Hunting Invisibles: Dark sectors, Dark matter and Neutrinos

Southampton School of Physics and Astronomy

Gravitational Waves as Probes of New Physics

STEVE KING 08/09/24 Dark Side of the Universe









Gravitational Waves enable us to look back to the earliest moments of the Universe (back to Inflation)

What powered the big bang?

Only gravitational waves can escape from the earliest moments of the Big Bang

Big Bang plus 10⁻⁴³ seconds

BIG BANG

Inflation (Big Bang plus 10⁻³⁵ seconds?)

> Cosmic microwave background, distorted by seeds of structure and gravitational waves

> > Light

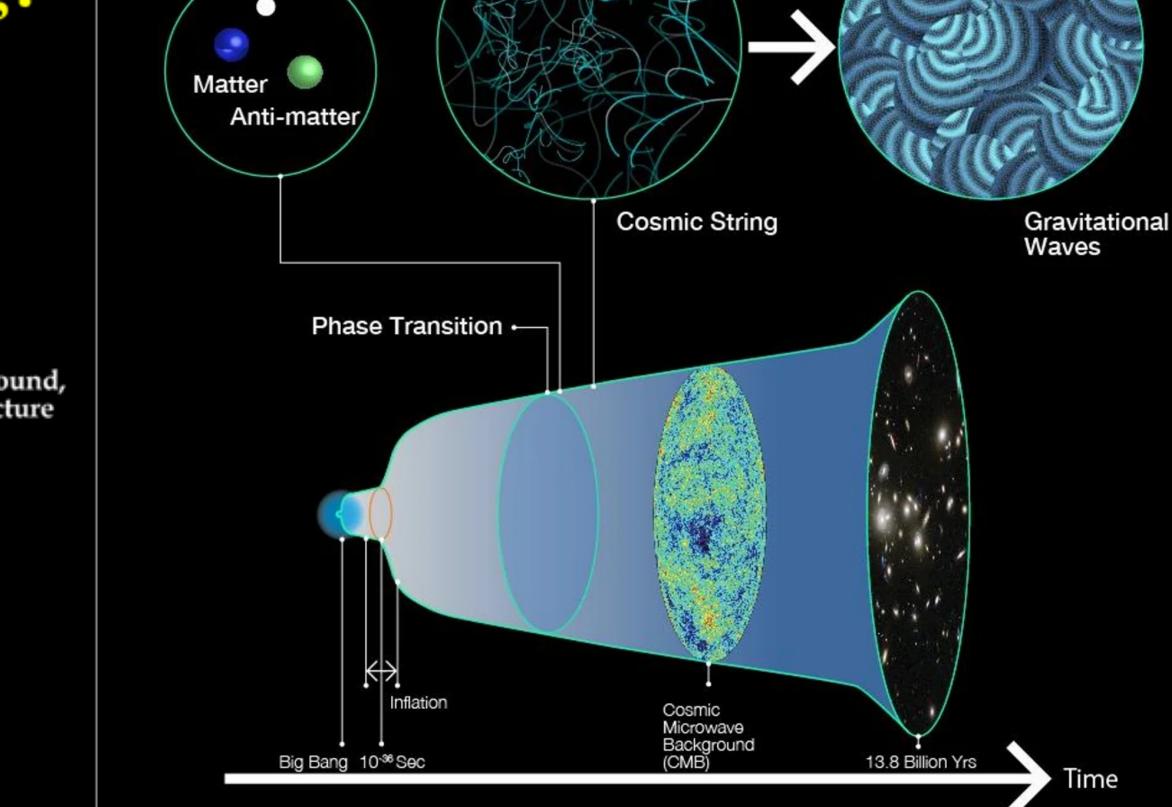
Now

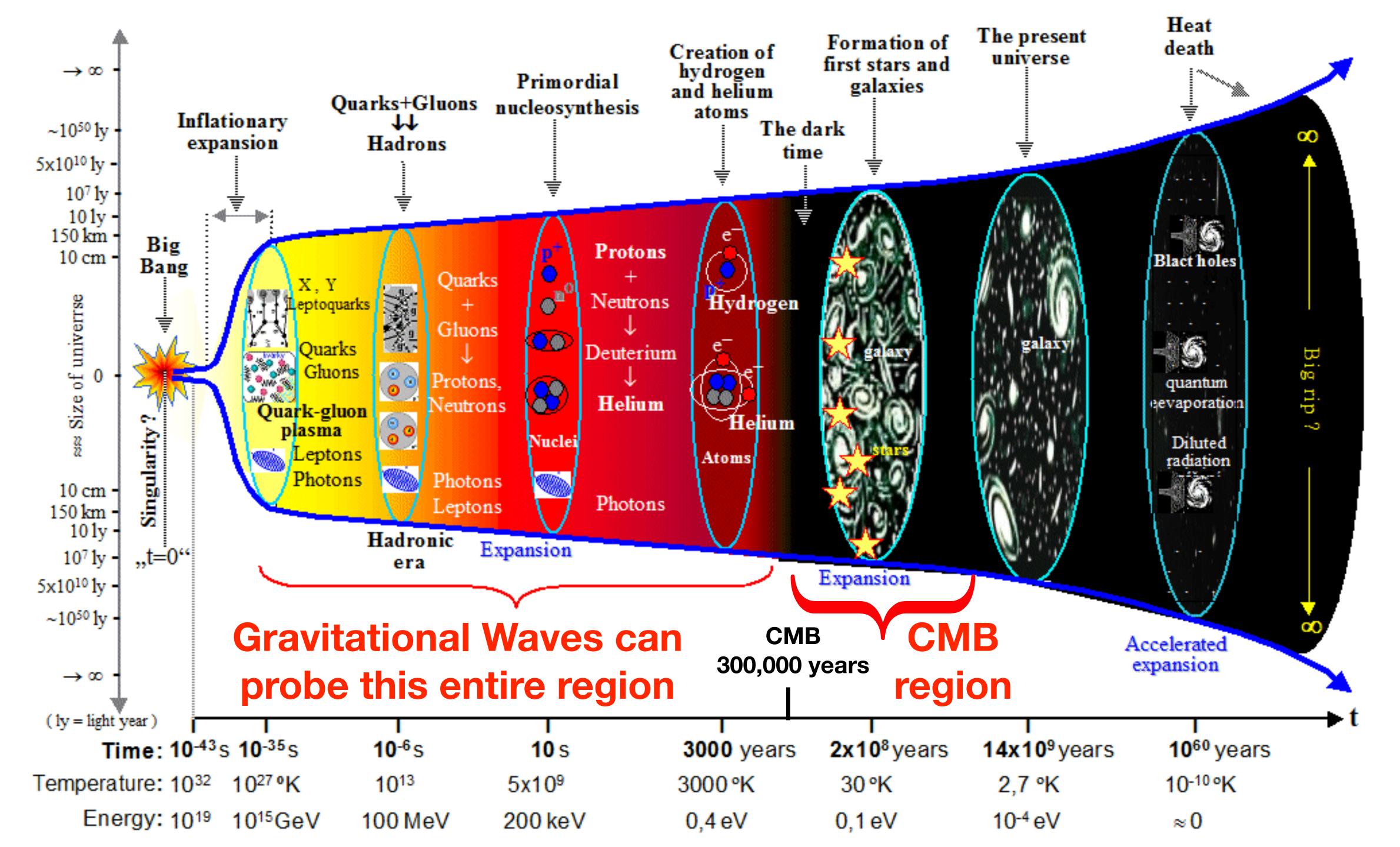
Big Bang plus 300,000 Years

Gravitational

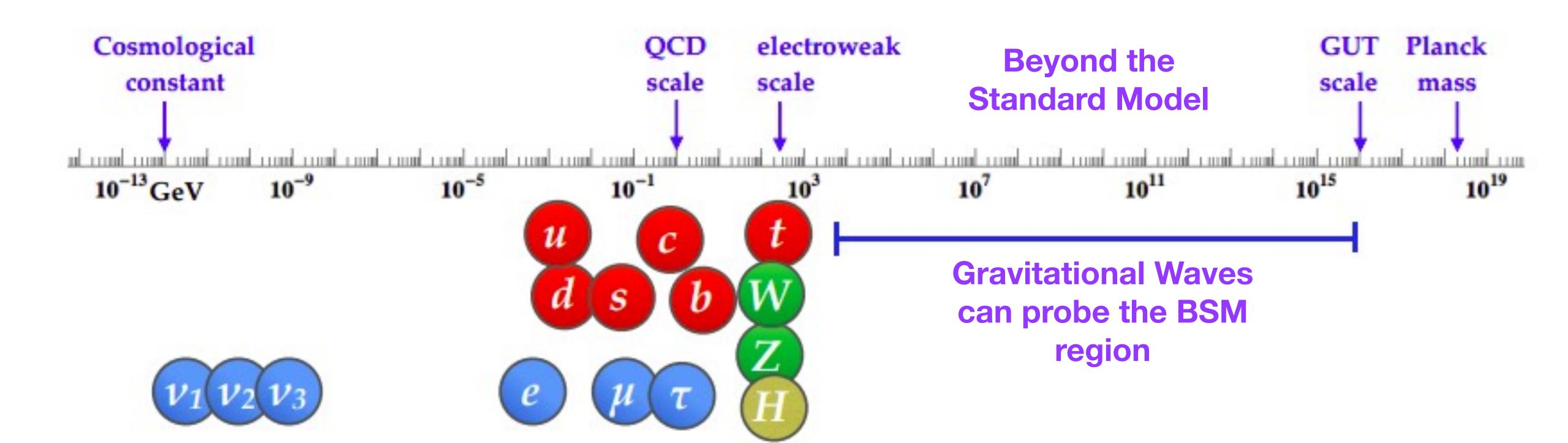
waves

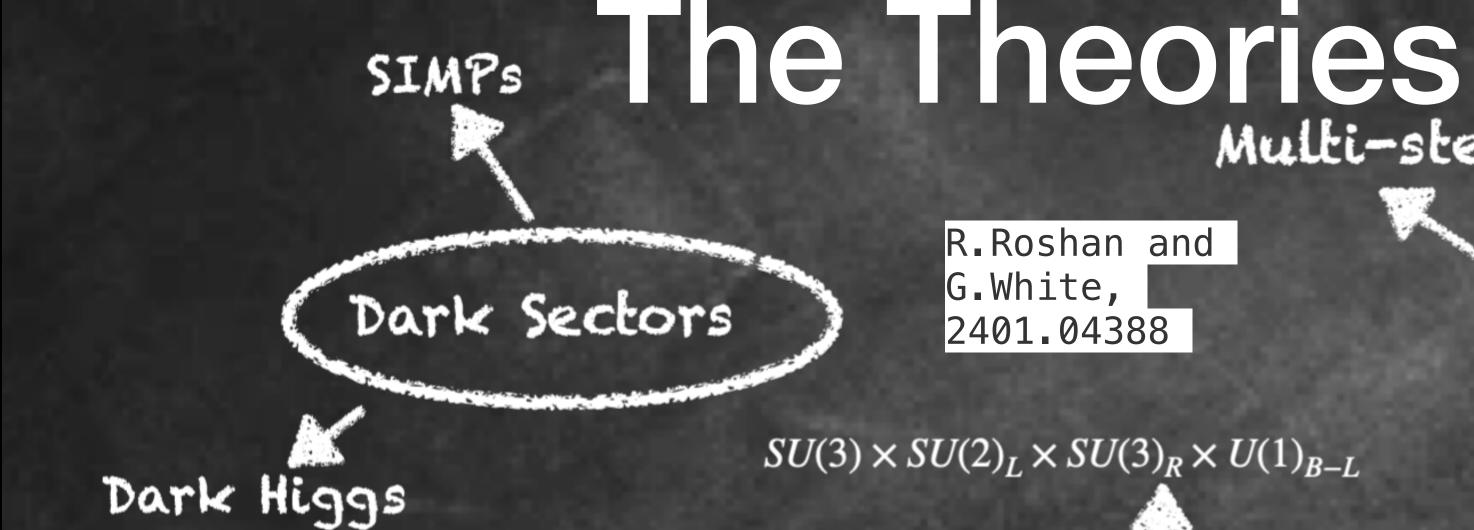
Big Bang plus 15 Billion Years





Gravitational Waves are sensitive to scales up to the Planck scale (well beyond the reach of colliders)





 $SU(4)_{\rm PS} \times SU(2)_L \times U(1)_R$



 $G_{SM} \times U(1)_{B-L}$

Top down

model building

Dynamical QCD Scale

Lepton asymmetry Dynamical quark

QCD

mass

Multi-step

Vector-Like Fermions

 $SU(5) \times U(1)_X$

New scalars

Supersymmetry

breaking

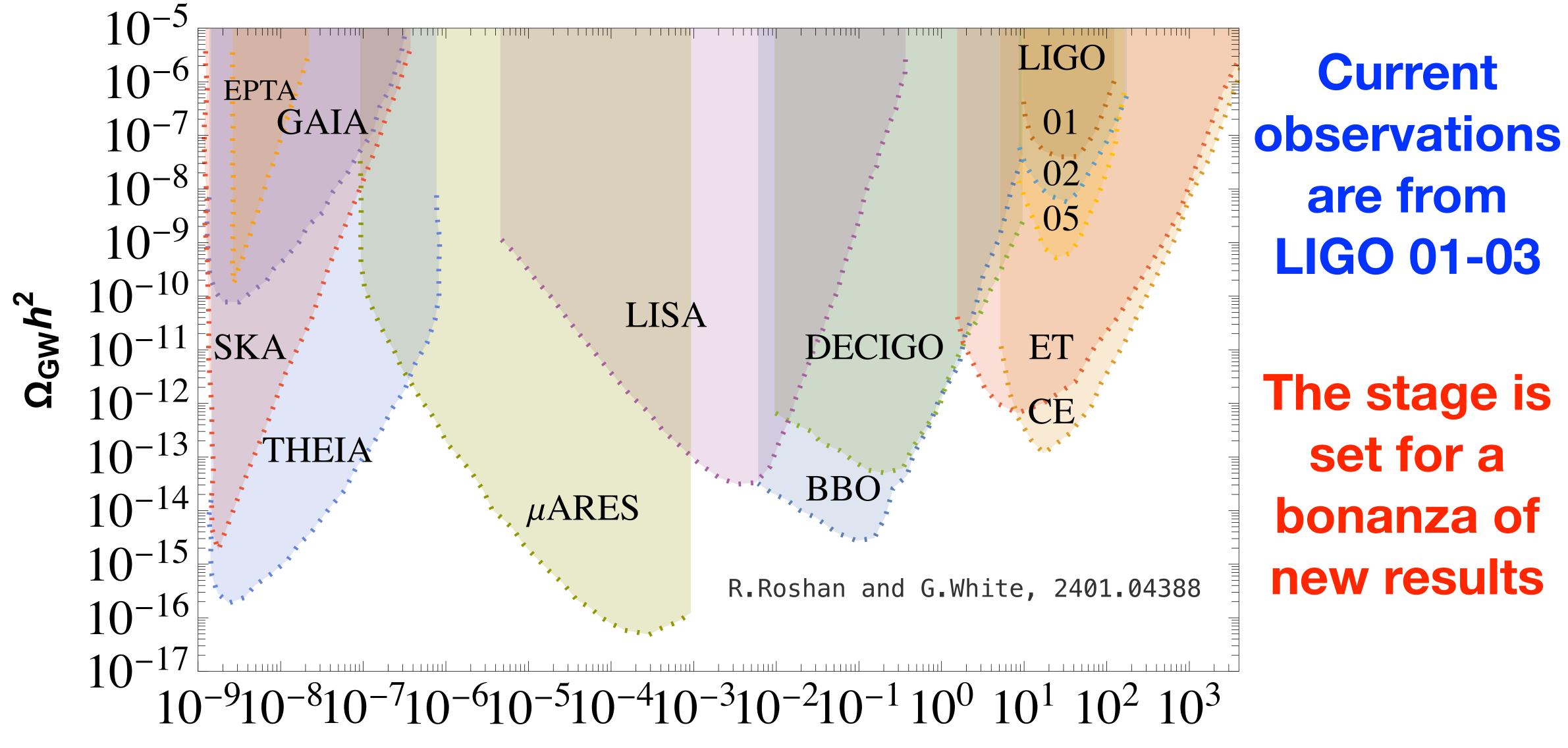
EWPT

Tree Level

Loop



The Detectors



f [Hz]



Gravitational Waves are sensitive to :

- First Order Phase transitions FOPT (e.g. QCD)
- Cosmic Strings CS (e.g. U(1) sym breaking)
- Domain Walls DW (e.g. Z₂ sym breaking)
- Inflation (e.g. with a kink or hybrid)
- Many other effects (e.g. PBHs,...)
- See other talks: Wang; Avgoustidis; Vikman; Papanikolaou; Goshal; Balazs; Zhou; Ota; Khalil; Mahapatra...

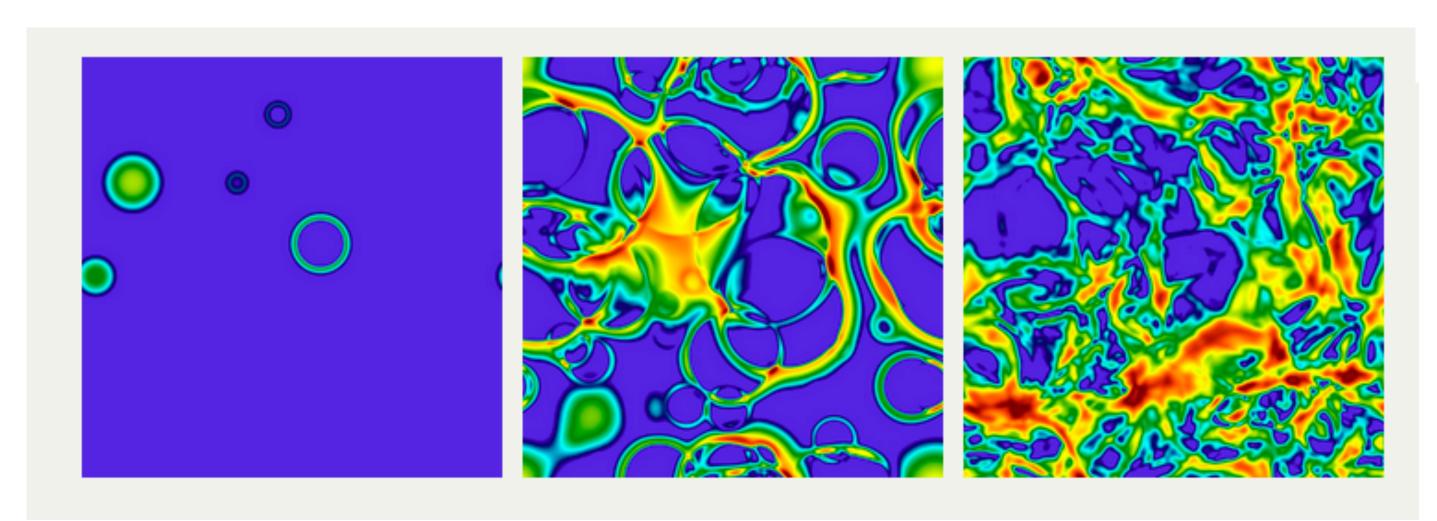
In this talk we are interested in a few **BSM examples**



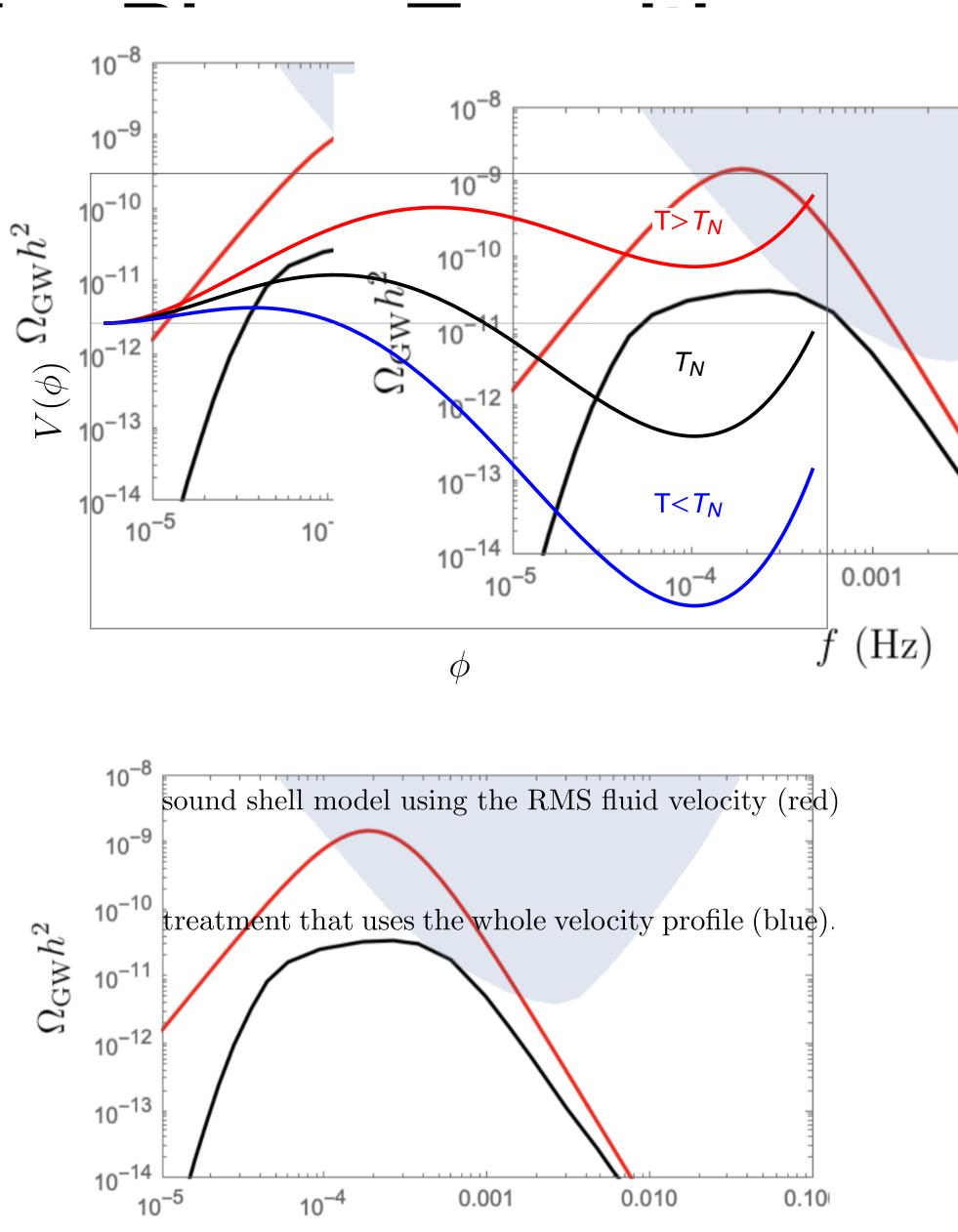
Gravitational Waves from First Ord

Phase Transitions:

- Bubbles nucleate and grow.
- Expand in plasma.
- Bubbles and fronts collide - violent process.
- Sound Waves left behind in thermal plasma.
- Turbulence, damping.

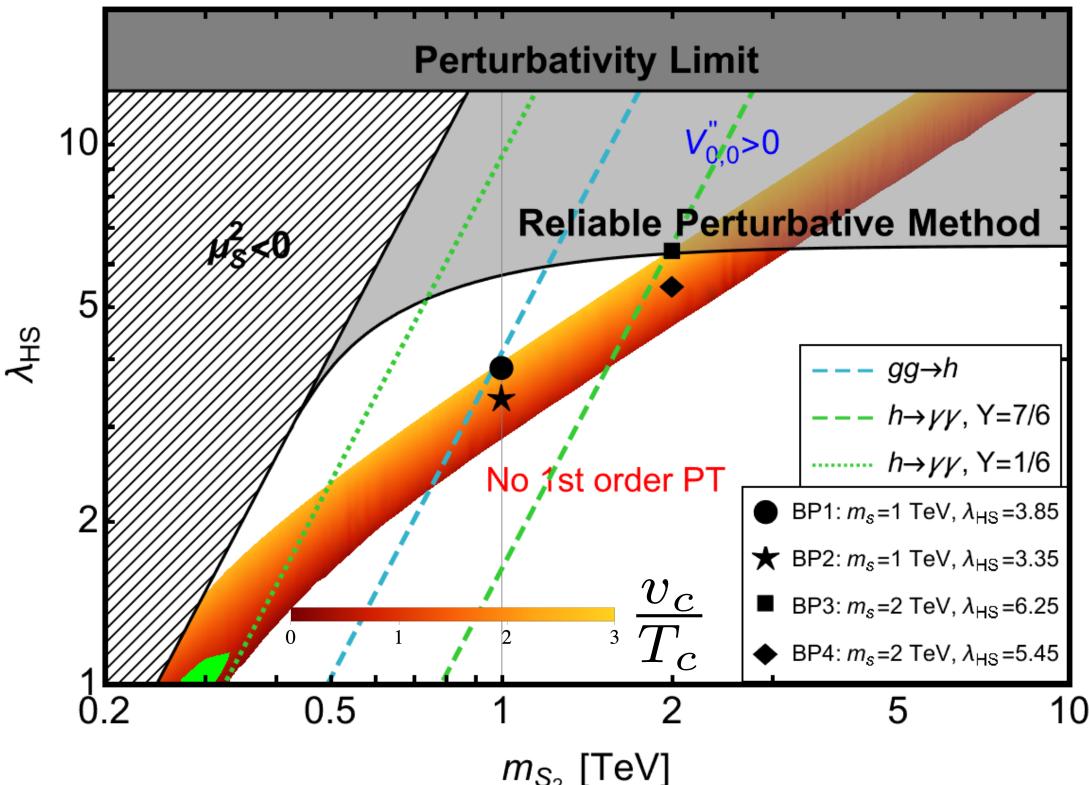


$\Omega_{\rm tot}(f) = \Omega_{\rm coll}(f) + \Omega_{\rm sw}(f) + \Omega_{\rm turb}(f)$



f (Hz)

B.Fu and S.F.K., 2209.14605 GW from leptoquark induced FOPT $V_0 = -\mu^2 |H|^2 + \lambda_H |H|^4 + \mu_S^2 |S_a|^2 + \lambda_S |S_a|^4 + 2\lambda_{HS} |H|^2 |S_a|^2$ Leptoquark singlet S₁, $V_{\text{eff}}(h,T) = V_0 + \Delta V_0^{1-\text{loop}}(h) + \Delta V_T^{1-\text{loop}}(h,T)$ doublet S₂ or triplet S₃ **Benchmark Point 1** μ Ares 10^{-6} **Perturbativity Limit** BBO 10^{-8} DECIGO V"/>0 10 LISA 10^{-10} **Reliable Perturbative Method** λ_{HS} gg→h 10^{-14} $---h \rightarrow \gamma \gamma$, Y=7/6 $\dots h \rightarrow \gamma \gamma, Y=1/6$ 10^{-16} No 1st order PT $\Omega_{\rm tot}h^2$ • BP1: m_s =1 TeV, λ_{HS} =3.85 $\Omega_{\rm sw}h^2$ **★** BP2: m_s =1 TeV, λ_{HS} =3.35 10^{-18} v_{c} BP3: m_s=2 TeV, λ_{HS}=6.25 $\frac{1}{3}\overline{T}_{c}$ ---- $\Omega_{\rm turb}h^2$ • BP4: m_s =2 TeV, λ_{HS} =5.45 10^{-20} 10^{-2} 10^{-4} 10^{-3} 10^{-1} 10^{1} 10^{-5} 0.2 0.5 5 10 m_{S_2} [TeV] f [Hz]









• NANOGrav Collaboration, arXiv:2306.16213

Pulsar Timing Arrays first observation of stochastic GW background



• NANOGrav Collaboration, arXiv:2306.16213

Pulsar Timing Arrays first observation of stochastic 10^{-1} GW background

 10^{-3} -10⁻⁵ -

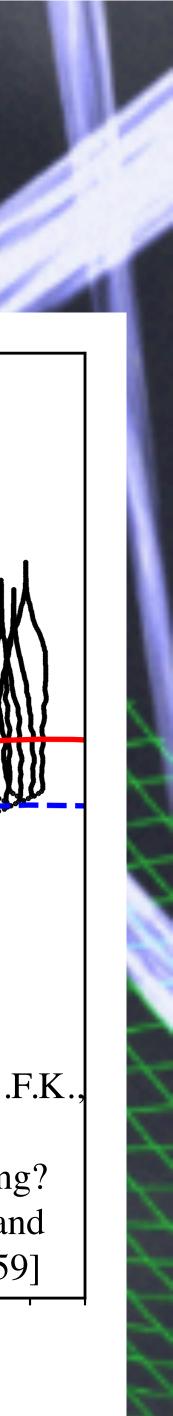
 10^{-7} $\Omega_{GW} h^2$ 10⁻⁹ ⊣

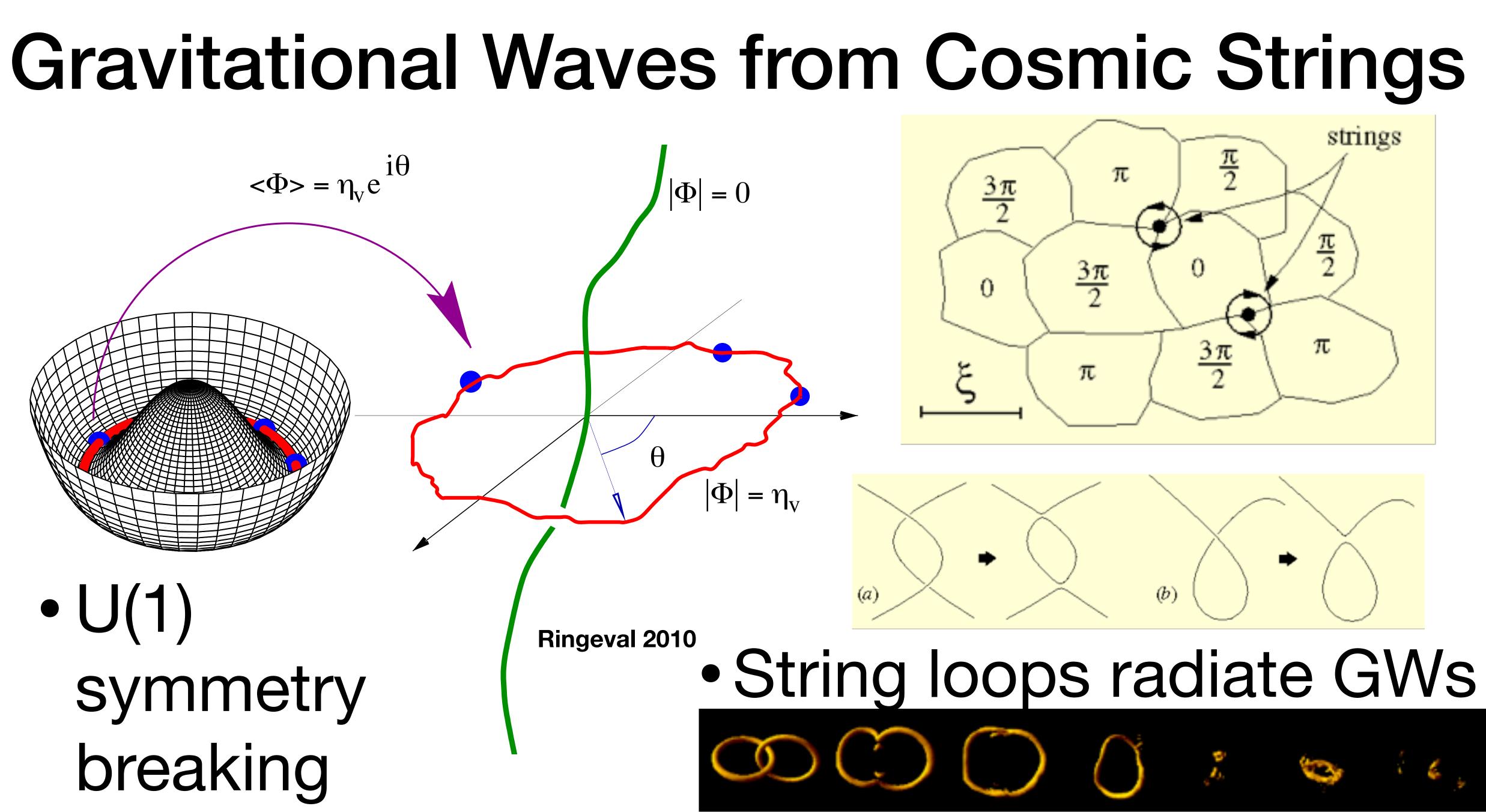
 10^{-11} -

 10^{-13} -

 10^{-15}

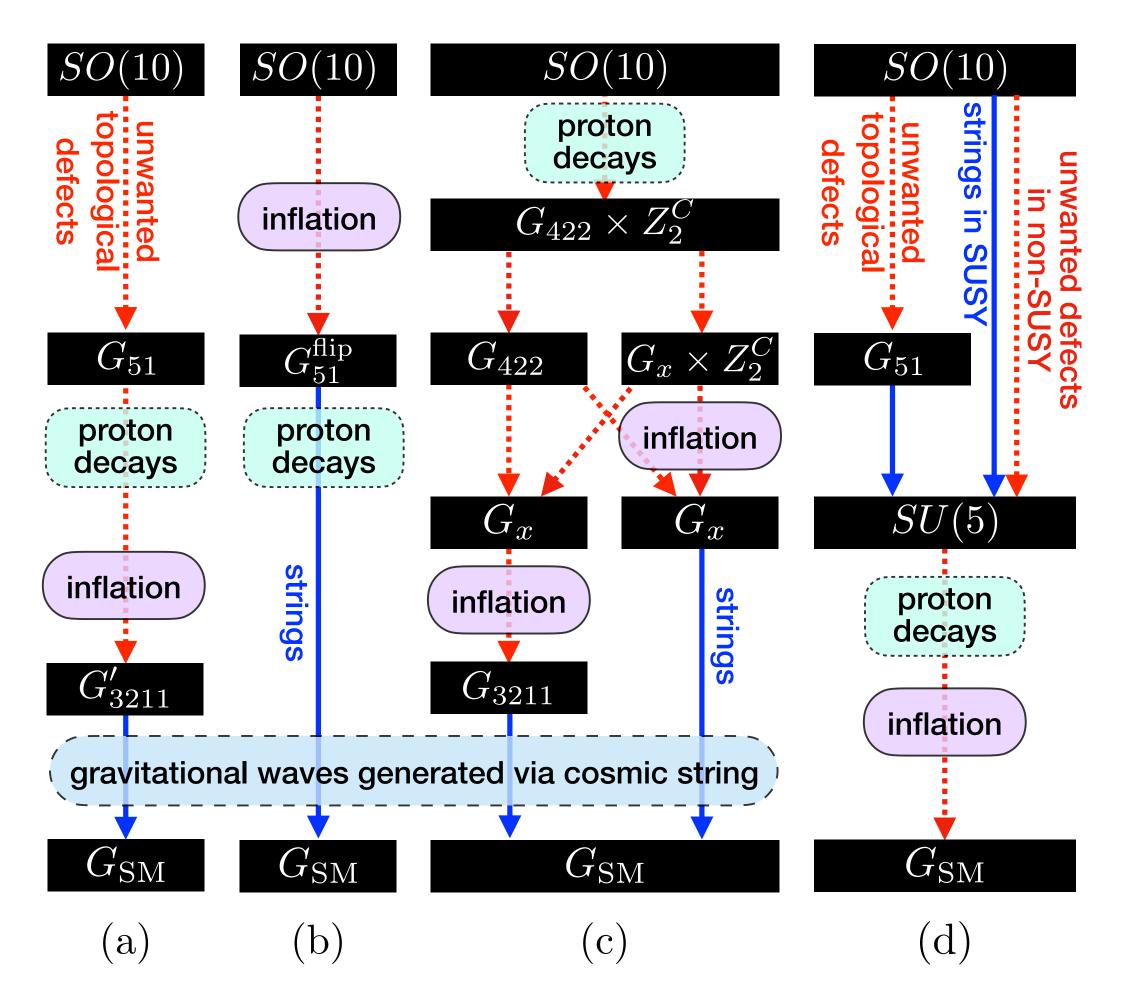
| | | nalytical Iumerical | | | ••••••••••••••••••••••••••••••••••••••• |
|----|----------------------------|------------------------|--------------------------|---------------|---|
| | Full Sound shell | | Broken p | ower law fit | |
| | Parameter Best fit value | | Parameter Best fit value | | |
| | $lpha_n$ | 0.85 | $lpha_n$ | 0.89 | |
| | eta/H_* | 42 | eta/H_* | 5.17 | |
| | T_n | 133 MeV | T_n | 142 MeV | |
| | v_w | 0.09 | v_{w} | 0.67 | |
| - | $\chi_{ m fit}$ | 1.4 | $\chi_{	ext{fit}}$ | 1.59 | |
| _ | | | | PHH H | |
| | | | T. | Ghosh, A.Gh | oshal, H.K.Guo, F.Hajkarim, S.I |
| | | | K K | .Sinha, X.Wa | ng and G.White, |
| - | | | | Did we hear t | he sound of the Universe boiling |
| | | | А | nalysis using | the full fluid velocity profiles ar |
| | | | | • • | -year data,'' [arXiv:2307.02259 |
| -1 | ľ | | | 10^{-8} | |
| | | | | f(Hz) | |

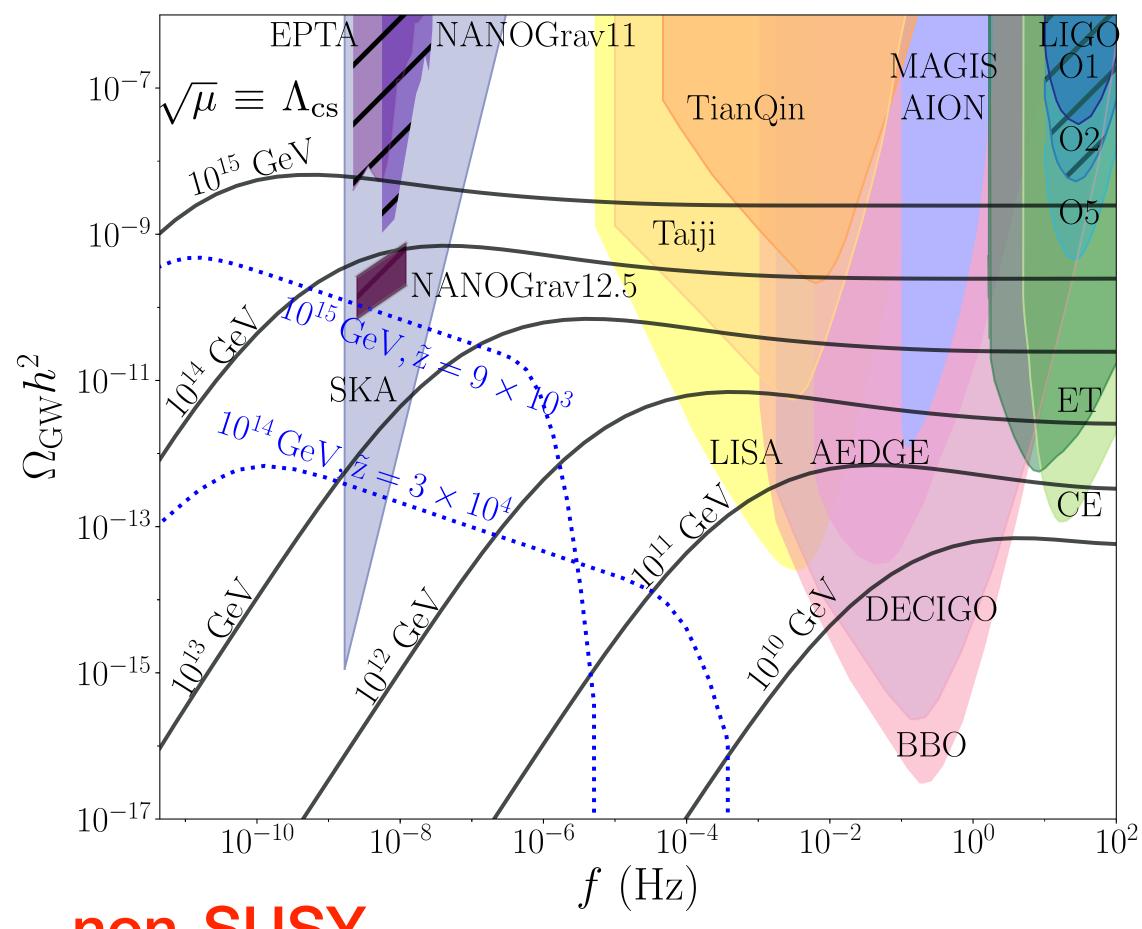




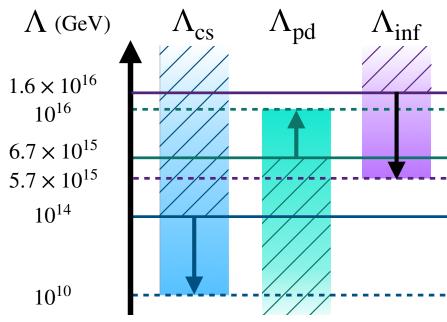
GWs via CSs from gauged U(1)_{B-L} in SO(10) GUTs

S.F.K., S.Pascoli, J.Turner and Y.L.Zhou, 2005.13549; 2106.15634; w/ Marsili 2209.00021; 2308.05799



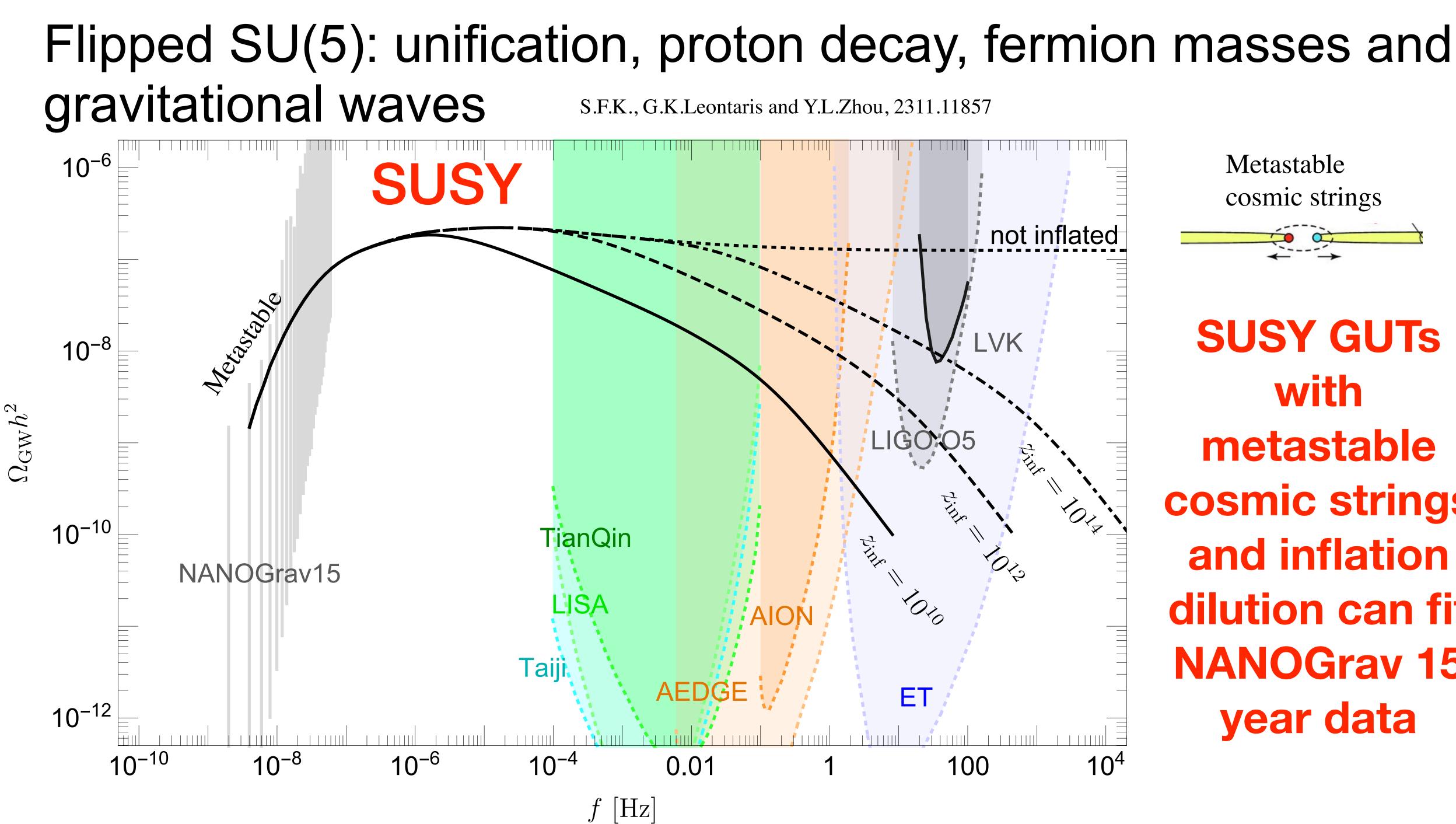


non-SUSY



| | Observables | | Proton decays | | | | | |
|--|-------------|----------|---|--|--|--|--|--|
| | | | $p \rightarrow \pi^0 e^+$ observed \Rightarrow non-SUSY contribution ind | | | | | |
| | | Observed | types (a) and (c) favoured types (b) and (d) excluded | | | | | |
| | GWs | Marginal | types (a) and (c) favoured type (d) excluded type (b) allowed if p → K⁺ν̄ not observed and Λ | | | | | |





Metastable cosmic strings

SUSY GUTs with metastable **cosmic strings** and inflation dilution can fit **NANOGrav 15** year data

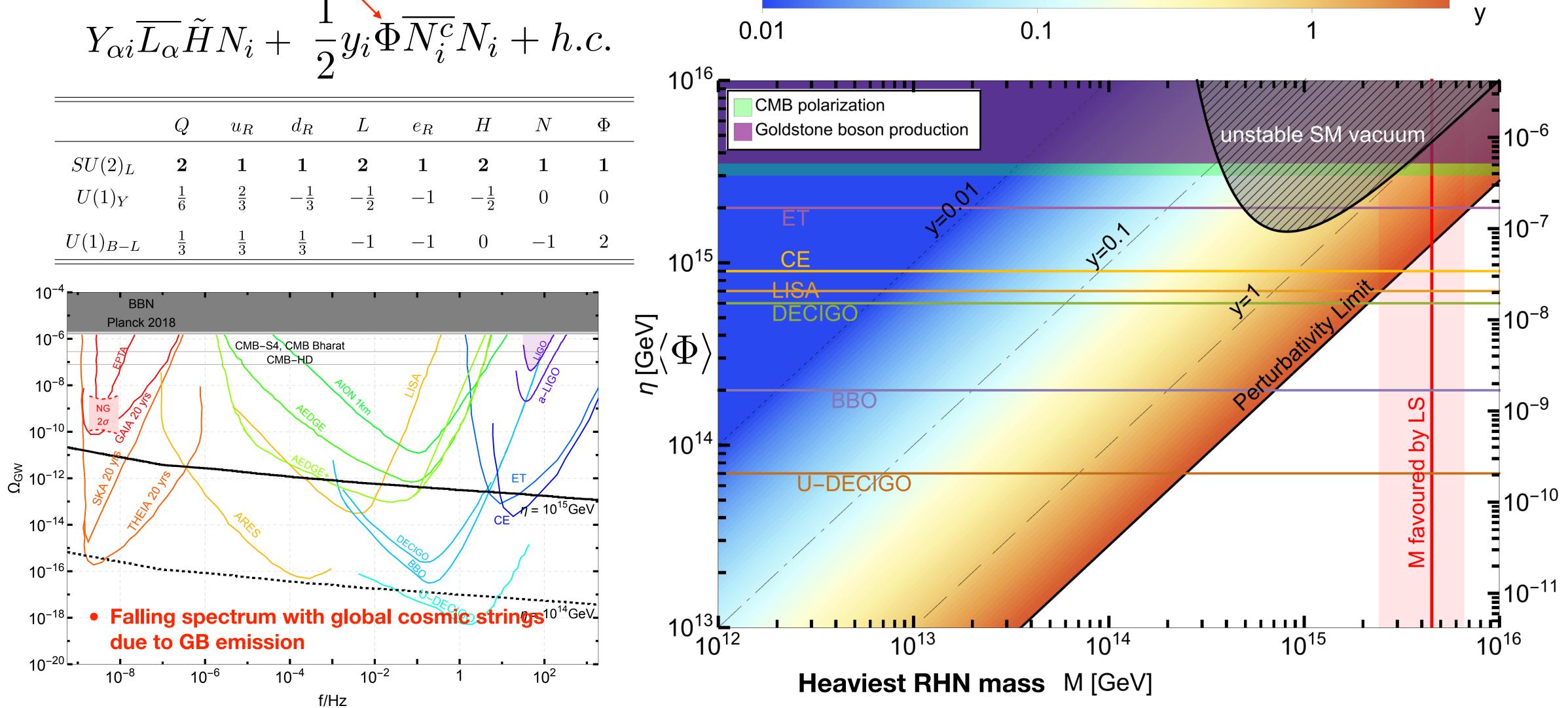






Cosmic string gravitational waves from global $U(1)_{B-L}$ symmetry breaking as a probe of the type I seesaw scale Majoron breaks U(1)B-L

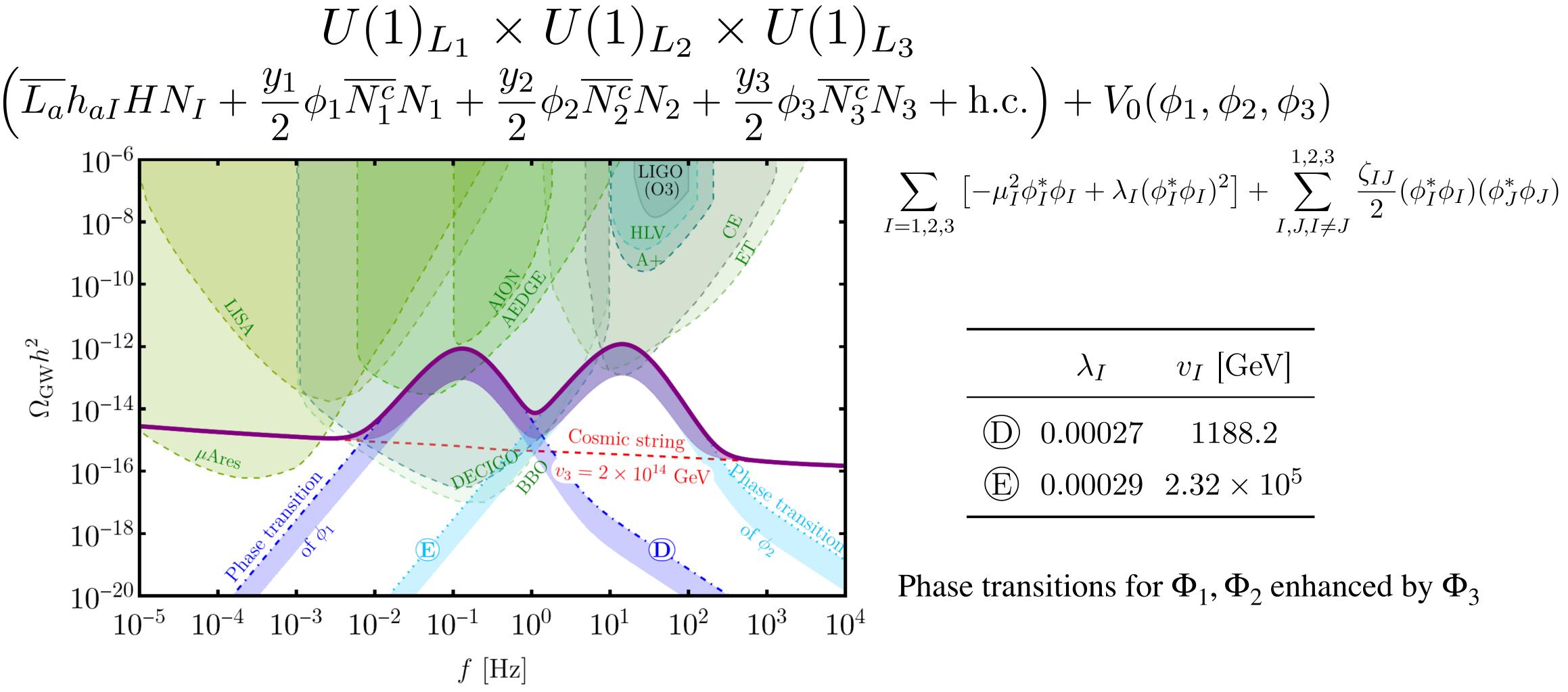
| | Q | u_R | d_R | L | e_R | H | N | Φ |
|--------------|---------------|---------------|----------------|----------------|-------|----------------|----|--------|
| $SU(2)_L$ | 2 | 1 | 1 | 2 | 1 | 2 | 1 | 1 |
| $U(1)_Y$ | $\frac{1}{6}$ | $\frac{2}{3}$ | $-\frac{1}{3}$ | $-\frac{1}{2}$ | -1 | $-\frac{1}{2}$ | 0 | 0 |
| $U(1)_{B-L}$ | $\frac{1}{3}$ | $\frac{1}{3}$ | $\frac{1}{3}$ | -1 | -1 | 0 | -1 | 2 |



B.Fu, A.Ghoshal and S.F.K.2306.07334



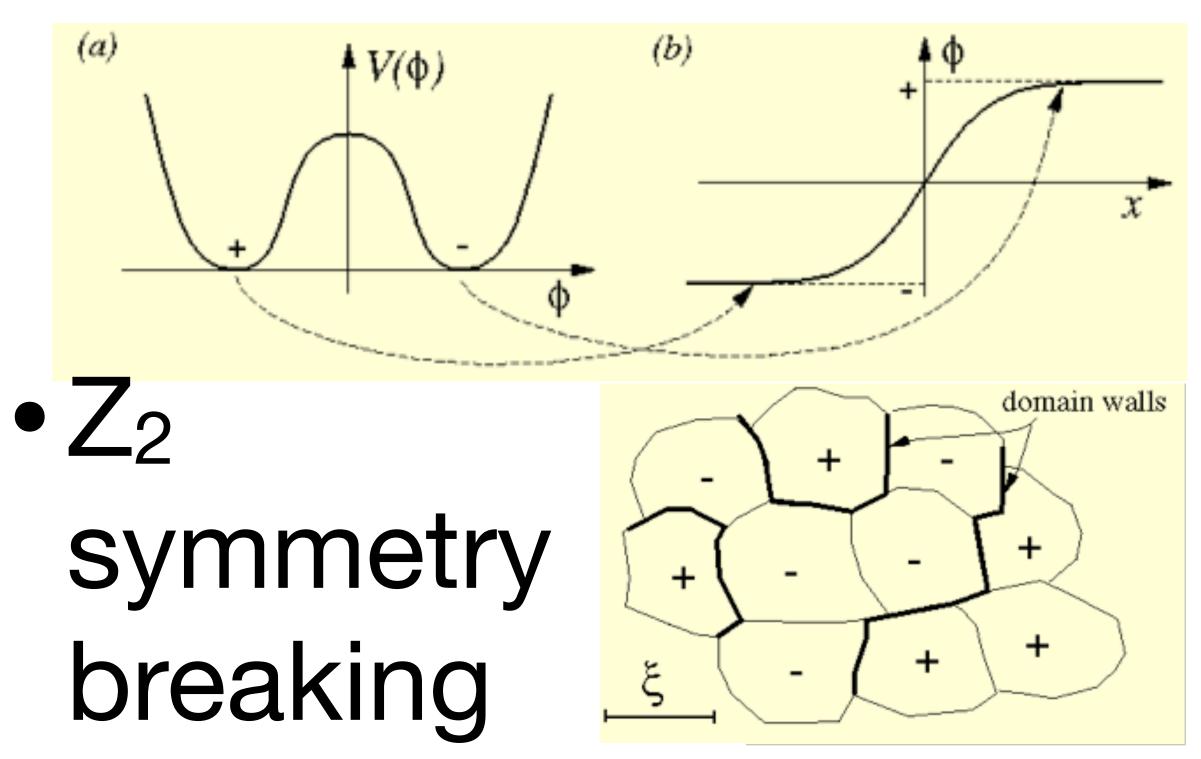
Ц С

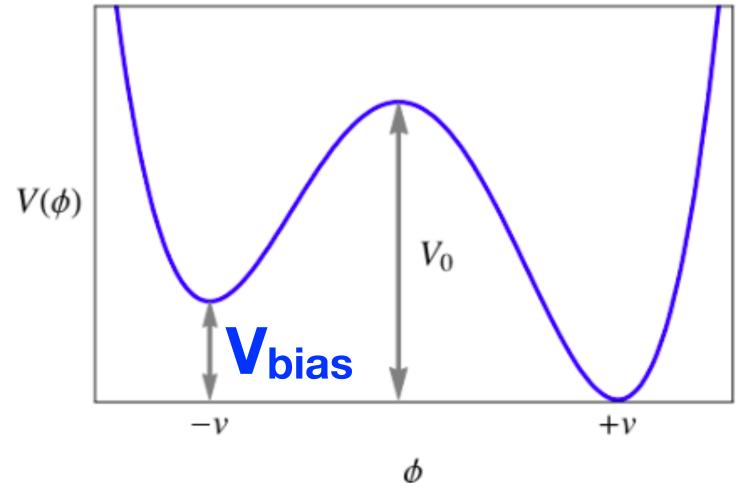


P.Di Bari, S.F.K. and M.H.Rahat, 2306.04680

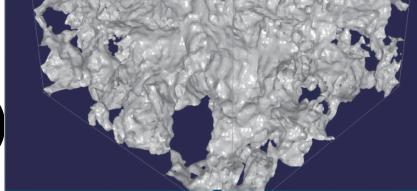
Gravitational waves from phase transitions and cosmic strings in neutrino mass models with multiple Majorons

Gravitational Waves fro

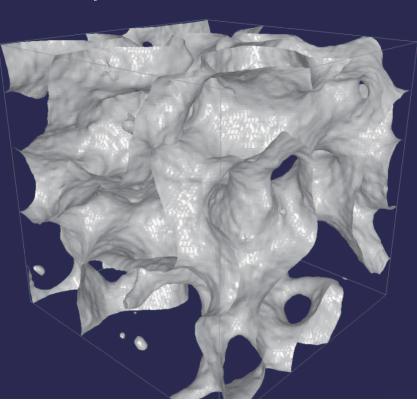




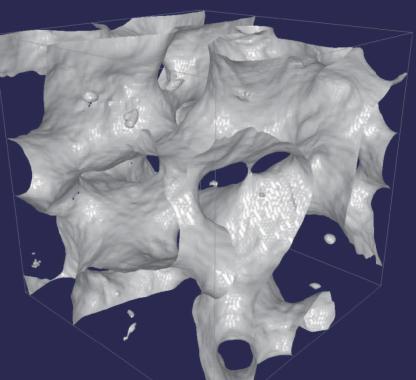
Introduce Vbias to allow domain walls to decay (otherwise dominate energy density of Universe)



 $t = 71t_i$

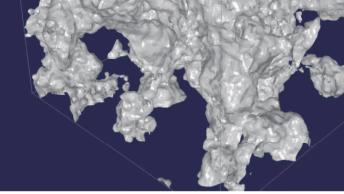


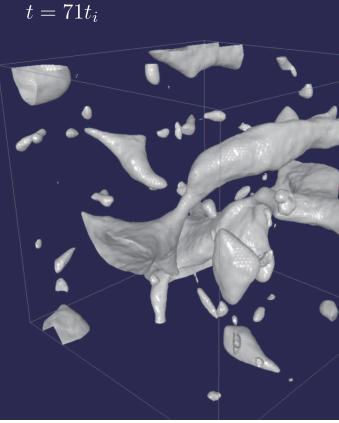
 $t = 111t_{i}$

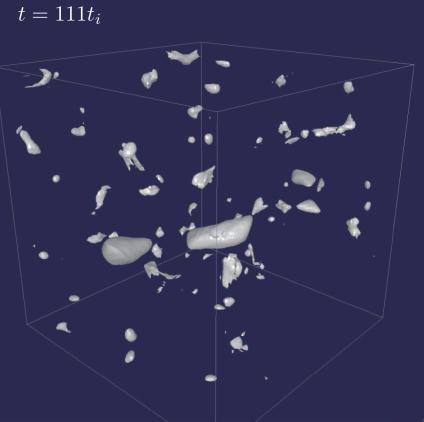


Stable DW on the left with V_{bias} =0

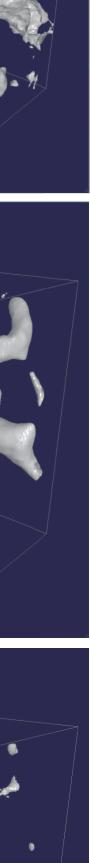
a







Unstable DW decay via GWs due to V_{bias}





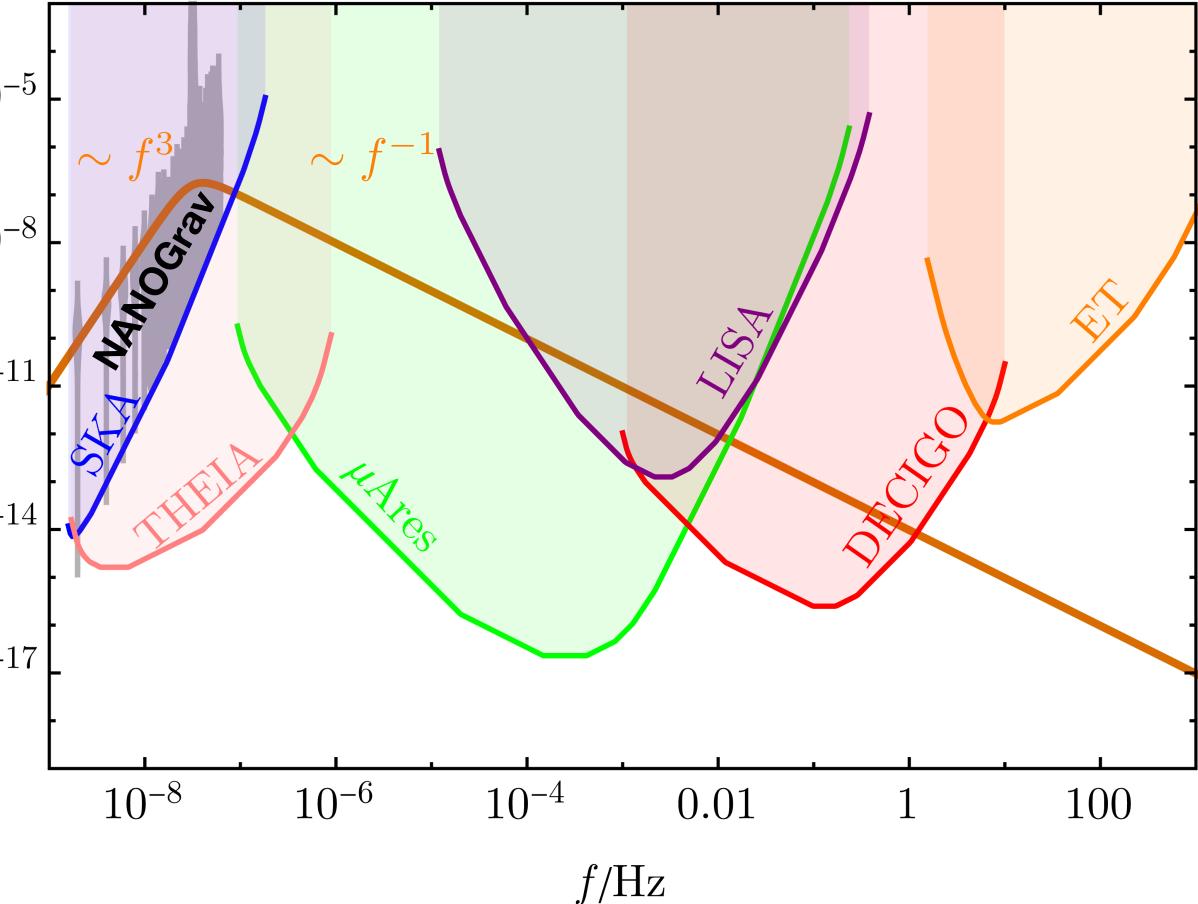


S.F.K., R.Roshan, X.Wang, G.White and M.Yamazaki, 2308.03724; 2311.12487

Quantum gravity effects on dark matter and gravitational waves **Peak occurs when volume pressure ~ V**_{bias} A tale of two Z₂'s 10^{-5} Stable Dark Sector $\mathbf{S_1}: \mathcal{Z}_2^{(\mathbf{DW})}$ **Quantum Gravity** 10^{-8} $\mathbf{S_2}: \mathcal{Z}_2^{(\mathbf{DM})}$ $h^2 \Omega_{ m GW}$ GW 10^{-11} Both broken by QG effects 10^{-14} $\mathcal{L}_{\mathbb{Z}_2} = \frac{1}{\Lambda_{OC}} \mathcal{O}_5 \quad \longrightarrow \text{Vbias}$ 10^{-17}

Due to instanton effects

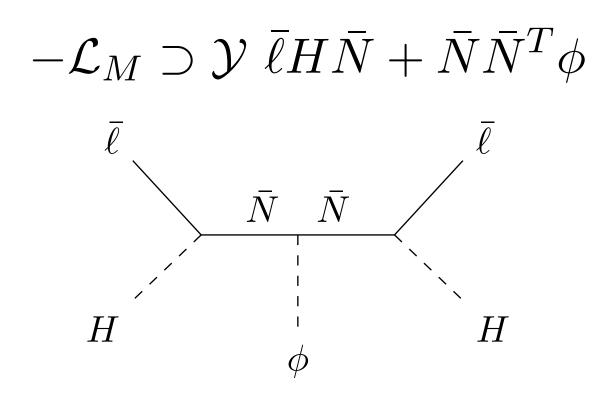
 $\Lambda_{\rm QG} \sim M_{\rm Pl} e^{\mathcal{S}} \gg M_{\rm Pl}$







mass with gravitational waves Majorana seesaw



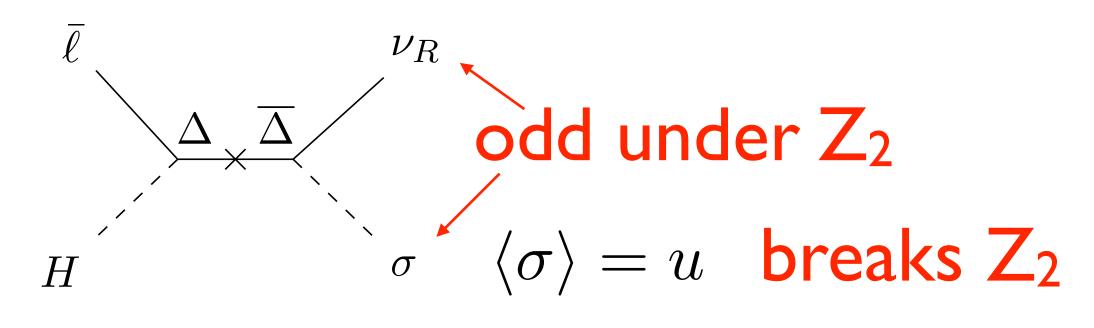
$$\mathcal{M}_M = \frac{1}{\sqrt{2}} v^2 \ \mathcal{Y} \ \mathcal{M}_N^{-1} \ \mathcal{Y}^T$$

Gauged $U(I)_{B-L}$ broken →Cosmic strings

S.F.K., D.Marfatia and M.H.Rahat, 2306.05389 Toward distinguishing Dirac from Majorana neutrino

Dirac seesaw

 $-\mathcal{L}_D \supset \mathcal{Y}_L \overline{\ell} H \Delta_R + \mathcal{Y}_R \overline{\Delta}_L \sigma \nu_R + \mathcal{M}_\Delta \overline{\Delta} \Delta$

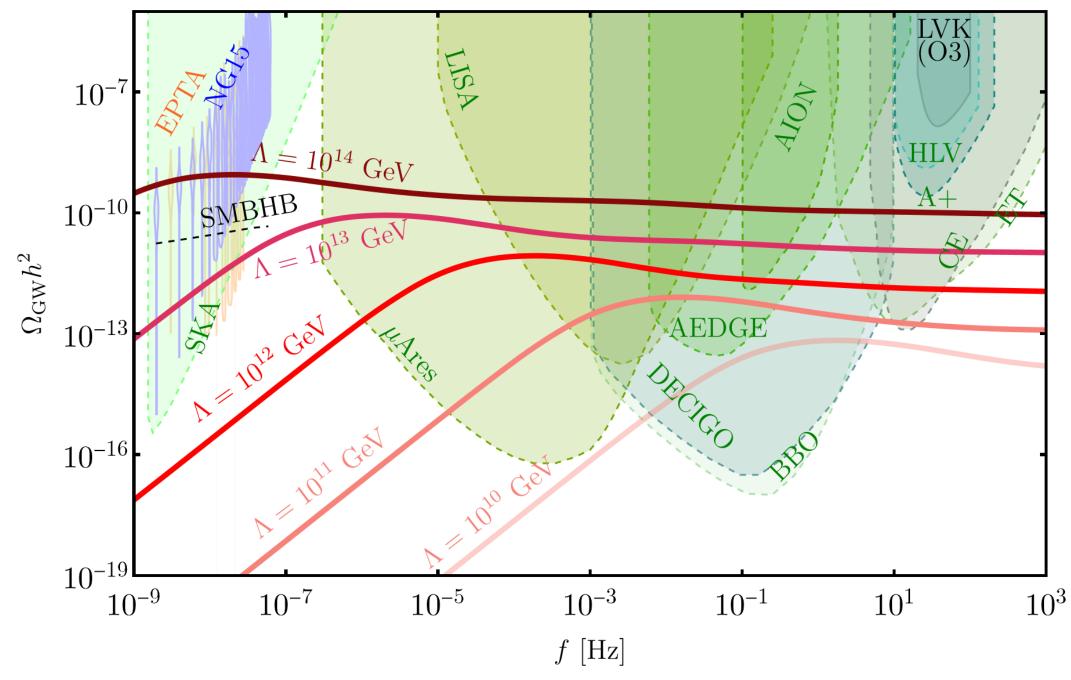


$$\mathcal{M}_D = \frac{1}{\sqrt{2}} v u \mathcal{Y}_L \mathcal{M}_\Delta^{-1} \mathcal{Y}_R$$

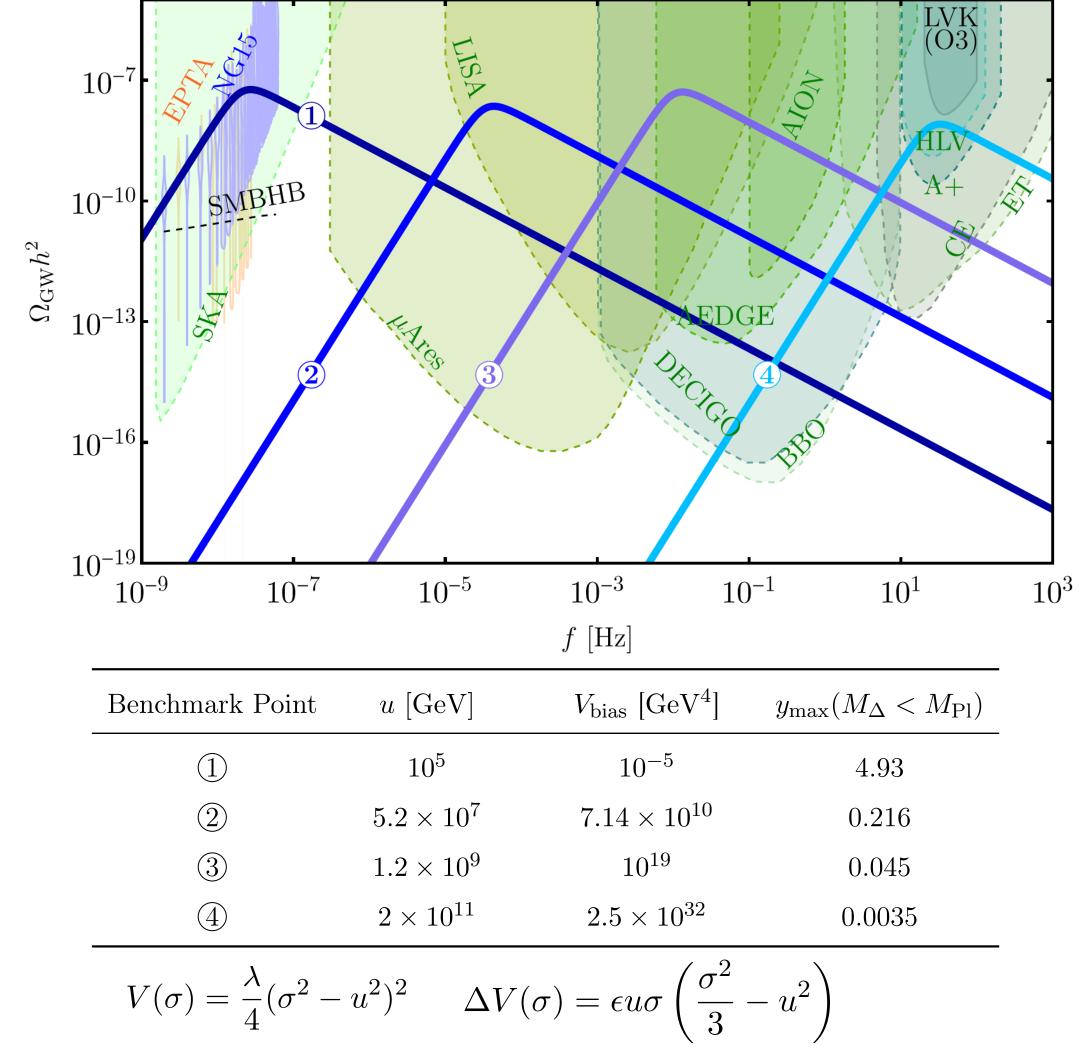
Global U(1)_{B-L} preserved Z_2 broken \rightarrow Domain Walls



S.F.K., D.Marfatia and M.H.Rahat, 2306.05389 Toward distinguishing Dirac from Majorana neutrino mass with gravitational waves Majorana seesaw Dirac seesaw LVK (O3)



Majorana vs Dirac can be distinguished from shape of GW spectrum - - Dirac is better fit to NANOGrav



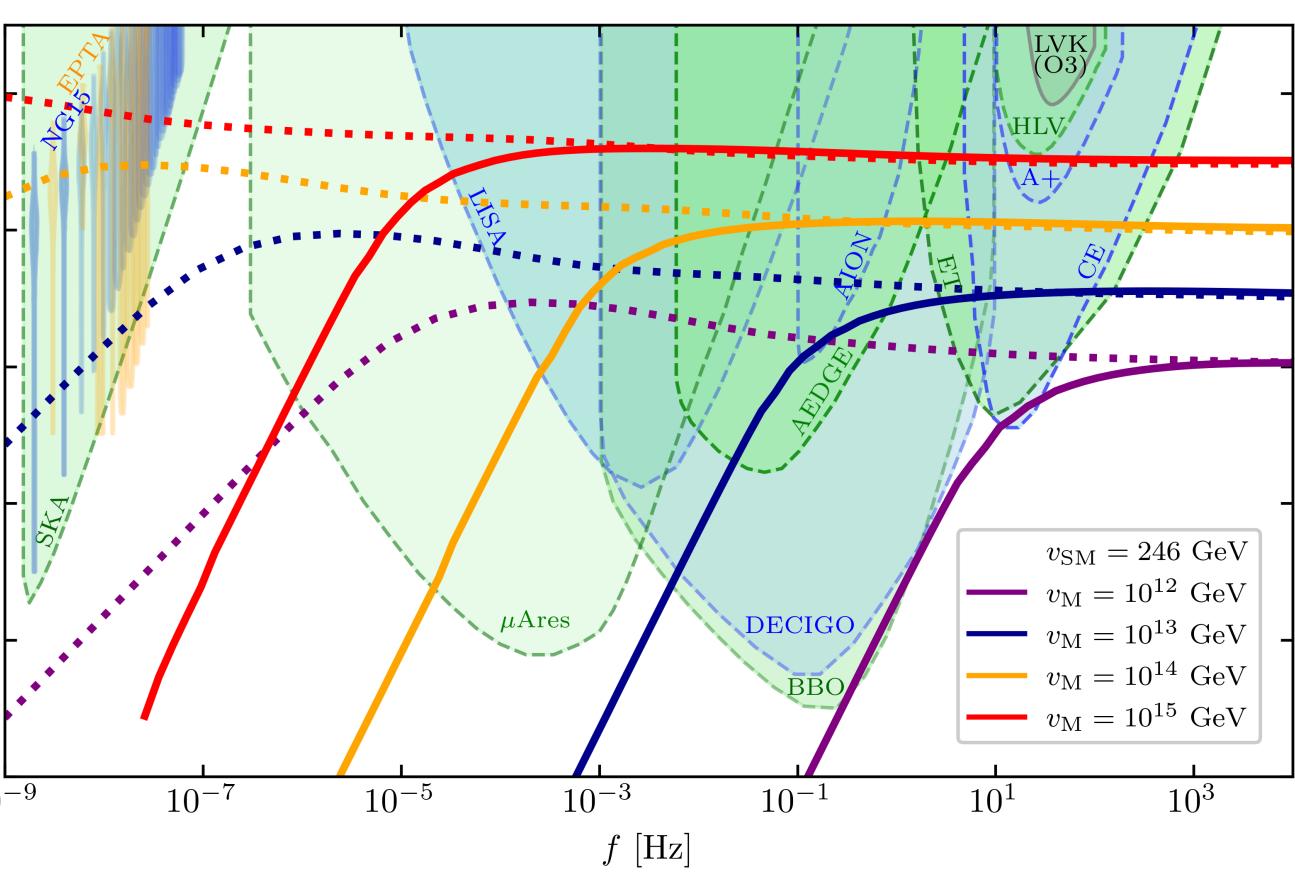


Type-I two-Higgs-doublet model and gravitational waves from domain walls bounded by strings

 $Y_u \overline{Q} \tilde{\Phi}_2 u_R + Y_d \overline{Q} \Phi_2 d_R + Y_e \overline{L} \Phi_2 e_R + Y_N \overline{L} \tilde{\Phi}_2 N_R + y_N \phi N_R N_R$ 10^{-8}

| | | | | | | | | | | 10 |
|-------------------------------------|---------------|------------------------|---------------|----------------|------------------------------------|--------------------------------|---------------|---------------|--------------|---|
| | $u_{R\beta}$ | d_{Reta} | Q_{lpha} | L_{lpha} | $e_{R\beta}$ | N_{Reta} | Φ_2 | Φ_1 | ϕ | |
| $SU(2)_L$ | 1 | 1 | 2 | 2 | 1 | 1 | 2 | 2 | 1 | 10^{-10} |
| $U(1)_Y$ | $\frac{2}{3}$ | $-\frac{1}{3}$ | $\frac{1}{6}$ | $-\frac{1}{2}$ | -1 | 0 | $\frac{1}{2}$ | $\frac{1}{2}$ | 0 | |
| $U(1)_R$ | 1 | -1 | 0 | 0 | -1 | 1 | 1 | -1 | -2 | $\overset{\mathbf{r}}{\prec}$ 10 ⁻¹² |
| residual Z_2 | _ | — | + | + | _ | _ | _ | _ | + | $^{\rm C}_{\rm C} h_{\rm D}^{\rm C}$ |
| | | | | | | | | | | 10^{-14} |
| U(1) | | $\langle \phi \rangle$ | . 7 | , | $\langle \Phi_i$ | $\stackrel{>}{\rightarrow}$ no | _+h | | \mathbf{r} | |
| | R | / | | 2 | | 7 110 | J [1] | 1115 | 5 | 10^{-16} |
| | | | | | | | | | | |
| | | string | θ+π | | | ç | | str | θ_{+} | wall 10^{-18} |
| | | | θ | | | | | | θ | 10 10 |
| | | | | V | Vilenk | in and | Ever | ett 19 | 982 | |
| (a) Before Z_2 symmetry breaking. | | | | | (b) After Z_2 symmetry breaking. | | | | king. | Sı |
| | | 000 | | •••• | th/ | + | \ /. | | | |
| DW decay without V bias! | | | | | | | | | | COIL |

B.Fu, A.Ghoshal, S.F.K. and M.H.Rahat, 2404.16931



urface tension in the walls causes the combined relic to decay earlier than strings





Conclusion

- FOPT at QCD scale can describe NANOGrav
- DW Z₂ w/QG bias; Majorana vs Dirac
- DW bounded by CS in 2HDM (type I)

• GWs can probe new physics BSM at HE, only a few examples here: FOPT, CS, DW (+combos)

CS U(1)_{B-L} gauged w/GUTs; global w/Majorons



