

# Gravitinos, Modulinos & Neutrinos

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Part based on:

L. Anchoredqui, I. Antoniadis, K.B., J. Cunnat, D. Lust, JHEP 2024

Corfu, September 5<sup>th</sup>, 2024

Let's consider a space-time

Minkowski  $\nearrow$   $\mathcal{M}_4 \otimes X_d \nwarrow$   $d$  extra-dimensions

SUGRA in  $d+4$   $\longrightarrow$   $N=1$  in 4 dimensions

We will discuss: gravitino & fermion masses

Phenomenological relevance  $\rightarrow$  SUSY must be broken

Supergravity : gravitino  $\xrightarrow{\text{SUSY}}$  mass  $m_{3/2}$

We assume that the supersymmetric theory has

massless fields singlet under the gauge symmetry

$\rightarrow$  moduli

SUSY breaking  $\rightarrow$  masses for the scalar moduli  
masses for the fermion **modulinos**

$$g = k + \ln |w|^2$$

Fermion masses

$$(M_f)_{\alpha\beta} = e^{G/2} \left[ g_{\alpha\beta} - g_{\alpha\beta\bar{\gamma}} g^{\bar{\gamma}} + \frac{1}{3} g_{\alpha} g_{\beta} \right]$$

$$g_{\alpha} = \frac{\partial g}{\partial z^{\alpha}}; \quad g_{\bar{\alpha}} = \frac{\partial g}{\partial \bar{z}^{\alpha}}; \quad g^{\alpha} = g_{\alpha} g^{\alpha\bar{\gamma}}; \quad g^{\alpha\bar{\beta}} g_{\bar{\beta}\alpha} = \delta^{\alpha}_{\alpha}$$

$$\langle e^{G/2} \rangle = m_{3/2} \quad \text{is the gravitino mass}$$

Typically, we get  $M_f \sim m_{3/2}$

# Gravitino Swampland Conjecture

$\Rightarrow m_{3/2}$  can NOT be arbitrarily small

$m_{3/2} \rightarrow 0 \Rightarrow$  tower of states have mass  $\rightarrow 0$

$\Downarrow$   
EFT no more valid

## Gravitino Swampland Conjecture:

$m_{3/2}$  Can NOT be arbitrarily small:

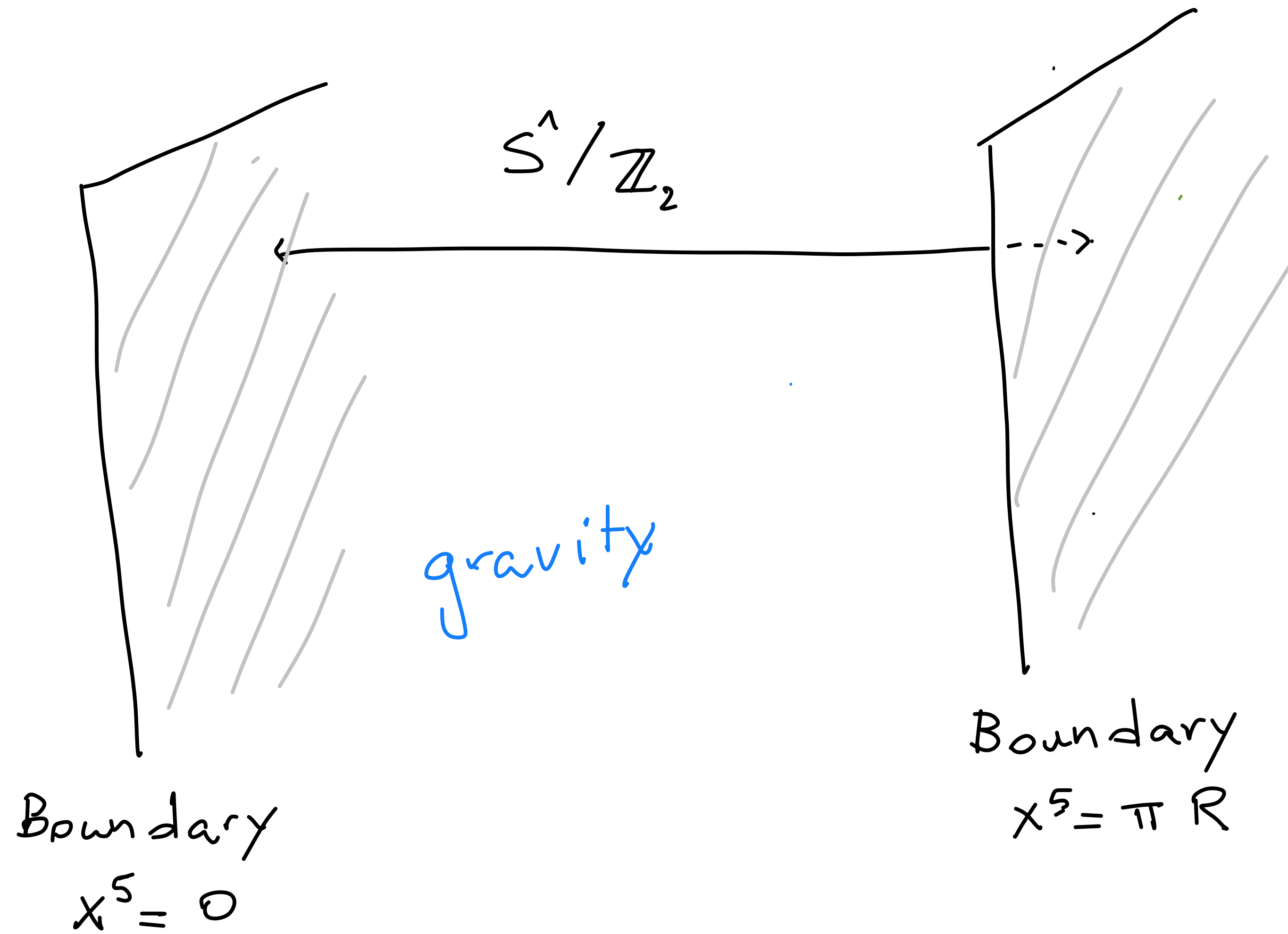
$m_{3/2} \rightarrow 0 \Rightarrow$  Infinite tower with masses  $\rightarrow 0$

EFT no more valid

Antoniadis-Bachas-Lewellen-Tomaras

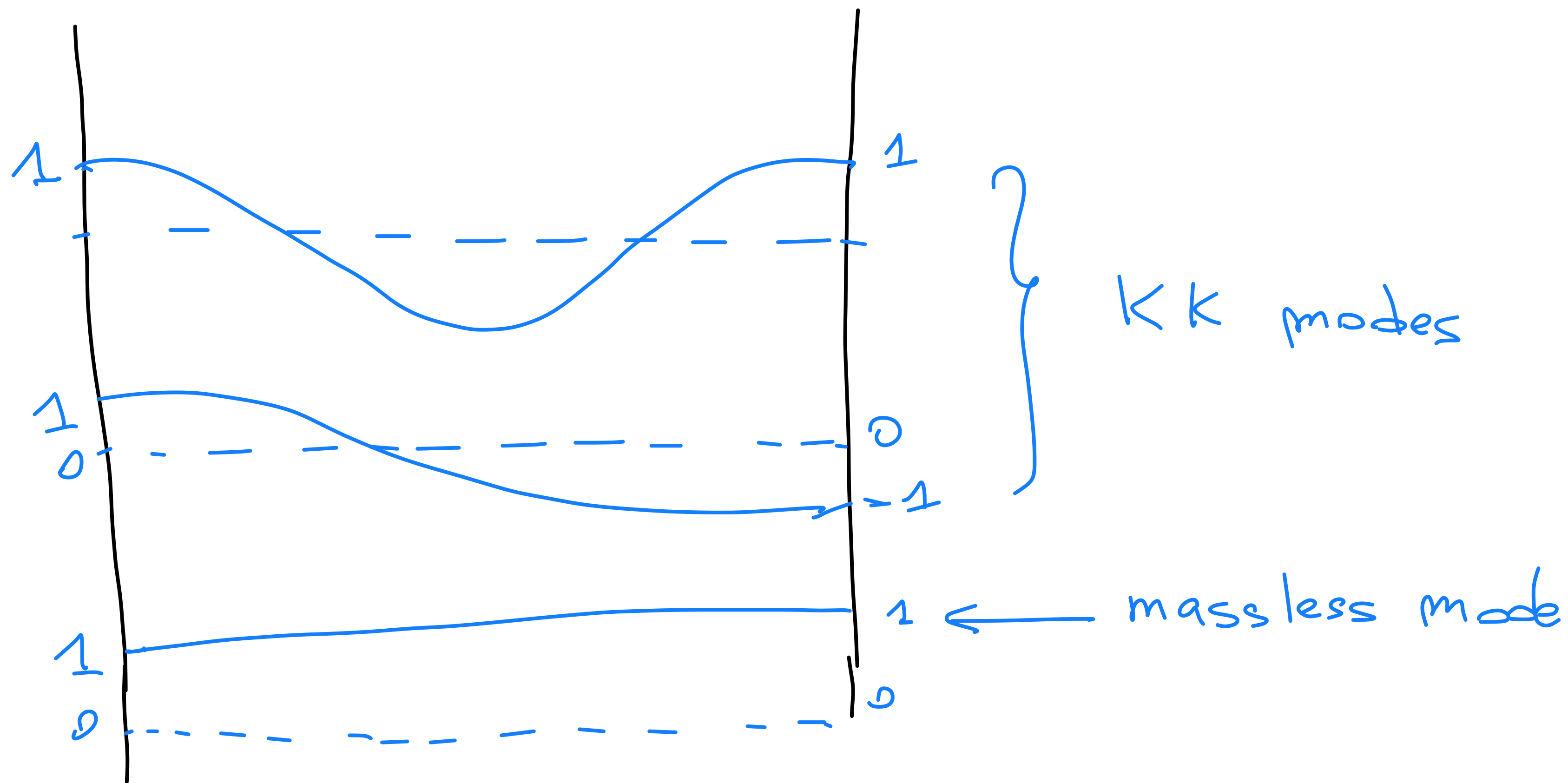
Cribiori-Lust-Scalisi

One extra-dimension



KK modes of the gravitino

SUSY





Let's break SUSY as a deformation of our previous set-up

two scenarios:

- ① Non-periodic boundary conditions
- ② Some dynamics on the boundary

Two usual scenarios to break SUSY:

- ① Non-trivial boundary conditions
- ② Dynamics on the brane

① Non-periodic condition for the gravitino

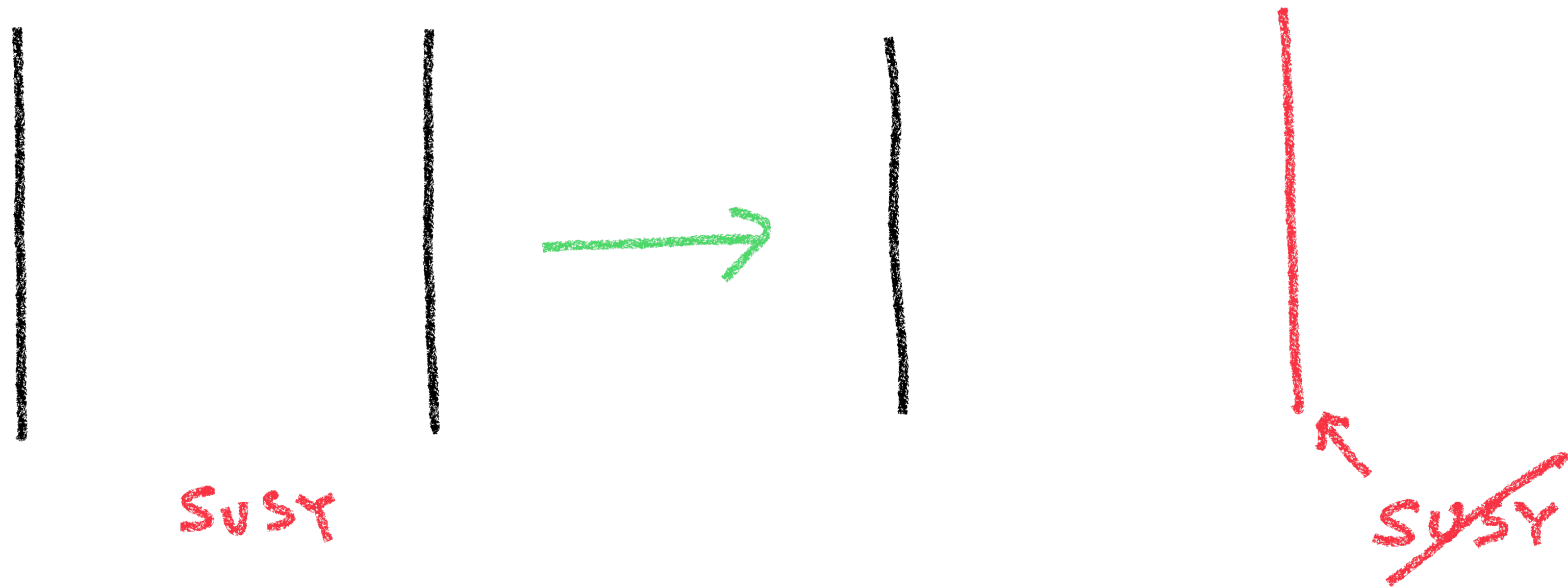
$$\rightarrow m_{3/2} = \frac{\alpha}{R} \quad 0 < \alpha \leq \frac{1}{2}$$

Scherk - Schwarz mechanism

Gravitino Swampland Conjecture  $\Rightarrow \alpha$  not arbitrarily small

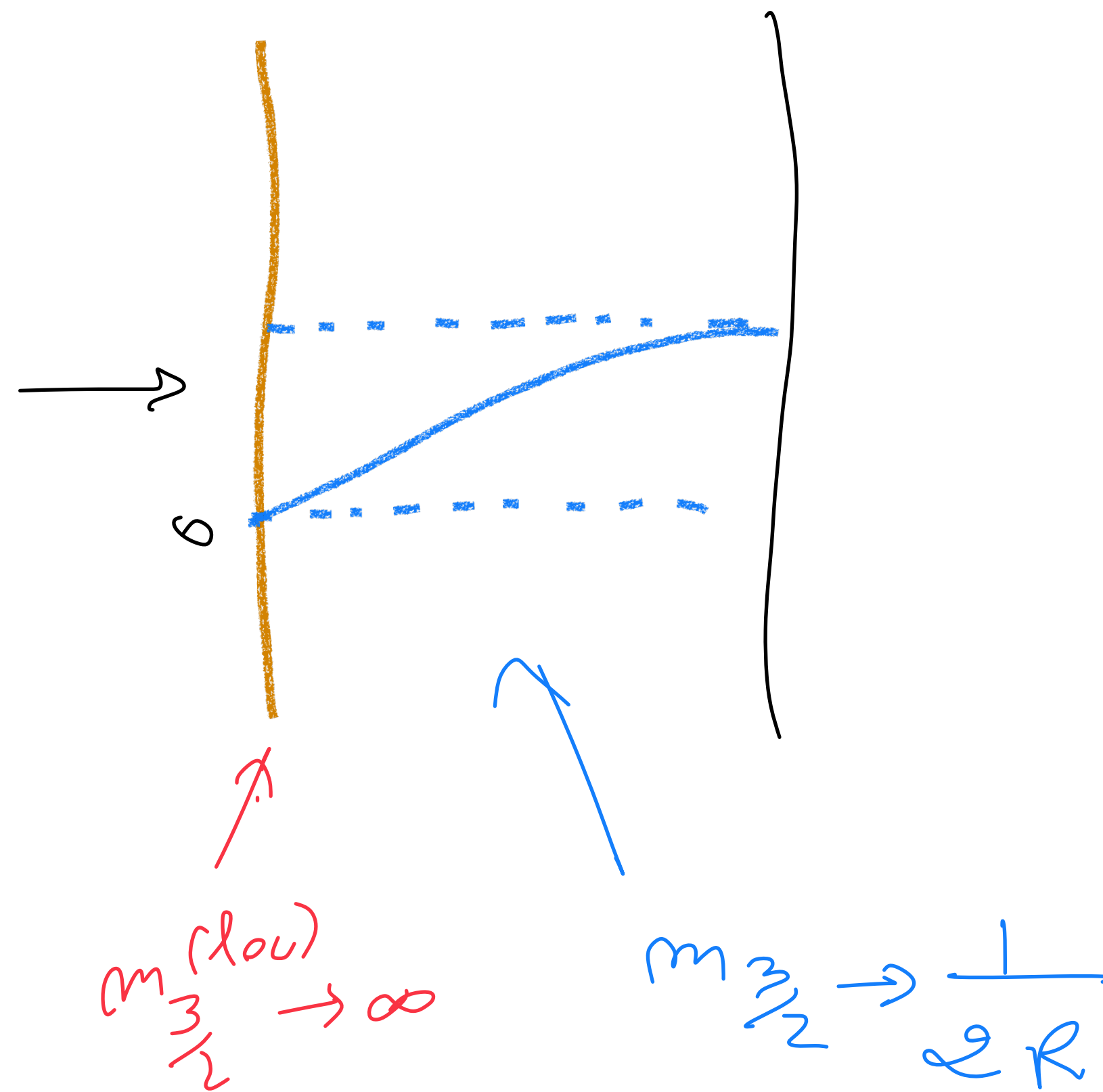
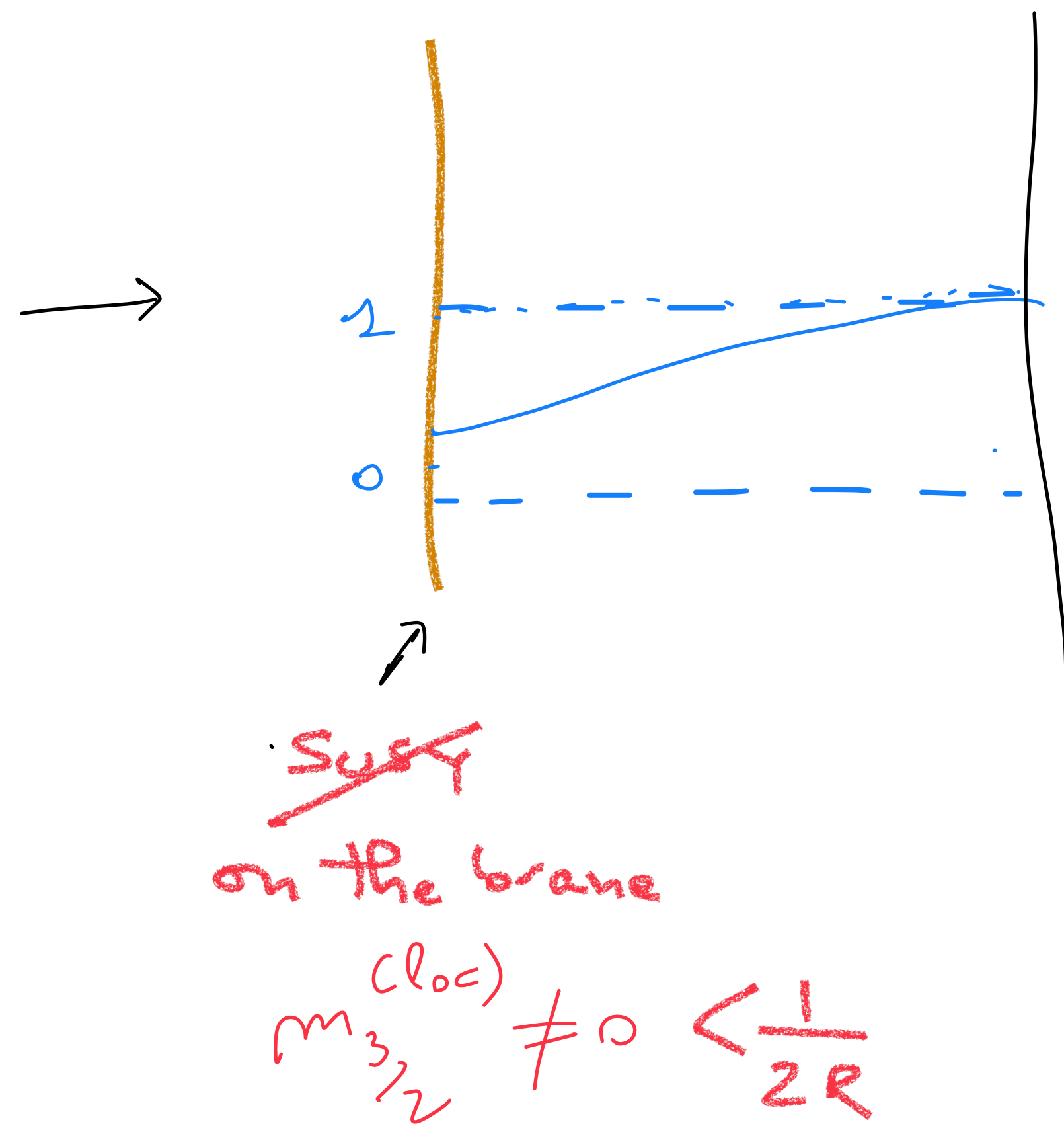
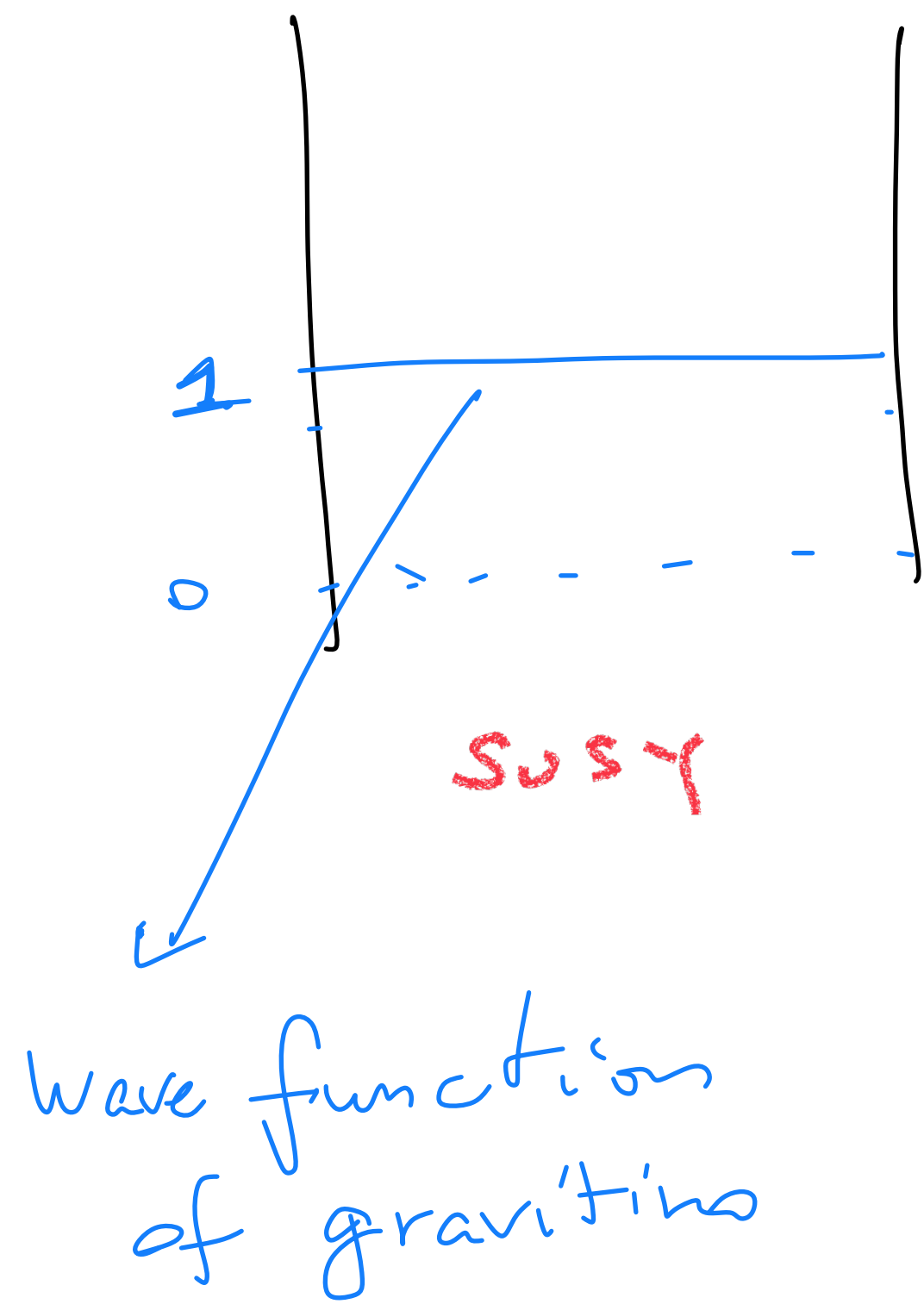
⊛ Note also that  $m_{3/2}$  is of order of  $\text{TeV}$  scale  
( $m_{3/2} < \frac{1}{R}$ )

② Breaking SUSY on the boundary:



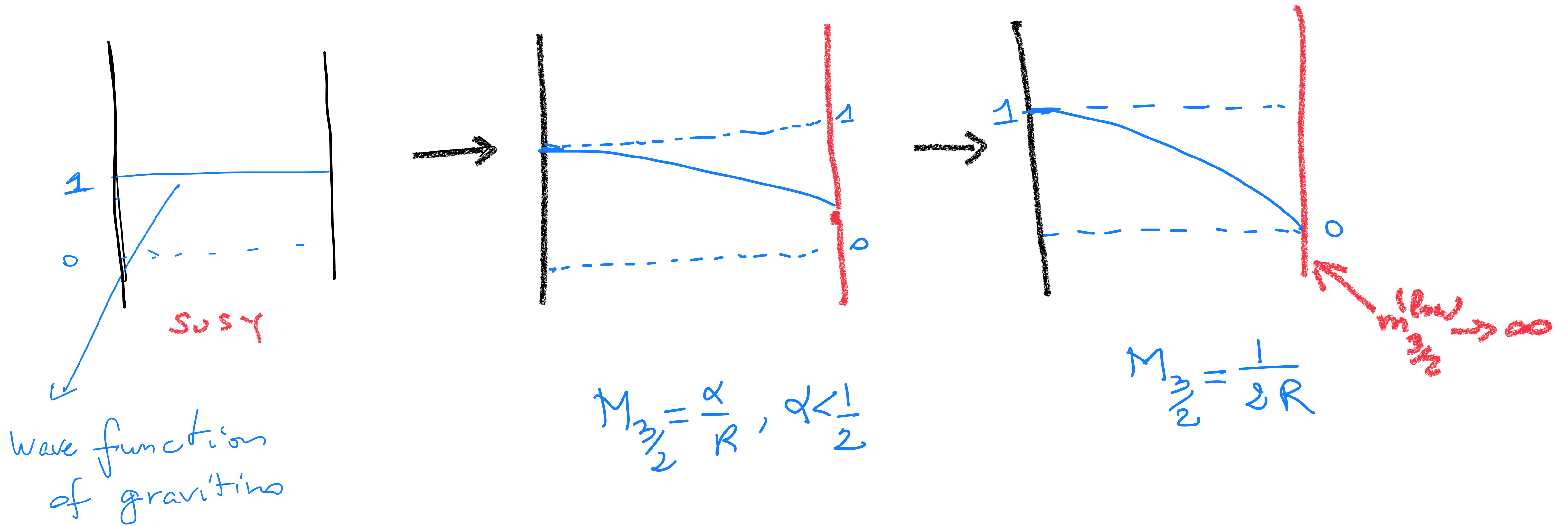
what is the mass of the gravitino?

$$\mathcal{L} \rightarrow \mathcal{L} + \delta(x^5) m_{3/2}^{(loc)} \bar{\Psi}_\mu \psi^\mu$$



Note gravitino mass is bounded by  $m_{KK}$

$$\mathcal{L} \rightarrow \mathcal{L} + \delta(x^5 - \pi R) m_{3/2}^{(loc)} \bar{\Psi}^\mu \Psi_\mu$$



Note: The gravitino mass is of order of KK scale

The gravitino mass formula:

$$M_{3/2} = \left[ n + \frac{1}{\pi} \arctan \left[ \frac{4\pi R m_{3/2}^{(loc)}}{\pi^2 R^2 m_{3/2}^{(loc)2} + 4} \right] \right] \frac{1}{R}$$

Note:  $M_{3/2}$  always smaller than the KK scale

For the case of many dimension, the generalization would be:

The gravitino wave-function spreads such that:

$$M_{3/2} < m_{KK}^{(1)}$$

$m_{KK}^{(1)}$  the lightest KK mode in the direction where

SUSY is broken



Modulino mass

- Let's stay with the  $S^1/\mathbb{Z}_2$  example
- We consider the case of Scherk-Schwarz breaking  
with  $\alpha = \frac{1}{2}$
- This can be realized in superstrings

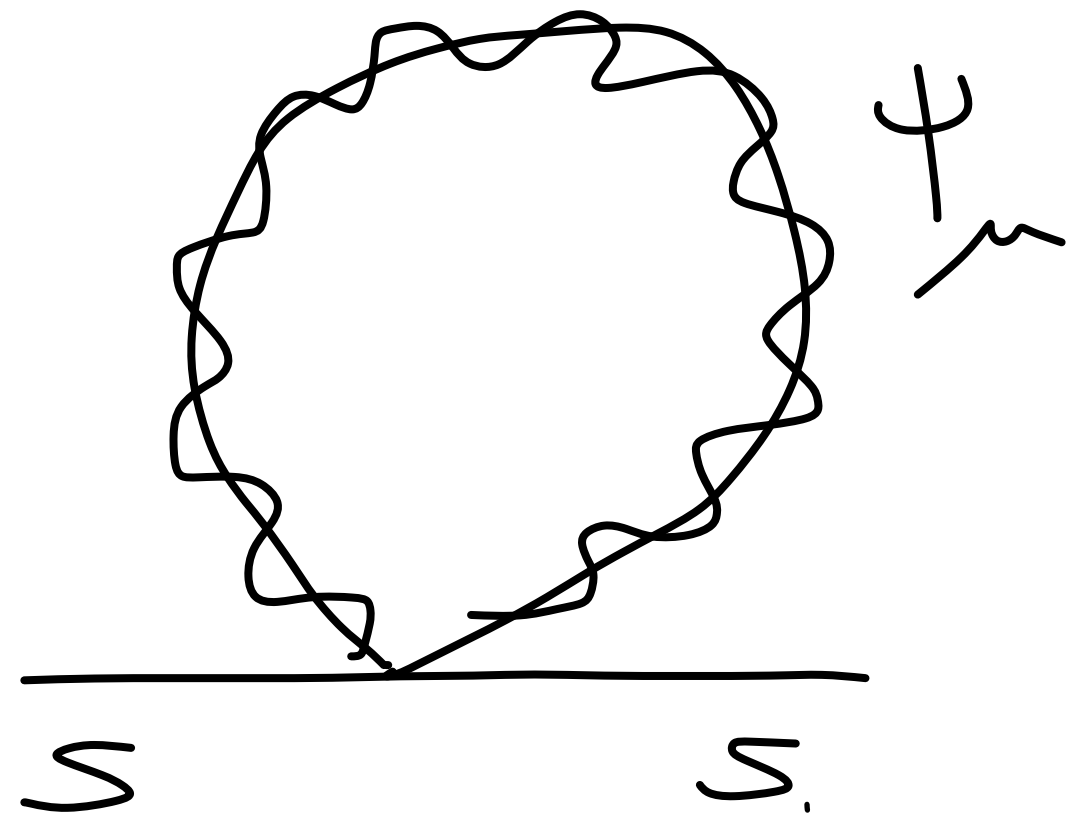
$$M_{3/2} = \frac{1}{2R}$$

Not described by 4D  $N=1$  SUGRA

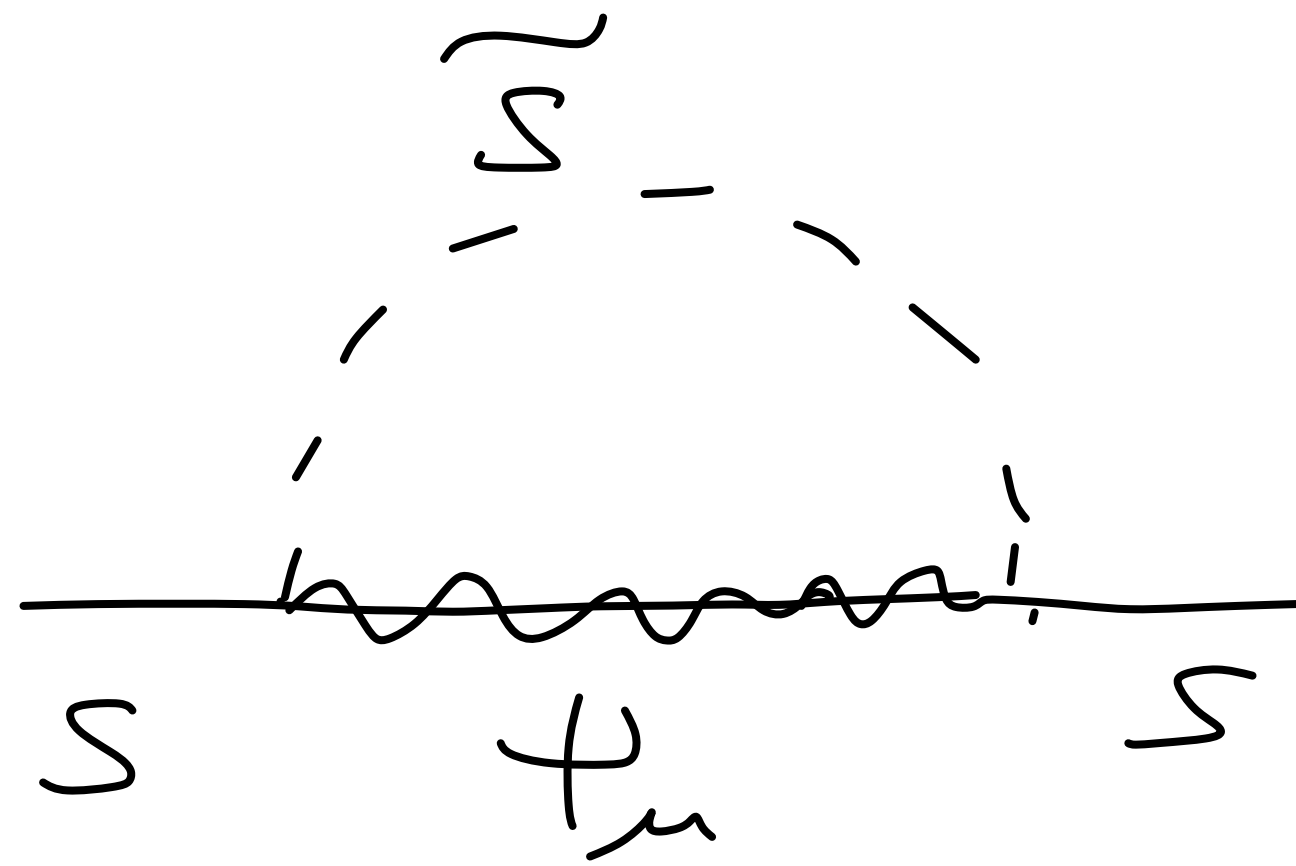
What is the expected size for fermion states localized on the boundary (modulino on boundary)

\* Bulk: gravitino  $M_{3/2} = \frac{1}{2R}$

\* Boundary modulino  $M_{\psi} = ?$



+



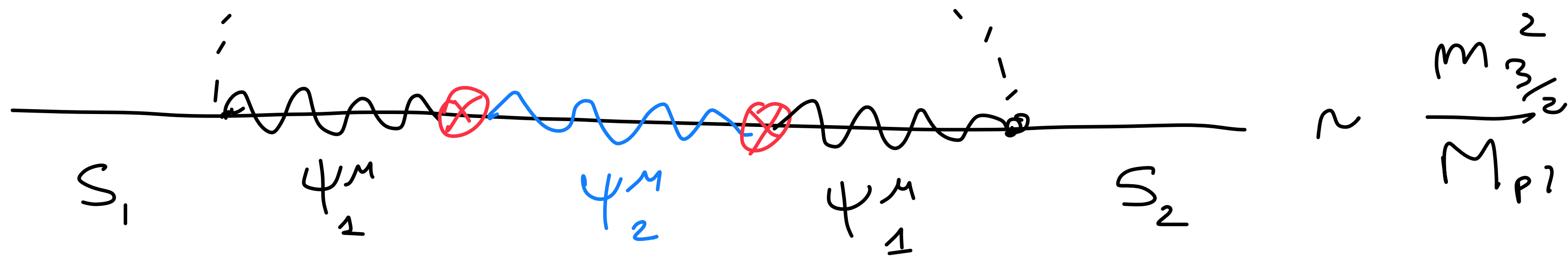
$\rightarrow 0$

Gherghetta - Riotto  
Antoniadis - KB - Quiros

Two issues:

① There is an R-symmetry restored @  $\alpha = \frac{1}{2}$   
 $\Rightarrow$  No Majorana masses

② There are two gravitini, so you can get a Dirac mass



We can get a modulus of mass:

$$M_f \sim m_{3/2}$$

or

$$M_f \sim \frac{m_{3/2}}{M_p}$$

What about coupling to matter

Two possibilities:

$$M_f \sim m_{3/2}$$

or

$$M_f \sim \frac{m_{3/2}^2}{M_p}$$

# Neutrino-modulino oscillations

K.B. Y. Smirnov

Anchor doqui - Antoniadis - K.B. Lunat - List

sterile neutrino of mass  $\sim 50$  eV can be detected  
by FPF @ CERN

For  $m_{3/2} = \frac{1}{2R} \sim 250$  TeV

$\rightarrow \frac{m_{3/2}^2}{M_{pl}} \sim 50$  eV  $\Rightarrow$  Modulino = sterile neutrino?

$\rightarrow \alpha m_{3/2} \rightarrow$  "MSSM" states @ few TeV  
"gauge coupling"  $\uparrow$   
(gaugino mediation from  $X^S \rightarrow \mathcal{N}_4$ )

Such a small  $\pi_{3/2}$  mass might also be obtained as

$$\sim \left( \frac{1}{2R} \right)^3 \frac{1}{M_s^2}$$

$$M_s \sim 3 \cdot 10^9 \text{ GeV}$$

$$R \sim 3 \cdot 10^8 \text{ GeV}$$



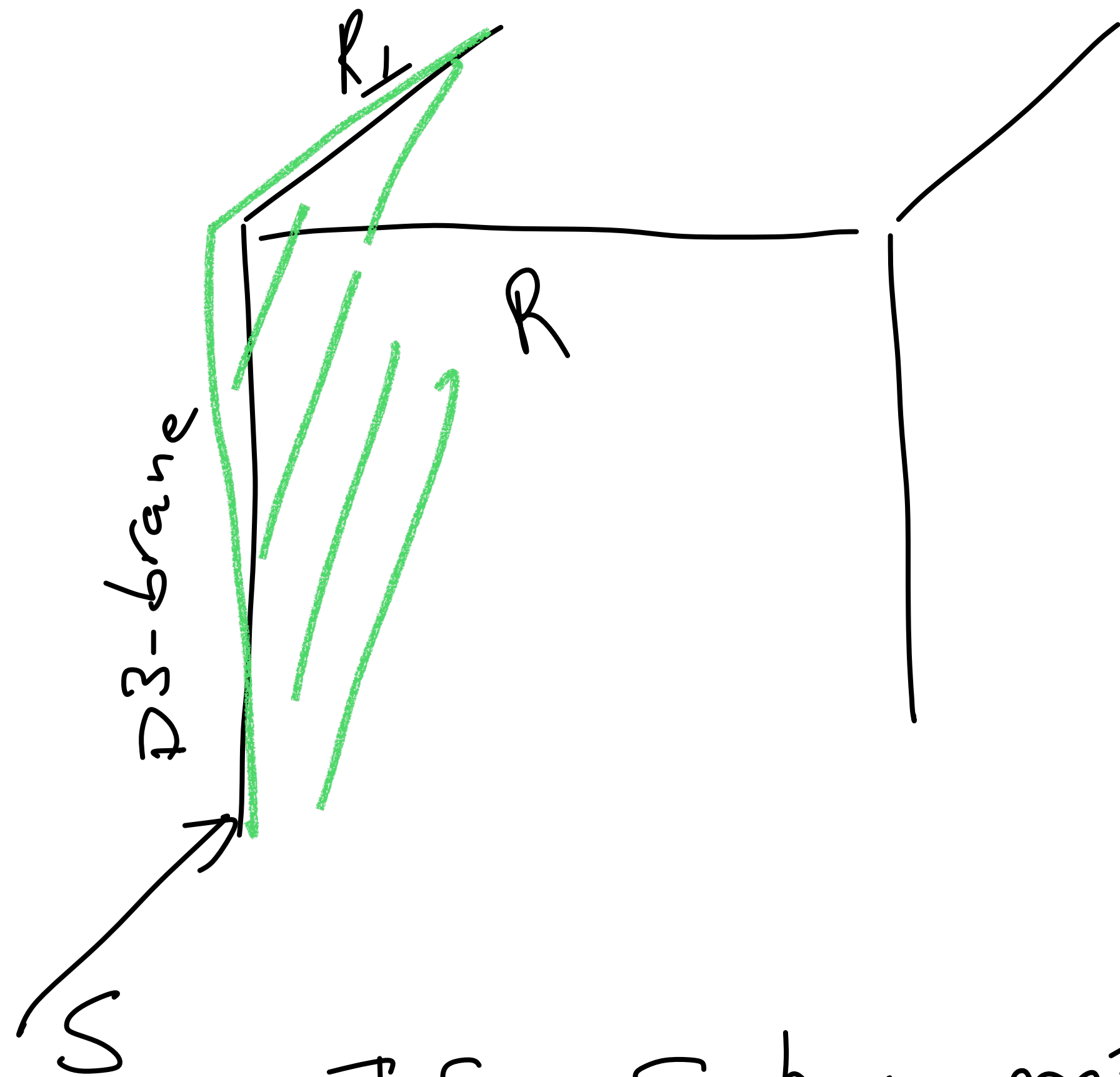
Modulino = sterile neutrino

IF

Modulino mixes with neutrino

Very model dependent

For sterile neutrino with mass  $< 10$  keV  
no strong constraints on mixing.



If  $S$  has mixing through

$$\frac{1}{R_{\perp}} \sim eV$$

Dark dimension

$$\frac{1}{R} \sim 250 \text{ TeV}$$

$$M_s \sim 10^{10} \text{ GeV}$$

Yukawa coupling  $\lambda S L H$  ?

1<sup>st</sup> possibility

R-parity  $\mu$ -term

$\mu L H$

$$\mu = m_{3/2}$$

Giudice - Masiero  
mechanism

$$\frac{\langle S \rangle}{M_s} \alpha$$

dependance  
on moduli

coupling

Note that

$$\alpha \frac{m^{3/2}}{M_{Pl}} \sim \alpha \frac{1}{R} \frac{1}{R R_{\perp} M_S^3}$$

$$\frac{1}{(R M_S)^2} \quad \frac{1}{R_{\perp} M_S}$$

coupling to  
state in the XS dimension

← Already there

⇒ ~~SUSY~~ not necessarily related to mixing

$$\alpha \frac{m_{3/2}}{M_{pl}} SLH \Rightarrow \alpha \frac{m_{3/2} \langle H \rangle}{M_{pl}} SL$$

$\Rightarrow$  mixing of order 0.1

$\Rightarrow$  The range of mass & mixing that

**FPE will probe**

## Conclusions

- \* Upper bound on gravitino masses?
  - Large volume scenario
- \* Modulinos are not subject of a lot of interest
  - "  $M_f \sim m_{3/2} \gg \text{masses of interest}$  "
- \* But  $M_f$  can be light and lead to new phenomena