim \frac{1}{M_t^{n-d}}$

+

# The Bottom-up Perspective: Comparison with S-Matrix Bootstrap bounds

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- 1) Absolute bounds in terms of EFT cutoff:

[Caron-Huot, Mazac, Rastelli, Simons-Duffin '21] [Bern, Kosmopoulos, Zhiboedov '21]  
2x[Caron-Huot, Li, Parra-Martinez, Simons-Duffin '22] [Albert, Knop, Rastelli '23] [Caron-Huot, Li '24]

Compatible with Double EFT bounds upon  
identification of EFT cutoff with  $M$  (or  $M_t$ )