Prospects for Gravitational Wave Signals from Higgs Inflation

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Based on R.I.P. by the Sinchon GW Group: Injun Cheon, JK, Seong Chan Park, Yeji Park, Stefano Scopel, Juhoon Son, Liliana Velasco Sevilla

A New Window to the Universe







Caltech/MIT/LIGO Lab

- Discovery of gravitational waves (GW) by LIGO & Virgo LIGO/Virgo, PRL 116 (2016), PRL 118 (2017)
- Source: mergers of black holes and neutron stars
- → GW astronomy and cosmology

A Window to the Very Early Universe?

- Pulsar Timing Arrays → Evidence for Stochastic GW Background NANOGrav, ApJL 951 (2023) EPTA, InPTA, A&A 678 (2023) Parkes PTA, ApJL 951 (2023) CPTA, RAA 23 (2023)
- Lower frequency than LIGO/Virgo events
 Mergers of supermassive black holes?
- More interesting: particle physics origin
 - First-Order Phase Transition (FOPT)
 - Cosmic strings



Champion/MPI for Radio Astronomy



Looking Forward to More Discoveries

THE GRAVITATIONAL WAVE SPECTRUM



https://www.astro.gla.ac.uk/users/martin/powersof60/images/gwspectrum.jpg

First-Order Phase Transitions

See talks by Merchand, Kowalska



Kinnunen et al., Rep. Prog. Phys. 81 (2018)

- High temperature: potential minimum at $\phi = 0$
- 2 $T < T_c$: deeper mimimum at $\phi = v$, separated by barrier
- Tunneling ~ bubbles of true vacuum
- GW from bubble collisions and turbulence

Gravitational Waves from Particle Physics

- SM: no FO electroweak PT for Higgs mass ≳ 70 GeV Kajantie et al., PRL 77 (1996); Karsch et al., NP Proc. Suppl. 53 (1997) Csikor et al., PRL 82 (1999)
- GW from thermal fluctuations in plasma Ghiglieri & Laine, JCAP 07 (2015)
- Peak at frequency ~ GHz
- Amplitude enhanced for
 - Very weakly interacting new particles Drewes et al., JCAP 06 (2024)
 - Gauss-Bonnet Cosmology Biswas, ..., Scopel, Velasco Sevilla, 2405.15998



SM Extensions with First-Order Phase Transitions

• Most famous: SUSY with light stop

Carena et al., PLB **380** (1996); Espinosa, NPB **475** (1996) Delepine et al., PLB **386** (1996); Cline & Kainulainen, NPB **482** (1996); ...

- Minimal from model building perspective: 2HDM Dorsch et al., JHEP 10 (2013); Basler et al., JHEP 02 (2017) Andersen et al., PRL 121 (2018); ...
- Minimal from EFT perspective: higher-dimensional operators Zhang, PRD 47 (1993); Grojean et al., PRD 71 (2005) Bödeker et al., JHEP 02 (2005); Chala et al., JHEP 07 (2018); ...
- Many recent works Review: Roshan & White, arXiv:2401.04388 See talks by Rubio, Merchand, Koutroulis, Gorbunov, Kowalska, Wang, King

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Bezrukov & Shaposhnikov, PLB 659 (2008)

$$\mathcal{L}_J = \sqrt{-g_J} \left[\frac{M_{\mathsf{P}}^2}{2} R_J + \xi \phi^{\dagger} \phi R_J - g_J^{\mu\nu} (D_{\mu} \phi)^{\dagger} (D_{\nu} \phi) - V_J(\phi) + \dots \right]$$

- J: Jordan frame
- R: Ricci scalar
- V_J: SM Higgs potential
- ξ : non-minimal coupling to gravity
 - Consistent with all symmetries
 - Required for renormalization in curved spacetime See talk by Rubio
 - Not necessarily small

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- Def.: conformal factor $\Omega^2(\phi) \equiv 1 + 2\xi \frac{\phi^T \phi}{M_p^2}$

$$\mathcal{L}_J = \sqrt{-g_J} \left[\frac{M_{\mathsf{P}}^2}{2} \, \Omega^2(\phi) \, R_J - g_J^{\mu\nu} (D_\mu \phi)^{\dagger} (D_\nu \phi) - V_J(\phi) + \dots \right]$$

• Unitary gauge:
$$\phi = \begin{pmatrix} 0 \\ (\varphi + v)/\sqrt{2} \end{pmatrix}$$

- Weyl transformation to Einstein frame: $g_J \rightarrow g_E \equiv \Omega^2(\varphi) g_J$
- Canonical normalization: $\chi = \int_{0}^{\varphi} d\varphi \sqrt{\frac{3}{2} \frac{(\Omega^{2}, \varphi)^{2}}{\Omega^{4}} + \frac{1}{\Omega^{2}}}$

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$$\mathcal{L}_{E} = \sqrt{-g_{E}} \left[\frac{M_{P}^{2}}{2} R_{E} - \frac{1}{2} g_{E}^{\mu\nu} (\partial_{\mu}\chi) (\partial_{\nu}\chi) - V_{E}(\chi) + \dots \right]$$

- Minimal coupling to gravity but higher-dimensional operators
- Gravity remains weak below cutoff ~ M_P Ω
- Potential flat for large field values ~> slow-roll



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FOPT and Inflation from the Higgs?

● SM + higher-dimensional operators ~> FO EWPT

$$V(\phi) = -\mu^2 \phi^{\dagger} \phi + \lambda (\phi^{\dagger} \phi)^2 + \frac{C_6}{f^2} (\phi^{\dagger} \phi)^3 + \frac{C_8}{f^4} (\phi^{\dagger} \phi)^4$$

- f: cutoff
- c₆, c₈: dimensionless couplings
- SM + non-minimal coupling to gravity ~> inflation

Can we have both?

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CMB measurements

 $\begin{array}{ll} \mbox{Scalar power spectrum amplitude} & \mbox{In}(10^{10}\,A_s) = 3.044 \pm 0.014 \\ \mbox{Scalar spectral index} & n_s = 0.9649 \pm 0.0042 \\ \mbox{Tensor-to-scalar power ratio} & r < 0.036 \ (95\% \ \mbox{CL}) \\ \mbox{Planck, A&A 641 (2020); BICEP/Keck, PRL 127 (2021)} \\ \end{array}$

- → Strong contraints on inflationary parameters
- → Strong bounds on higher-dimensional operators

Effects of Higher-Dimensional Operators



Attractor point for $\xi \gg 1 \rightarrow$ effects of c_6 and c_8 reduced

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CMB Constraints



- Very strong constraints for $\xi \lesssim 1$
- Room for new physics at high scales ($f \leq M_P$) for $\xi \gg 1$
- No way to affect EWPT (requires *f* ~ TeV)

- $U(1)_X$ gauge symmetry
- SM singlet scalar with $U(1)_X$ charge
 - Spontaneously breaks $U(1)_X \rightsquigarrow \text{FOPT} \rightsquigarrow \text{GW}$
 - Non-minimal coupling to gravity ~> Higgs-inflation-like inflation
- Optional: fermion(s) with mass from Yukawa coupling
- Very weak coupling to SM (details t.b.d.) ~ reheating

Effective Scalar Potential at Finite Temperature

$$V(h_{c}, T) = V_{\text{tree}}(h_{c}) + V_{1-\text{loop}}(h_{c}) + V_{\text{th}}(h_{c}, T)$$

$$V_{\text{tree}}(h_{c}) = -\frac{\mu^{2}}{2}h_{c}^{2} + \frac{\lambda}{4}h_{c}^{4}$$

$$V_{1-\text{loop}}(h_{c}) = \sum_{i=h,\chi,g,f} \frac{n_{i}}{64\pi^{2}}m_{i}^{4}(h_{c})\left[\ln\frac{|m_{i}^{2}(h_{c})|}{v^{2}} - C_{i}\right]$$

$$V_{\text{th}}(h_{c}, T) = \sum_{i=h,\chi,g} \frac{n_{i}}{2\pi^{2}}T^{4}J_{b}\left(\frac{m_{i}^{2}(h_{c})}{T^{2}}\right) + \frac{n_{f}}{2\pi^{2}}T^{4}J_{f}\left(\frac{m_{i}^{2}(h_{c})}{T^{2}}\right)$$

$$m_{h}^{2}(h_{c}) = -\mu^{2} + 3\lambda h_{c}^{2}$$

$$m_{\chi}^{2}(h_{c}) = -\mu^{2} + \lambda h_{c}^{2}$$

$$m_{g}^{2}(h_{c}) = \frac{g^{2}}{4}h_{c}^{2}$$

2 options for generating potential barrier around $T \sim v$ (so far)

- No fermions, $\lambda \ll 1$ See talk by Kowalska
- 1 (Dirac) fermion, $\lambda \sim 0.1$

→ FOPT possible

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GW spectrum determined by

- Nucleation temperature T_n
- $\alpha \leftrightarrow \text{strength of PT}$
- $\beta \nleftrightarrow duration$

See talks by Merchand, Wang

Calculated with help from CosmoTransitions Wainwright, Comput. Phys. Commun. 183 (2012)

Example: 1 Fermion

Case 3. U(1) model

- VEV: $\Lambda=10^{15}\;\text{GeV}$
- · D.O.F of particles: (scalar, gauge boson, fermion) =(1+1, 3, 4)
- coupling constants

0.10

- scalar self-coupling: λ = 0.09
- gauge coupling: g = 1
- Yukawa coupling: y = 1





Slide by Yeji Park

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Resulting Gravitational Wave Spectrum



Plot by Liliana Velasco Sevilla

- Higgs inflation doesn't like higher-dimension operators in potential
- Dark $U(1)_X \sim$ gravitational waves and inflation from same scalar
- Ongoing work
 - Constraints on and from inflation
 - Map parameter space
 - Calculation of GW predictions
- Future directions
 - Inflaton coupling to SM ~ reheating
 - Lower scale of symmetry breaking
 - Different gauge groups
 - Dark Matter