

MAX PLANCK INSTITUTE
FOR PHYSICS



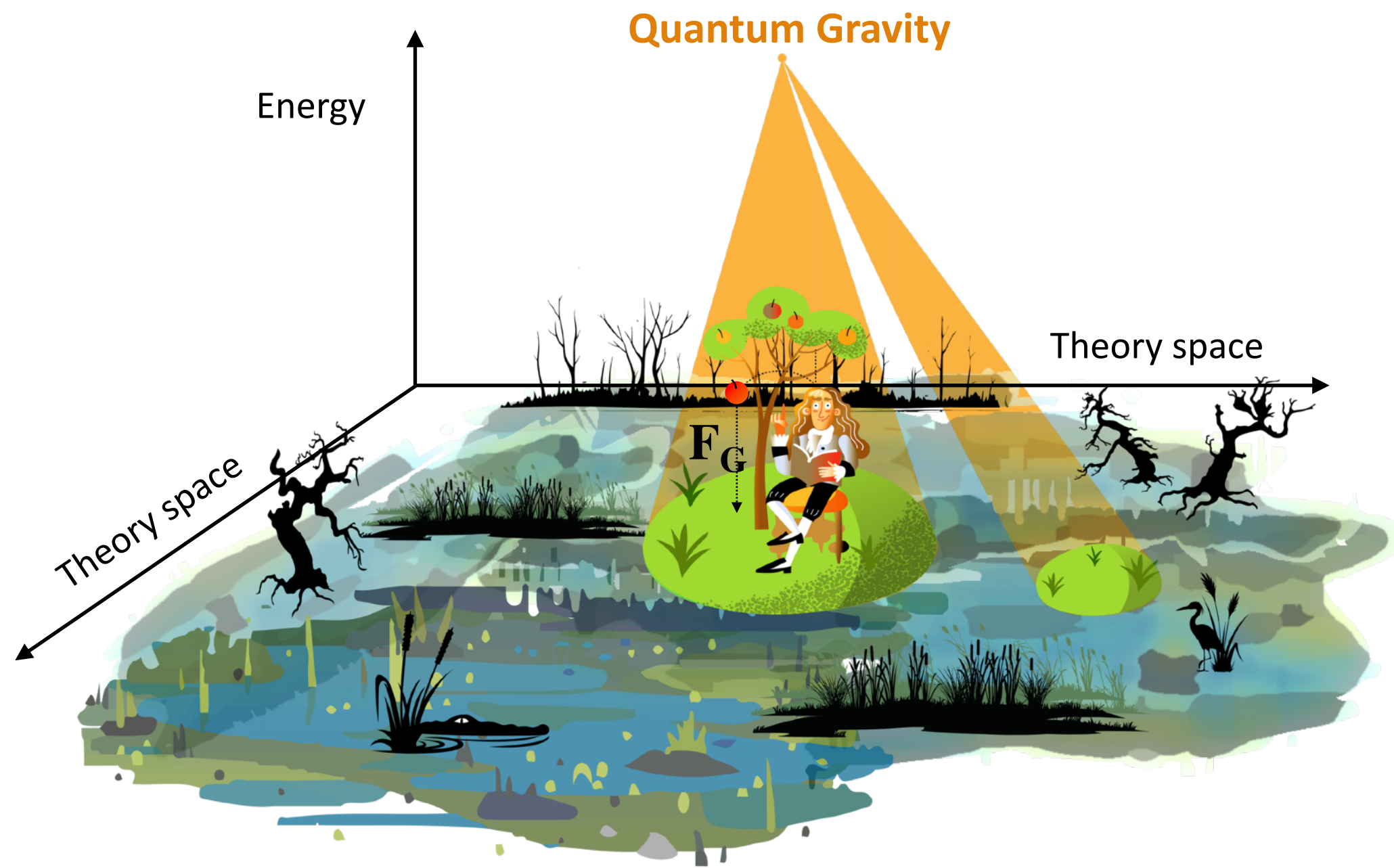
Infinite distances, the scalar potential and Ricci flow

[Saskia Demulder, Dieter Lust, TR 2312.07674 & ongoing]

24TH HELLENIC SCHOOL AND WORKSHOPS ON ELEMENTARY PARTICLE PHYSICS AND GRAVITY
Quantum Gravity, Strings and the Swampland
Corfu, 5.09.2024

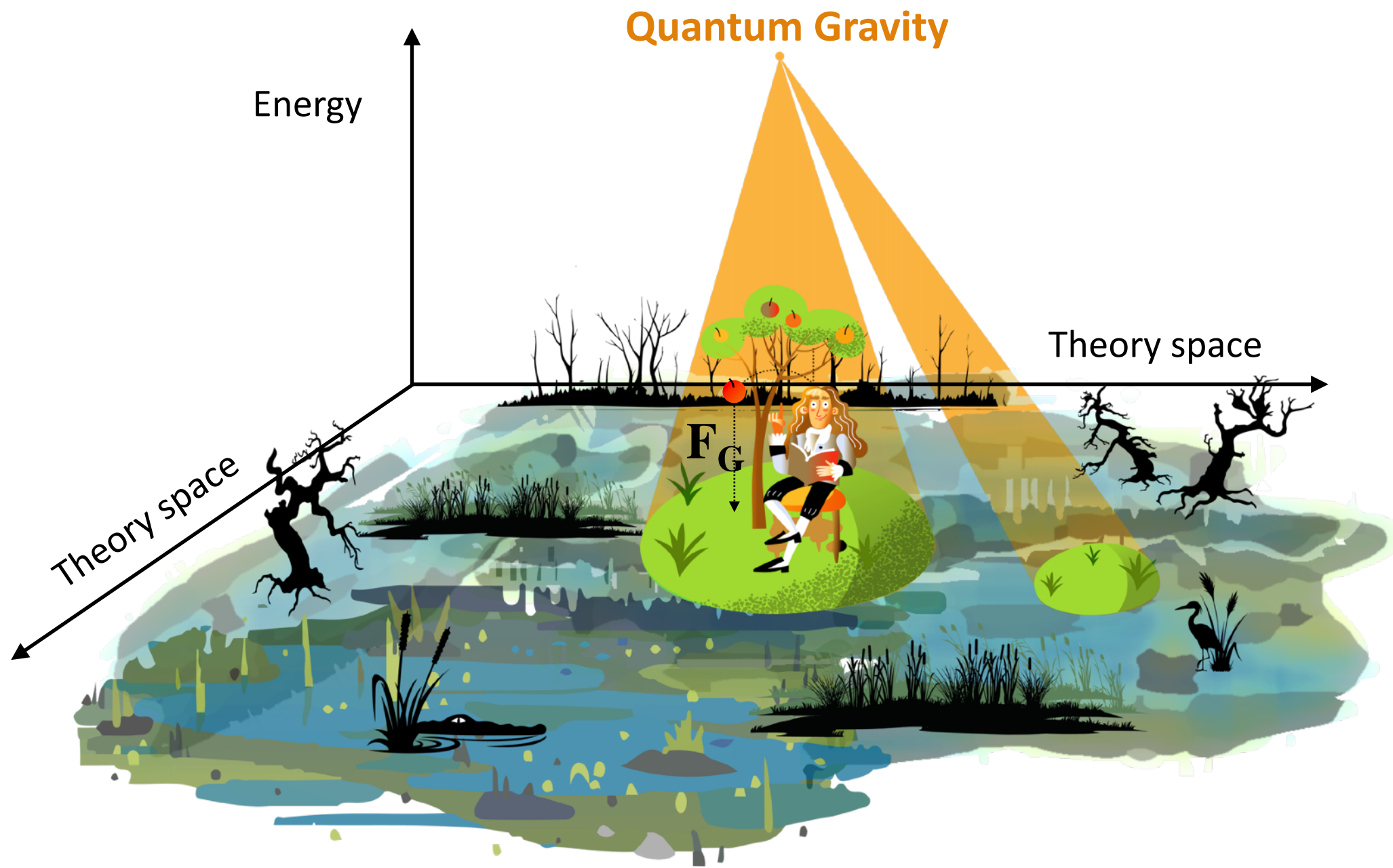
Thomas Raml

Quantum Gravity and the Swampland



Quantum field theories that emerge in the low energy limit of a quantum gravity theory are **very special**

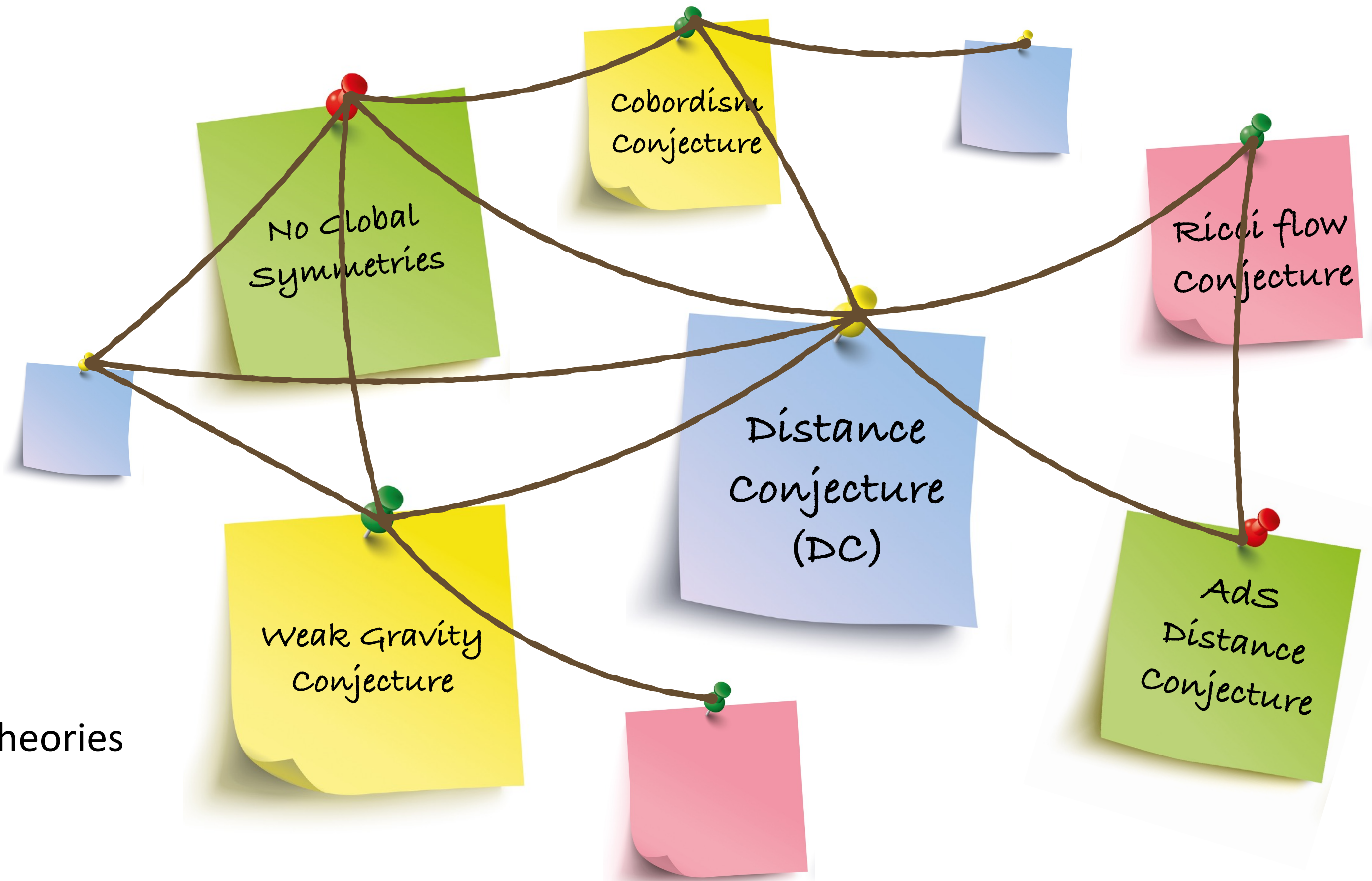
Quantum Gravity and the Swampland



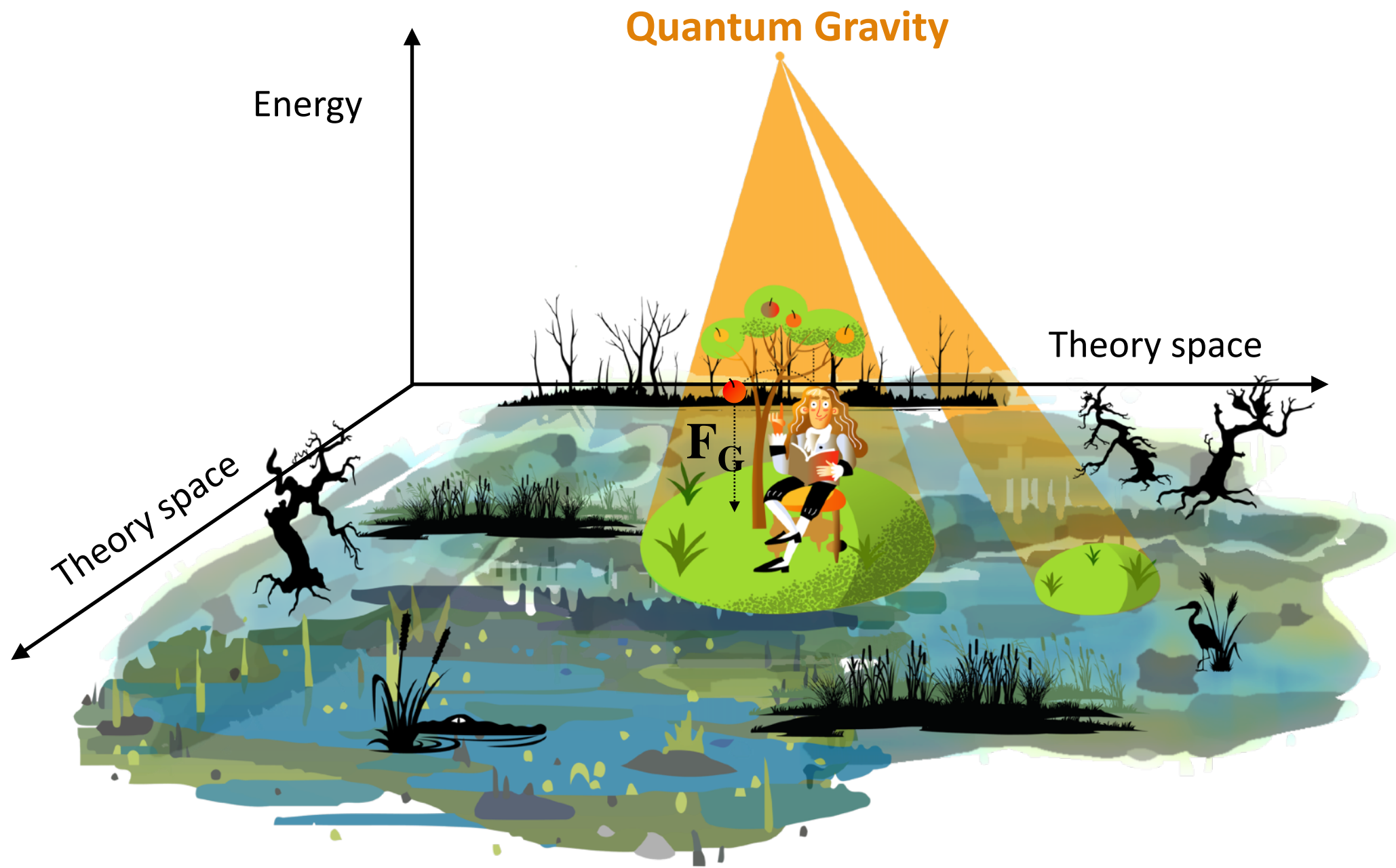
Swampland program:

formulate criteria that identify consistent quantum gravity theories by defining a set of constraining **conjectures**

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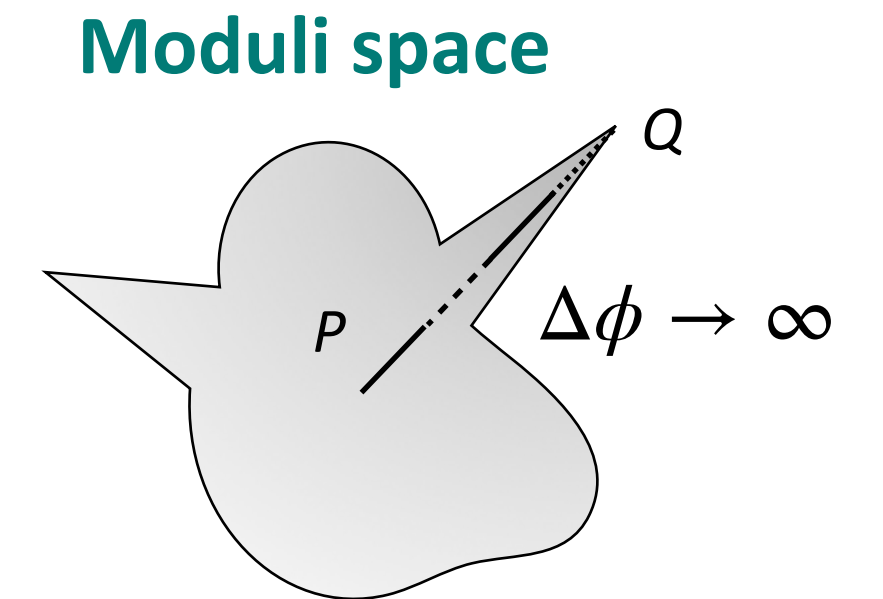


Recap: Distance Conjecture & S^1

In any consistent theory of quantum gravity: [Ooguri, Vafa '06]

when going to **large distances in its moduli space**,
we encounter an **infinite tower of states** which **become light** exponentially

$$M(Q) \sim M(P)e^{-\lambda\Delta\phi} \quad \text{when} \quad \Delta\phi \rightarrow \infty, \quad \Delta\phi \equiv d(P, Q)$$



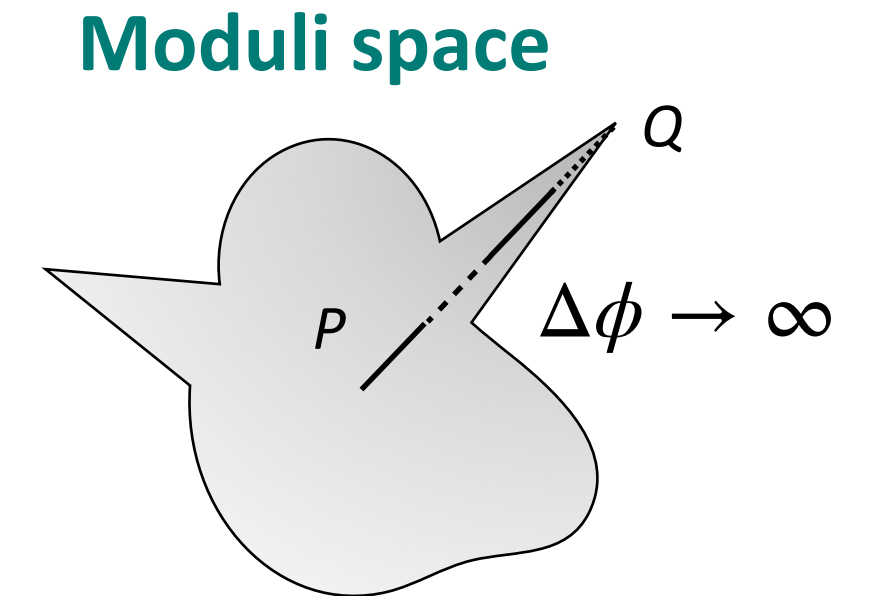
It describes the parameters
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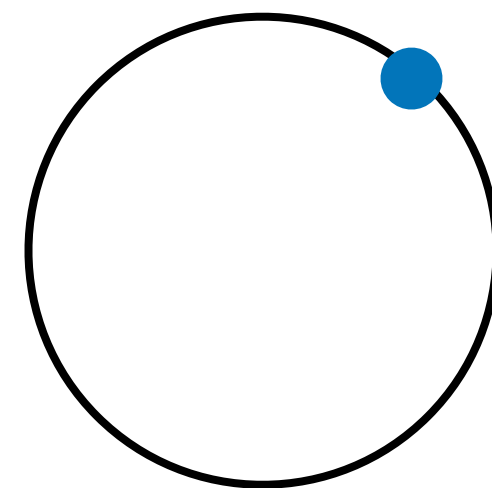
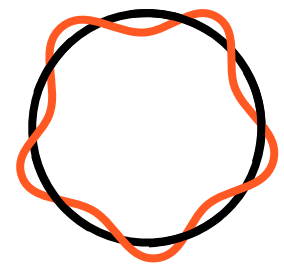
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Example: Circle compactification

$$S_{\text{EH}} \sim \int d^{D-1}x \sqrt{-g} \left(\mathcal{R}(g) - \frac{c}{R^2}(\partial R)^2 \right)$$



For $R \rightarrow \infty$
Infinite tower of
massless **KK**-modes

$$m_{\text{KK}}^2 \sim \frac{1}{R^2}$$

&

For $R \rightarrow 0$
Infinite tower of
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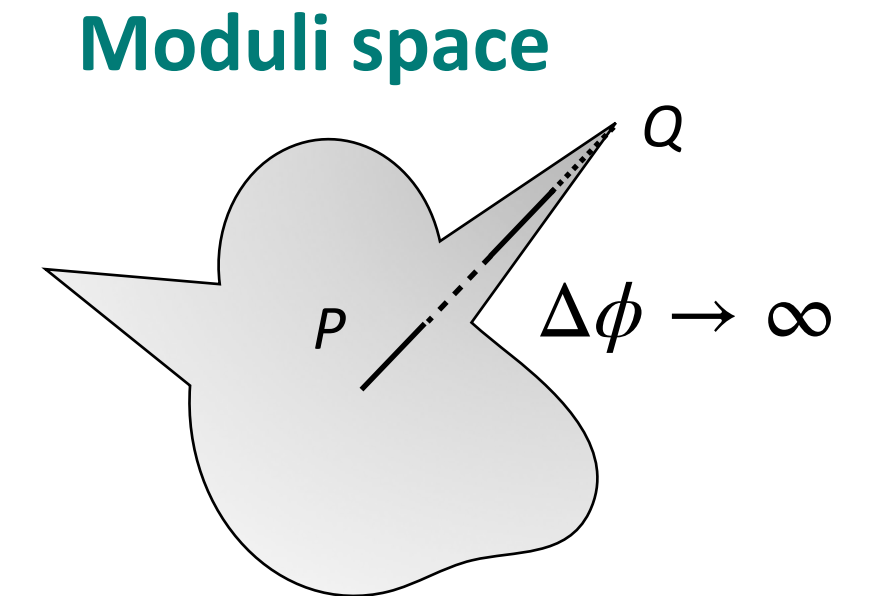
$$m_w^2 \sim R^2$$

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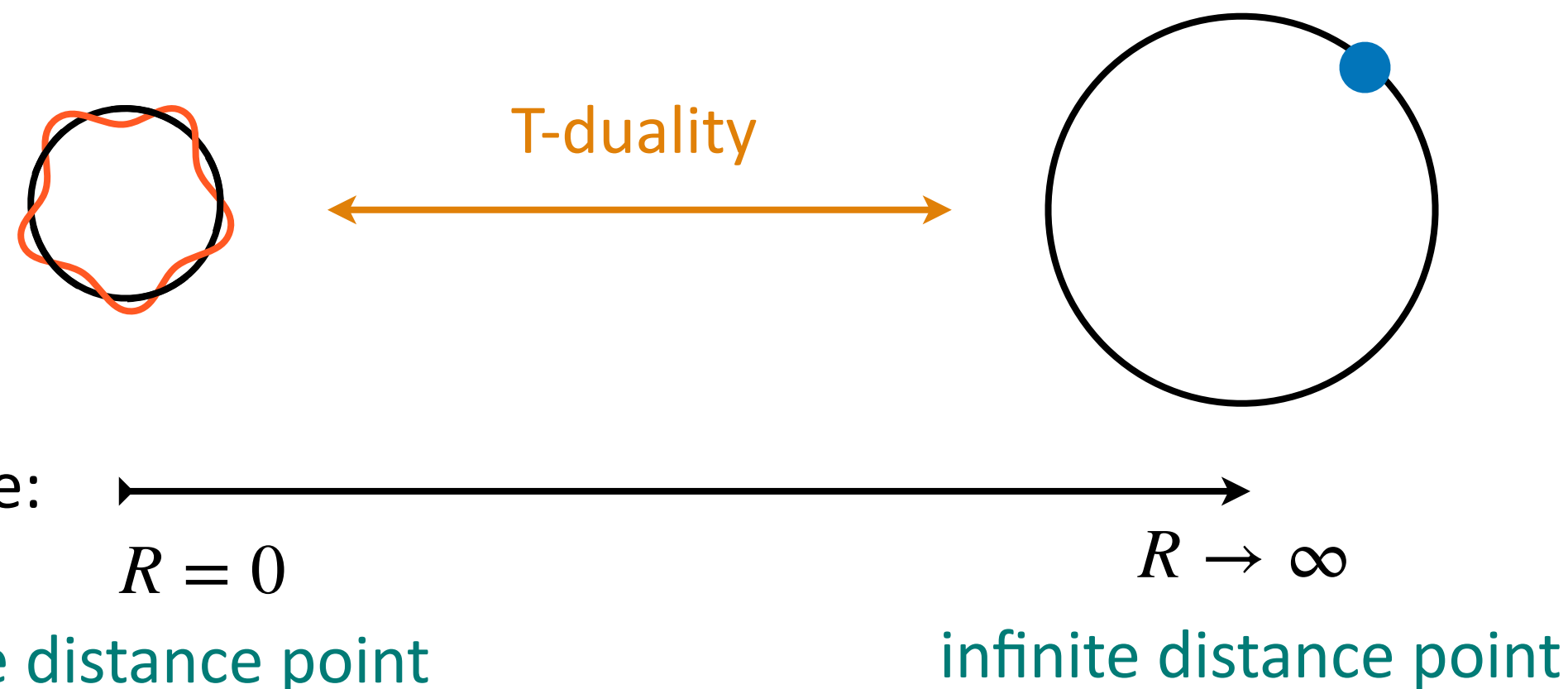
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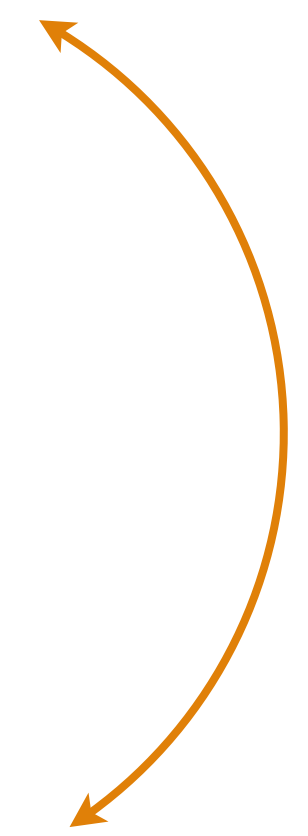
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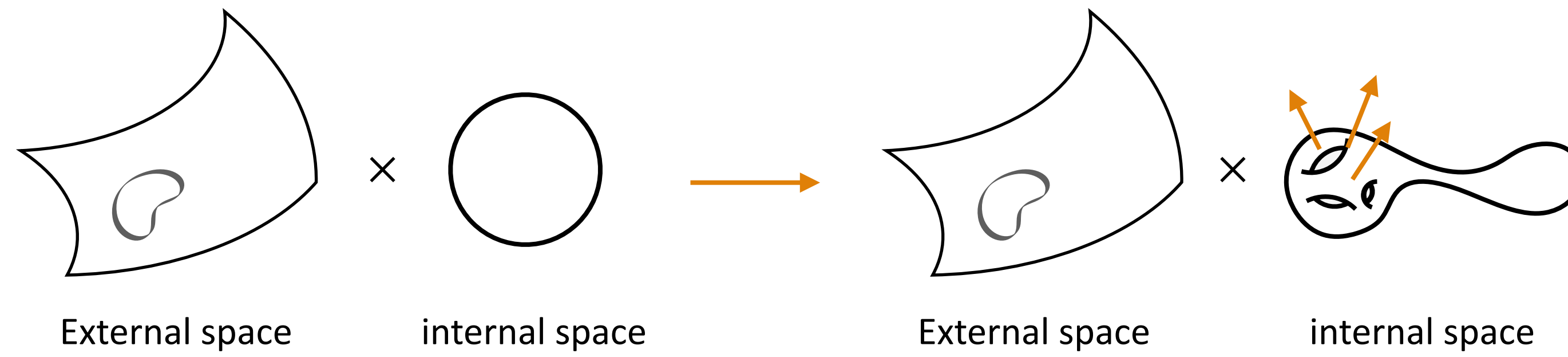
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T-duality

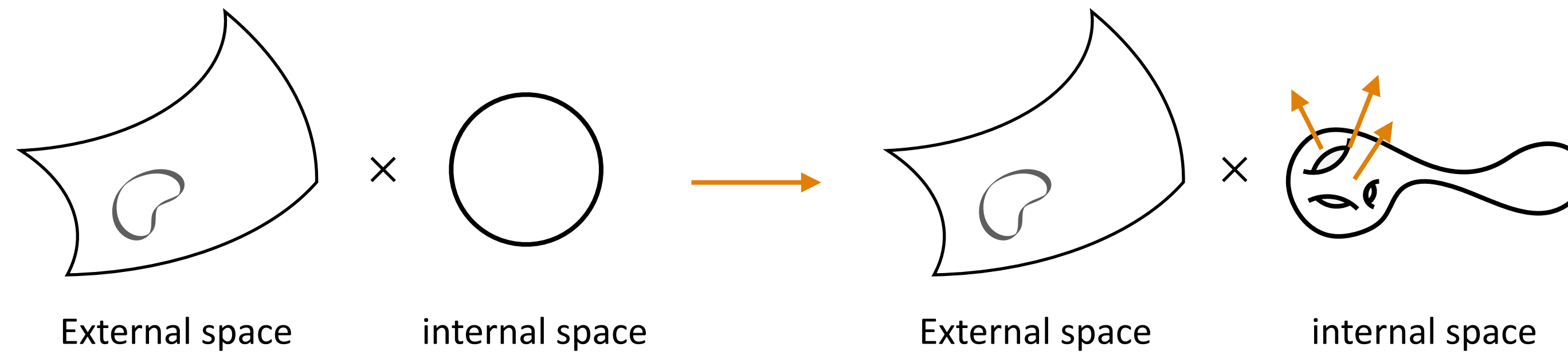


What about more complicated compact geometries?



A much more challenging question...

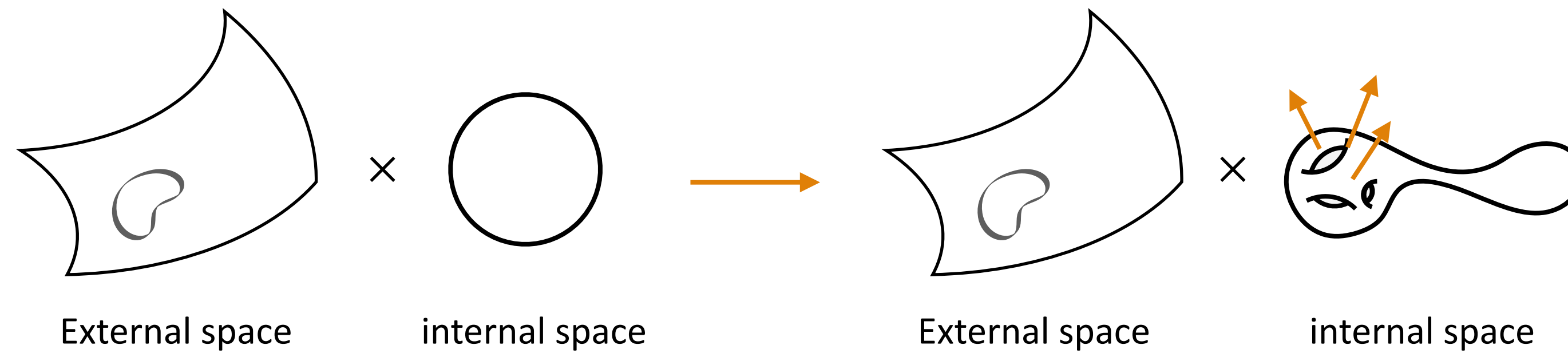
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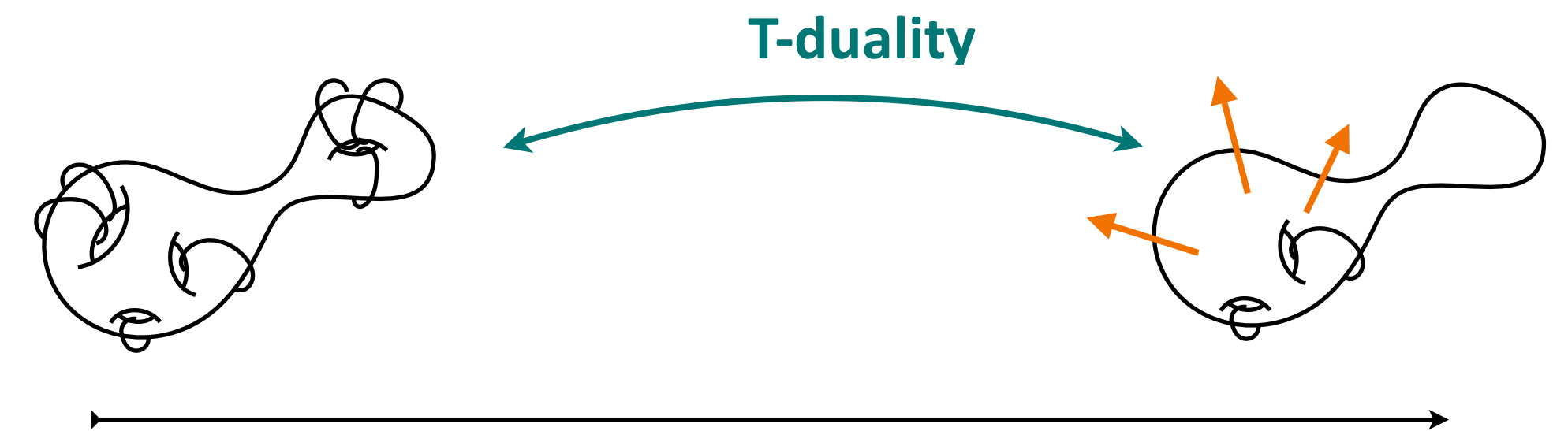


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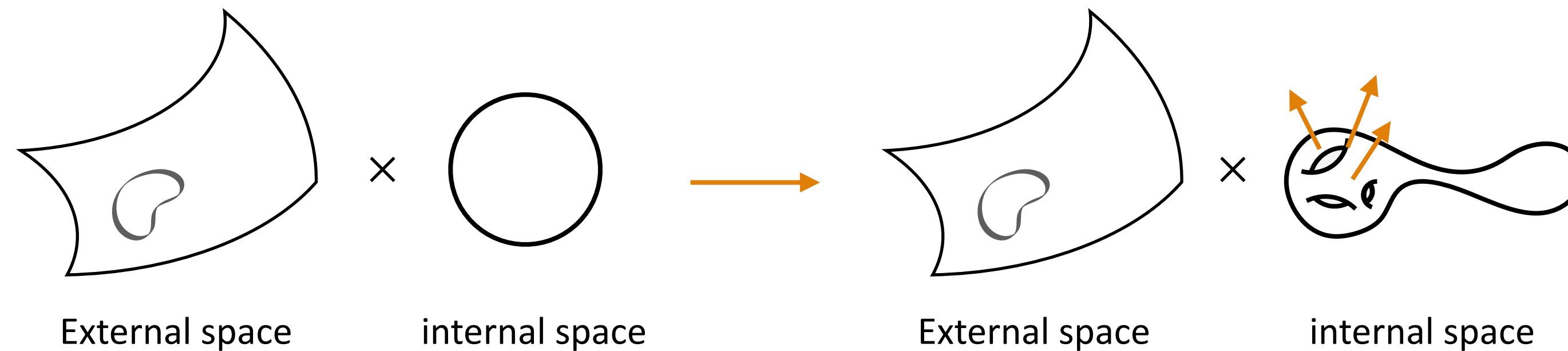
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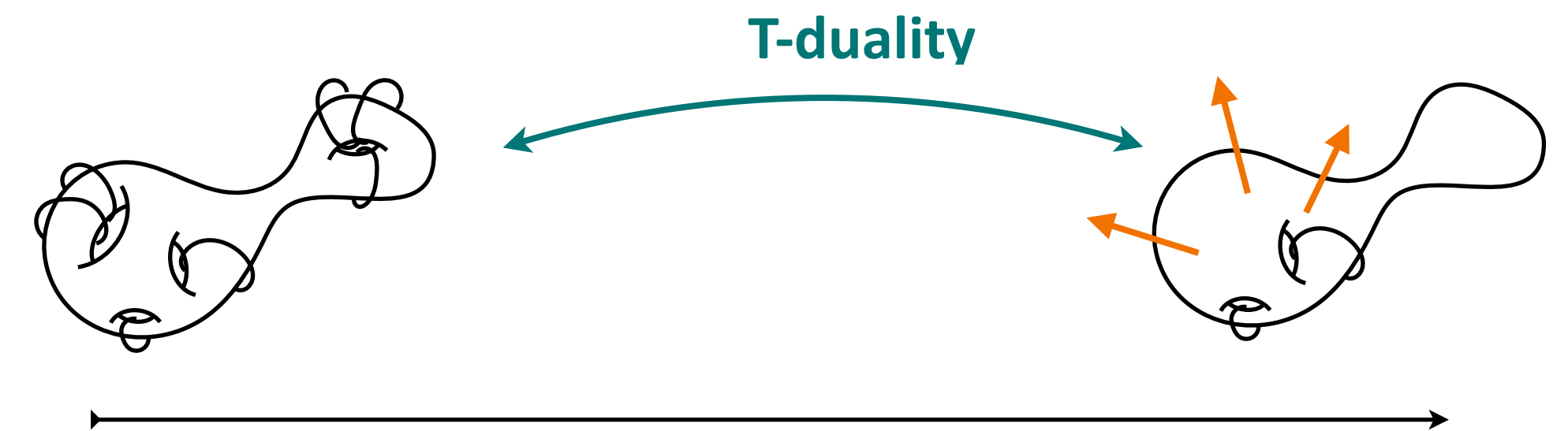
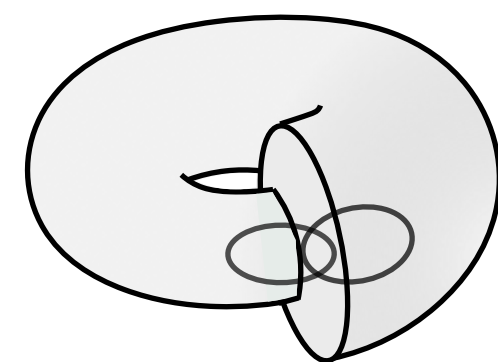
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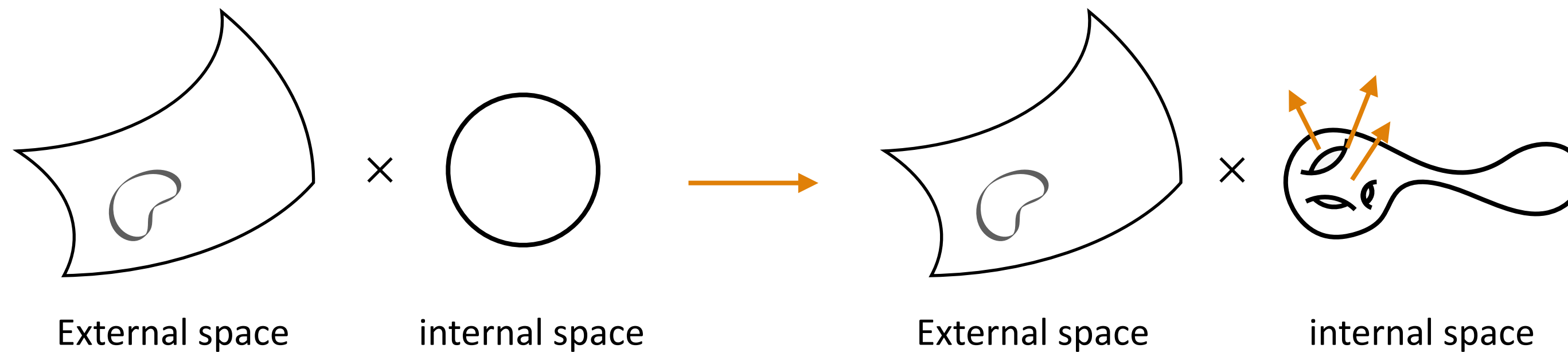
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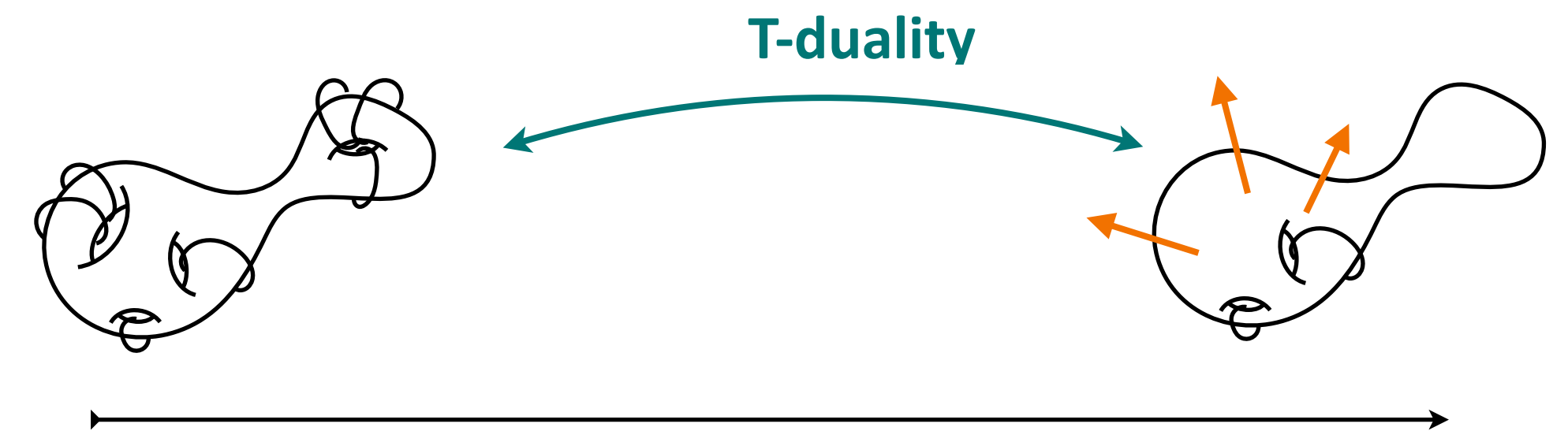
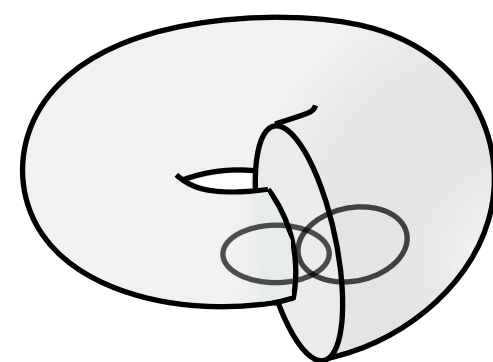
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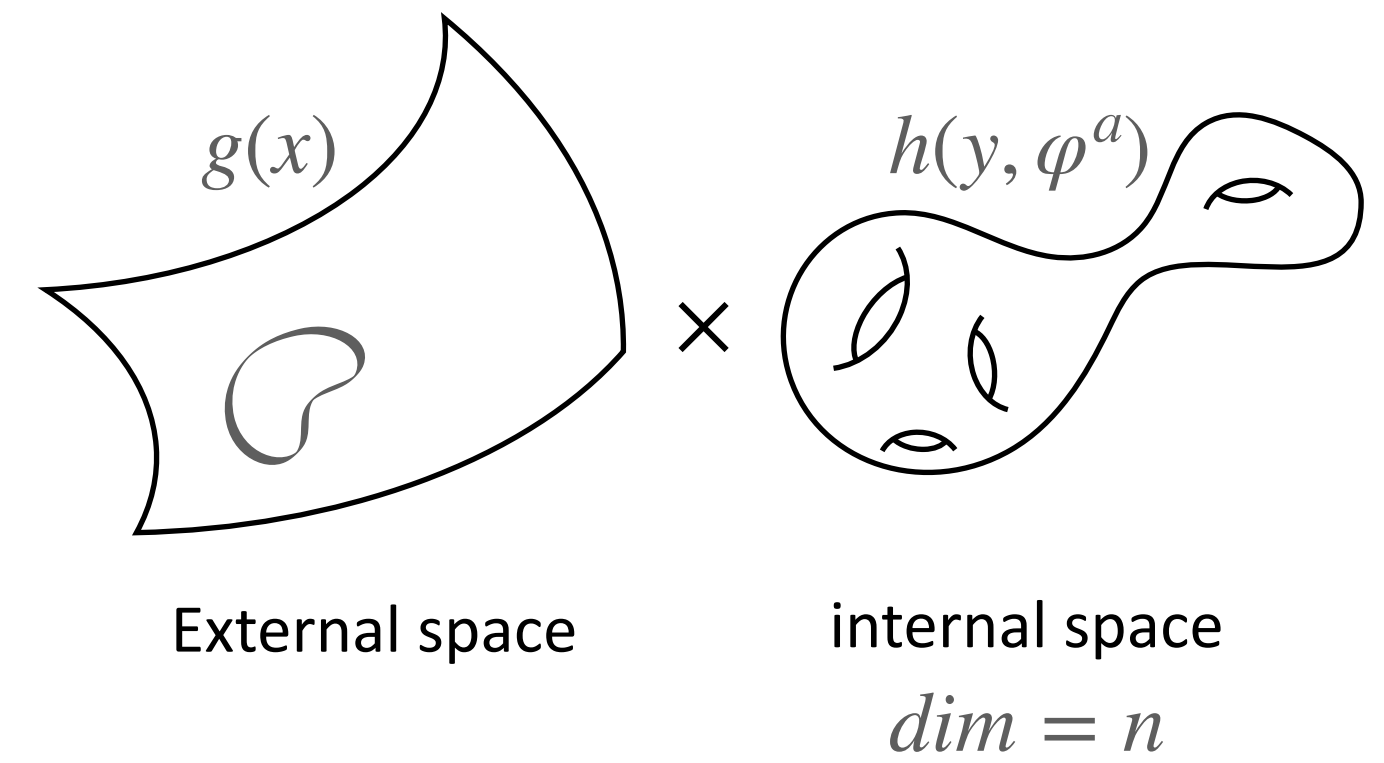
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Do these properties modify the Swampland Distance Conjecture?

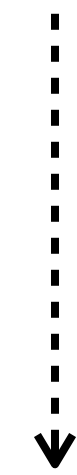
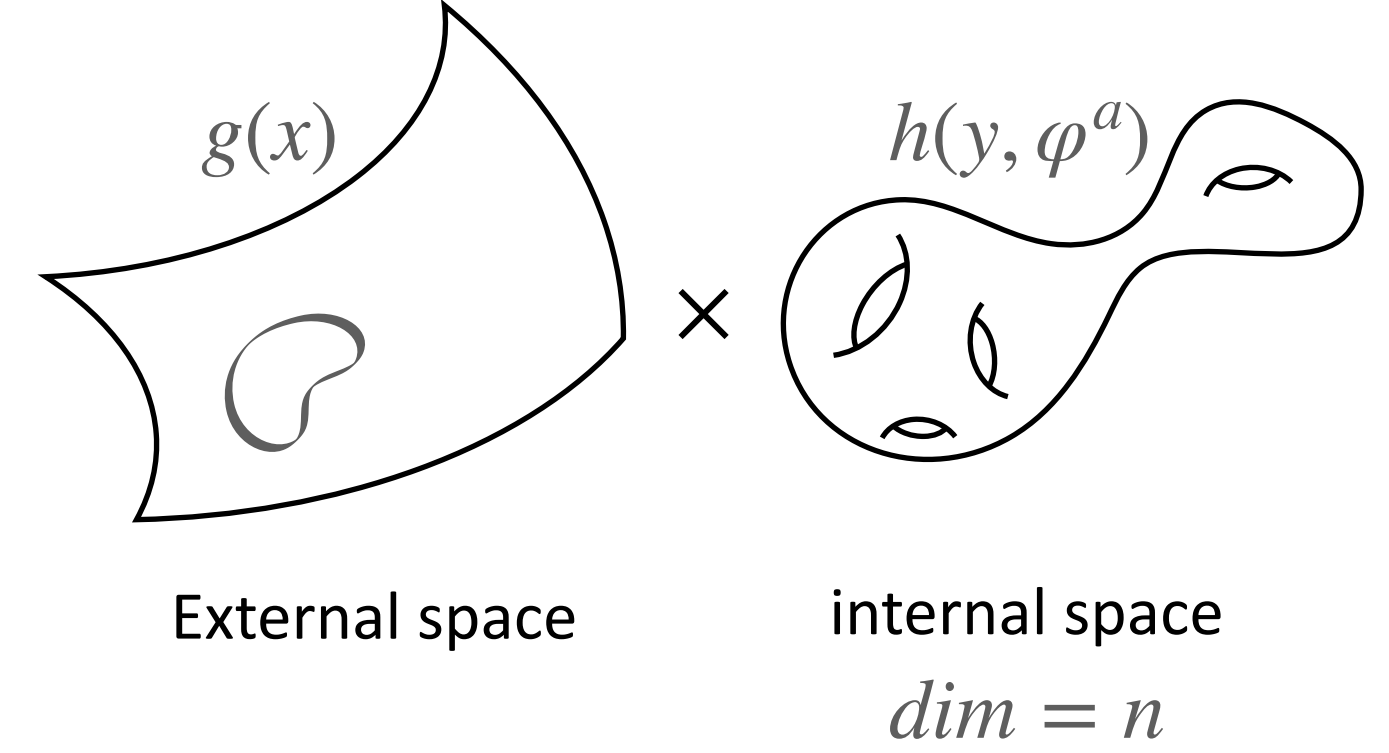
Generic setup & reduction

$$S = \frac{1}{2\kappa_0^2} \int d^D X \sqrt{-G} e^{-2\Phi} \left(\mathcal{R}(G) - \frac{1}{12} H_{IJK} H^{IJK} + 4\partial_I \Phi \partial^I \Phi \right)$$



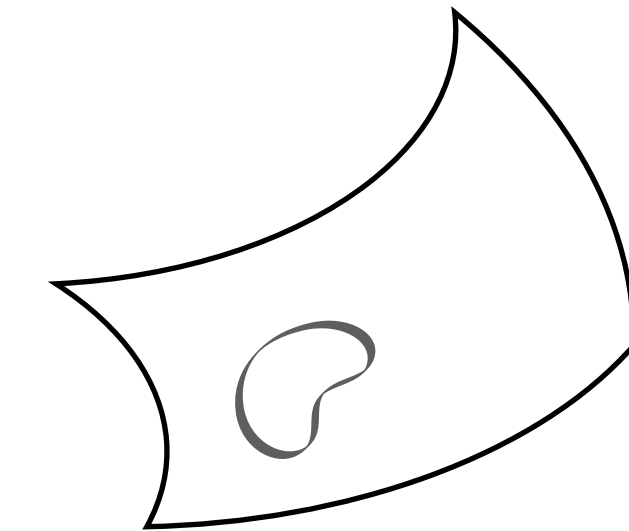
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$$G(x, y) = g(x) \oplus h(y, \varphi^a(x))$$

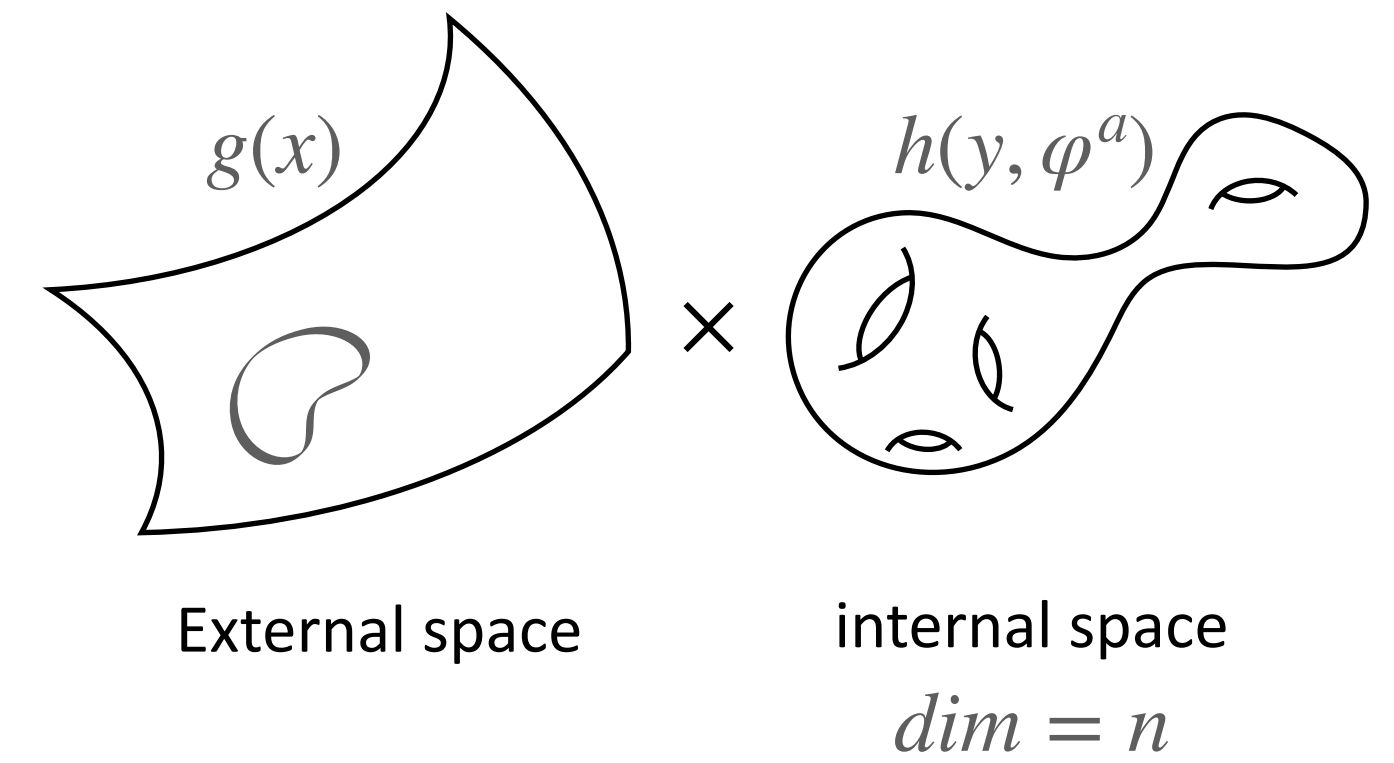
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Low-energy EFT on "External manifold"

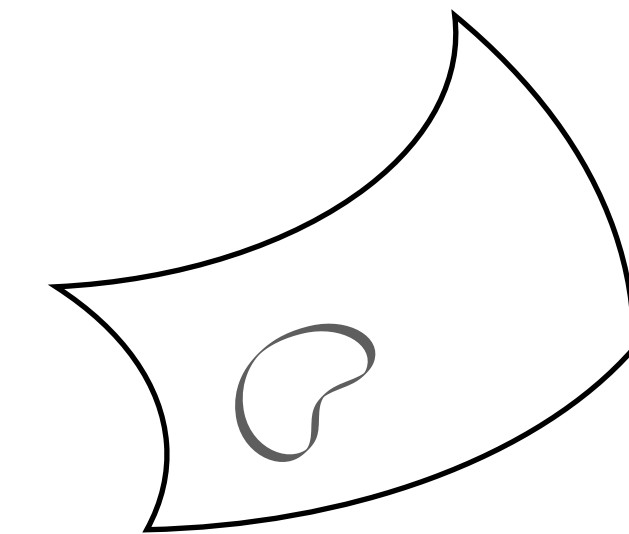
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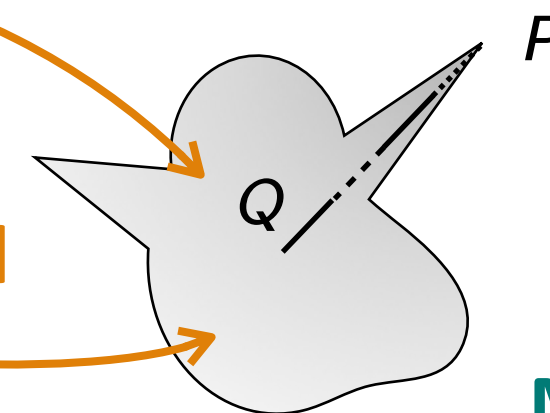
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metric

$$V(\varphi^i) \sim \int d^n y \sqrt{h} \left(\mathcal{R}(h) - \frac{1}{12} H_{ijk} H^{ijk} + 4\partial_i \Phi \partial^i \Phi \right)$$

potential



Moduli space

Example: S^3 with H -flux

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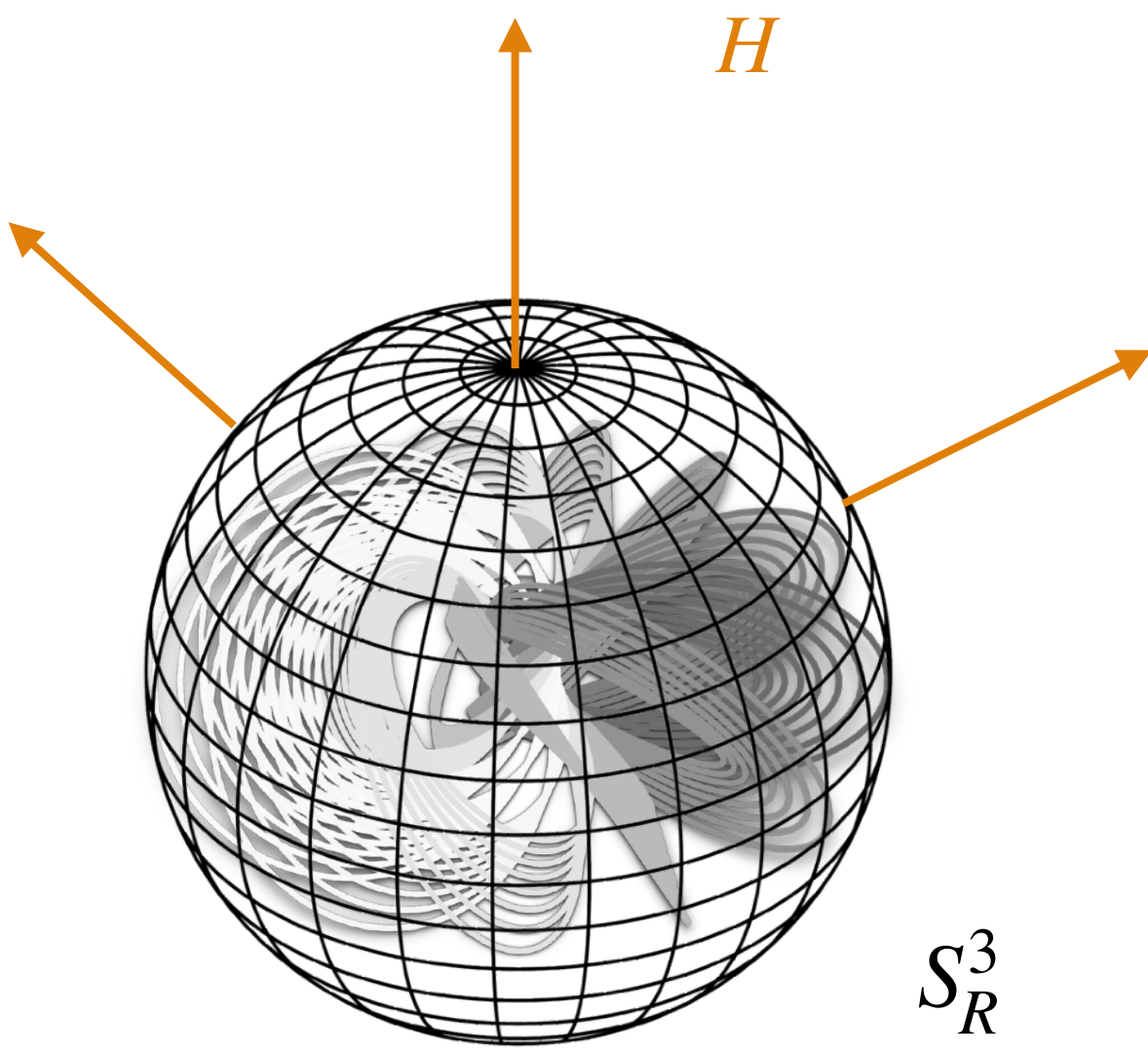


$$ds^2 = R^2(d\eta^2 + d\xi_1^2 + d\xi_2^2 + 2 \cos(\eta) d\xi_1 d\xi_2)$$

$$H = k \sin(\eta) d\eta \wedge d\xi_1 \wedge d\xi_2$$

$$\gamma_{RR} = \frac{3}{R^2}$$

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~~winding~~

$$\pi_1(S^3) = 0$$

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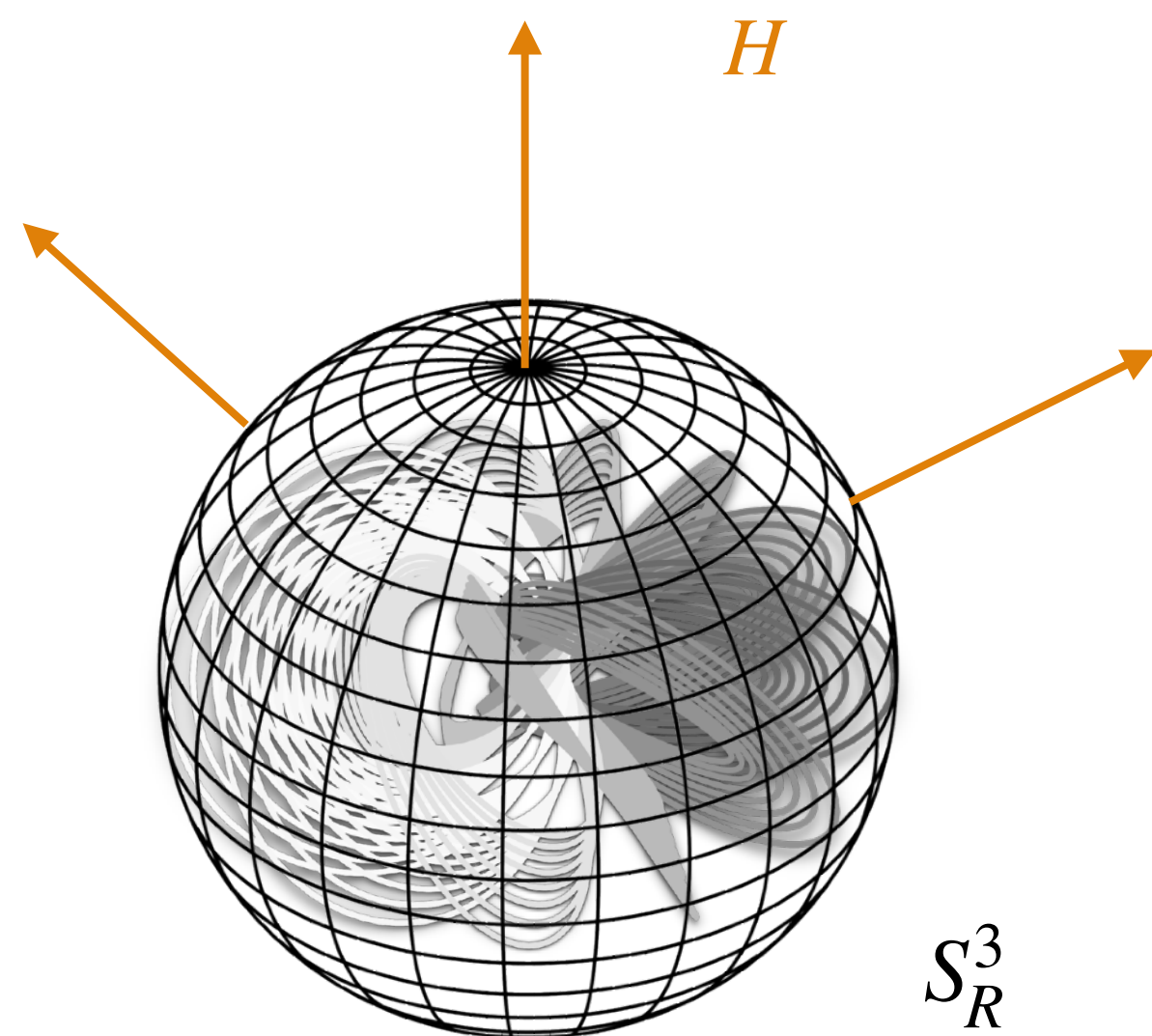
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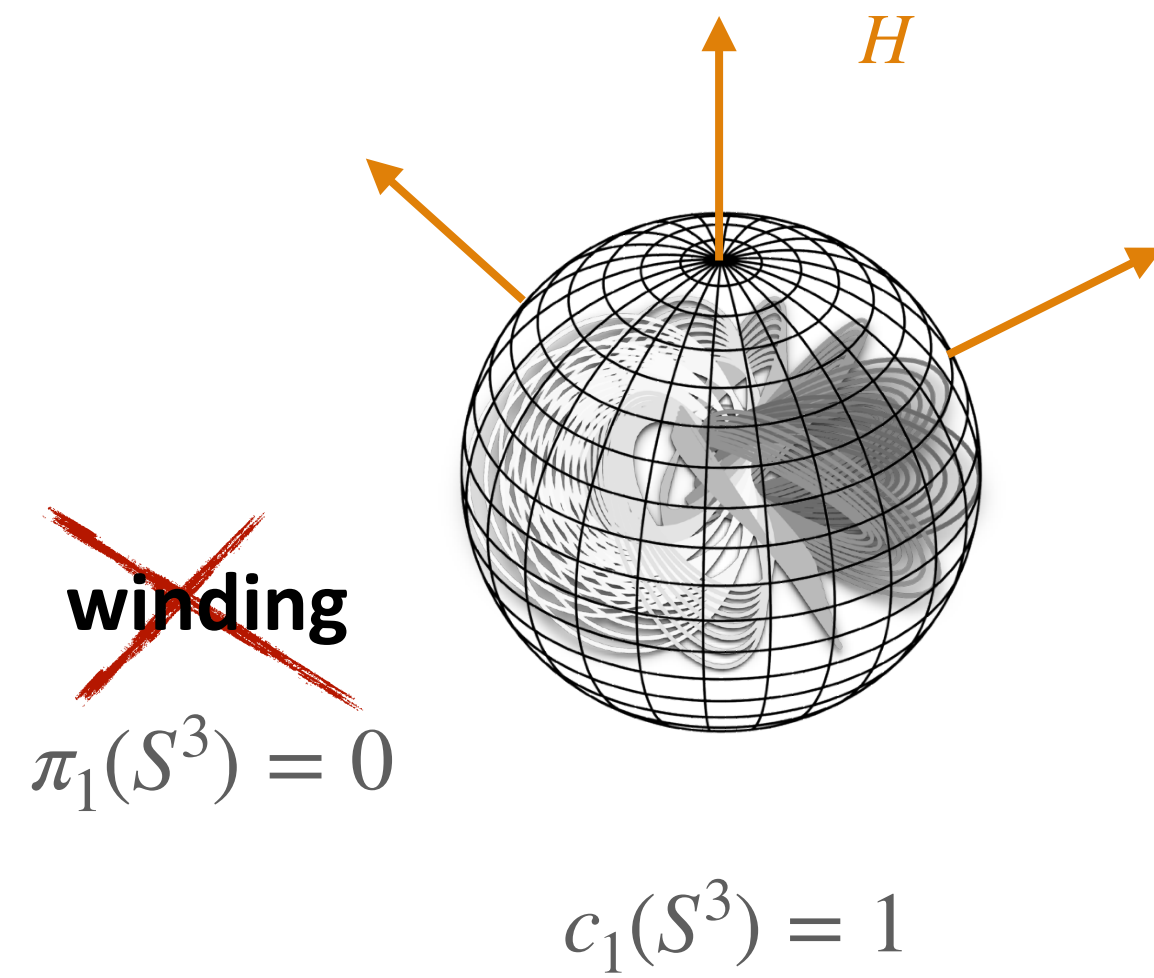
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How is **absence of winding modes** compatible with **T-duality**?

What does this mean for the **Swampland Distance Conjecture**?

► T-duality:

S^3_R with $[\mathbf{H}] = k$

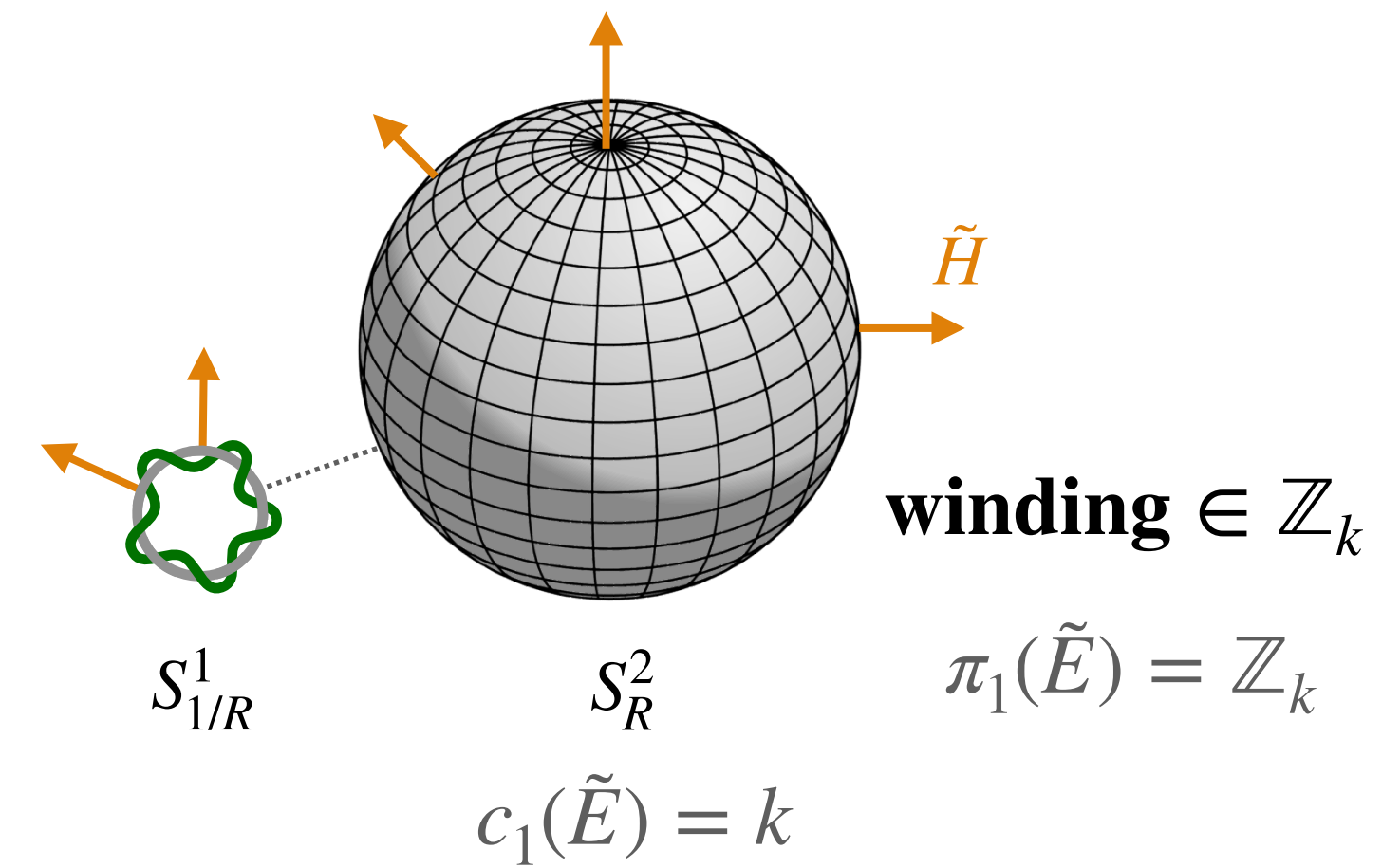


\longleftrightarrow
 T-duality
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 along Hopf-fiber

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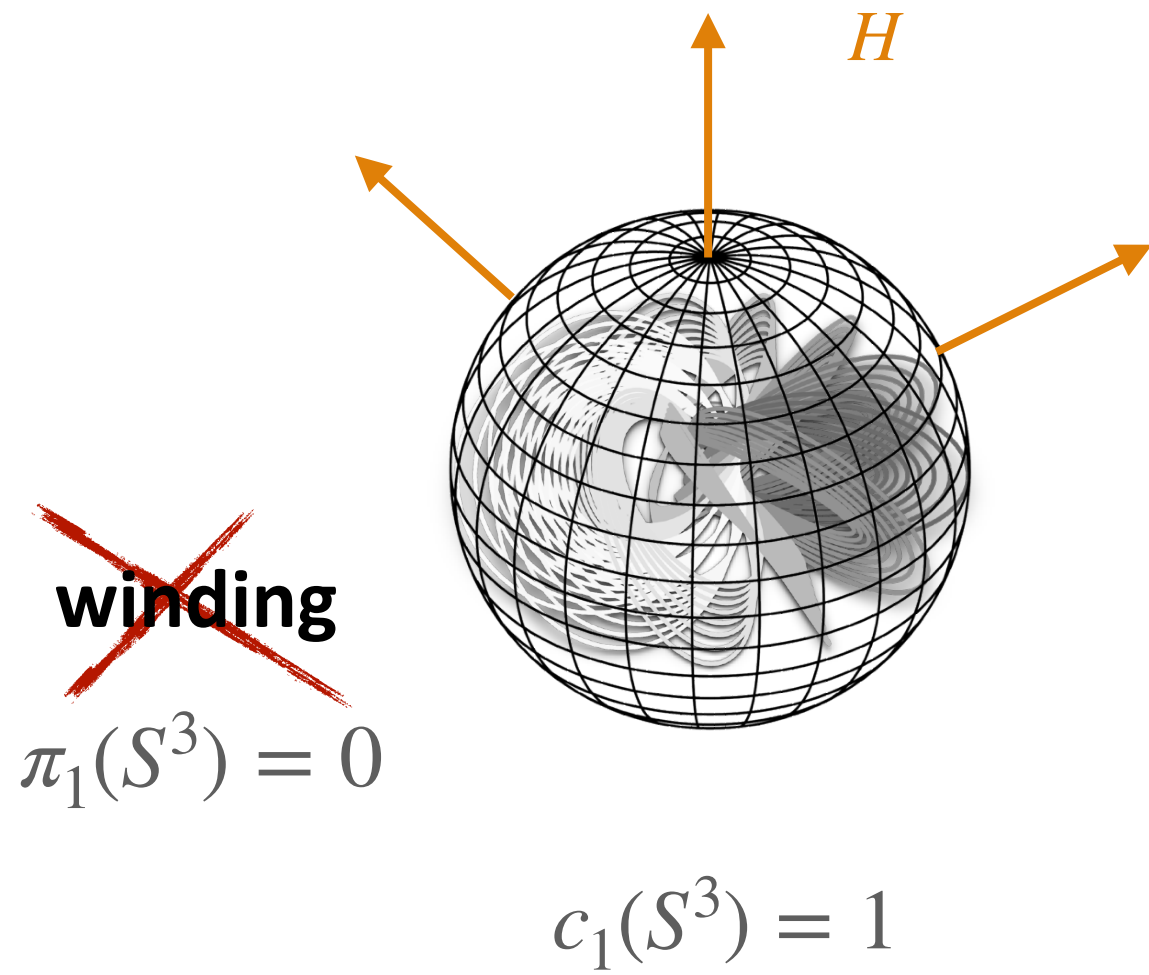
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$S^1_{1/R} \hookrightarrow \tilde{E}$
 \downarrow with $[\tilde{\mathbf{H}}] = 1$
 S^2_R



► T-duality:

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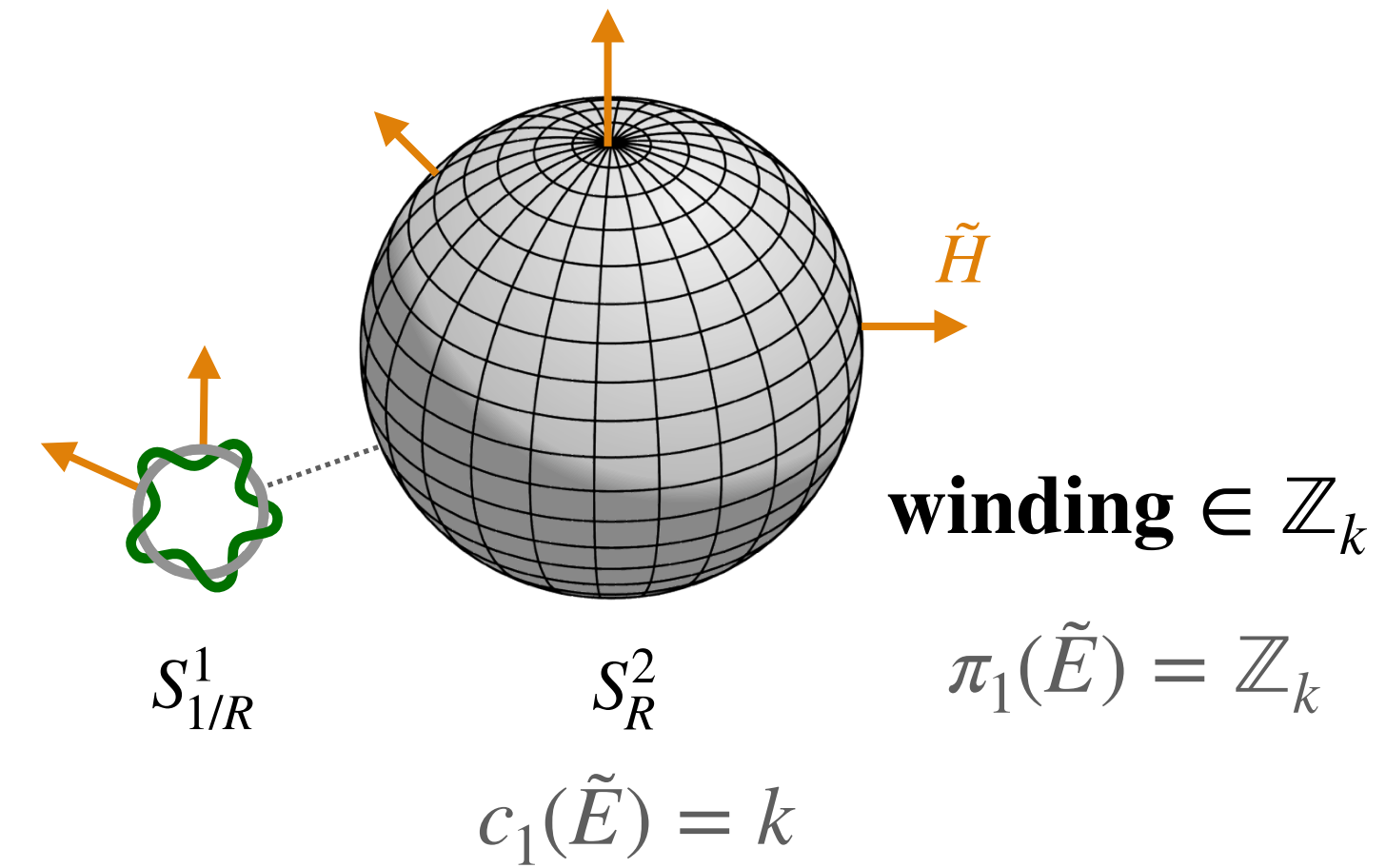
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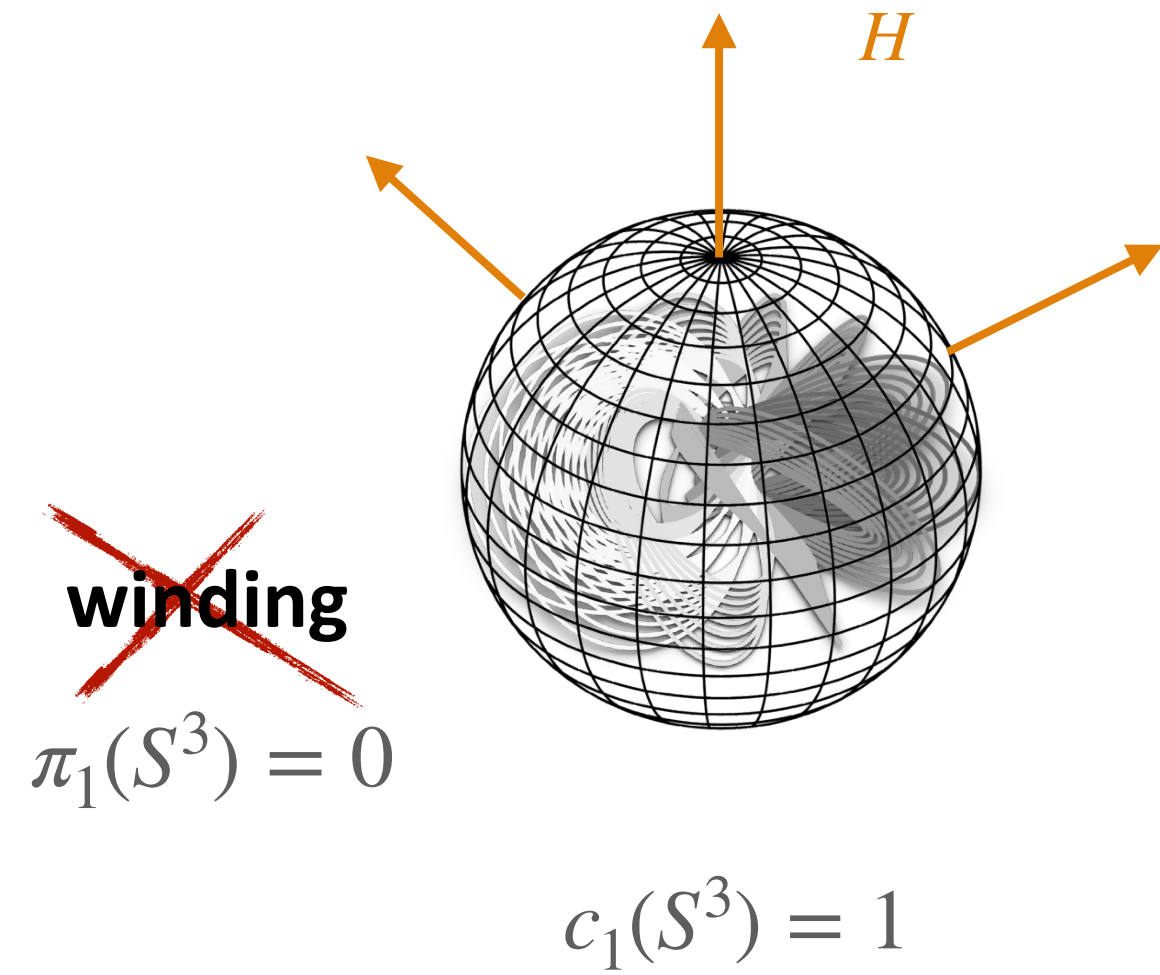
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	$R = 0$	$R \rightarrow \infty$
S^3 { winding momentum	\emptyset heavy	\emptyset light
\tilde{E} { winding momentum	\mathbb{Z}_k (heavy) heavy/ non-conserved	\mathbb{Z}_k (light) light
	no modes becoming light	tower of light states

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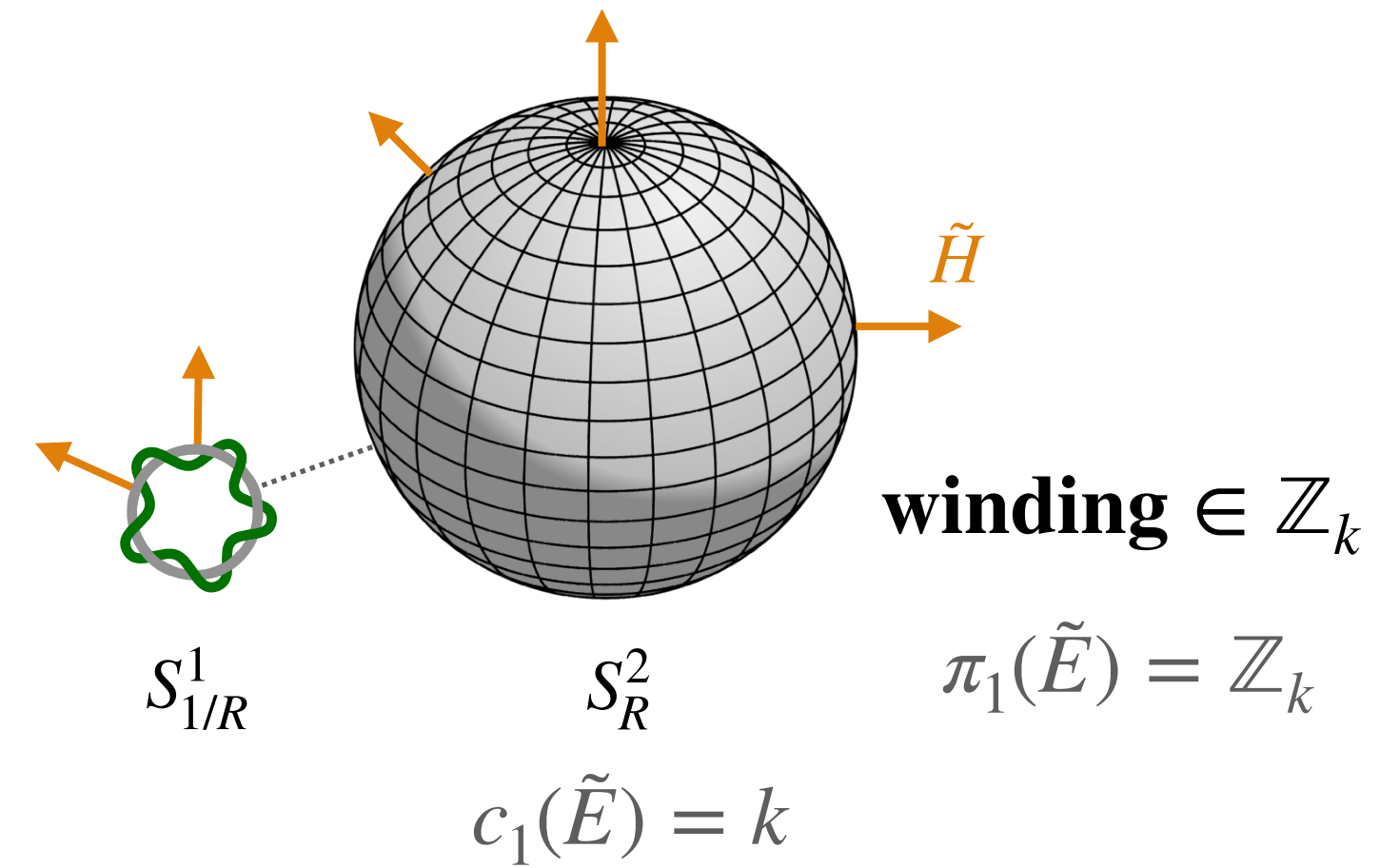
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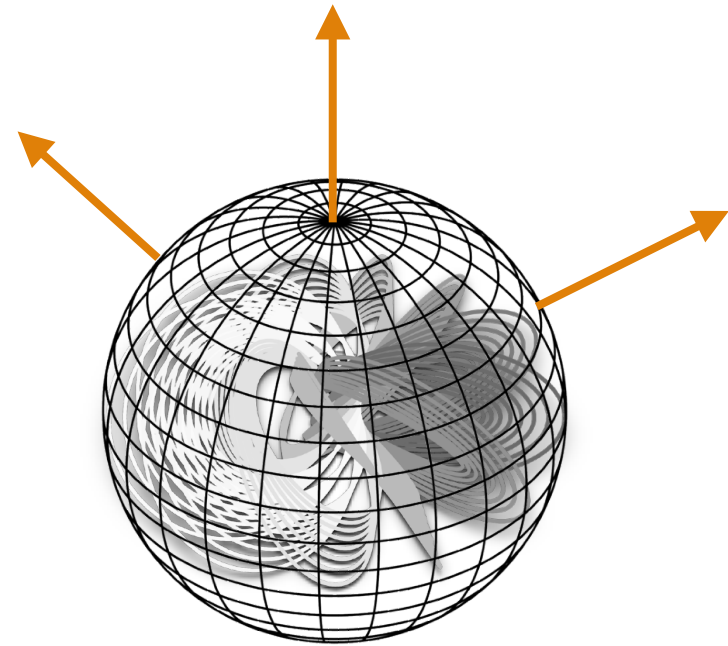
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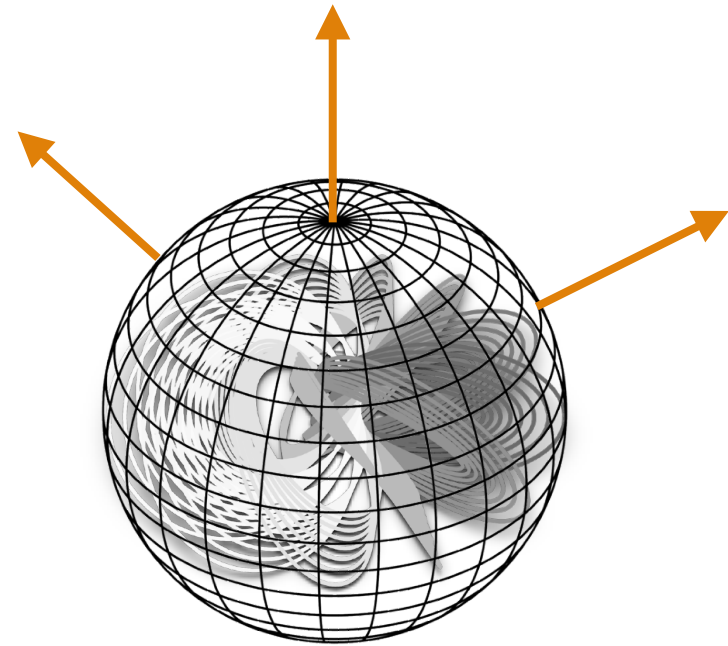
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Apparent inconsistency: S^3 with appropriately tuned H-flux is **valid string background** and therefore should be in the **Landscape**

However there is **no tower of light states** for $R \rightarrow 0$ which is an infinite distance limit

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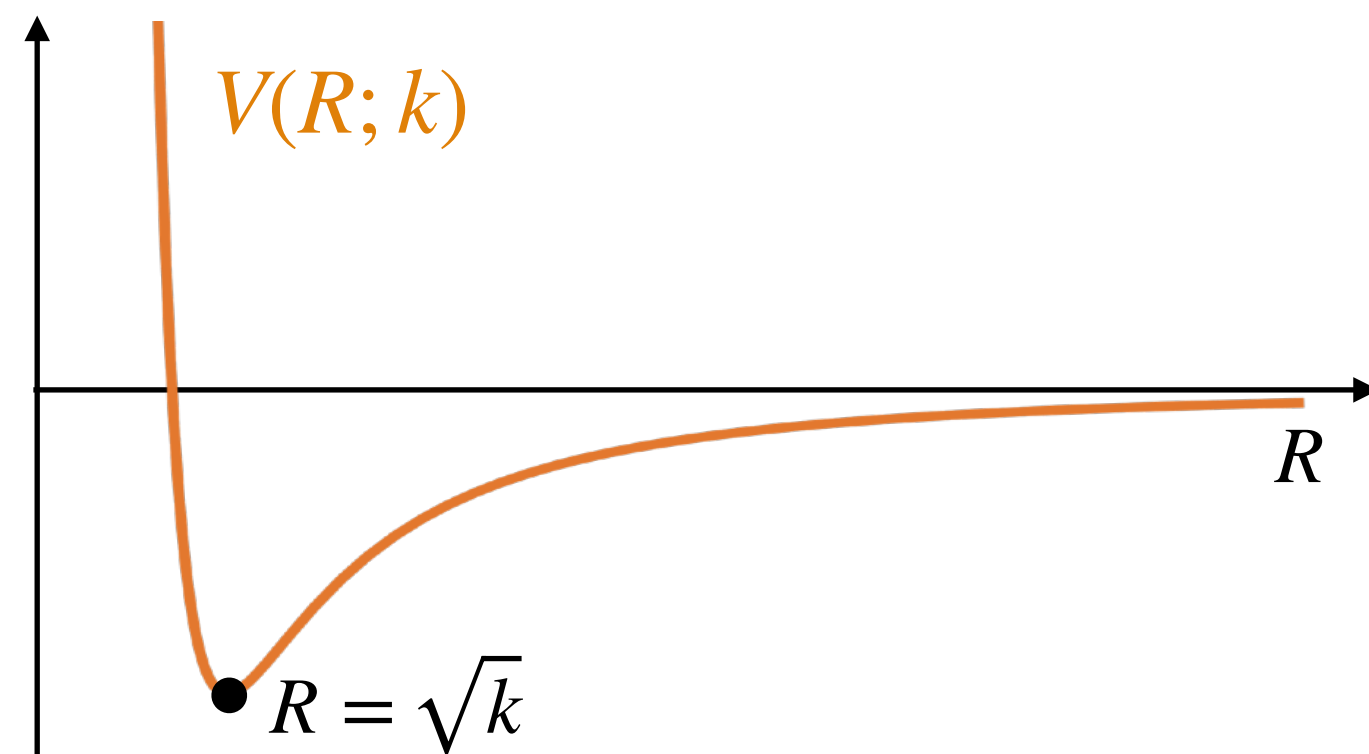


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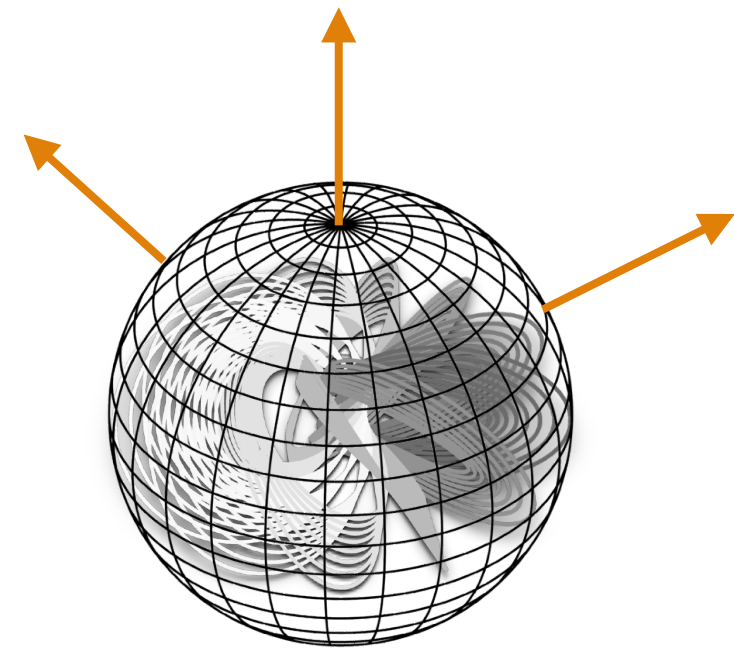
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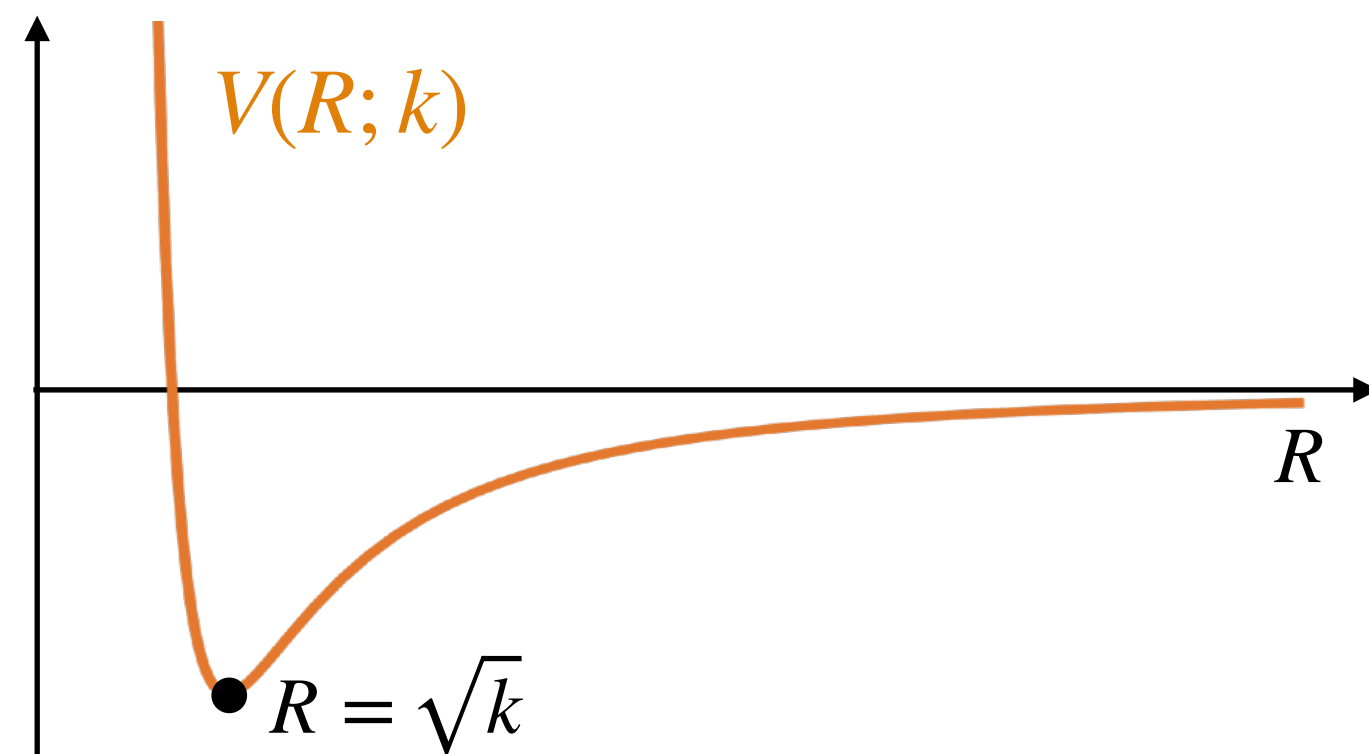
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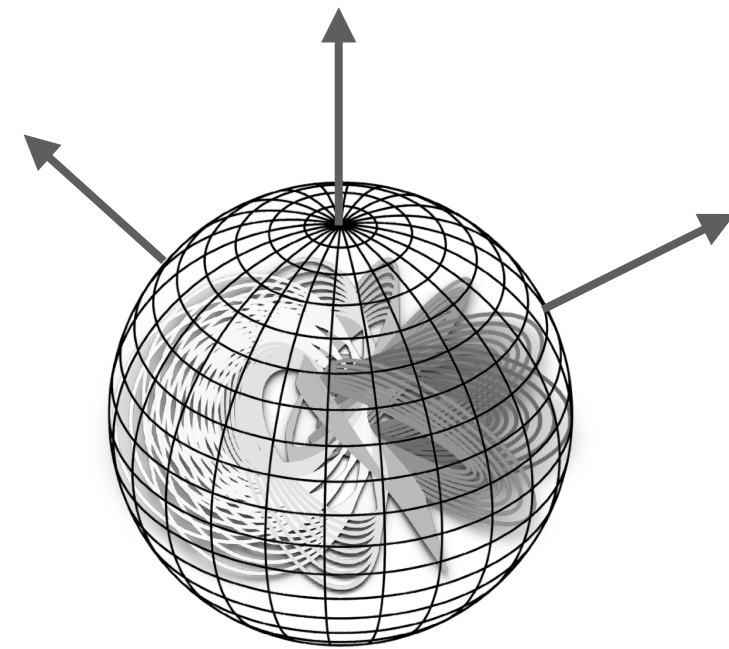
[Demulder, Lüster, TR '23]

In **effective field theories** that **can be lifted to a theory of quantum gravity** in the UV, a **divergence in the scalar potential** emerges when approaching an **infinite locus point** for which the target space geometry **cannot give rise to a light tower of states**.

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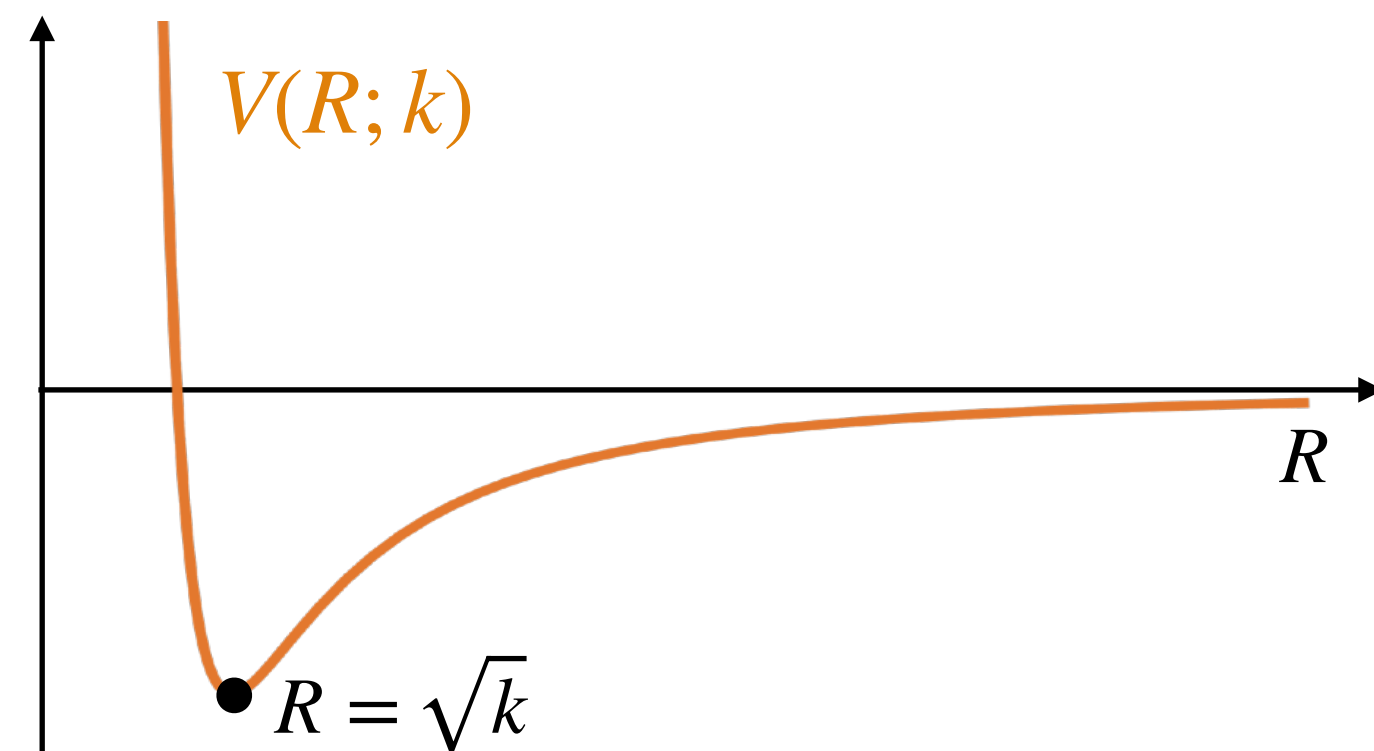
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► Invariance of metric & flux variations:

Metric on moduli space given by

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So by $O(d, d)$ invariance γ_{ab} is invariant under (abelian) **T-duality**.

[Demulder, Lüst, TR '23]

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► Example: S^3 with $k(x)=R^2(x)$ $\xleftrightarrow{\text{T-duality}}$ \tilde{E}

\tilde{E} ...modulus only in spacetime metric h

γ_{RR} obtained in standard way from “deWitt” metric



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$$H = R^2 \sin(\eta) d\eta \wedge d\xi_1 \wedge d\xi_2$$

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\tilde{E} ...modulus only in spacetime metric h } $S^3_{\mathbf{R}}$ with $[\mathbf{H}] = \mathbf{k} = \mathbf{R}^2$...modulus in h and B
 γ_{RR} obtained in standard way from "deWitt" metric } also contribution $\text{tr}(h^{-1} \partial_{\varphi_a} B h^{-1} \partial_{\varphi_b} B) \neq 0 \subset \gamma_{RR}$

$\tilde{\gamma}_{RR} = \gamma_{RR}$ only if flux variation are taken into account

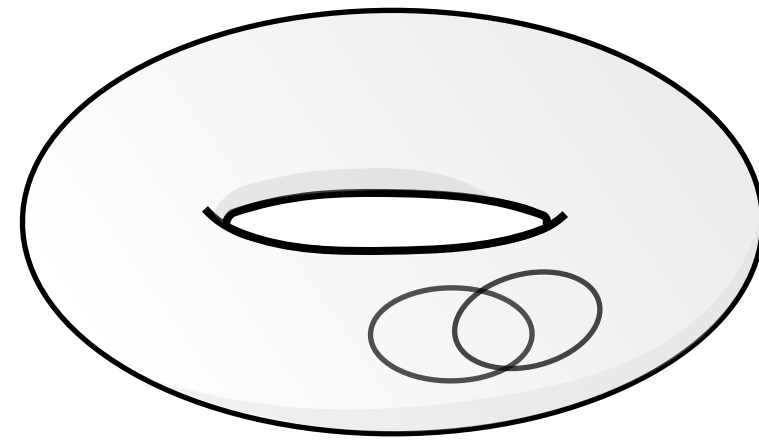
c.f also [Li,Palti,Petri '23][Palti,Petri '24]

Non-geometric backgrounds

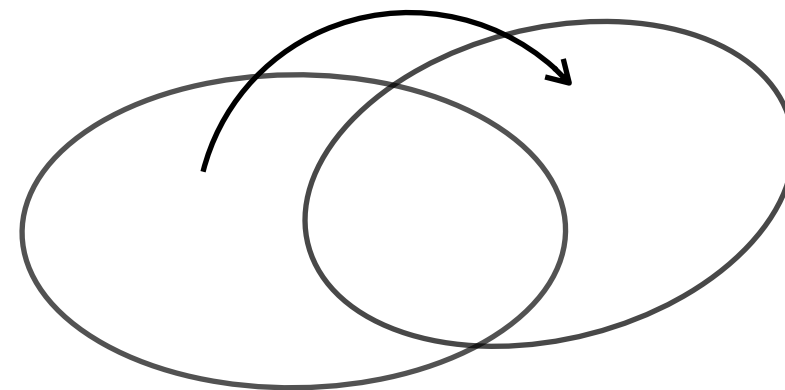
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[Dabholkar, Hull '02&'05; Flournoy, Wecht, Williams '04...]

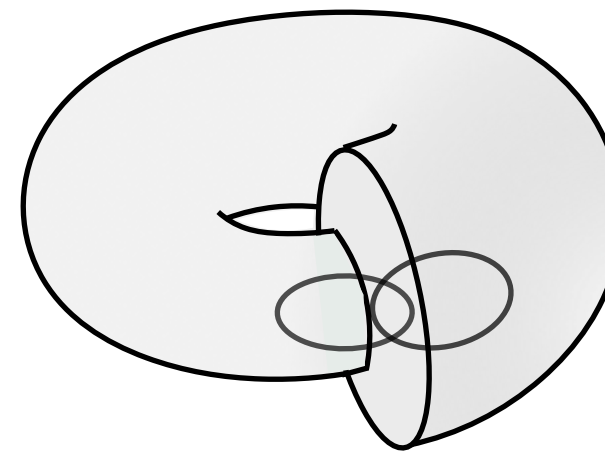
Manifold, e.g. a torus



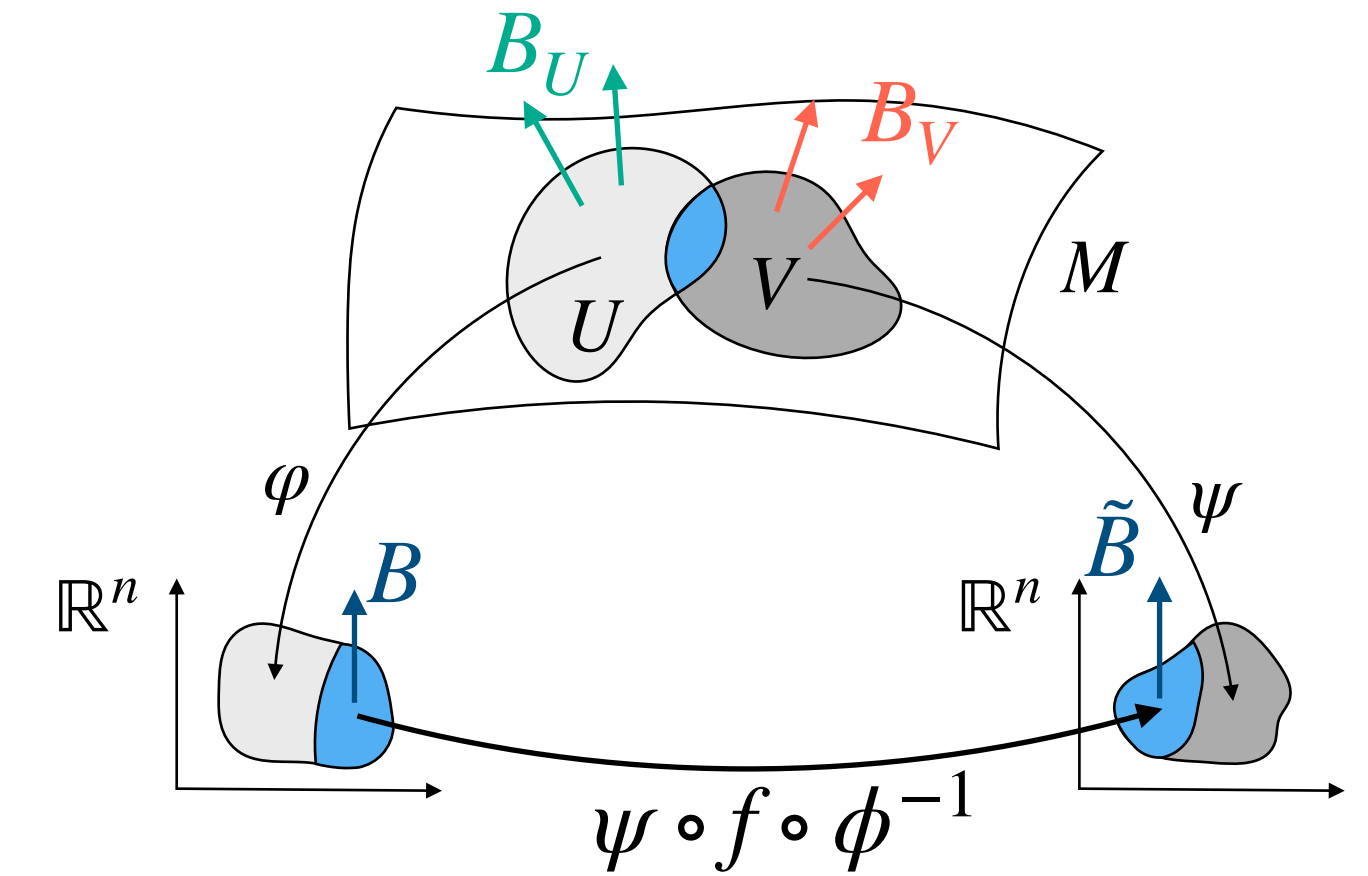
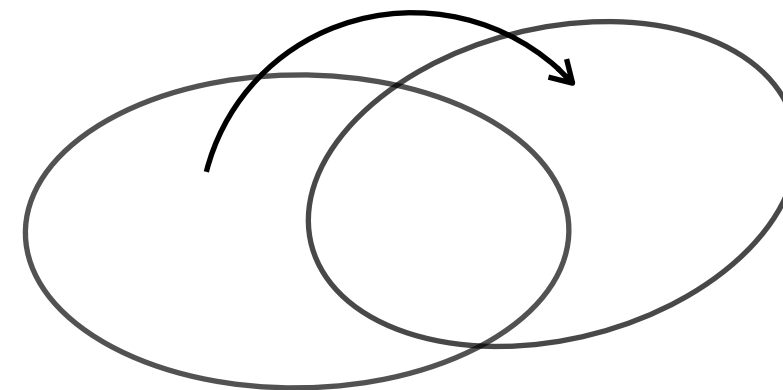
diffeomorphism



vs. non-geometric background



... and T-duality



$$f \in \begin{cases} Diff(M) & : \text{Riemannian} \\ Diff(M) \cup \text{T-duality} & : \text{non-geometric} \end{cases}$$

▷ **“Inevitable”** in string theory

▷ **Moduli stabilisation**

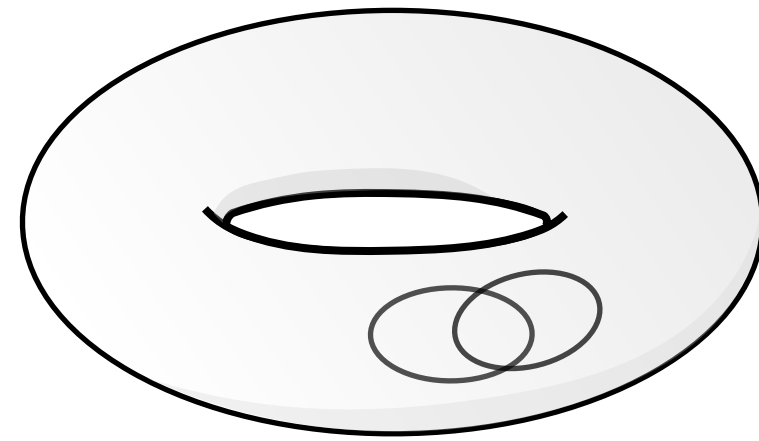
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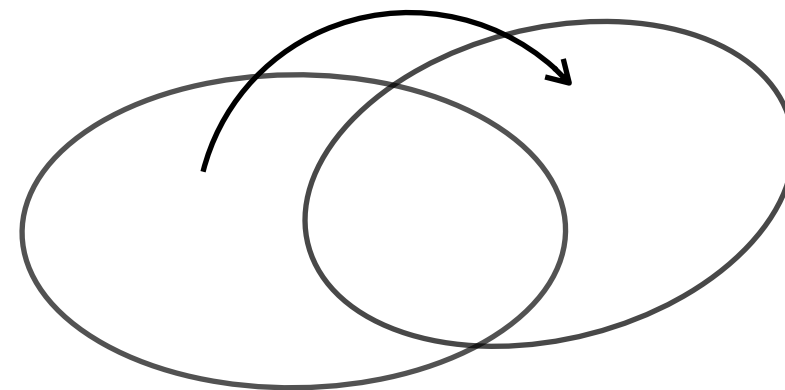
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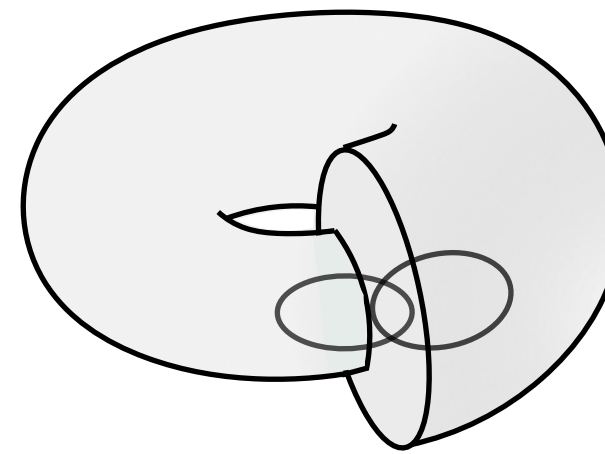
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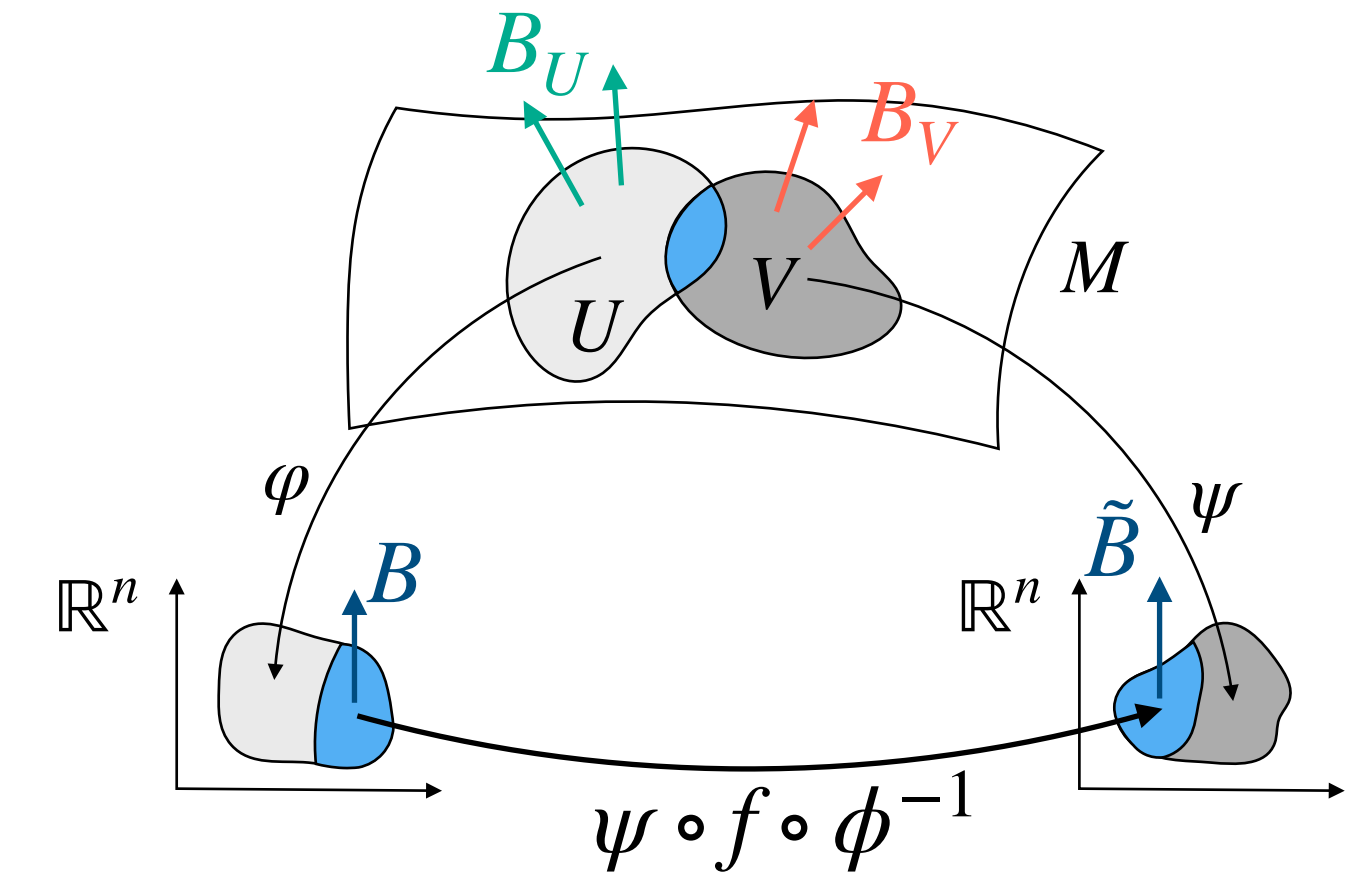
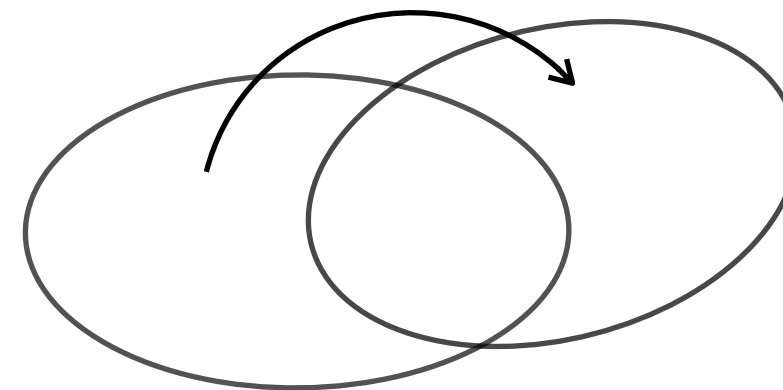
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[Becker, Becker, Blumenhagen, Lüst, Plauschinn, Shelton, Taylor, Vafa, Wecht, Walcher, ...]

In addition to standard fluxes, there are **non-geometric fluxes**, called Q - and R -fluxes.

Basic example: T-duality chain of T^3 with H -flux $H_{ijk} \rightarrow f_{ij}^k \rightarrow Q_i^{jk} \rightarrow R^{ijk}$

Are these valid backgrounds for quantum gravity?

Non-geometric backgrounds

Generically the action of non-geometric backgrounds is **ill-defined** in standard **NSNS frame**

- ▷ How to obtain previous reduction procedure?
- ▷ How to obtain **metric & potential** on moduli space
- ▷ Consistent picture under **T-duality**?

$$\theta \simeq \theta + 2\pi$$
$$g_{\mu\nu}(\theta + 2\pi) \neq g_{\mu\nu}(\theta)$$

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$$\tilde{V}(\varphi^a) \sim \int d^n y \sqrt{\tilde{h}_0} \left(\mathcal{R}(\tilde{h}) - \frac{1}{4} Q^2 + 4(\partial\tilde{\Phi}_y)^2 \right)$$

Perform the **field redefinition**: $(h + B)^{-1} = (\tilde{h}^{-1} + \beta)$ \longrightarrow $\mathcal{L}_\beta = \mathcal{L}_{NSNS} + \partial(\dots)$

... **β -supergravity action**

[Andriot, Larfors, Lüst, Patalong '11]

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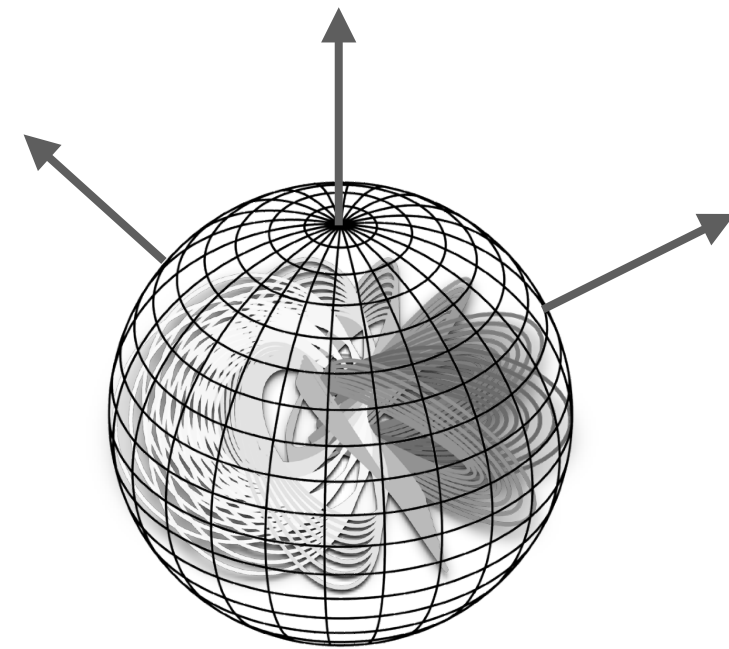
Crucial to use **β -supergravity** for consistency of **non-geometry backgrounds & geometric duals**:

A **consistent picture** between a (globally) non-geometric space and its geometric dual
 - i.e. matching moduli spaces, potentials and towers of states -
 can be established **only after moving to the β -frame**, where the background is well-defined.

[Demulder, Lüst, TR '23]

► Distance Conjecture:

S^3_R with $[H] = k$



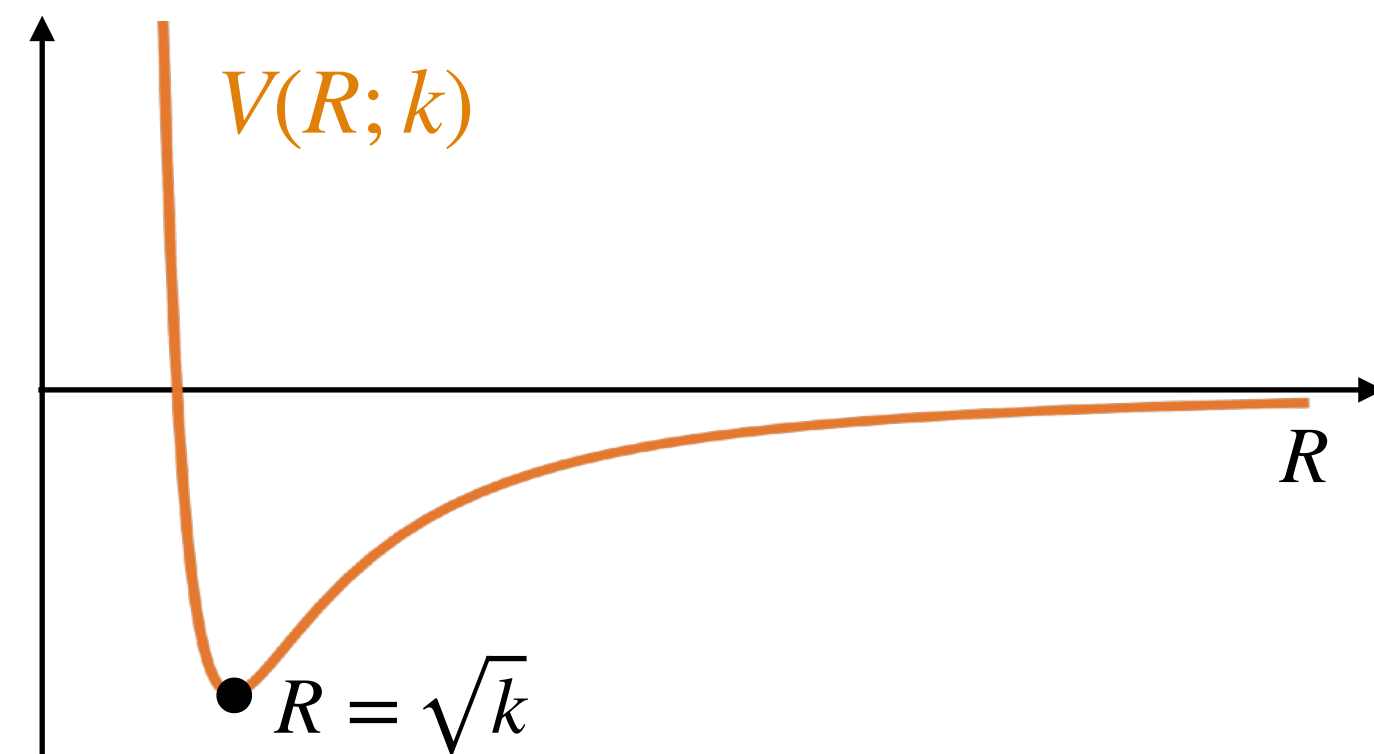
$$\gamma_{RR} = \frac{3}{R^2}$$

$$V(R; k) = -\frac{3}{2R^2} + \frac{k^2}{R^6}$$

Apparent inconsistency: S^3 with appropriately tuned H-flux is **valid string background** and therefore should be in the **Landscape**

However there is **no tower of light states** for $R \rightarrow 0$ which is an infinite distance limit

...need to **take into account scalar potential**



[Demulder, Lüster, TR '23]

In **effective field theories** that can be lifted to a theory of quantum gravity in the UV, a **divergence in the scalar potential** emerges when approaching an **infinite locus point** for which the target space geometry **cannot give rise to a light tower of states**.

That is, the **potential signals pathological infinite distance loci** in the scalar field space.

A different perspective



In which sense is $R = 0$ at **infinite distance** in the presence of a **(divergent) scalar potential**?

- ▷ Is the point “**reachable**” in the presence of a divergent **potential**?
- ▷ What is a “**good**” notion of **distance**?

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There is a plethora of works, examining the notion of distance in various settings.

- Very recently taking into account a scalar potential:
- ▷ [Mohseni, Montero, Vafa, Valenzuela '24]
 - ▷ [Debusschere, Tonioni, van Riet '24]
 - ▷ [De Biasio '22]

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...another somewhat **alternative approach**: **Ricci flow**

[Kehagias, Lüst, Lüst '19]

Ricci flow & Ricci flow Conjecture

Ricci flow equation: $\frac{\partial}{\partial t} g_{\mu\nu}(t) = -2\mathcal{R}_{\mu\nu}(t)$ [Hamilton '82]

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$$\Delta_g \sim \int_{\tau_i}^{\tau_f} \left(\frac{1}{V_M} \int_M \sqrt{|g|} g^{mn} g^{op} \frac{\partial g_{mo}}{\partial \tau} \frac{\partial g_{np}}{\partial \tau} \right)^{1/2} d\tau$$

AdS: $\mathcal{R}_{\mu\nu}(0) = \Lambda_0 \hat{g}_{\mu\nu} \longrightarrow \Lambda(t) = \frac{\Lambda_0}{(1 - 2\Lambda_0 t)}$
 fixed point with $\Lambda = 0$ as $t \rightarrow \infty$

$$\left\{ \begin{array}{l} \text{I: } \Delta_g \simeq \log(1 - 2\Lambda_0 t) \rightarrow \infty \\ \text{II: } \Delta_R \equiv \log\left(\frac{\mathcal{R}(0)}{\mathcal{R}(t)}\right) \sim \log(1 - 2\Lambda_0 t) \rightarrow \infty \end{array} \right.$$

[Kehagias, Lüst, Lüst '19]

Ricci flow Conjecture \implies AdS Distance Conjecture

Distance from Entropy functionals

Ricci flow equation: $\frac{\partial}{\partial t} g_{\mu\nu}(t) = -2\mathcal{R}_{\mu\nu}(t)$ [Hamilton '82]

Including the dilation leads to **Perelmans combined flow** [Perelman '02]

$$\begin{cases} \partial_t g_{\mu\nu}(t) &= -2\mathcal{R}_{\mu\nu}(t) \\ \partial_t \phi(t) &= -\frac{1}{2}\mathcal{R}(t) - \Delta\phi(t) + 2(\nabla\phi(t))^2 \end{cases}$$

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...gradient with respect to “Entropy Functional”

$$\mathcal{F}(g, f) = \int_M d^d x \sqrt{-g} e^{-2\phi} \left(\mathcal{R} + 4(\nabla\phi)^2 \right)$$

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Conjecture: The **distance in field space** $\Delta_{\mathcal{F}}$ along the combined flow is determined by the **entropy functional** $\mathcal{F}(g, f)$.

$\mathcal{F} = 0$ lies at **infinite distance** and is accompanied by an **infinite tower of massless states** and

$$\Delta_{\mathcal{F}} \simeq \log\left(\frac{\mathcal{F}_i}{\mathcal{F}_f}\right) \qquad \text{[Kehagias, Lüst, Lüst '19]}$$

Generalized Ricci flow & Distance Conjecture

[Demulder, Lüst, TR] work in progress



Can we define **notion of distance** for generic internal manifold using **generalized Ricci flow**?
What are the implications for the **Swampland Distance Conjecture** in presence of a **potential**?

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Metric & potential on moduli space both derived from **background data** on internal space

→ what are the implications of defining **distance in presence of potential** according to $\Delta_{\mathcal{F}^H} \simeq \log \left(\frac{\mathcal{F}_i^H}{\mathcal{F}_f^H} \right)$

[Demulder, Lüst, TR; ongoing]

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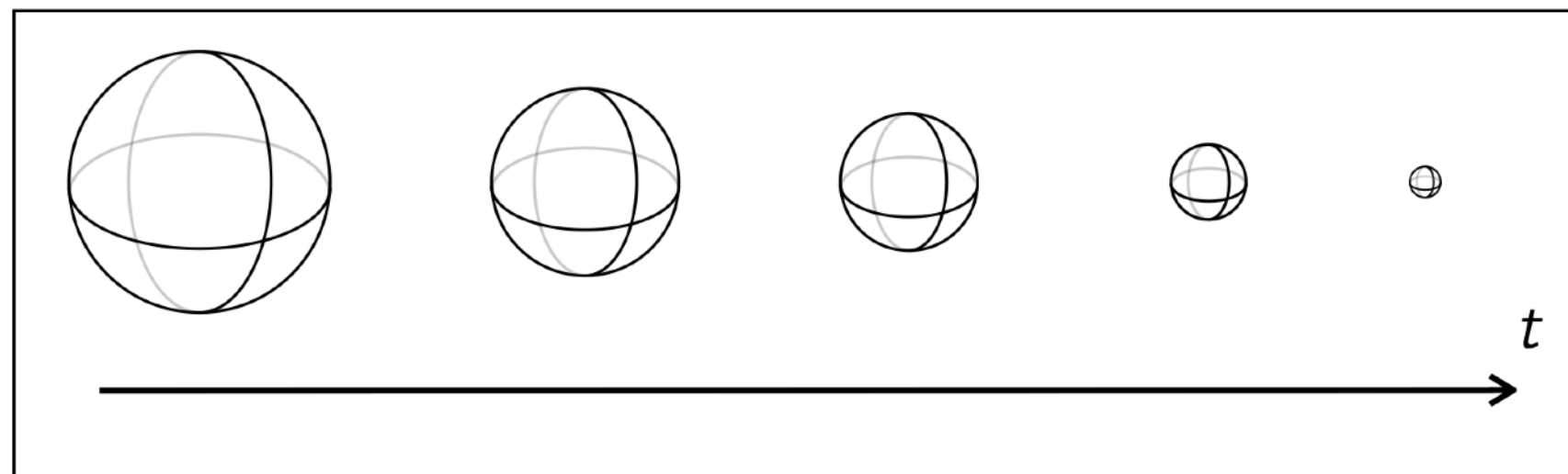
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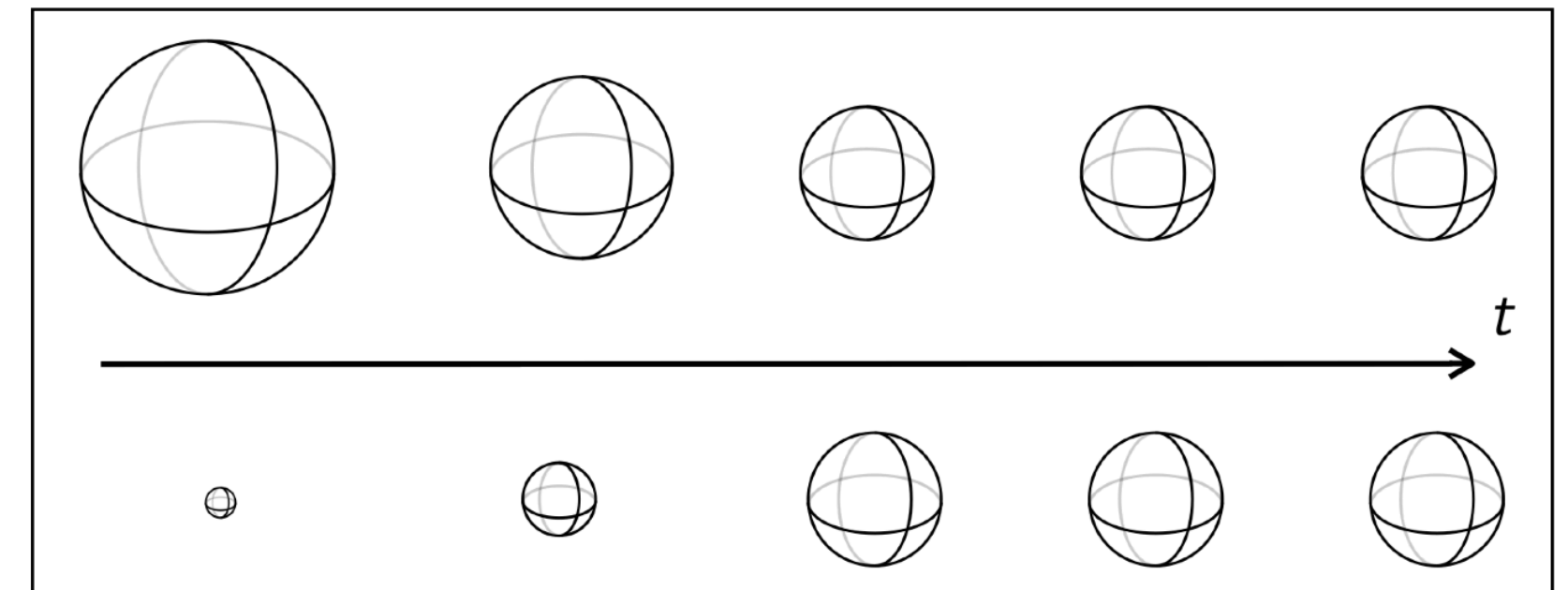
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[Demulder, Lüst, TR; ongoing]

$$\partial_t r(t) = - \frac{1}{2r(t)} \quad S_r^3$$



$$\partial_t r(t) = - \frac{1}{2r(t)} + \frac{h^2}{2r(t)^5} \quad S_r^3 \text{ with H-flux [h]}$$



Generalized Ricci flow & Distance Conjecture

[Demulder, Lüst, TR] work in progress



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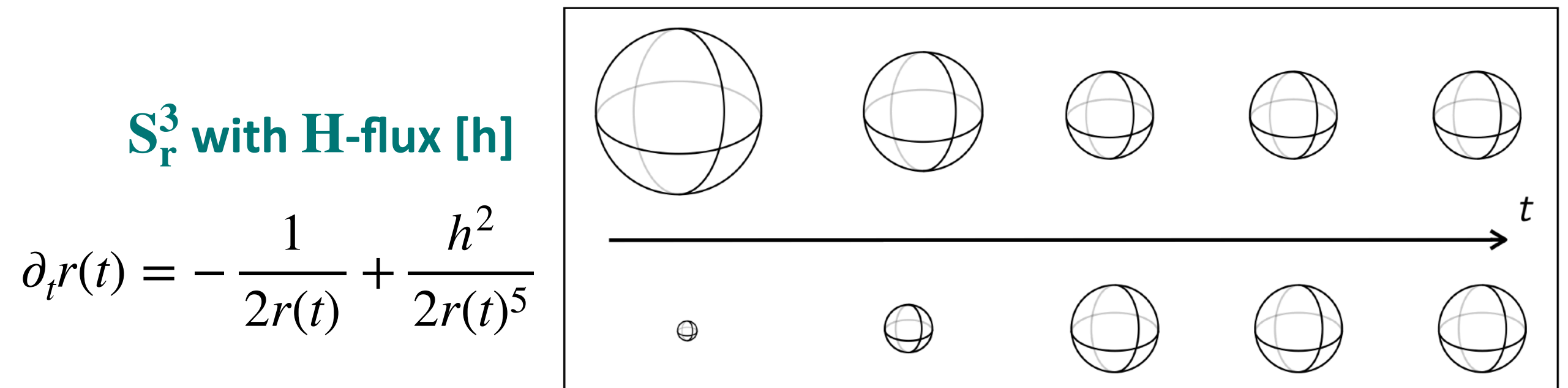
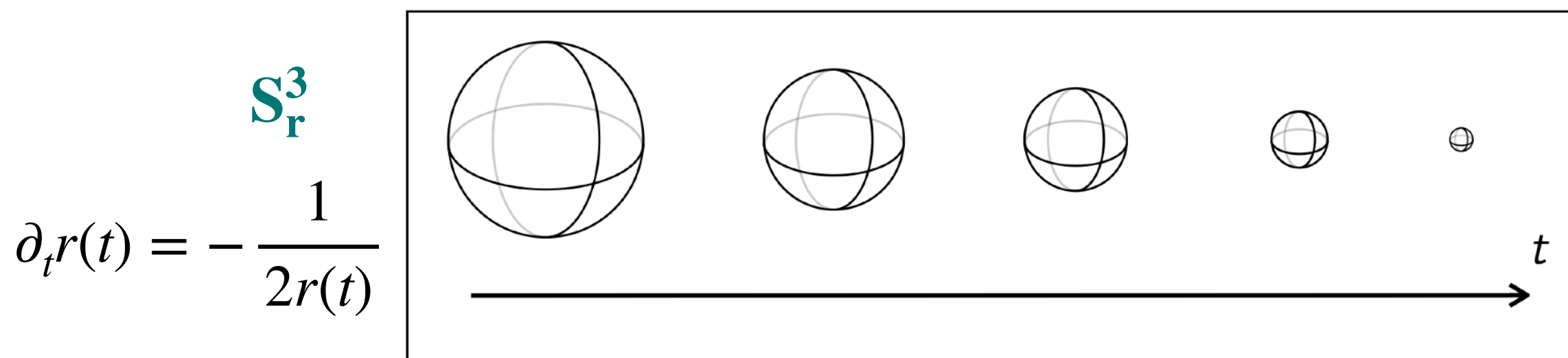
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[Demulder, Lüst, TR; ongoing]



- ▷ Adding **RR-fluxes**? [Demulder, Lüst, TR; ongoing]
- ▷ **Non-geometric flow**?

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- ▷ **Invariance** of the metric on moduli space **under (abelian) T-duality**
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- ▷ More **realistic setups: full 10d backgrounds**, e.g. $AdS_5 \times S^5$, $AdS_4 \times T^6$ with fluxes,...
- ▷ Flux variations & potential: **on-shell vs off-shell** [Li,Palti,Petri '23][Palti,Petri '24]
- ▷ **Deformations** and **generalized T-duality** (Poisson-Lie T-duality)
- ▷ **Distance** in presence of **potential** and **generalized Ricci flow** [Kehagias, Lüst, Lüst '19]
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— work in progress —

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Thank you!

— work in progress —

BACKUP SLIDES

T-duality and change in topology

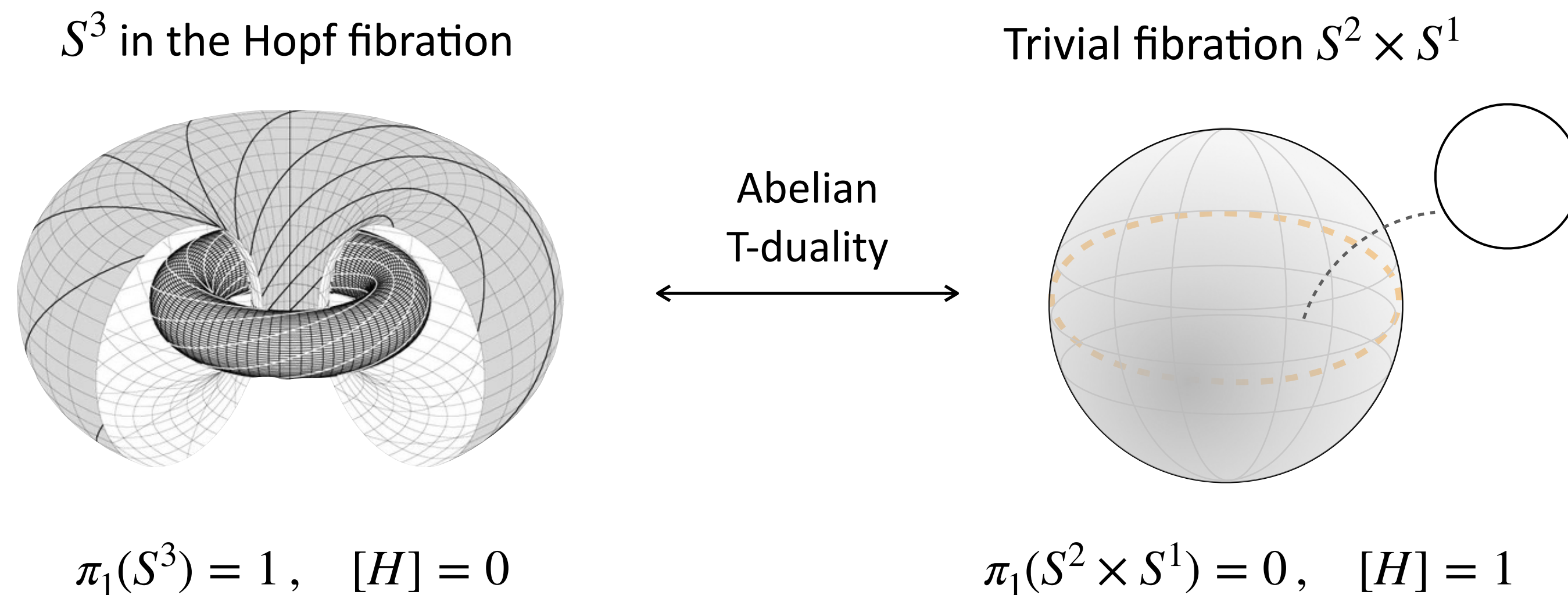
[Giveon, Kiritsis '94; Bouwknegt, Evslin, Mathai '04; Alvarez, Alvarez-Gaume, Lozano '05, ...]

T-duality does not only affect the geometry locally but **can also affect its global structure**

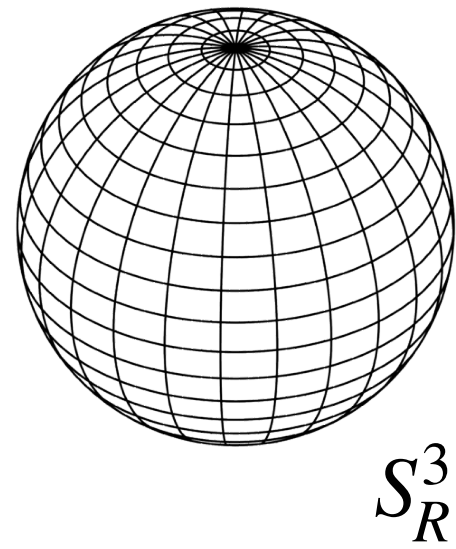
T-duality exchanges **H-flux** and the **first Chern class**

$$\text{fluxes : } \nu_k \mathcal{H} \quad \longleftrightarrow \quad \text{topology : } c_1(E_k)$$

Concrete example:



NATD of S^3



S^3_R

NATD

$$V(\varphi^a) \sim \int d^n y \sqrt{h_0} \left(\mathcal{R}(h) - \frac{1}{12} H^2 + 4(\partial\Phi_y)^2 \right)$$

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$$h = \frac{1}{R^2 \chi} \begin{pmatrix} R^4 + \phi^2 & \phi\psi & \theta\phi \\ \phi\psi & R^4 + \psi^2 & \theta\psi \\ \theta\phi & \theta\psi & R^4 + \theta^2 \end{pmatrix}, \quad B = \frac{1}{\chi} \begin{pmatrix} 0 & -\theta & \psi \\ \theta & 0 & -\phi \\ -\psi & \phi & 0 \end{pmatrix}$$

$\chi = R^4 + \phi^2 + \psi^2 + \theta^2$, and $\{\phi, \psi, \theta\}$ periodic angular coordinates

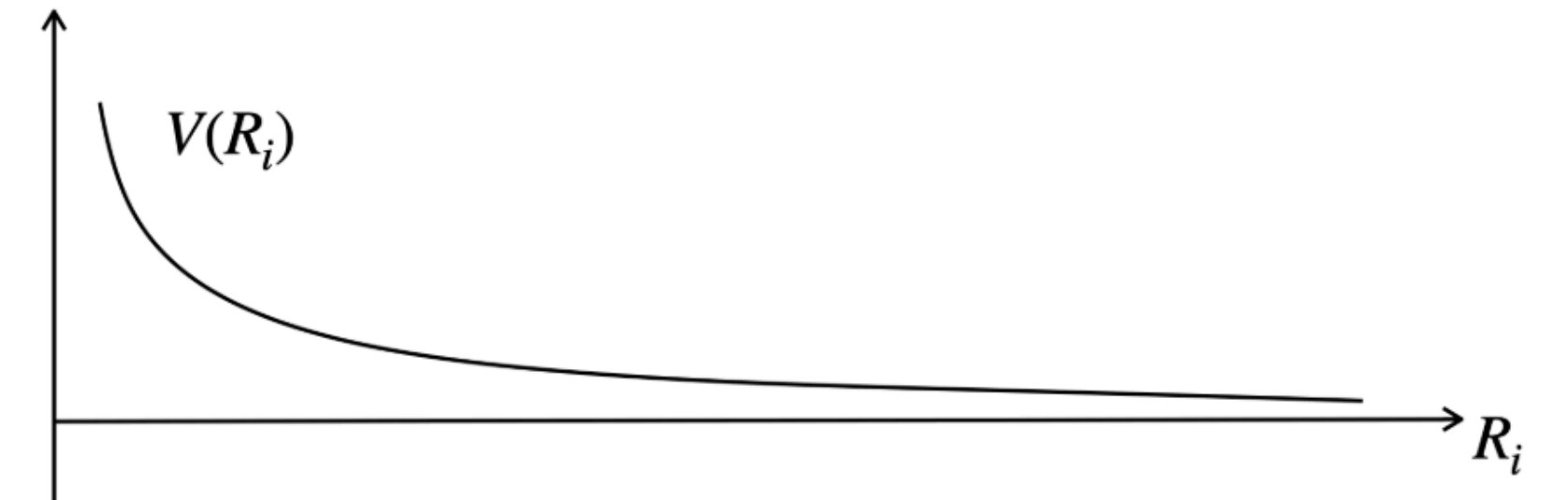
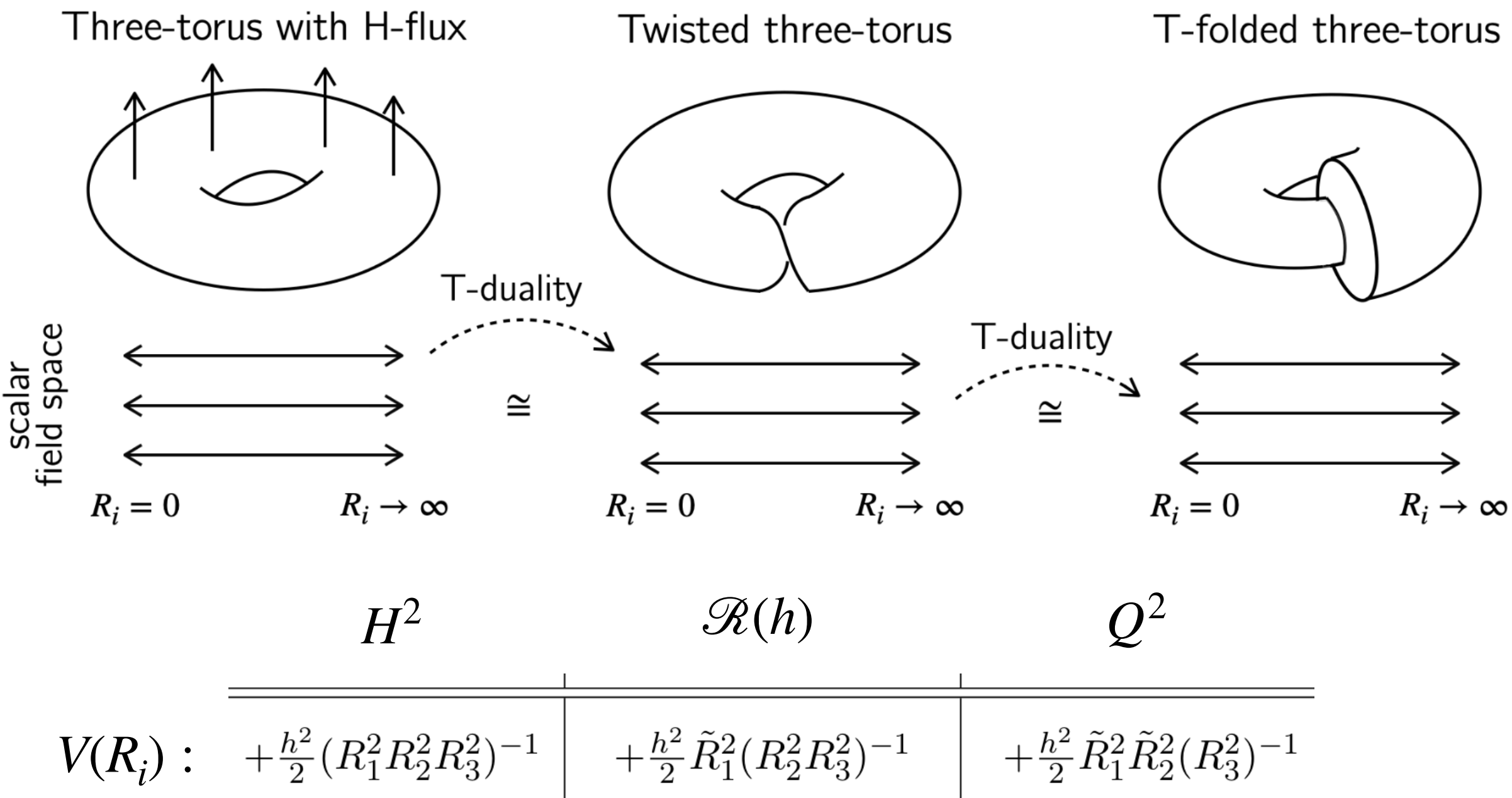
β -gravity

[Andriot, Larfors, Lüst, Patalong '11]

$$\hat{h} = R^{-2} I_3, \quad \beta = \begin{pmatrix} 0 & -\theta & \psi \\ \theta & 0 & -\phi \\ -\psi & \phi & 0 \end{pmatrix}$$

Q -flux: $Q_i^{jk} = \partial_i \beta^{jk}$

T-duality chain



	modes	$R_i \rightarrow 0$	$R_i \rightarrow \infty$
$T_{\mathcal{H}}^3, [\mathcal{H}] = k$ $\{R_1, R_2, R_3\}$	$w : \mathbb{Z} \oplus \mathbb{Z} \oplus \mathbb{Z}$ $p : nc \oplus nc \oplus nc$	$w : (\text{light}, \text{light}, \text{light})$ $p : (\overline{\text{heavy}}, \overline{\text{heavy}}, \overline{\text{heavy}})$	$w : (\text{heavy}, \text{heavy}, \text{heavy})$ $p : (\underline{\text{light}}, \underline{\text{light}}, \underline{\text{light}})$
$T_{\text{tw}}^3, [f] = k$ $\{R_1, R_2, R_3^{-1}\}$	$w : \mathbb{Z} \oplus \mathbb{Z} \oplus \mathbb{Z}_k$ $p : nc \oplus nc \oplus \mathbb{Z}$	$w : (\text{light}, \text{light}, \overline{\text{heavy}})$ $p : (\overline{\text{heavy}}, \overline{\text{heavy}}, \text{light})$	$w : (\text{heavy}, \text{heavy}, \underline{\text{light}})$ $p : (\underline{\text{light}}, \underline{\text{light}}, \text{heavy})$
$T_Q^3, [Q] = k$ $\{R_1, R_2^{-1}, R_3^{-1}\}$	$w : \mathbb{Z} \oplus \mathbb{Z}_k \oplus \mathbb{Z}_k$ $p : nc \oplus \mathbb{Z} \oplus \mathbb{Z}$	$w : (\text{light}, \overline{\text{heavy}}, \overline{\text{heavy}})$ $p : (\overline{\text{heavy}}, \text{light}, \text{light})$	$w : (\text{heavy}, \underline{\text{light}}, \underline{\text{light}})$ $p : (\underline{\text{light}}, \text{heavy}, \text{heavy})$

T-duality chain of T^3 with H -flux is not a proper string background. Careful analysis reveals a lack of towers of states for some infinite distance points and therefore:

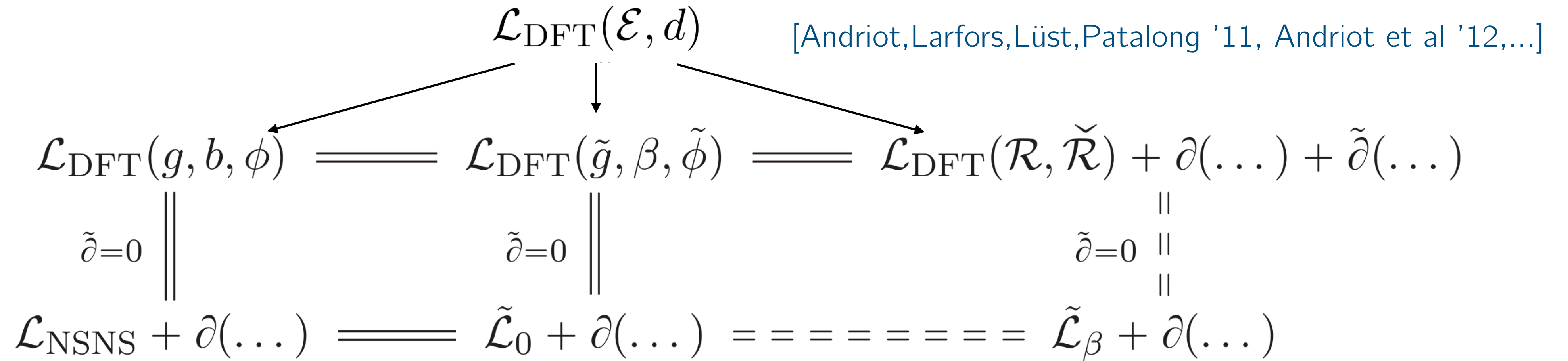
Without a completion, e.g. additional fluxes, the backgrounds of the **T-duality chain** lie in the **Swampland**.

[Demulder, Lüst, TR '23]

$$S_{\text{DFT}} = \int dx d\tilde{x} e^{-2d} \left(\frac{1}{8} \mathcal{H}^{MN} \partial_M \mathcal{H}^{KL} \partial_N \mathcal{H}_{KL} - \frac{1}{2} \mathcal{H}^{MN} \partial_N \mathcal{H}^{KL} \partial_L \mathcal{H}_{MK} - 2 \partial_M d \partial_N \mathcal{H}^{MN} + 4 \mathcal{H}^{MN} \partial_M d \partial_N d \right).$$

$$\begin{pmatrix} h - Bh^{-1}B & Bh^{-1} \\ -h^{-1}B & h^{-1} \end{pmatrix} = \mathcal{H} = \begin{pmatrix} \tilde{h} & \tilde{h}\beta \\ -\beta\tilde{h} & \tilde{h}^{-1} - \beta\tilde{h}\beta \end{pmatrix}$$

β -gravity



$$\tilde{\mathcal{L}}_\beta = e^{-2\tilde{\phi}} \sqrt{|\tilde{g}|} \left(\mathcal{R}(\tilde{g}) + \check{\mathcal{R}}(\tilde{g}) + 4(\partial\tilde{\phi})^2 - \frac{1}{2}R^2 + 4(\beta^{mp} \partial_p \tilde{\phi} - \mathcal{T}^m)^2 \right)$$

$$\simeq \mathcal{R}(g) - \frac{1}{4}Q^2 - \frac{1}{2}\tilde{g}_{ij} Q_k^{jl} Q_l^{ki} + 4(\partial\tilde{\phi})^2 + \dots$$

= $\partial(\dots)$

$$\mathcal{L}_{\text{NSNS}} = \mathcal{R}(h) - \frac{1}{12} \mathcal{H}^2 + 4(\partial\Phi_y)^2 = \frac{3}{2R^2} + 4R^2 \chi^{-2} (3R^4 + \phi^2 + \psi^2 + \theta^2) \quad \mathcal{L}_\beta = \frac{3}{2R^2}$$