

QCD in higher orders

Peter Marquard

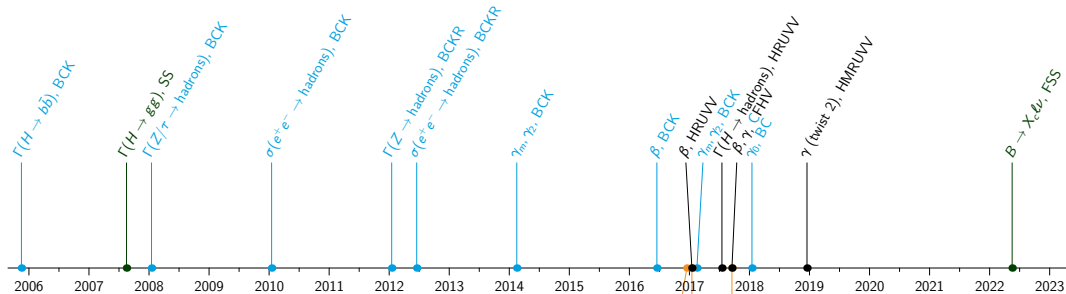
in collaboration with

A. Maier, Y. Schröder



Corfu, August 2024

Towards QCD at five loops



BC(K)(R): [Baikov, Chetyrkin, (Kühn), (Rittinger)]

SS: [Schreck, Steinhauser], FSS: [Fael, Schönwald, Steinhauser]

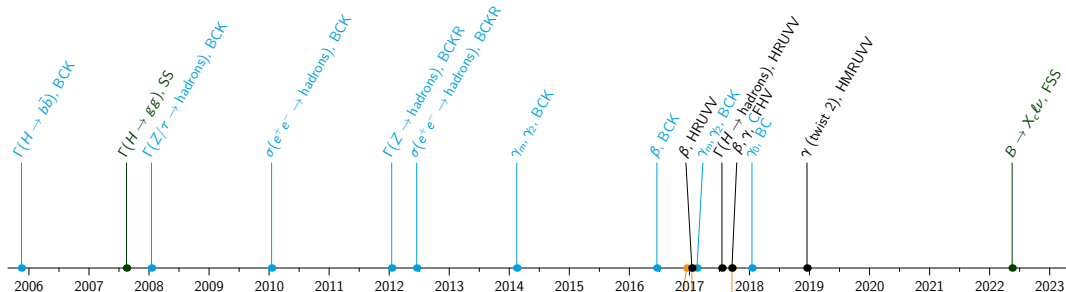
H(M)RUVV: [Herzog, (Moch), Ruijl, Ueda, Vermaseren, Vogt]

CFHV: [Chetyrkin, Falcioni, Herzog, Vermaseren]

LMMS: [Luthe, Maier, Marquard, Schröder]

$\gamma_m \gamma_2$, LMMS
 γ , LMMS
 β, γ , LMMS

Towards QCD at five loops



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} at most four loops
 not quite QCD: massive gluons

What can we learn at five loops?

- Total decay widths, $\sigma(e^+e^- \rightarrow \text{hadrons})$ at order α_s^5
↔ α_s **determinations**, mostly limited by experiment

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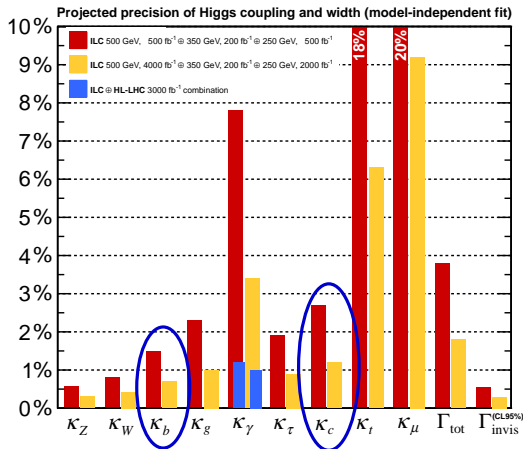
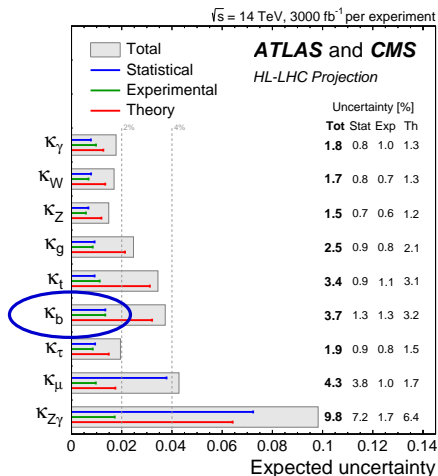
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- Renormalisation group evolution at six loops,
including five-loop decoupling

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- Renormalisation group evolution at six loops,
including five-loop decoupling
- Moments of $\sigma(e^+e^- \rightarrow \text{hadrons})$ at order α_s^5
 \hookrightarrow **heavy quark mass determinations**
- ...

Why precise quark masses?

Higgs coupling measurements



Heavy quark masses at four loops

[Chetyrkin, Kühn, Maier, Maierhöfer, Marquard, Steinhauser, Sturm 2009 + 2017]

$$m_b(10 \text{ GeV}) = (3610 \pm 10(\text{exp}) \pm 12(\alpha_s) \pm 3(\mu)) \text{ MeV}$$

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$$m_c(3 \text{ GeV}) = (993 \pm 7(\text{exp}) \pm 4(\alpha_s) \pm 2(\mu) \pm 1(\text{np})) \text{ MeV}$$

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[Dehnadi, Hoang, Mateu, Zebarjad 2013; Dehnadi, Hoang, Mateu 2015]

$$m_b(m_b) = (4176 \pm 4(\text{stat}) \pm 19(\text{sys}) \pm 7(\alpha_s) \pm 10(\mu)) \text{ MeV}$$

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Need masses at five loops

Where to start?

- 1 Generate diagrams (qgraf)
- 2 Identify integral families (nauty and Traces)
- 3 Insert Feynman rules (FORM)
- 4 Reduce to master integrals
- 5 Calculate & insert master integrals

[Nogueira 1991]

[McKay, Piperno 2014]

[Vermaseren et al.]

Where to start?

One momentum or one mass?



• Dedicated methods

- Triangle / diamond reduction rules
[Chetyrkin, Tkachov 1981; Ruijl, Ueda, Vermaseren 2015]
- Graphical functions
- ...
- **64** integral families
- **15** propagators
- **20** possible scalar products

- **Standard reduction methods:**
Integration-by-parts reduction
over finite fields

[Chetyrkin, Tkachov 1981; Laporta 2000; von Manteuffel, Schabinger 2014; Peraro 2016; ...]

- **34** integral families
- **12** propagators
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Where to start?

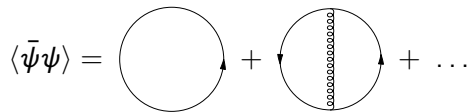
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Heavy-quark condensate

$$\langle \bar{\psi}\psi \rangle = \text{circle} + \text{circle with vertical line} + \dots$$
The equation shows the expansion of the heavy-quark condensate $\langle \bar{\psi}\psi \rangle$. It consists of three terms: a simple circle with an arrow on the right side, a circle with a vertical line of small circles inside and an arrow on the right side, and an ellipsis indicating further terms.

- Leading non-analytic contribution in
 - Operator Product Expansion ($m \sim \Lambda_{\text{QCD}}$)
 - Asymptotic small-mass expansion ($m \gg \Lambda_{\text{QCD}}$)

Extrapolation from heavy to light quarks via
Renormalisation Group Optimised Perturbation Theory

[Kneur, Neveu 2010-2020]

Heavy-quark condensate

$$\langle \bar{\psi}\psi \rangle = \text{[Diagram 1]} + \text{[Diagram 2]} + \dots$$

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Extrapolation from heavy to light quarks via
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- Direct relation to vacuum anomalous dimension:

$$\mu^2 \frac{d}{d\mu^2} m \langle \bar{\psi}\psi \rangle = -4m^4 \gamma_0$$

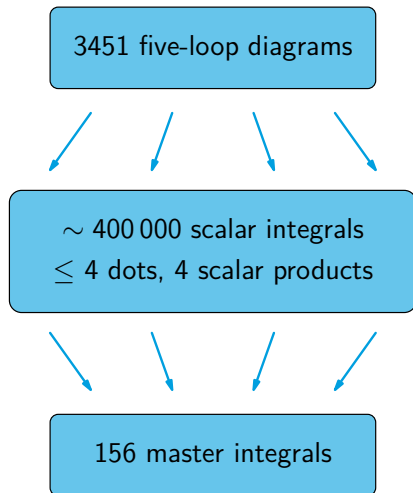
↔ independent check of five-loop result [Baikov, Chetyrkin 2018]

[Kneur, Neveu 2010-2020]

[Spiridonov, Chetyrkin 1988]

Heavy-quark condensate

Calculation



Master integrals

Sector decomposition [Hepp 1966; Speer 1963; Dlapa, Heinrich 2000]

Compute master integrals with FIESTA:

[Smirnov et al. 2008-2021]

$$\langle \bar{\psi}\psi \rangle \Big|_{\left(\frac{\alpha_s}{\pi}\right)^4} = \frac{m^3}{16\pi^2} \left[\frac{(-3.5 \pm 3.0) \times 10^{-8}}{\epsilon^{11}} + \frac{(-2.1 \pm 7.6) \times 10^{-6}}{\epsilon^{10}} \right. \\ \left. + \frac{(0.2 \pm 2.1) \times 10^{-4}}{\epsilon^9} + \frac{(-0.5 \pm 7.5) \times 10^{-2}}{\epsilon^8} + \dots \right]$$

Master integrals

Sector decomposition [Hepp 1966; Speer 1963; Dineth, Heinrich 2000]

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- ~ 2 digits lost per order in $\epsilon \Rightarrow$ **need better than double precision**
- Found no significant improvement with
 - Quasi-Monte Carlo methods [Borowka, Heinrich, Jahn, Jones, Kerner, Schlenk 2018]
 - Quasi-finite integral basis [von Manteuffel, Panzer, Schabinger 2014]
 - Above plus direct Feynman parameter integration
 - Tropical integration [Borinsky 2020]
 - pySecDec [Heinrich, Jones, Kerner, Magerya, Olsson, Schlenk 2023]

Master integrals

Difference equations [Laporta 2000; Luthe, Schröder]

- Raise one denominator to symbolic power x :

$$M_i(x) = \int \frac{1}{D_1^x \cdots D_{15}^{\alpha_{15}}}$$

- IBP reduction \rightarrow coupled difference equations in x

$$\sum_{j,k} q_{jk}(d, x) M_j(x+k) = 0$$

- Factorial series ansatz \rightarrow recurrence relations for $a_{i,s}$

$$M_i(x) = \sum_{s=0}^{\infty} a_{i,s} \frac{\Gamma(x+1)}{\Gamma(x+s-d/2+1)}$$

- Boundary conditions from $x \rightarrow \infty$

High-precision numerical solution

Master integrals

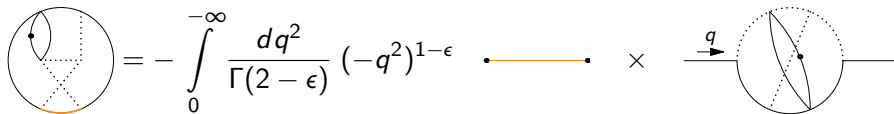
Direct numerical integration



The diagram shows an equation between two Feynman diagrams. On the left is a circle with a solid line on the left, a solid line on the bottom, and a dashed line on the right. A dashed line also connects the top-left and bottom-right vertices. A solid line connects the top-left and bottom-left vertices. On the right is a circle with a solid line on the left, a solid line on the right, and a dashed line on the top. A dashed line also connects the top-left and bottom-right vertices. A solid line connects the top-left and bottom-left vertices. An arrow labeled q points to the right from the left solid line. The equation is:
$$\text{Left Diagram} = \int \frac{d^d q}{i\pi^{\frac{d}{2}}} \text{Right Diagram}$$

Master integrals

Direct numerical integration


$$\text{Diagram} = - \int_0^{\infty} \frac{dq^2}{\Gamma(2 - \epsilon)} (-q^2)^{1 - \epsilon} \times \text{Diagram}$$

Integral is


- **IR finite:** q only flows through massive lines
- **UV finite:** otherwise raise **tree-level propagator** to higher power

Alternative: introduce counterterms

Expansion in ϵ straightforward

Master integrals

Direct numerical integration

$P_0 = \text{---} \xrightarrow{q} \text{---}$  from differential equations for propagator master integrals P_i :

$$q^2 \frac{\partial}{\partial q^2} P_i = Q_{ij} P_j$$

[Kotikov 1991, Remiddi 1997,...]

Expansions around $q^2 \rightarrow 0$, $q^2 \rightarrow -\infty$ from power series ansätze:

$$P_i = \sum_{k=0}^{N-1} c_{ik} q^{2k} + \mathcal{O}(q^{2N}), \quad P_i = \sum_{n=0}^{\#\text{loops}} \sum_{k=k_0}^{N-1} d_{ikn} \left(-\frac{1}{q^2}\right)^{k+n\epsilon} + \mathcal{O}\left(\frac{1}{(-q^2)^N}\right)$$

Boundary conditions $c_{i,0}$, d_{ik_0n} :

products of massive vacuum diagrams & massless propagators with ≤ 4 loops

[Schröder, Vuorinen 2005; Faisst, Maierhöfer, Sturm 2006; Baikov, Chetyrkin 2010; Lee, Smirnov, Smirnov 2011]

Master integrals

Padé approximation

Accelerate series convergence through Padé approximation $[n/m]$:

$$P_i = \sum_{k=0}^{N-1} c_{ik} q^{2k} + \mathcal{O}(q^{2N})$$

\Downarrow

$$P_i = [n/m] + \mathcal{O}(q^{2N}) \equiv \frac{a_0 + a_1 q^2 + \dots + a_n q^{2n}}{1 + b_1 q^2 + \dots + b_m q^{2m}} + \mathcal{O}(q^{2N})$$

Further improvements:

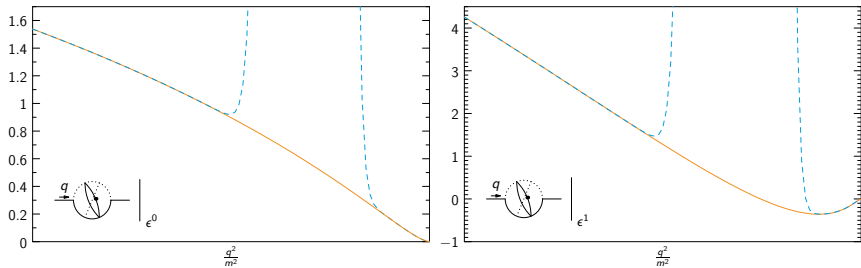
[Baikov, Broadhurst 1995]

- Combine low- and high-energy expansions into single Padé approximation
- Subtract high-energy logarithms $\ln^n(-q^2)$
- Conformal mapping $q^2 \rightarrow \frac{4\omega}{(1+\omega)^2} m^2$

See also [Faisst, Chetyrkin, Kühn 2004]

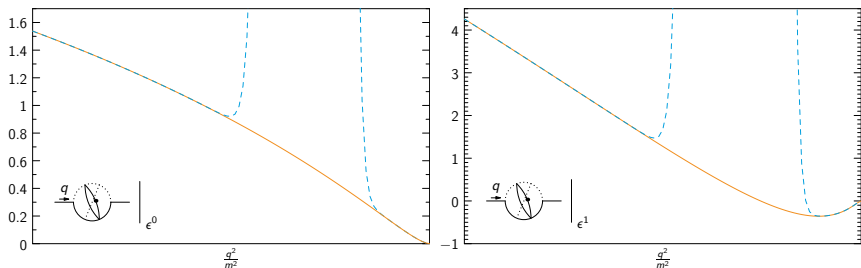
Master integrals

Result for $N = 20$ expansion terms



Master integrals

Result for $N = 20$ expansion terms



$$\text{Diagram} = 5.8125309358416596949 - 31.572349480122869826\epsilon + \mathcal{O}(\epsilon^2)$$

Checks:

- ✓ Change integration contour to $0 \rightarrow -i\infty \Rightarrow$ agrees within $10^{-19} + 2 \times 10^{-18}\epsilon$
- ✓ Agrees with result for $N = 25$ within uncertainty
- ✓ FIESTA: $5.81409 - 31.58070(87)\epsilon$
- ⚠ Compare with numerical solution of differential equations

Next Steps

Decoupling

Decoupling at five loops

Motivation

Top quarks are problematic for QCD processes with $E < m_t$:

- Diagrams with massive internal lines \Rightarrow hard to calculate
- Large logarithms $\ln\left(\frac{E}{m_t}\right)$ spoil perturbative convergence

Decoupling at five loops

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Solution: effective 5-flavour theory with coupling

$$\alpha_s^{(5)} = \alpha_s^{(6)} \times \left[\frac{\left(\text{triangle diagram} \right)^2}{\left(\text{self-energy diagram} \right)^2 \cdot \left(\text{vacuum polarization diagram} \right)} \right]$$

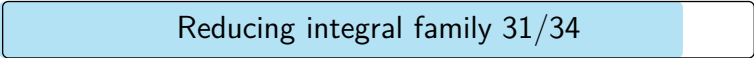
$q=0$
 $m_t \neq 0$
 $m_Q=0$

Relation known to four loops [\[Schröder, Steinhauser 2005\]](#)

Decoupling at five loops

Status

- 131 860 803 scalar integrals
- Up to 7 dots, 6 scalar products
- Reduction to master integrals with `crusher` + `tinbox` [Maier, Marquard, Seidel]



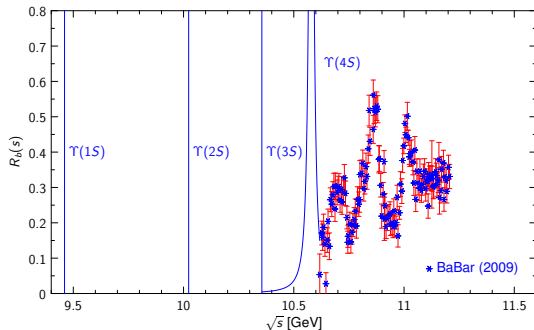
Reducing integral family 31/34

A horizontal progress bar with a light blue fill and a black border. The bar is approximately 90% full, with the remaining 10% being white. The text "Reducing integral family 31/34" is centered within the bar.

Next Steps

Quark Masses

Heavy Quark Mass Determination



Consider *moments* \mathcal{M}_n of $R_Q(s) = \frac{\sigma(e^+e^- \rightarrow Q\bar{Q})}{\sigma(e^+e^- \rightarrow \mu^+\mu^-)}$

$$\mathcal{M}_n = \int_{s_0}^{\infty} ds \frac{R_Q(s)}{s^{n+1}} = \frac{12\pi^2}{n!} \underbrace{\left[\left(\frac{d}{dq^2} \right)^n \Pi_Q(q^2) \right]_{q^2=0}}_{\text{massive vacuum diagrams}}$$

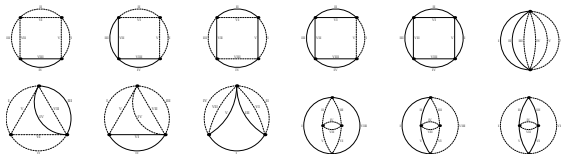
The first moment at five loops

First results

n_h closed heavy (massive) quark loops + n_l closed light (massless) quark loops

	n_l^0	n_l^1	n_l^2	n_l^3
n_h^1	✗	✗	✓	✓
n_h^2	✗	✗	✓	
n_h^3	✗	✓		
n_h^4	✓			

Master integrals:



$$\mathcal{M}_1^{5 \text{ loop}} = \frac{3\pi^2}{m_Q^2} \left(\frac{\alpha_s}{\pi} \right)^4 n_h C_F \left[0.6 T_F^3 n_l^3 + 1.2 T_F^3 n_l^2 n_h + 0.9 T_F^3 n_l n_h^2 + 0.2 T_F^3 n_h^3 + (C_F - 5C_A) T_F^2 n_l^2 + \dots \right]$$

Conclusion

Five-loop QCD work in progress

- Quark condensate
- Decoupling
- Charm & bottom mass determinations

Combination of methods for computing five-loop vacuum diagrams:

- Sector decomposition
- Recurrence equations
- Direct integration