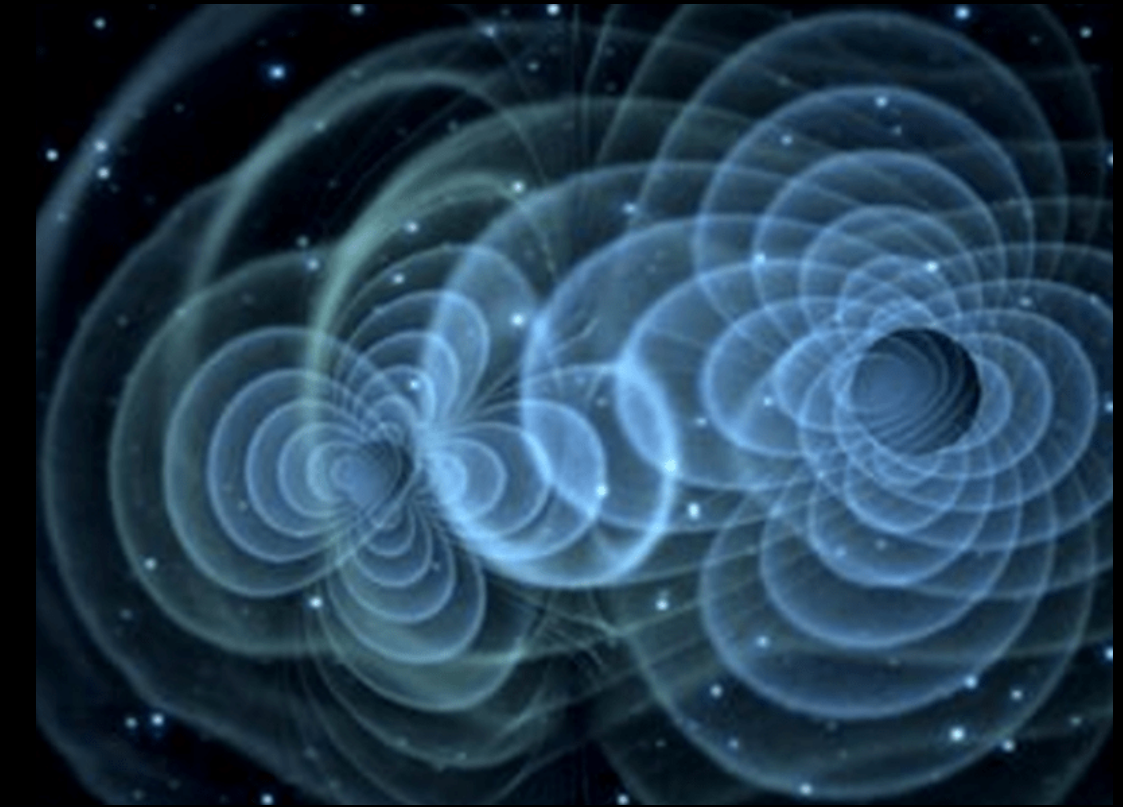




Gravitational Waves as Probes of New Physics



STEVE KING 02/09/24

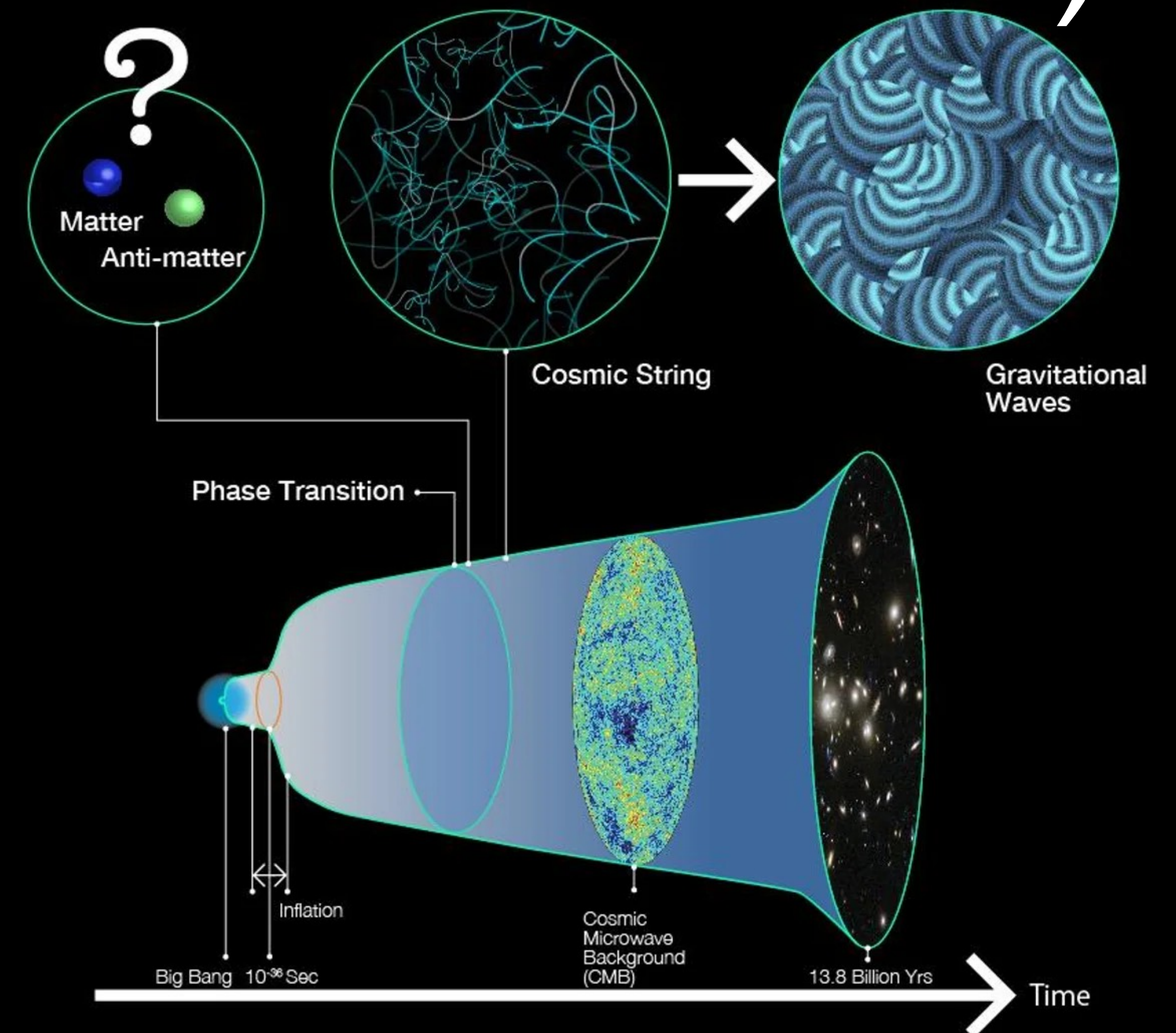
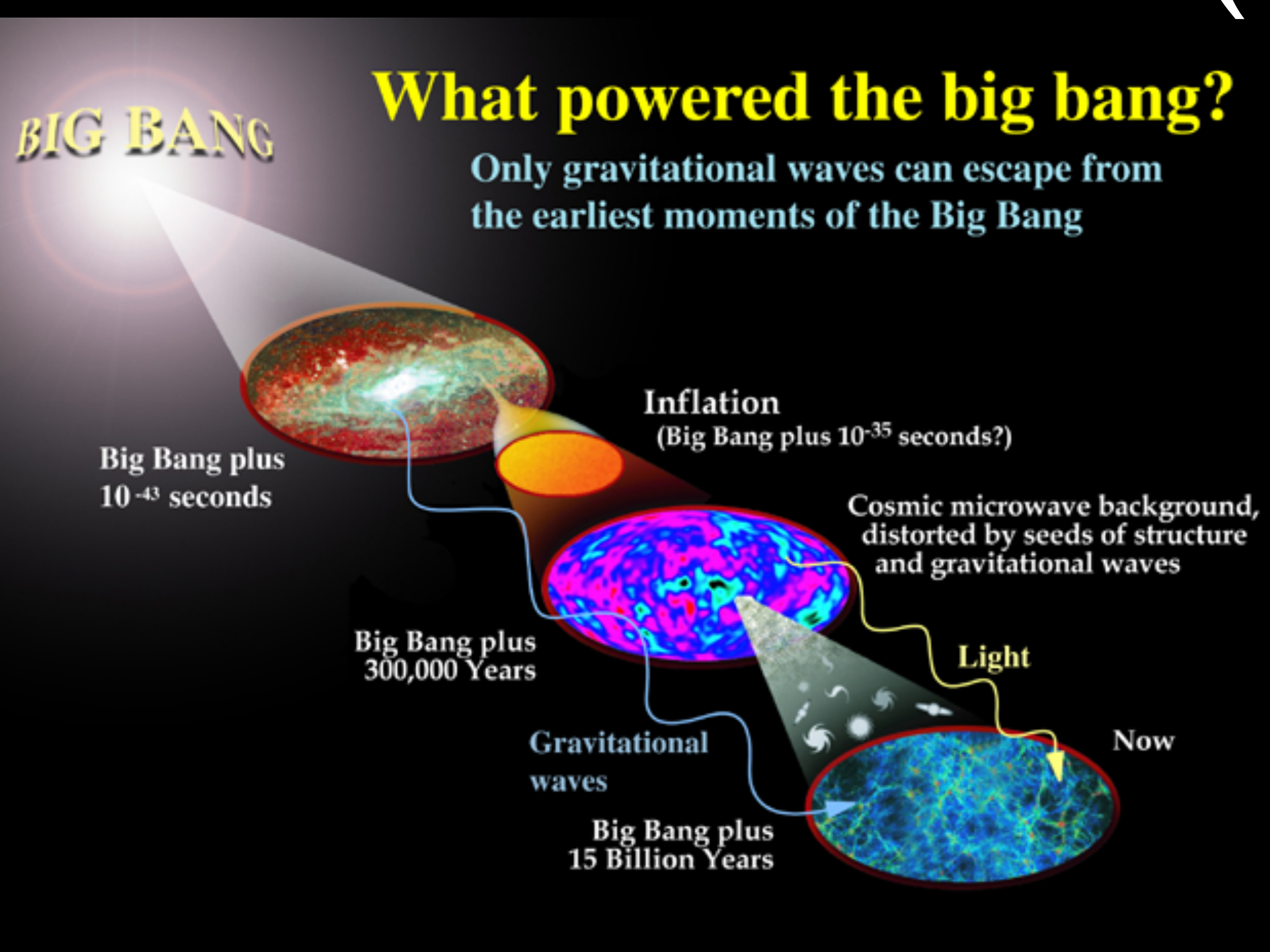
**Workshop on the Standard Model
and Beyond**

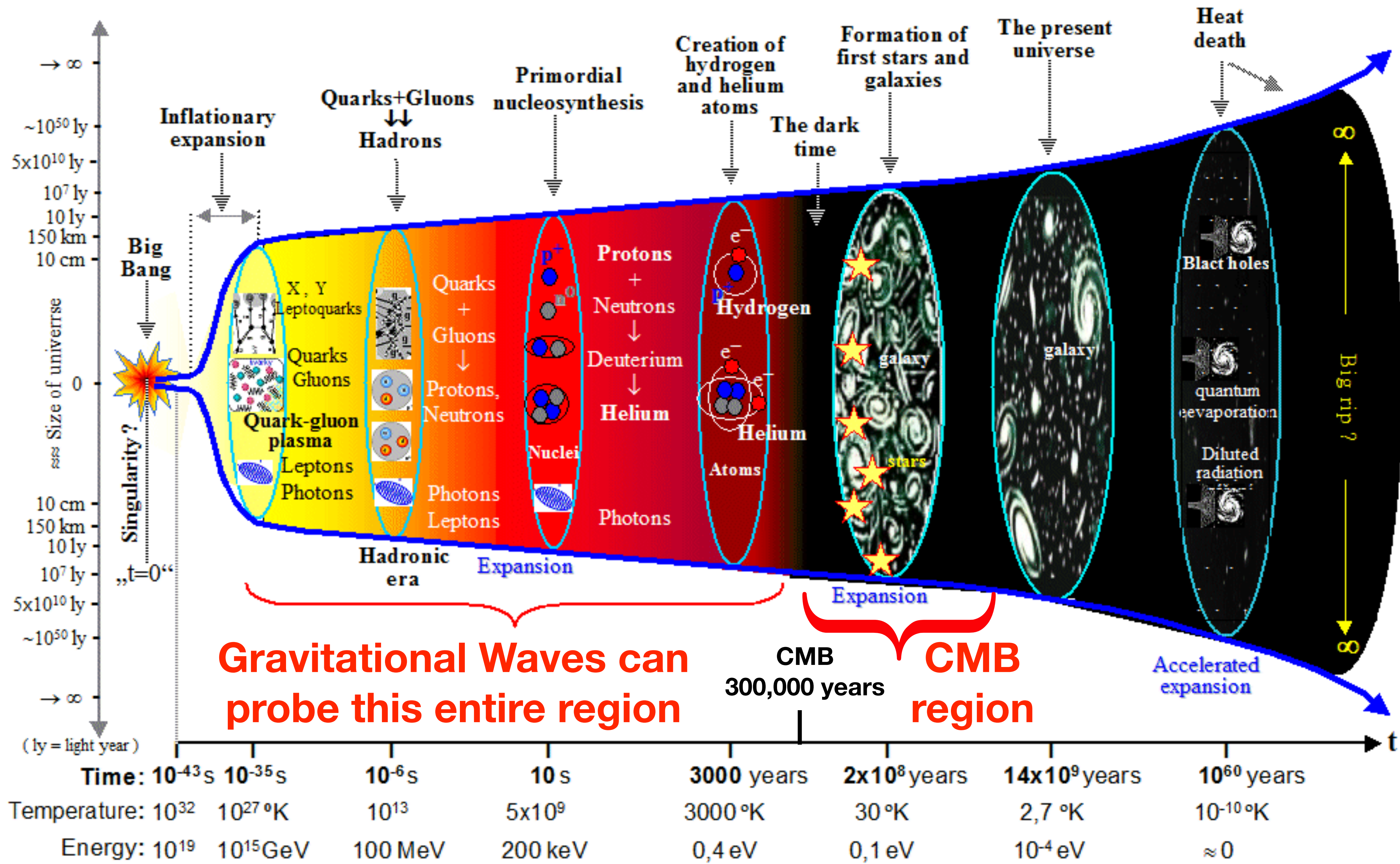


Gravitational Waves enable us to look back to the earliest moments of the Universe (back to Inflation)

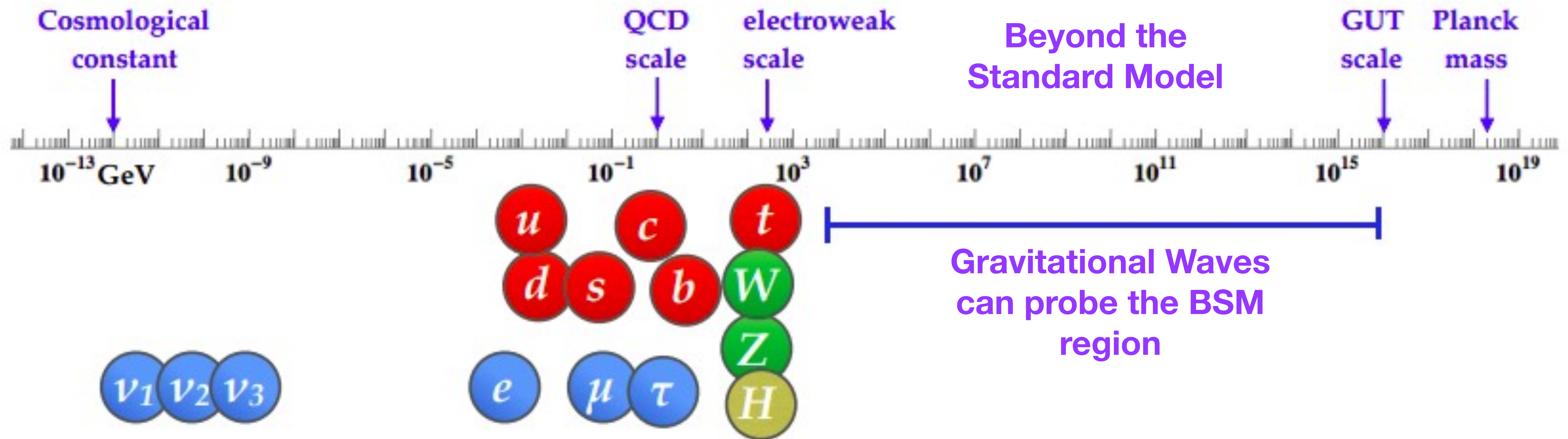
What powered the big bang?

Only gravitational waves can escape from the earliest moments of the Big Bang



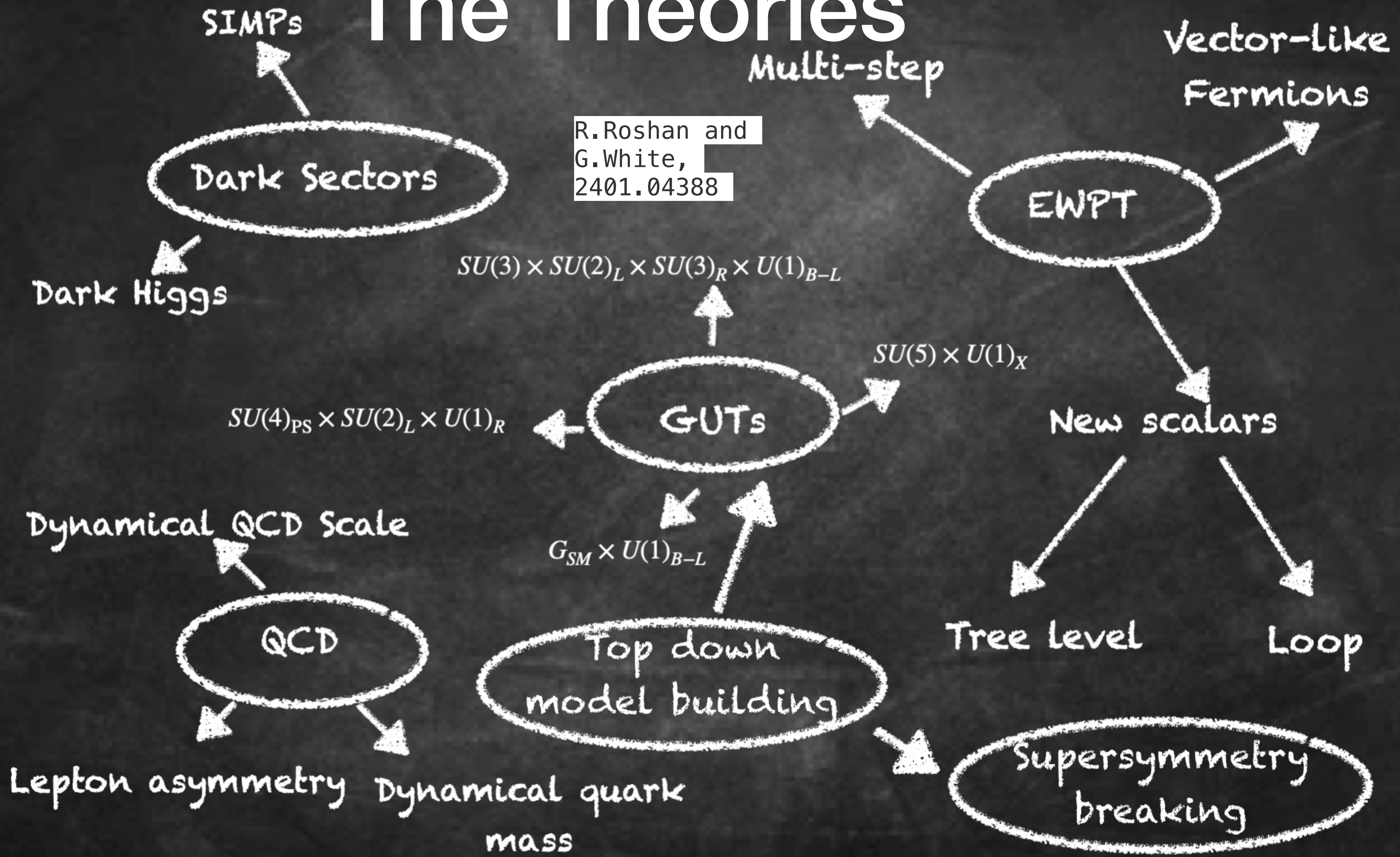


Gravitational Waves are sensitive to scales up to the Planck scale (well beyond the reach of colliders)

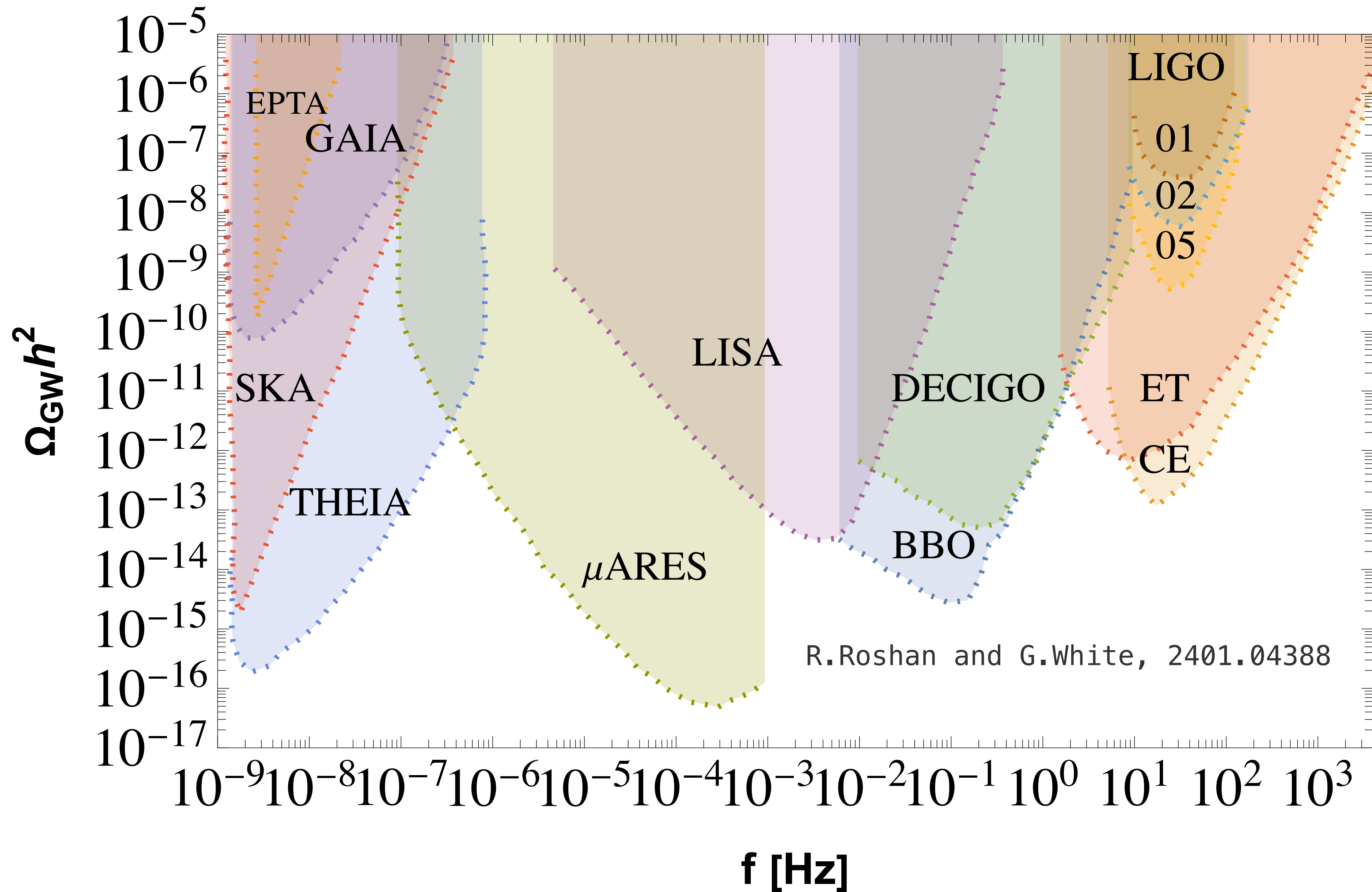


The Theories

R. Roshan and
G. White,
2401.04388



The Detectors



**Current
observations
are from
LIGO 01-03**

**The stage is
set for a
bonanza of
new results**

Gravitational Waves are sensitive to :

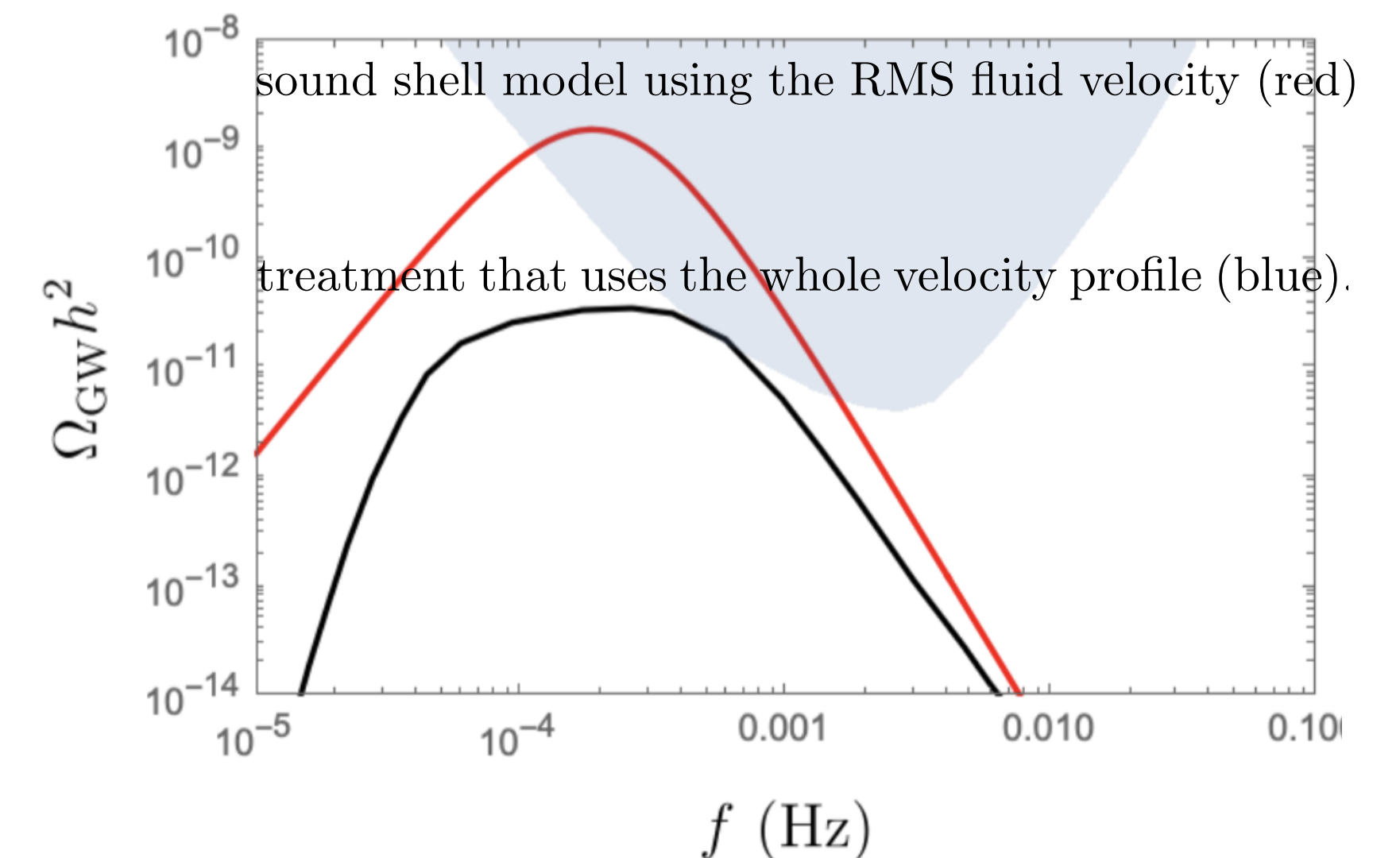
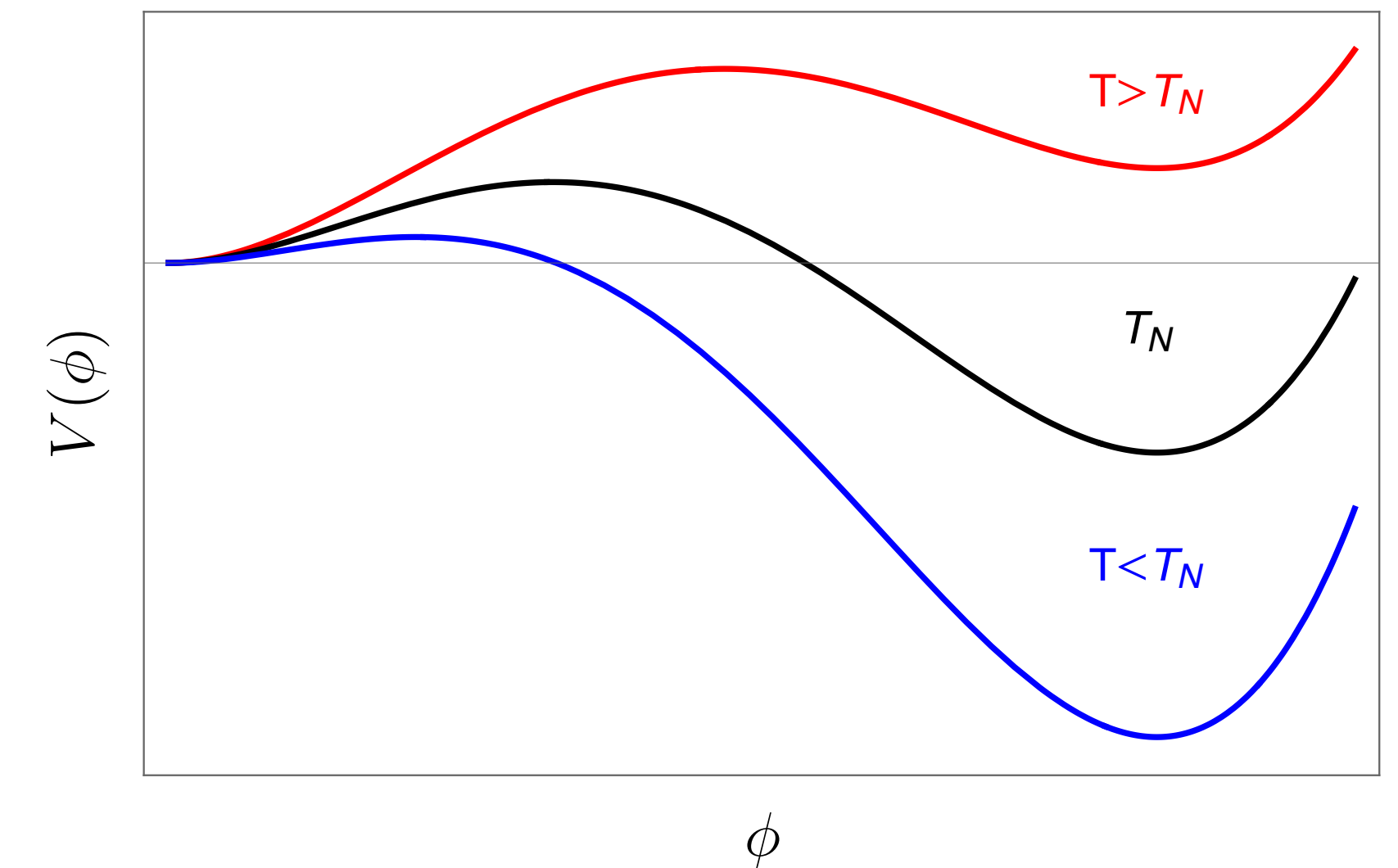
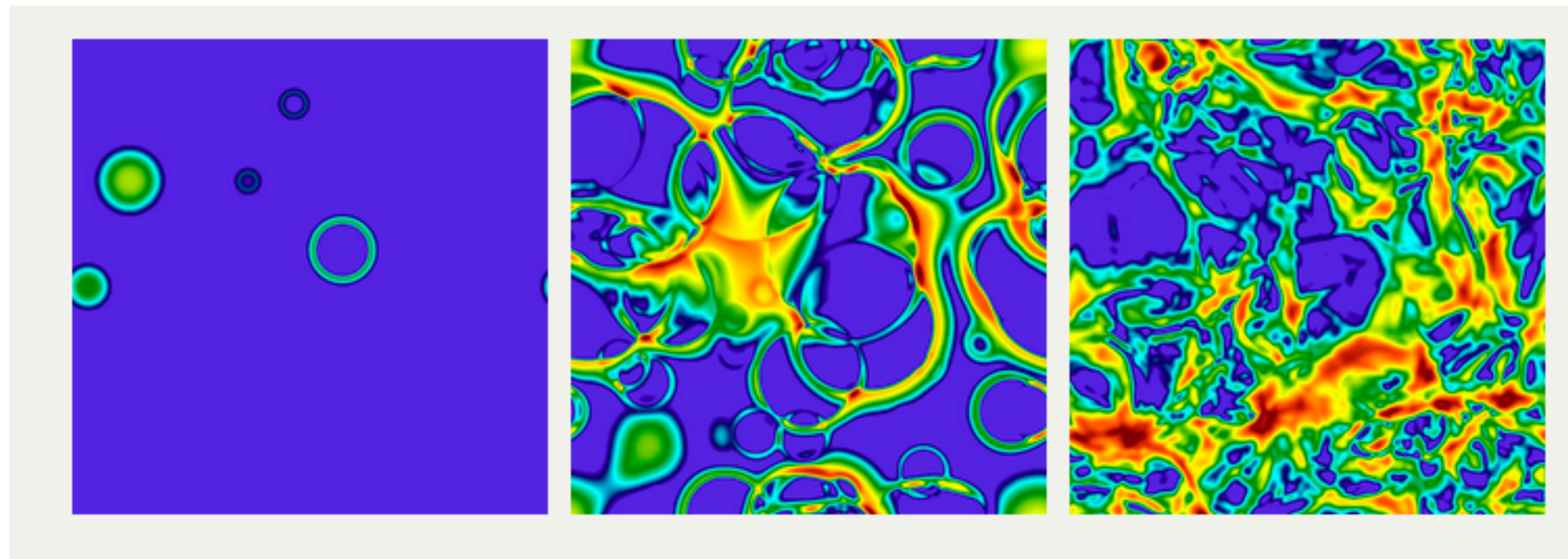
- First Order Phase transitions FOPT (e.g. QCD)
- Cosmic Strings CS (e.g. U(1) sym breaking)
- Domain Walls DW (e.g. Z_2 sym breaking)
- Inflation (e.g. with a kink or hybrid)
- Many other effects (e.g. PBHs,...)

**In this talk we
are interested
in a few BSM
examples**

Gravitational Waves from First Order Phase Transitions

Phase Transitions:

- Bubbles nucleate and grow.
- Expand in plasma.
- Bubbles and fronts collide - - violent process.
- Sound Waves left behind in thermal plasma.
- Turbulence, damping.



$$\Omega_{\text{tot}}(f) = \Omega_{\text{coll}}(f) + \Omega_{\text{sw}}(f) + \Omega_{\text{turb}}(f)$$

GW from leptoquark induced FOPT

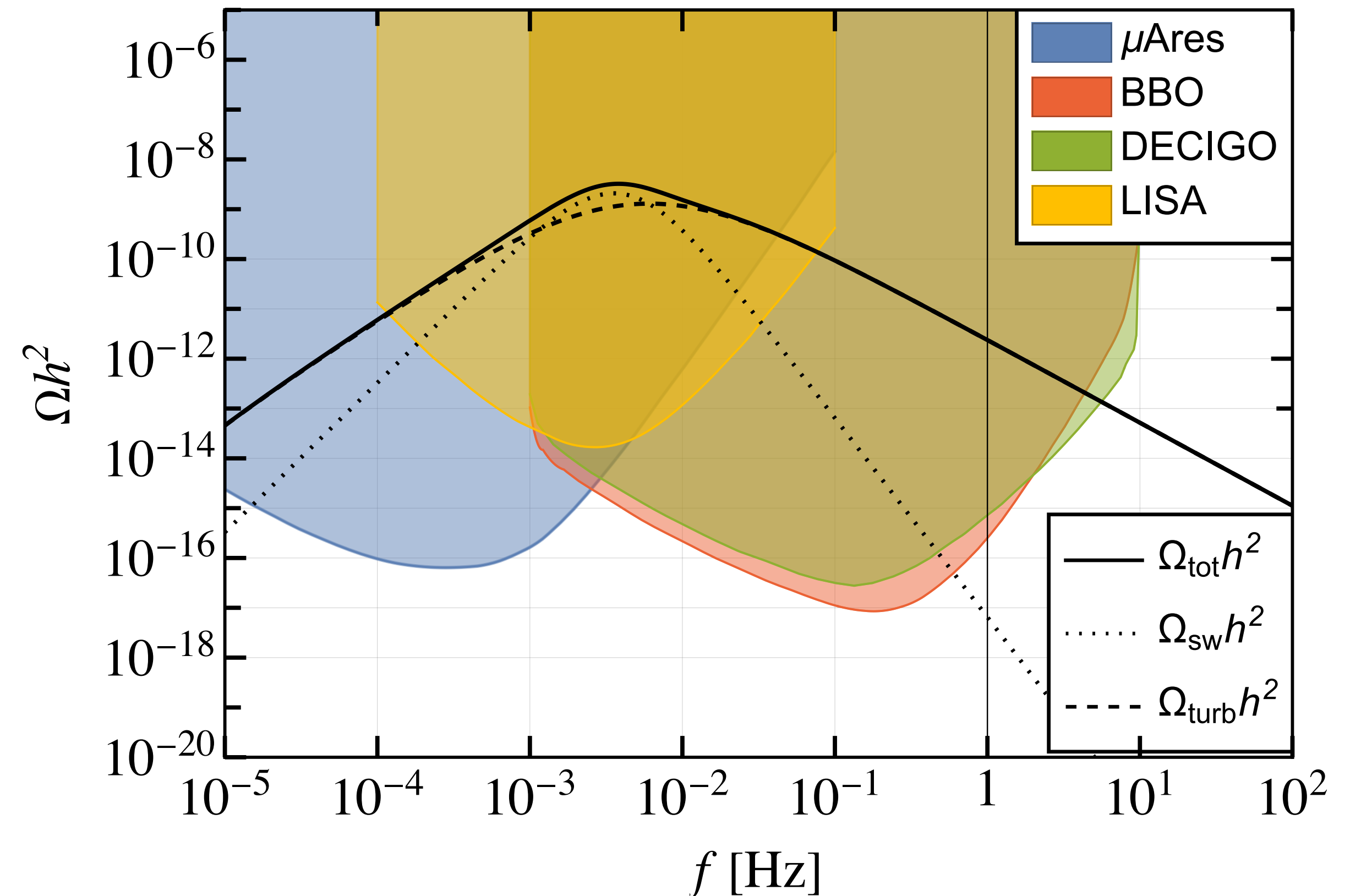
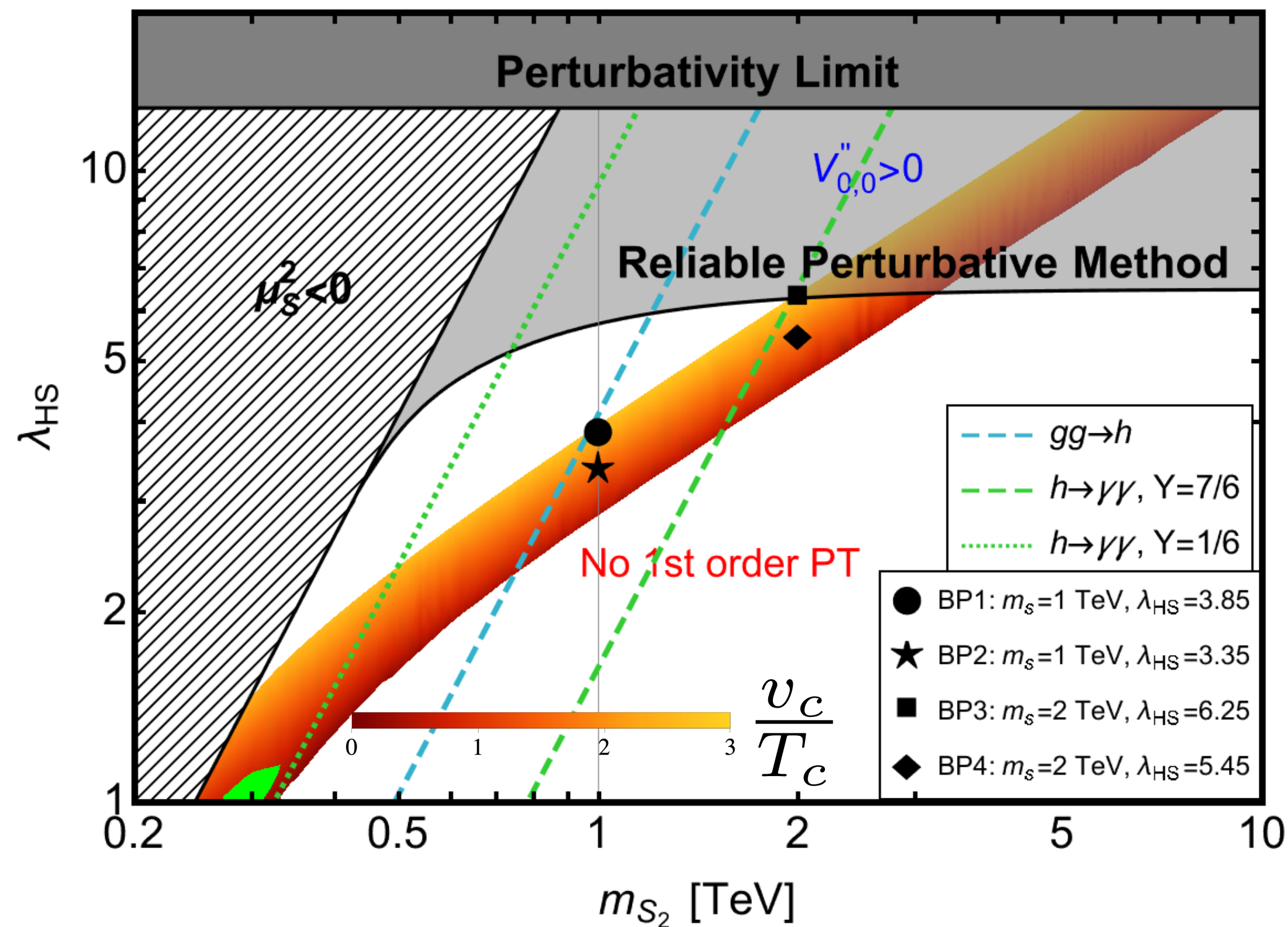
Kamila Kowalska talk

$$V_0 = -\mu^2 |H|^2 + \lambda_H |H|^4 + \mu_S^2 |S_a|^2 + \lambda_S |S_a|^4 + 2\lambda_{HS} |H|^2 |S_a|^2$$

**Leptoquark singlet S_1 ,
doublet S_2 or triplet S_3**

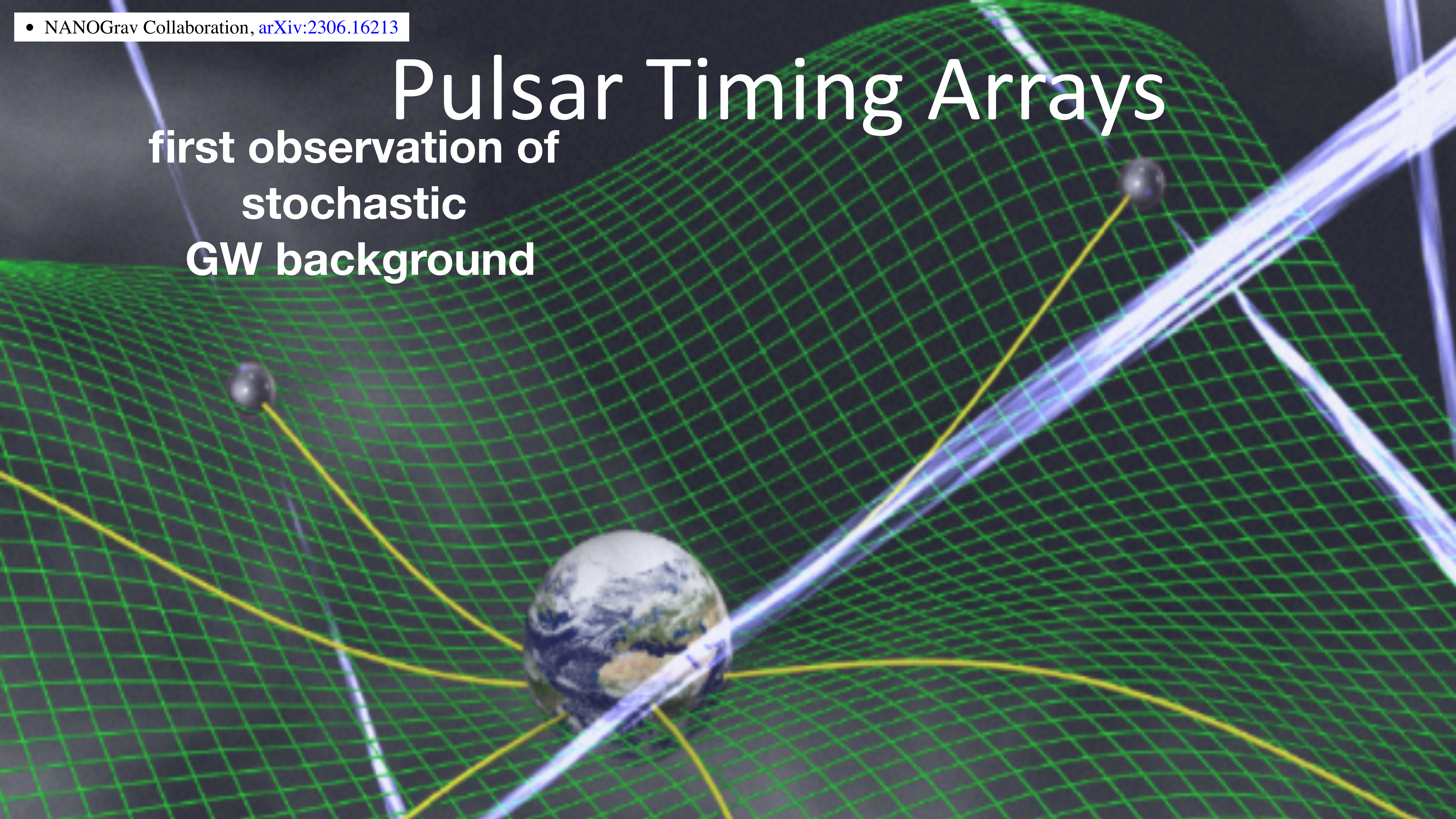
$$V_{\text{eff}}(h, T) = V_0 + \Delta V_0^{1\text{-loop}}(h) + \Delta V_T^{1\text{-loop}}(h, T)$$

Benchmark Point 1



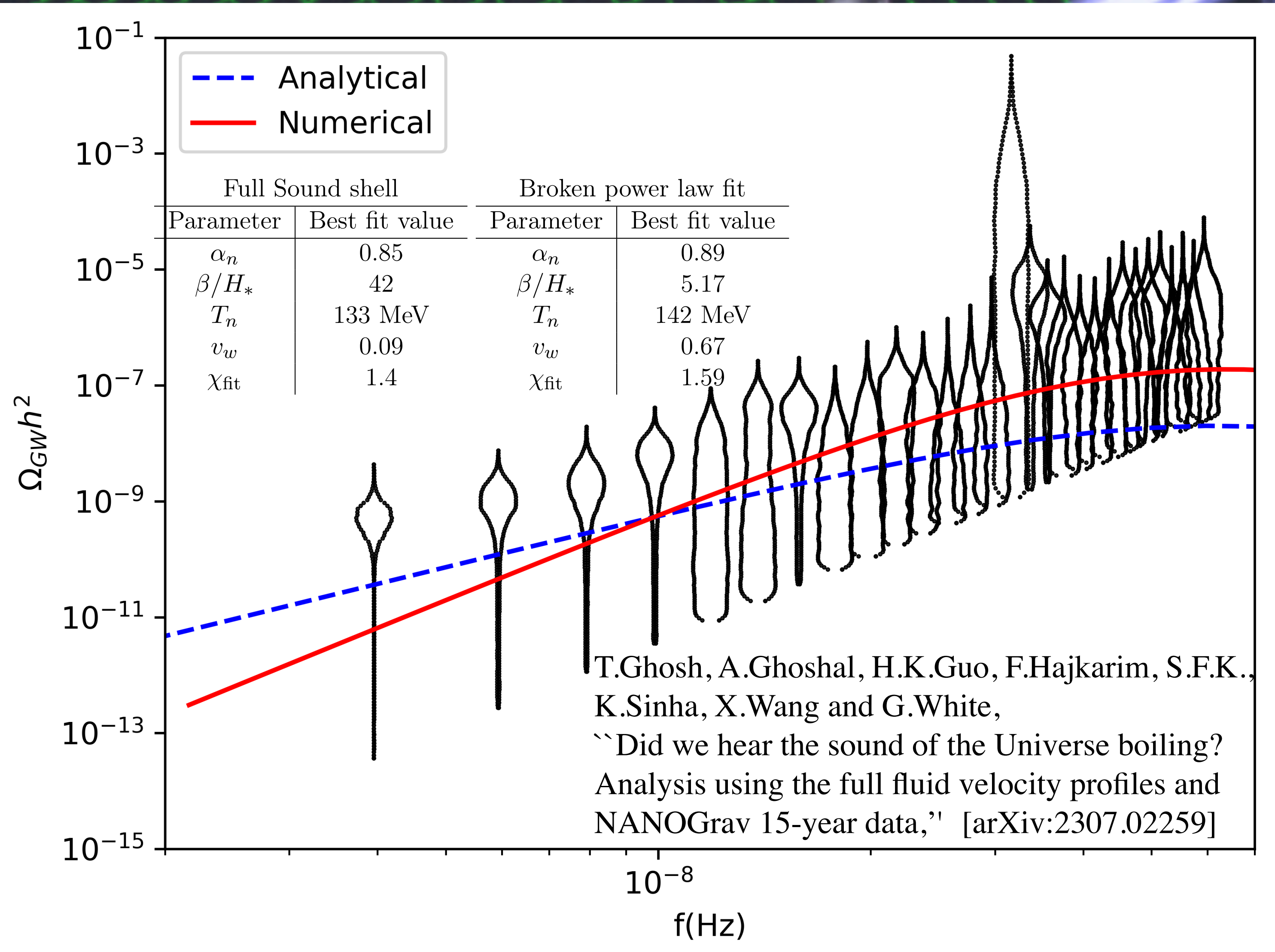
Pulsar Timing Arrays

first observation of
stochastic
GW background

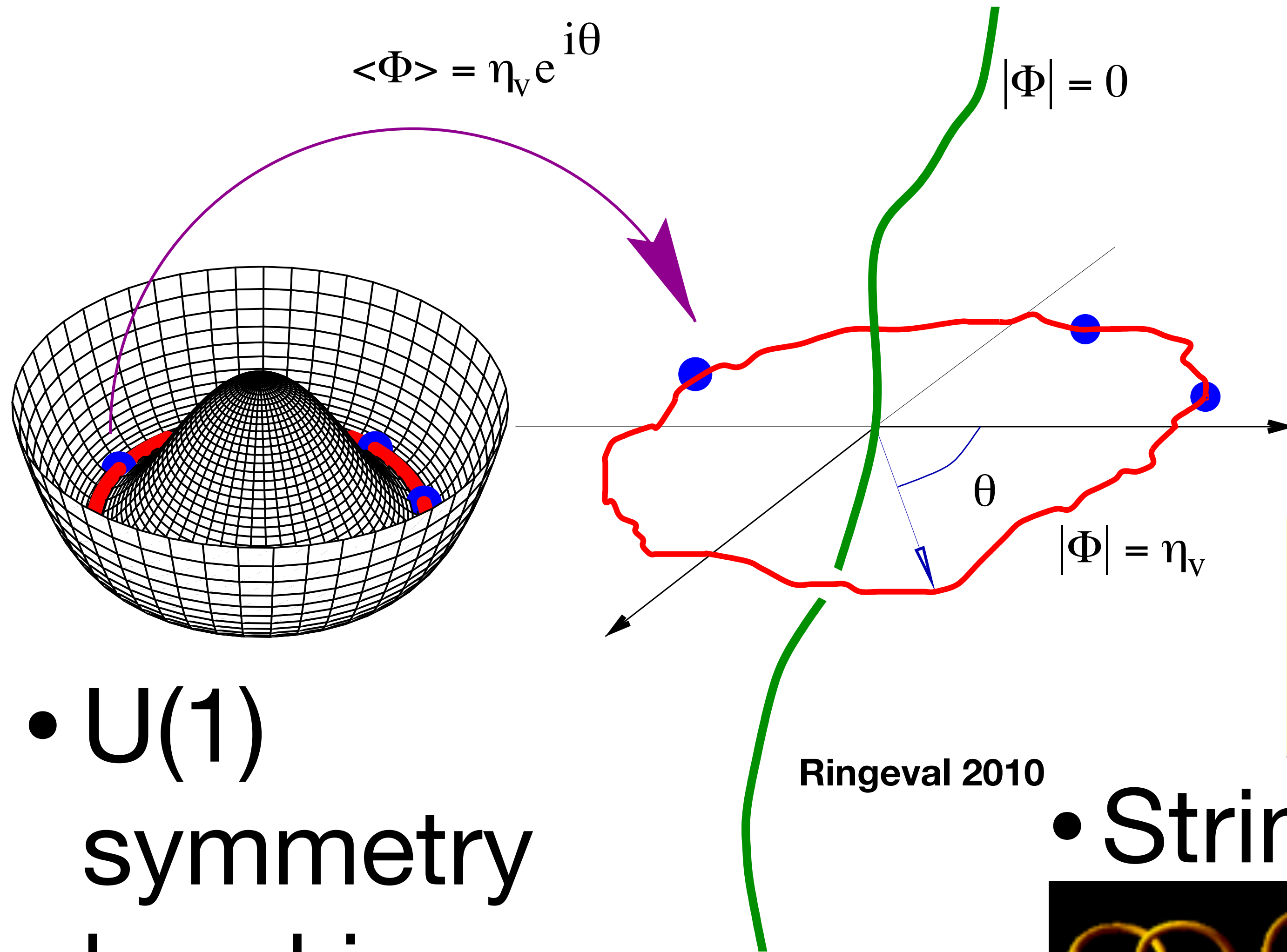


Pulsar Timing Arrays

first observation of
stochastic
GW background

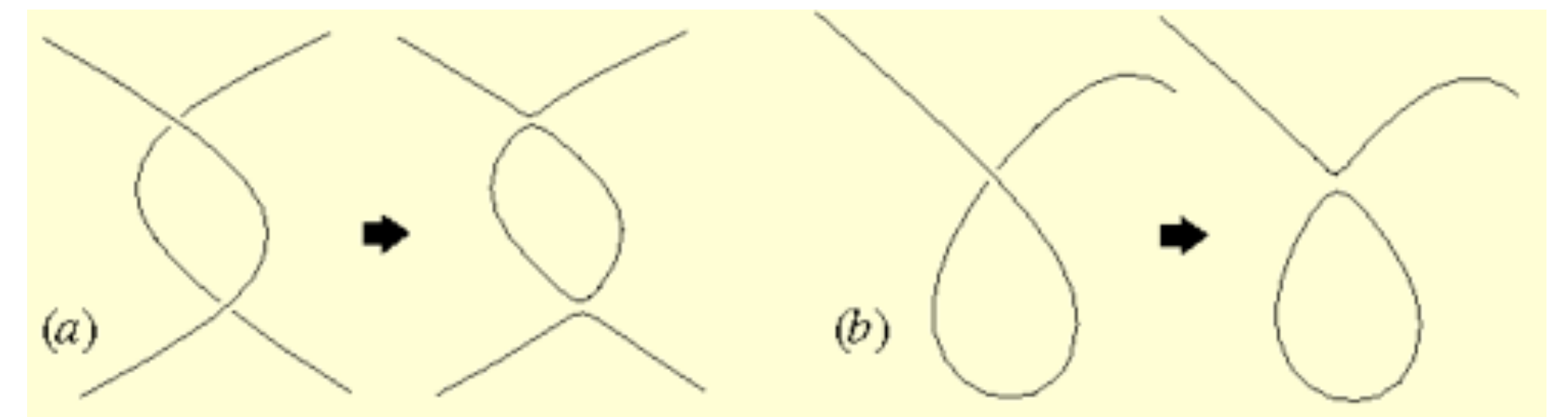
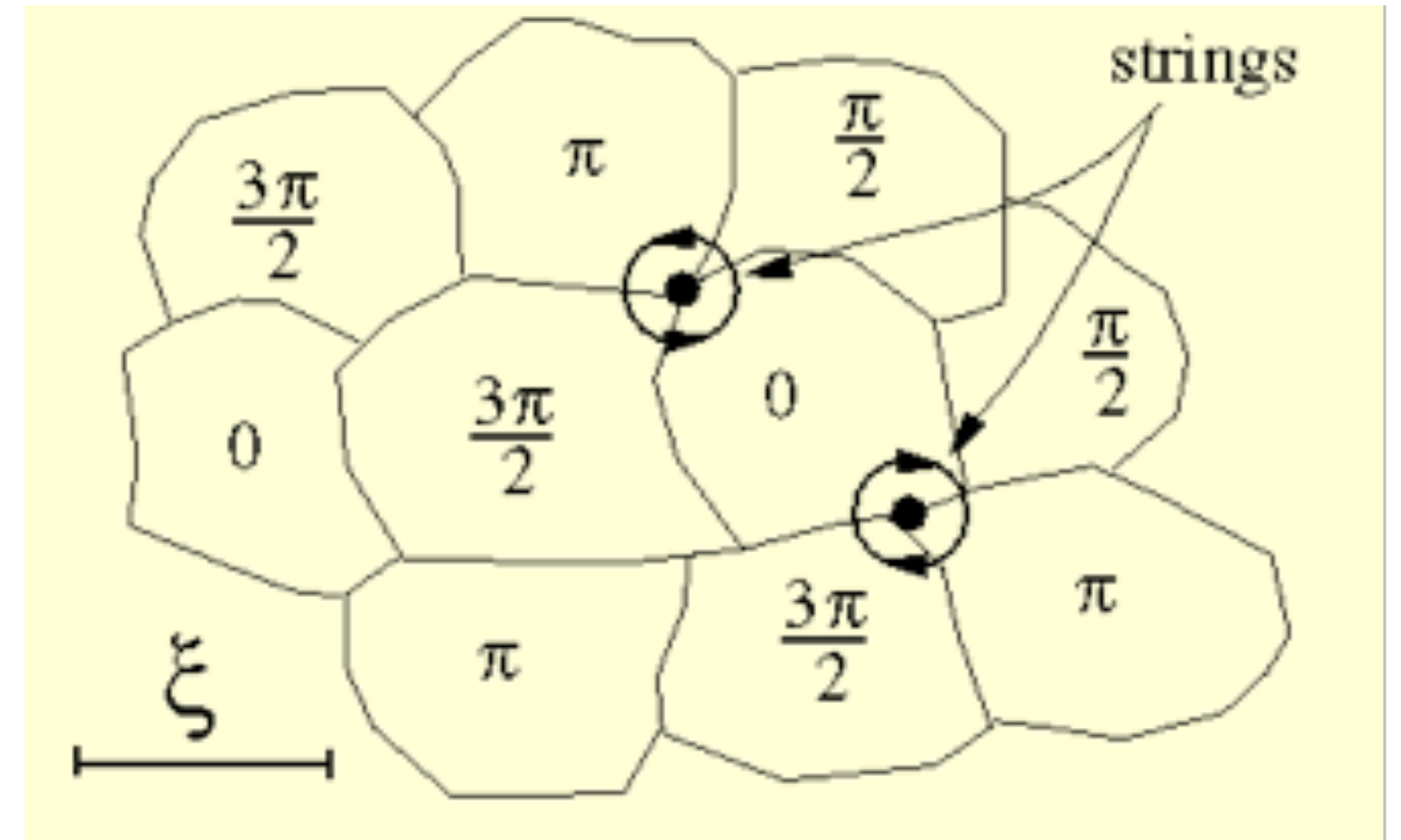


Gravitational Waves from Cosmic Strings

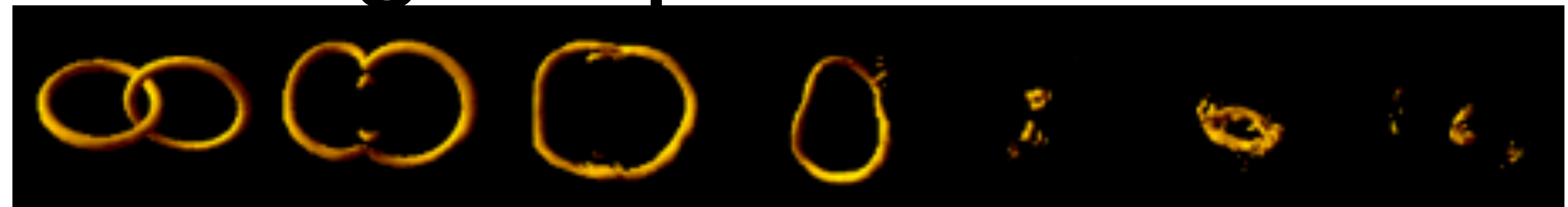


- U(1) symmetry breaking

Ringeval 2010

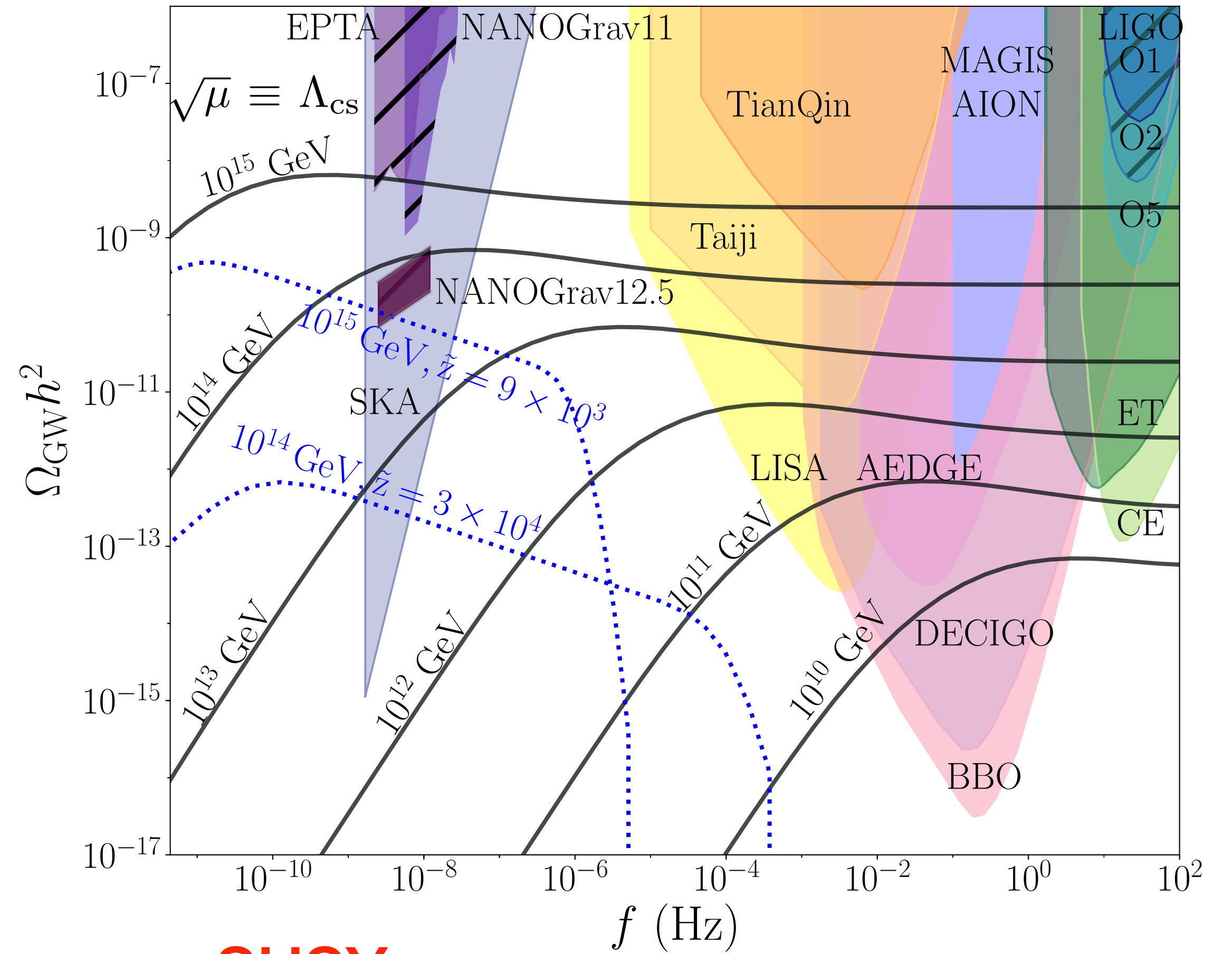
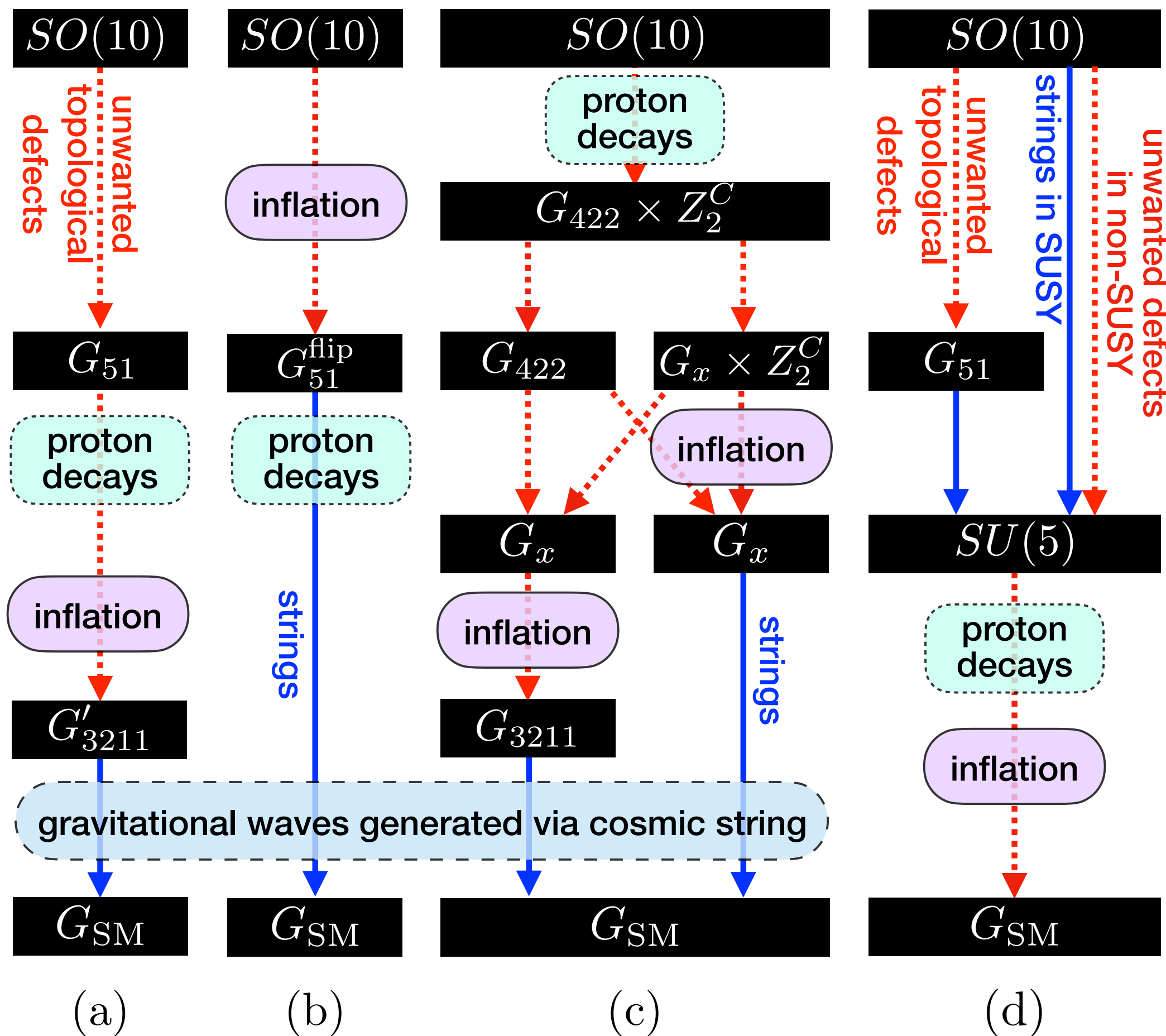


- String loops radiate GWs

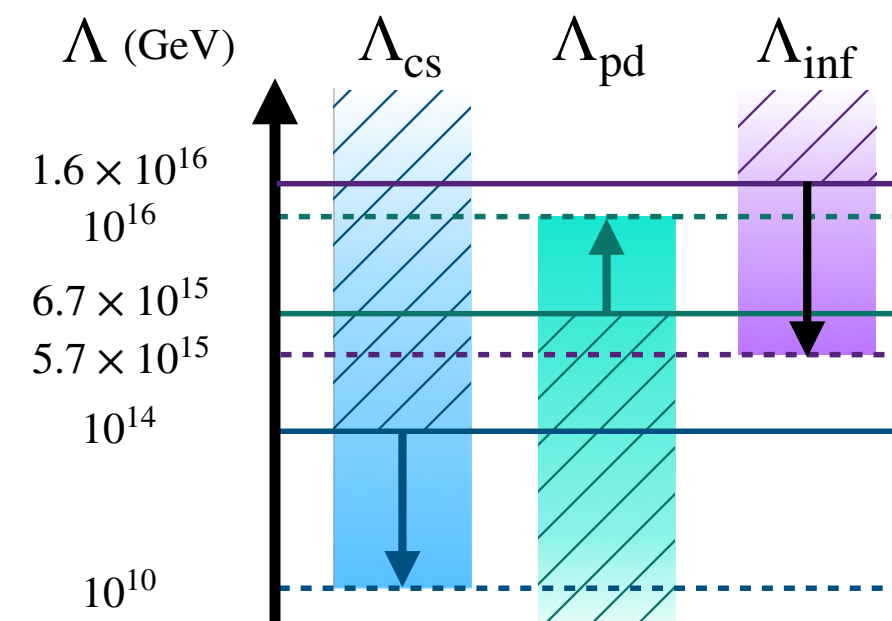


GWs via CSs from gauged $U(1)_{B-L}$ in $SO(10)$ GUTs

S.F.K., S.Pascoli, J.Turner and Y.L.Zhou, 2005.13549;
2106.15634; w/ Marsili 2209.00021; 2308.05799



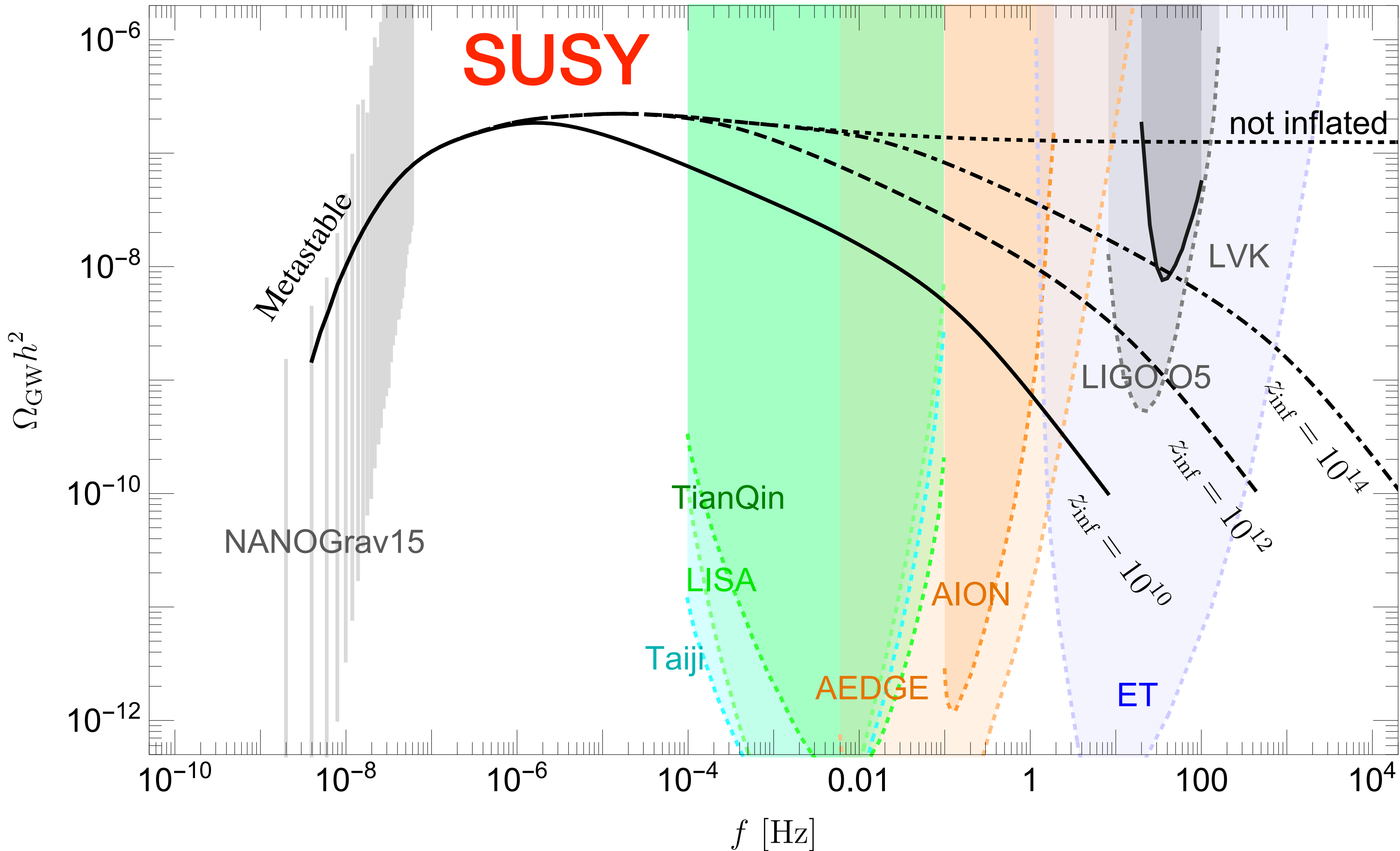
non-SUSY



Observables		Proton decays	
		$p \rightarrow \pi^0 e^+$ observed \Rightarrow non-SUSY contribution indicated	
GWs	Observed	• types (a) and (c) favoured	• types (b) and (d) excluded
	Marginal	• types (a) and (c) favoured	• type (d) excluded • type (b) allowed if $p \rightarrow K^+ \bar{\nu}$ not observed and $\Lambda_{pd} \sim \Lambda_{cs}$

Flipped SU(5): unification, proton decay, fermion masses and gravitational waves

S.F.K., G.K.Leontaris and Y.L.Zhou, 2311.11857



Metastable
cosmic strings

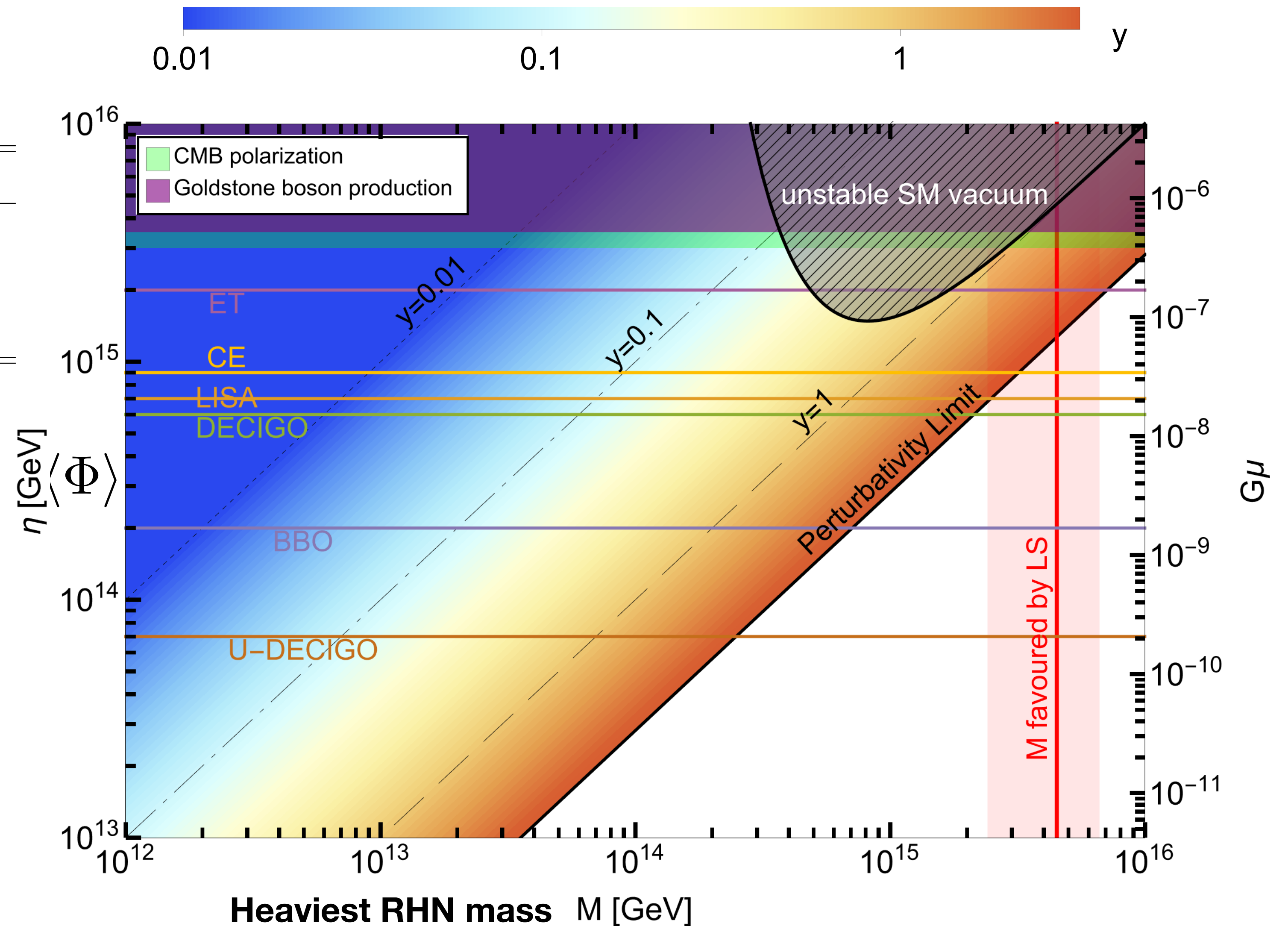
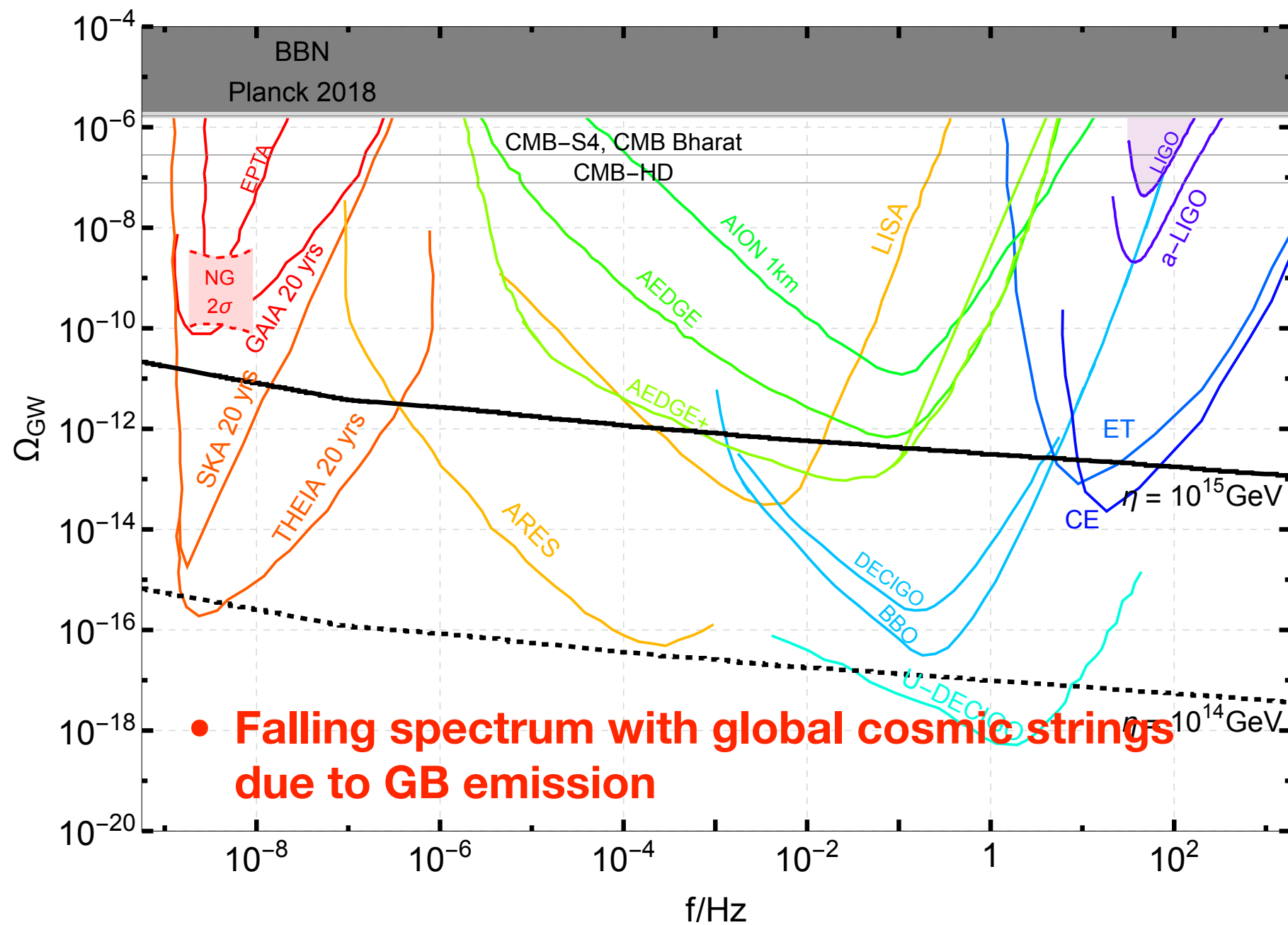
**SUSY GUTs
with
metastable
cosmic strings
and inflation
dilution can fit
NANOGrav 15
year data**

Cosmic string gravitational waves from **global** $U(1)_{B-L}$ symmetry breaking as a probe of the type I seesaw scale

Majoron breaks $U(1)_{B-L}$

$$Y_{\alpha i} \bar{L}_{\alpha} \tilde{H} N_i + \frac{1}{2} y_i \Phi \bar{N}_i^c N_i + h.c.$$

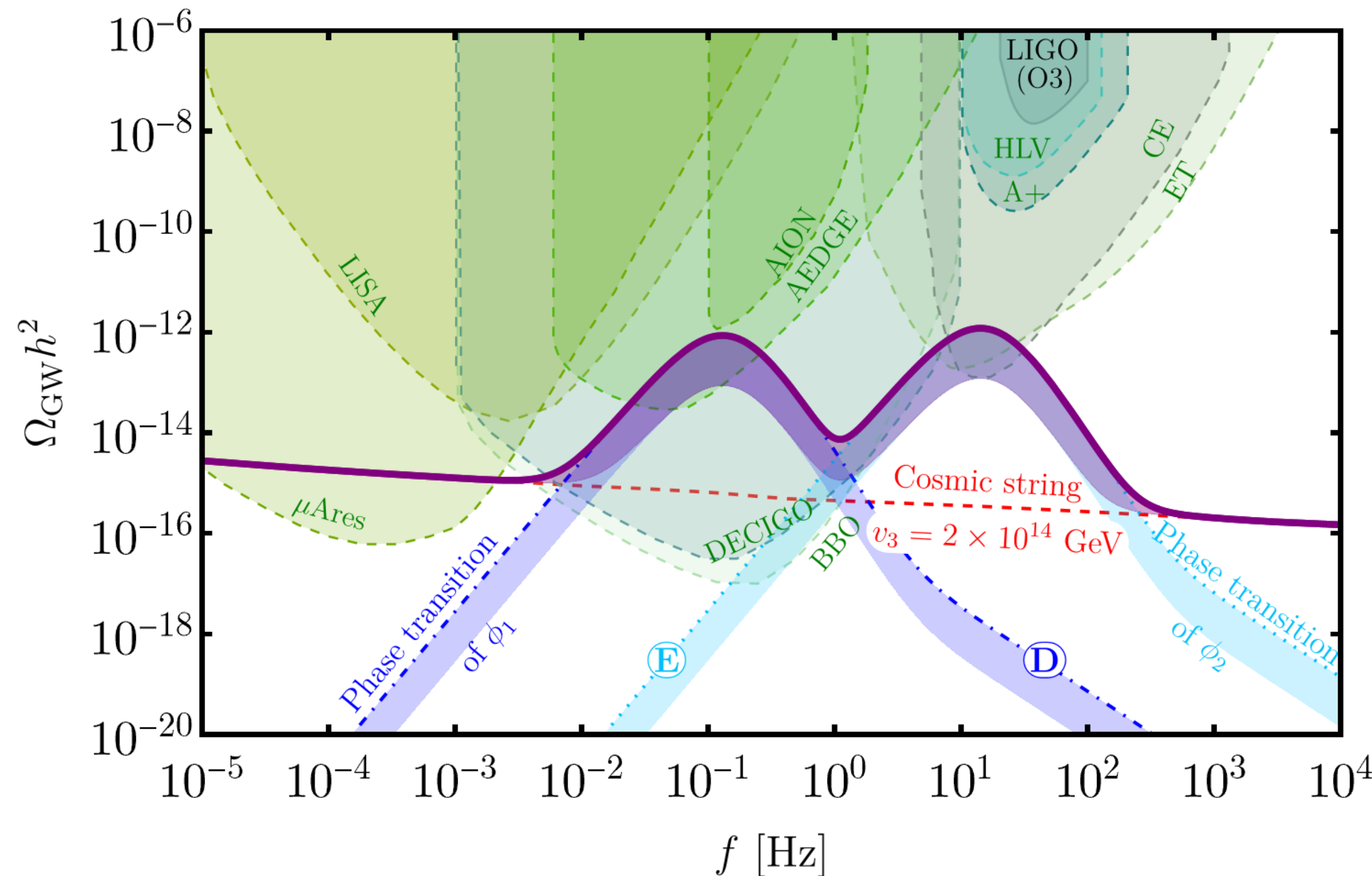
	Q	u_R	d_R	L	e_R	H	N	Φ
$SU(2)_L$	2	1	1	2	1	2	1	1
$U(1)_Y$	$\frac{1}{6}$	$\frac{2}{3}$	$-\frac{1}{3}$	$-\frac{1}{2}$	-1	$-\frac{1}{2}$	0	0
$U(1)_{B-L}$	$\frac{1}{3}$	$\frac{1}{3}$	$\frac{1}{3}$	-1	-1	0	-1	2



Gravitational waves from **phase transitions** and **cosmic strings** in neutrino mass models with **multiple Majorons**

$$U(1)_{L_1} \times U(1)_{L_2} \times U(1)_{L_3}$$

$$\left(\overline{L}_a h_{aI} H N_I + \frac{y_1}{2} \phi_1 \overline{N}_1^c N_1 + \frac{y_2}{2} \phi_2 \overline{N}_2^c N_2 + \frac{y_3}{2} \phi_3 \overline{N}_3^c N_3 + \text{h.c.} \right) + V_0(\phi_1, \phi_2, \phi_3)$$



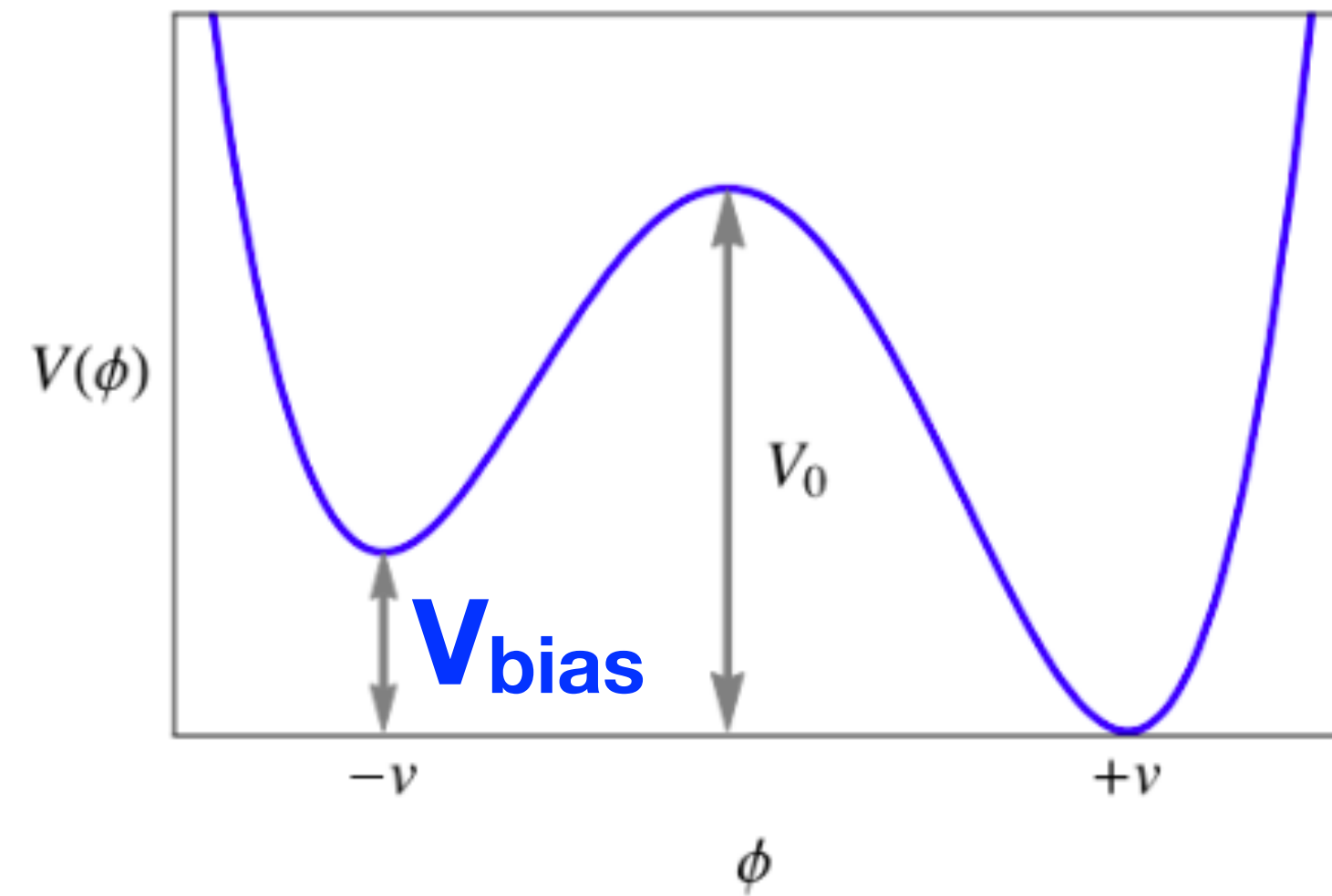
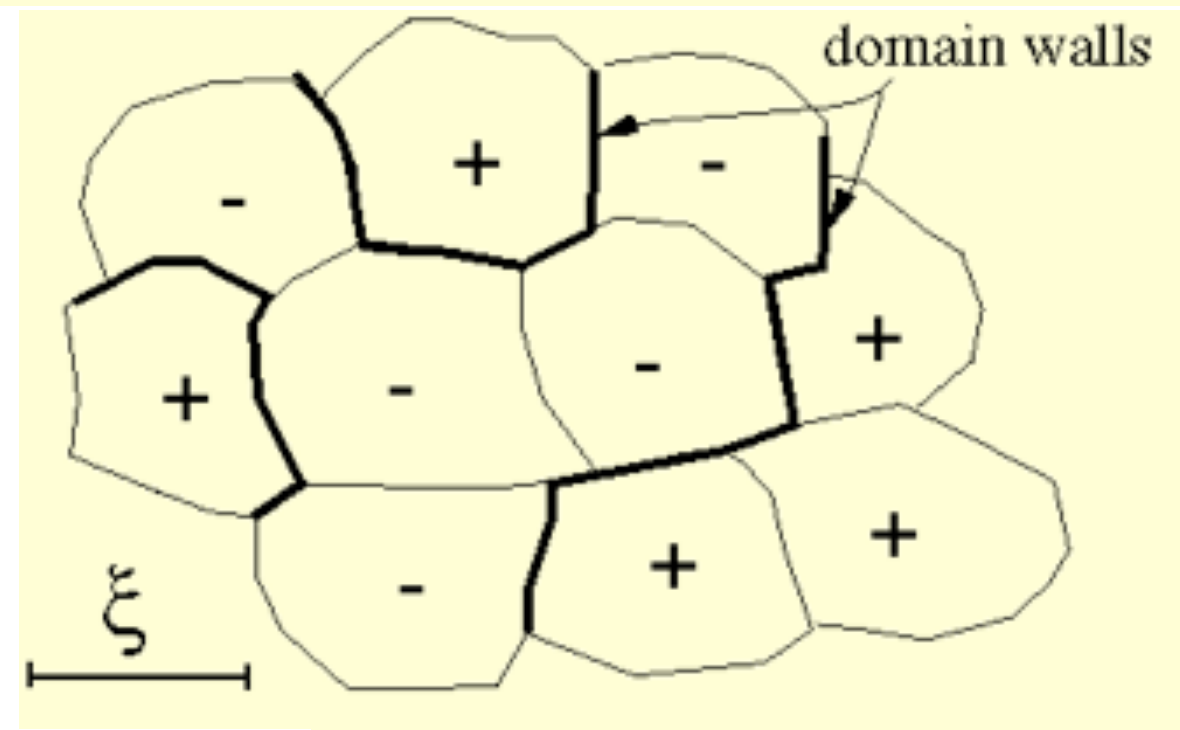
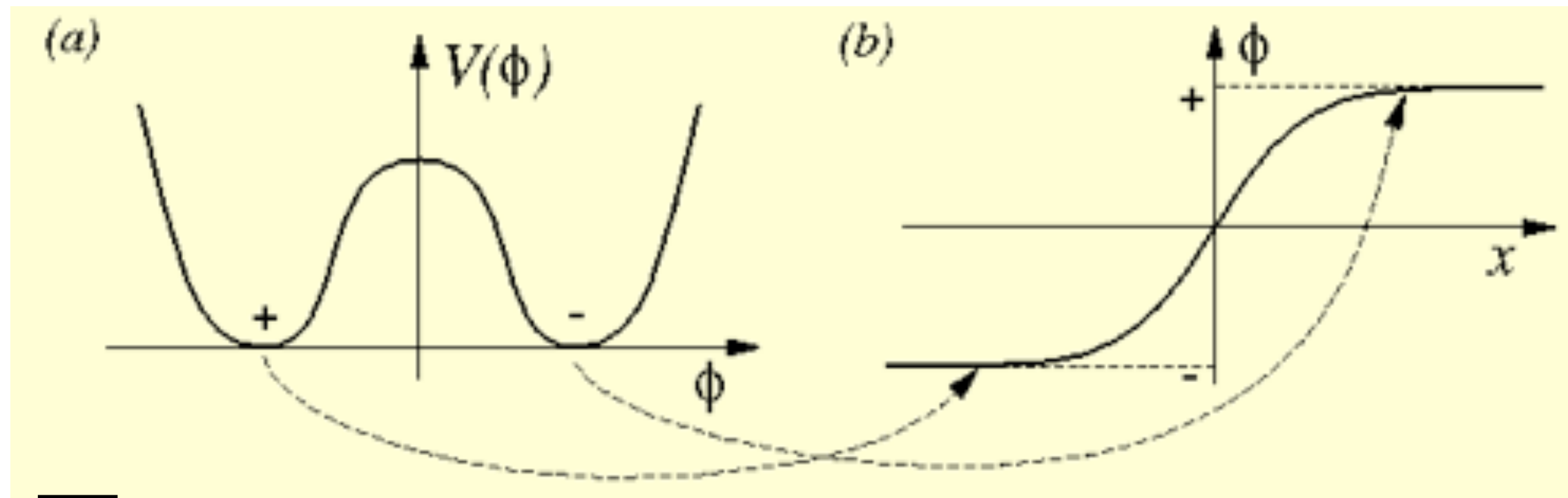
$$\sum_{I=1,2,3} [-\mu_I^2 \phi_I^* \phi_I + \lambda_I (\phi_I^* \phi_I)^2] + \sum_{I,J,I \neq J}^{1,2,3} \frac{\zeta_{IJ}}{2} (\phi_I^* \phi_I) (\phi_J^* \phi_J)$$

	λ_I	v_I [GeV]
(D)	0.00027	1188.2
(E)	0.00029	2.32×10^5

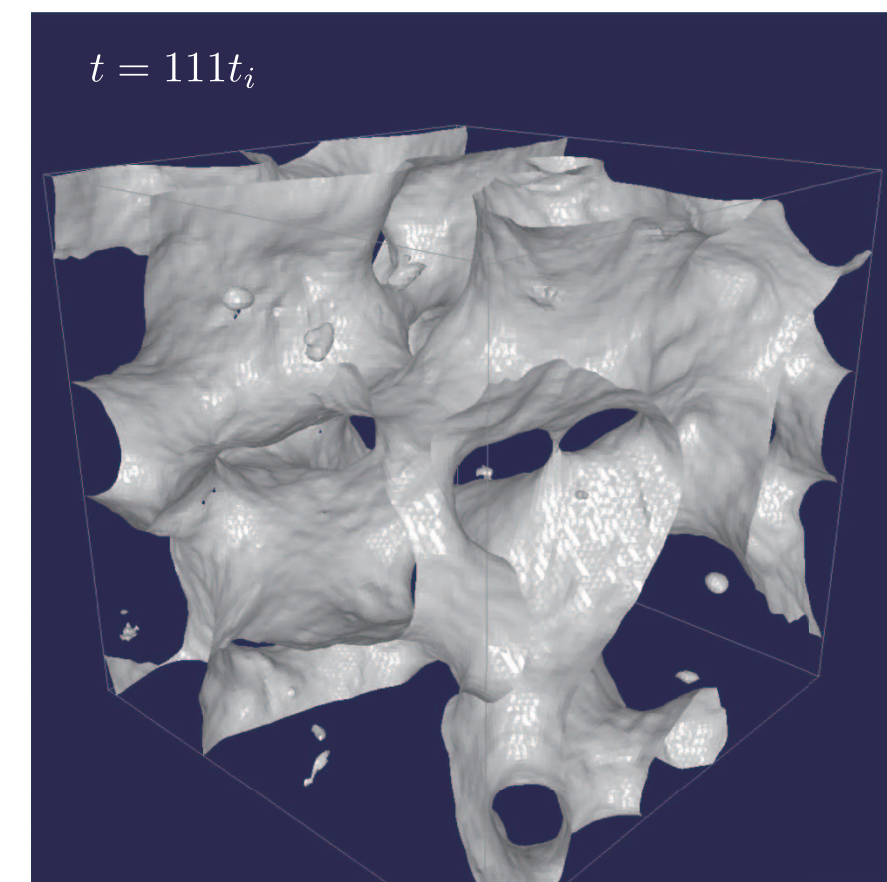
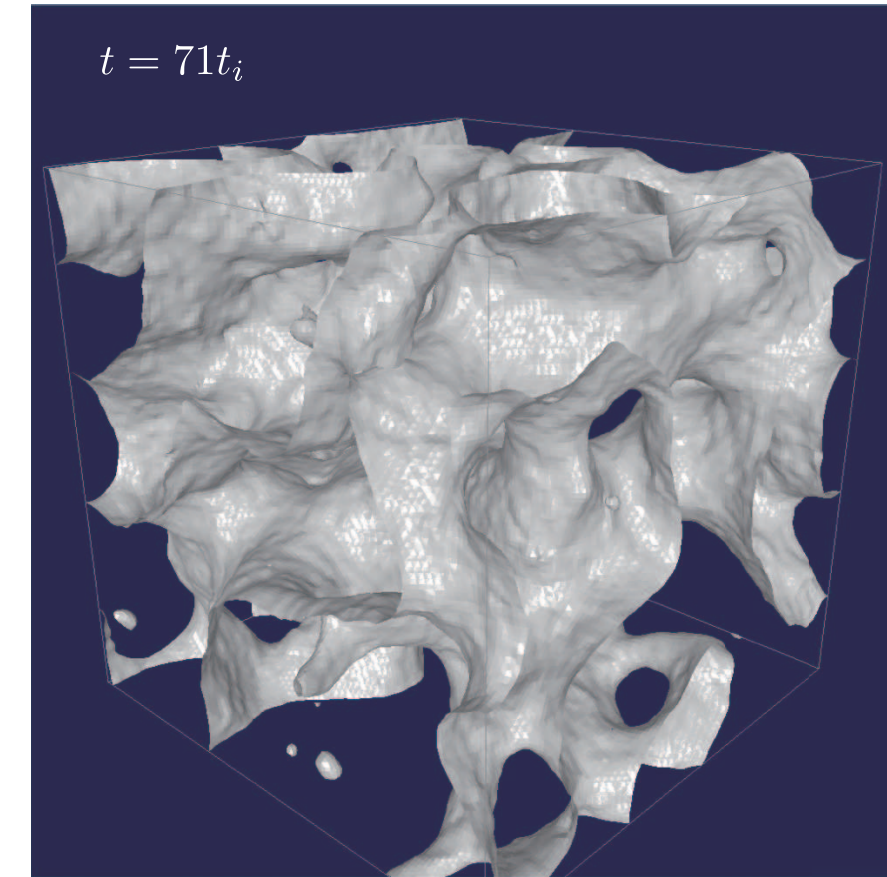
Phase transitions for Φ_1, Φ_2 enhanced by Φ_3

Gravitational Waves from Domain Walls

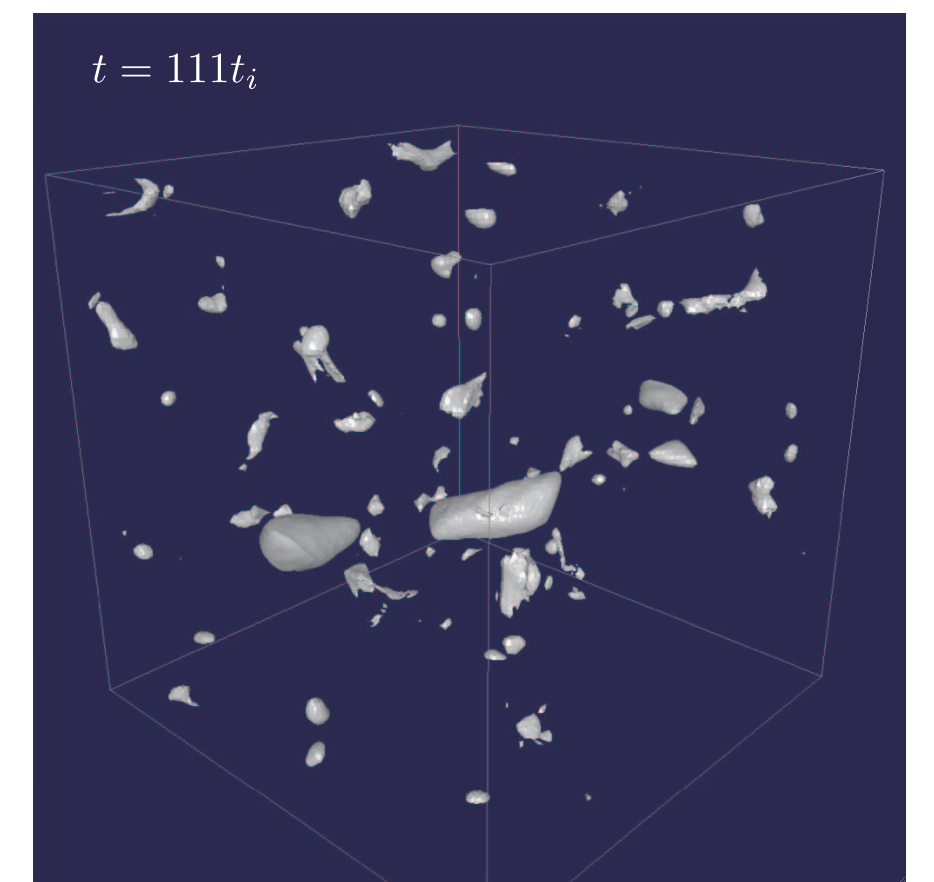
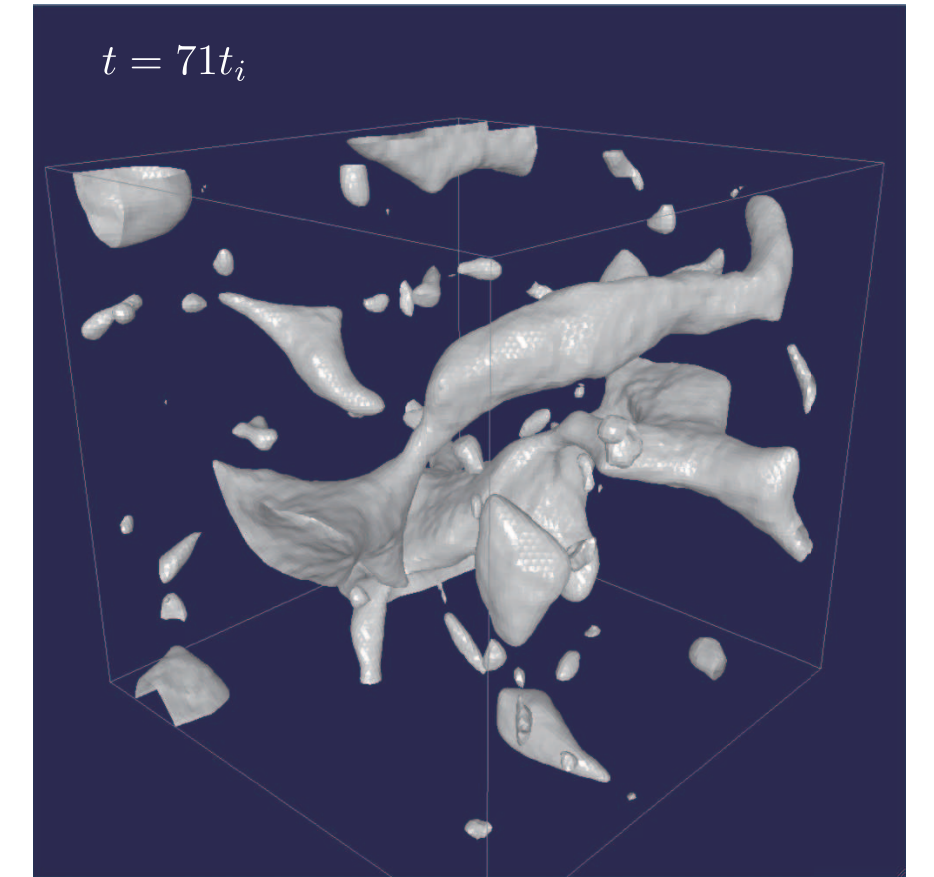
- Z_2 symmetry breaking



Introduce V_{bias} to allow domain walls to decay (otherwise dominate energy density of Universe)



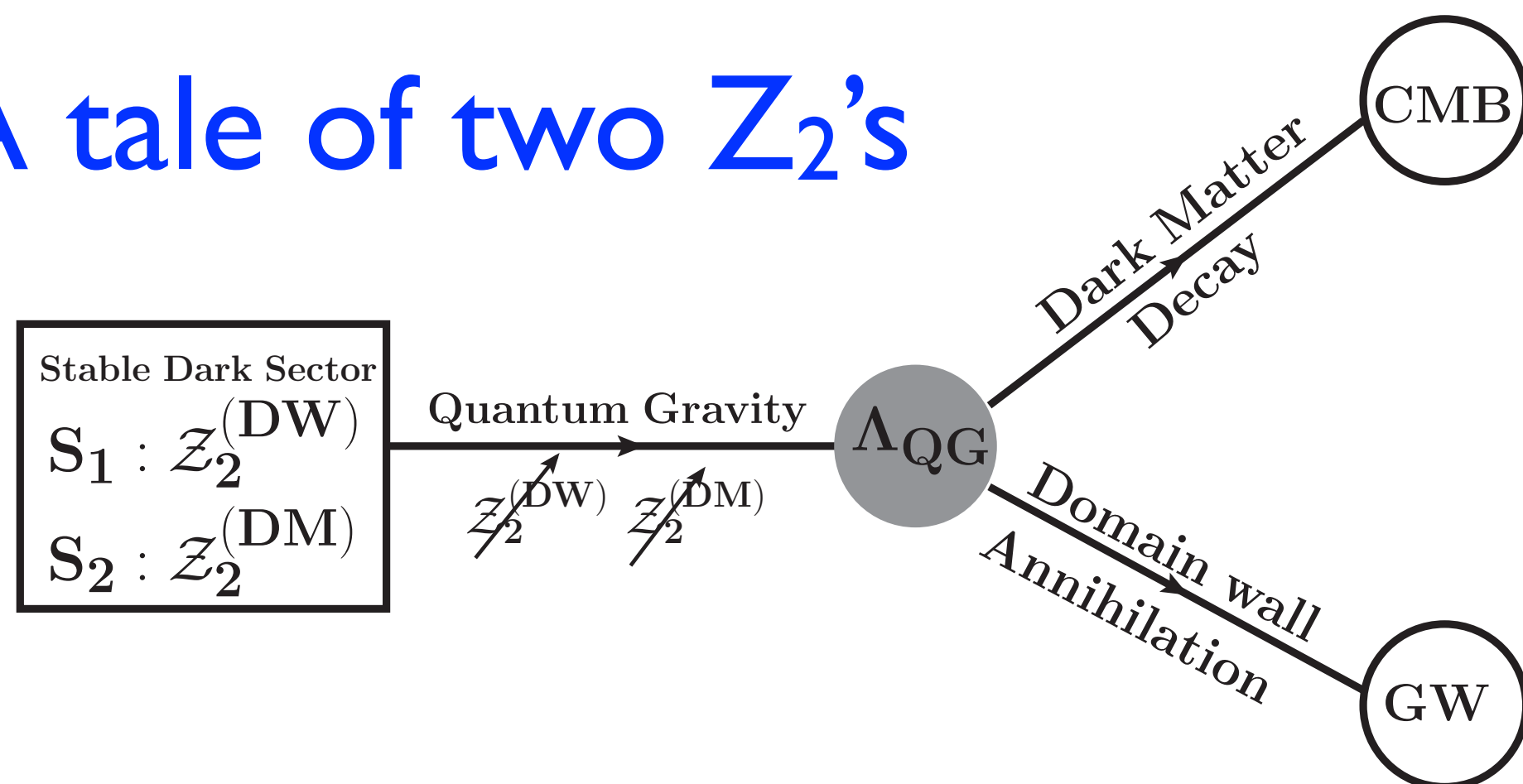
Stable DW on the left with $V_{\text{bias}} = 0$



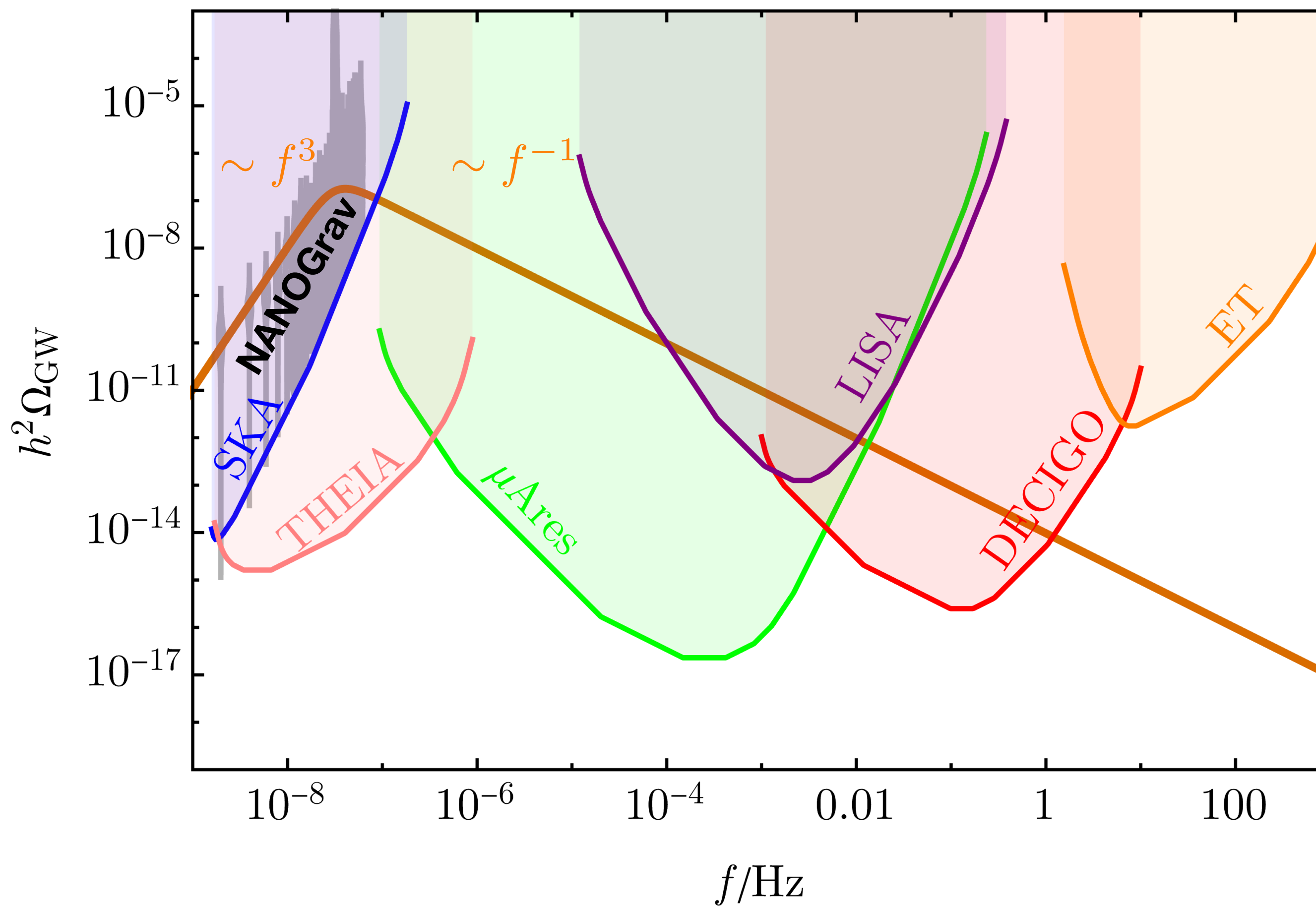
Unstable DW decay via GWs due to V_{bias}

Quantum gravity effects on dark matter and gravitational waves

A tale of two Z_2 's



Peak occurs when volume pressure $\sim V_{bias}$



Both broken by QG effects

$$\mathcal{L}_{Z_2} = \frac{1}{\Lambda_{QG}} \mathcal{O}_5 \quad \rightarrow \quad \mathbf{V}_{bias}$$

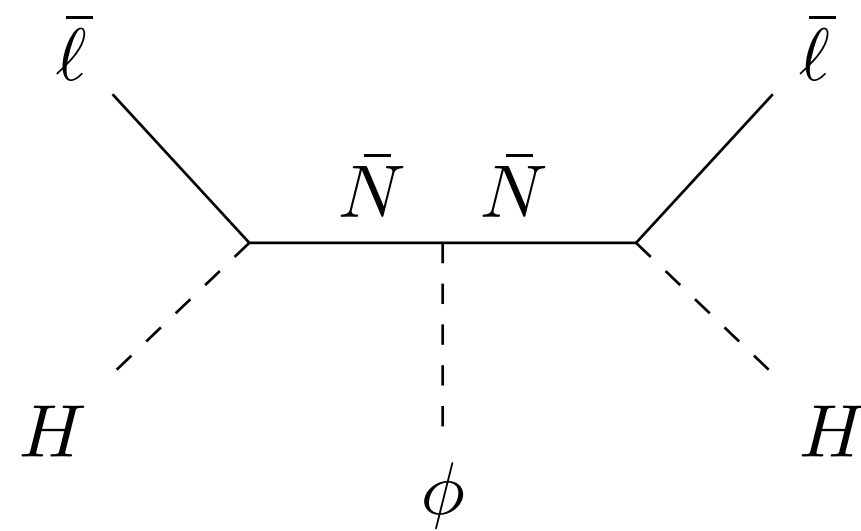
Due to instanton effects

$$\Lambda_{QG} \sim M_{Pl} e^S \gg M_{Pl}$$

Toward distinguishing **Dirac** from **Majorana** neutrino mass with gravitational waves

Majorana seesaw

$$-\mathcal{L}_M \supset \mathcal{Y} \bar{\ell} H \bar{N} + \bar{N} \bar{N}^T \phi$$

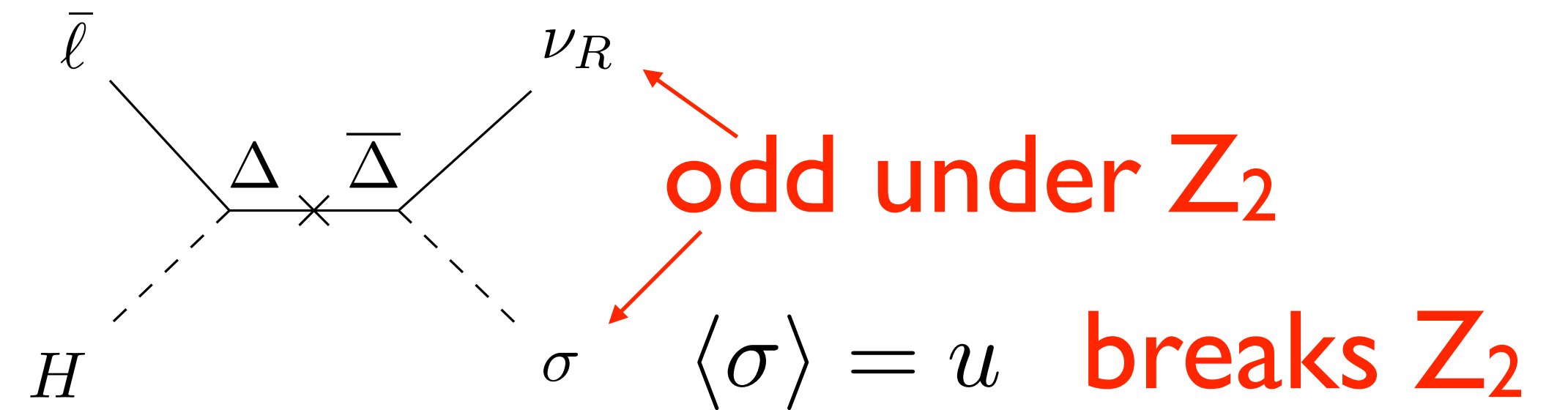


$$\mathcal{M}_M = \frac{1}{\sqrt{2}} v^2 \mathcal{Y} \mathcal{M}_N^{-1} \mathcal{Y}^T$$

Gauged $U(1)_{B-L}$ broken
→ Cosmic strings

Dirac seesaw

$$-\mathcal{L}_D \supset \mathcal{Y}_L \bar{\ell} H \Delta_R + \mathcal{Y}_R \bar{\Delta}_L \sigma \nu_R + \mathcal{M}_\Delta \bar{\Delta} \Delta$$

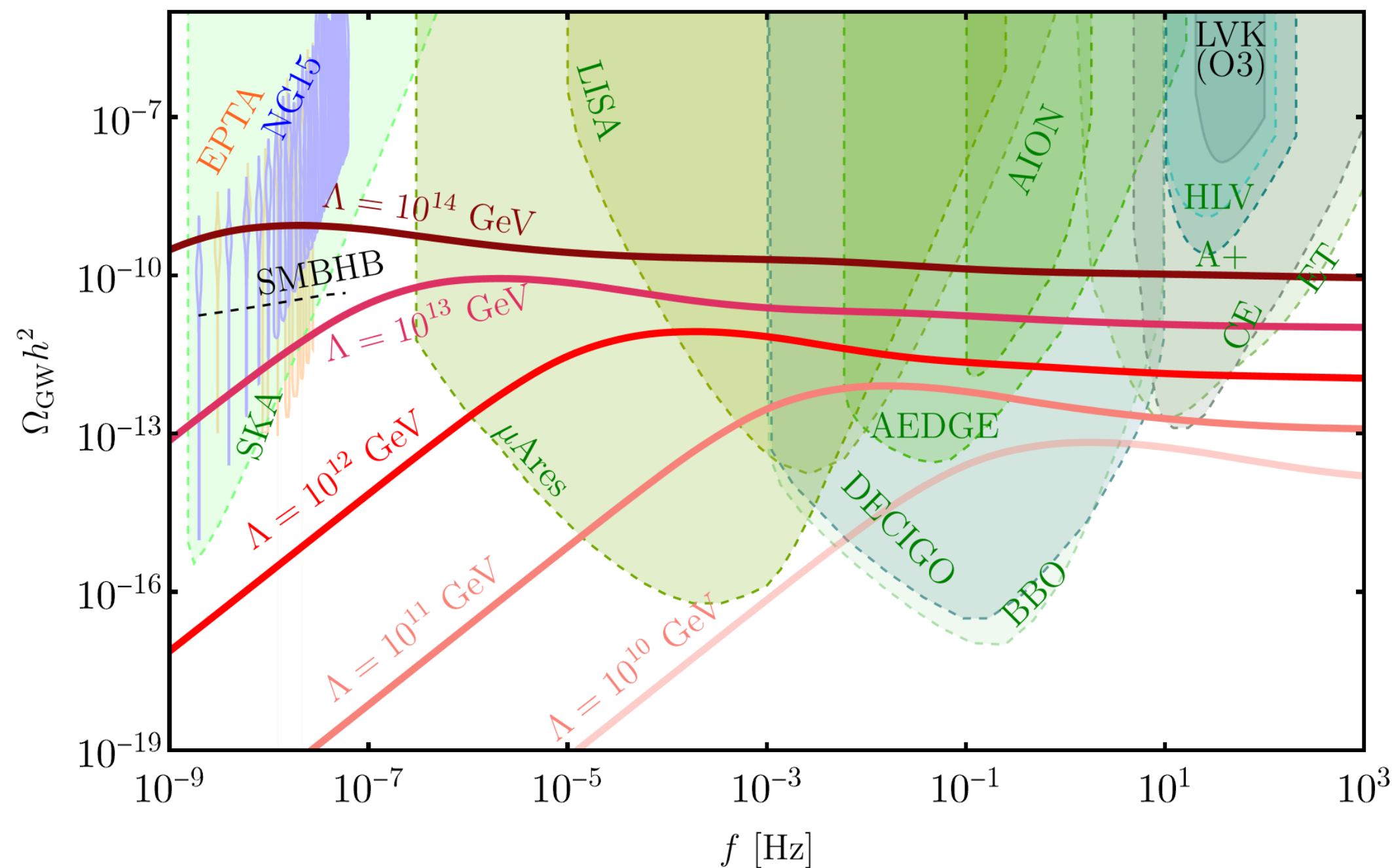


$$\mathcal{M}_D = \frac{1}{\sqrt{2}} v u \mathcal{Y}_L \mathcal{M}_\Delta^{-1} \mathcal{Y}_R$$

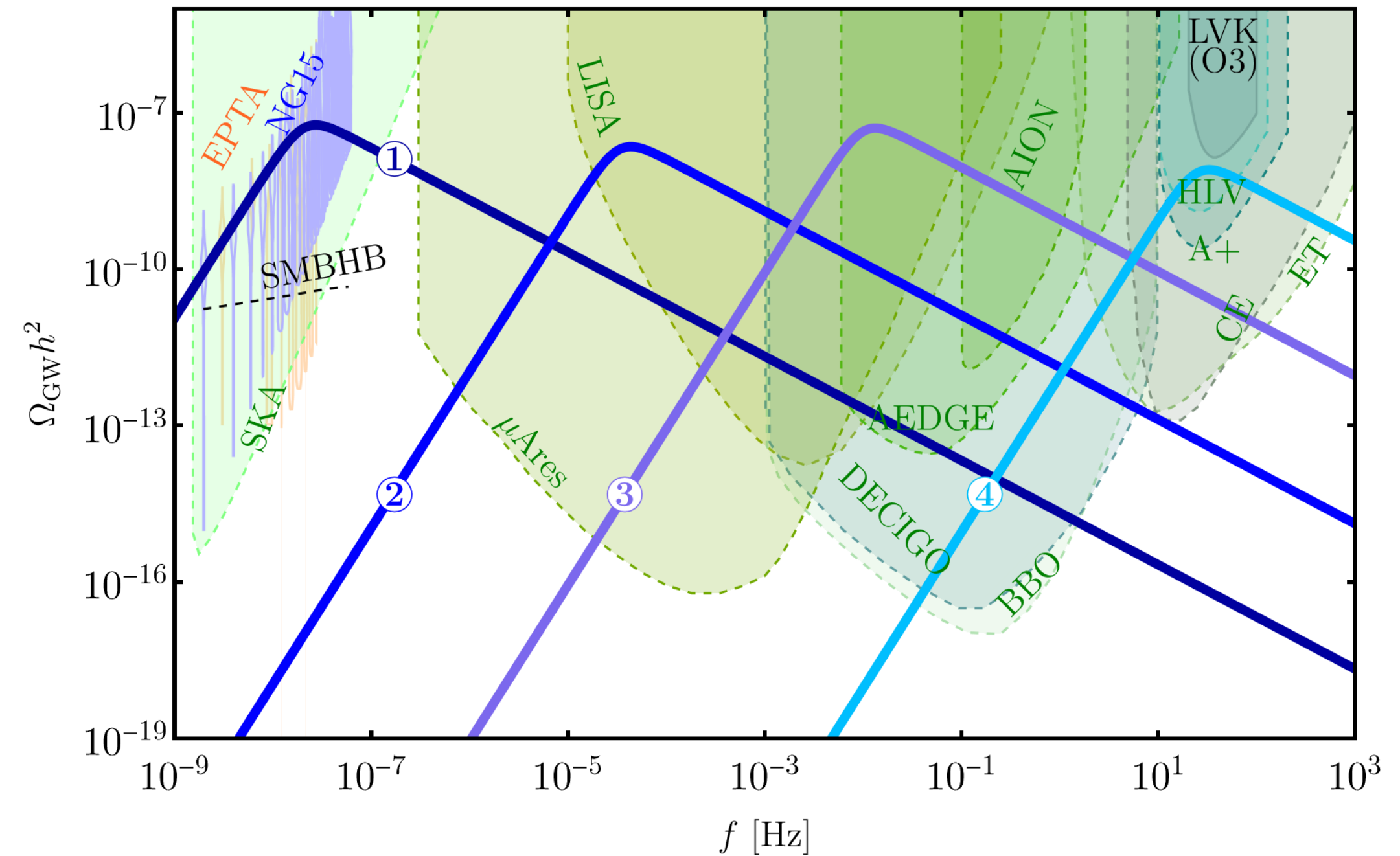
Global $U(1)_{B-L}$ preserved
 Z_2 broken → Domain Walls

Toward distinguishing **Dirac** from **Majorana** neutrino mass with gravitational waves

Majorana seesaw



Dirac seesaw



Majorana vs Dirac can be distinguished from shape of GW spectrum

- - Dirac is better fit to NANOGrav

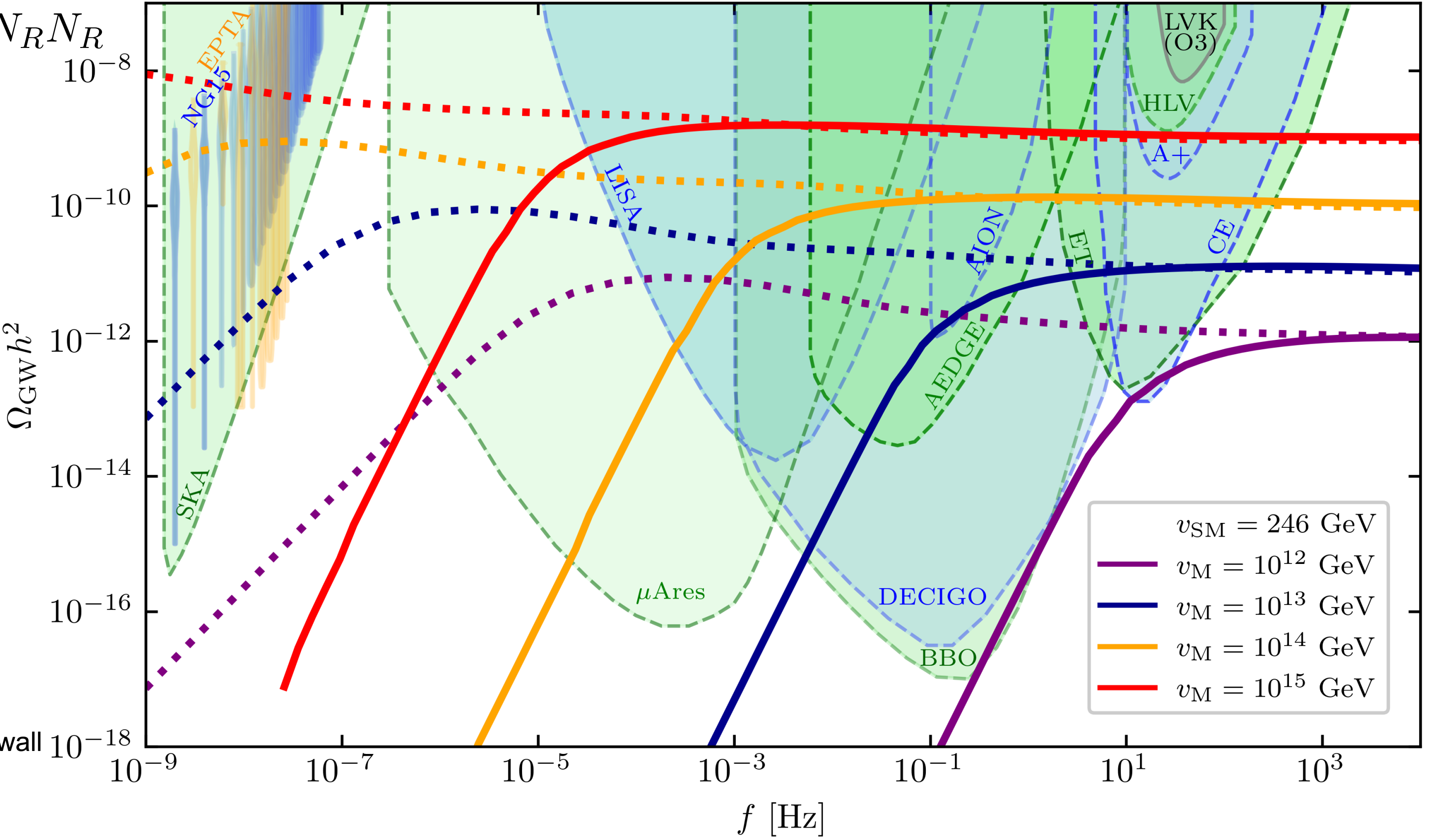
Benchmark Point	u [GeV]	V_{bias} [GeV ⁴]	$y_{\text{max}}(M_{\Delta} < M_{\text{Pl}})$
①	10^5	10^{-5}	4.93
②	5.2×10^7	7.14×10^{10}	0.216
③	1.2×10^9	10^{19}	0.045
④	2×10^{11}	2.5×10^{32}	0.0035

$$V(\sigma) = \frac{\lambda}{4}(\sigma^2 - u^2)^2 \quad \Delta V(\sigma) = \epsilon u \sigma \left(\frac{\sigma^2}{3} - u^2 \right)$$

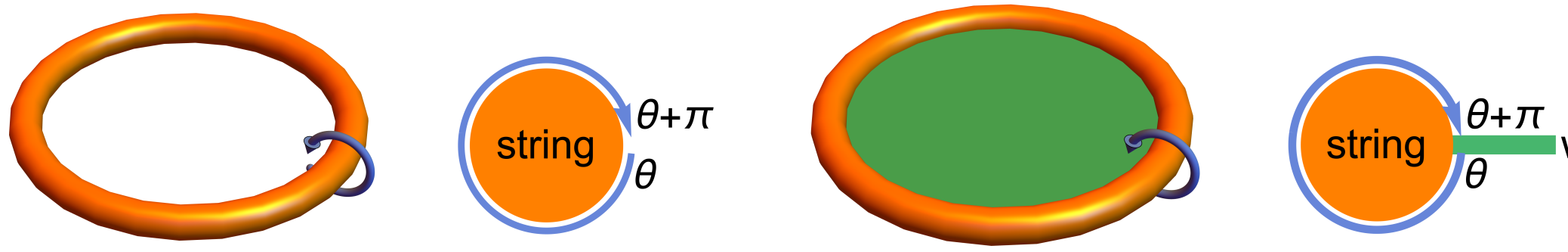
Type-I two-Higgs-doublet model and gravitational waves from domain walls bounded by strings

$$Y_u \bar{Q} \tilde{\Phi}_2 u_R + Y_d \bar{Q} \Phi_2 d_R + Y_e \bar{L} \Phi_2 e_R + Y_N \bar{L} \tilde{\Phi}_2 N_R + y_N \phi N_R N_R$$

	$u_{R\beta}$	$d_{R\beta}$	Q_α	L_α	$e_{R\beta}$	$N_{R\beta}$	Φ_2	Φ_1	ϕ
$SU(2)_L$	1	1	2	2	1	1	2	2	1
$U(1)_Y$	$\frac{2}{3}$	$-\frac{1}{3}$	$\frac{1}{6}$	$-\frac{1}{2}$	-1	0	$\frac{1}{2}$	$\frac{1}{2}$	0
$U(1)_R$	1	-1	0	0	-1	1	1	-1	-2
residual Z_2	-	-	+	+	-	-	-	-	+



$$U(1)_R \xrightarrow{\langle \phi \rangle} Z_2 \xrightarrow{\langle \Phi_i \rangle} \text{nothing}$$



Vilenkin and Everett 1982

(a) Before Z_2 symmetry breaking.

(b) After Z_2 symmetry breaking.

Surface tension in the walls causes the combined relic to decay earlier than strings

DW decay without V_{bias} !

Conclusion

- GWs can probe new physics BSM at HE, only a few examples here: FOPT, CS, DW (+combos)
- FOPT at QCD scale can describe NANOGrav
- CS $U(1)_{B-L}$ gauged w/GUTs; global w/Majorons
- DW Z_2 w/QG bias; Majorana vs Dirac
- DW bounded by CS in 2HDM (type I)