Towards spectral evolution Andrzej Sitarz

Supported by the Polish National Science Centre grant 2020/37/B/ST1/0154

Jagiellonian University

Workshop on Noncommutative and Generalized Geometry in String Theory, Gauge Theory and Related Physical Models **CORFU 2024**

ALGEBRA, TOPOLOGY, ANALYSIS, GEOMETRY

ALGEBRAIC APPROACH

- algebras (generators and relations): nice, but?
- depends on the choice of the (polynomial) basis (!)

ALGEBRA, TOPOLOGY, ANALYSIS, GEOMETRY

ALGEBRAIC APPROACH

- algebras (generators and relations): nice, but?
- depends on the choice of the (polynomial) basis (!)

TOPOLOGICAL APPROACH

- C* algebras (normed, complete algebras): nice, but?
- topology does not see any smoothness (!)

ALGEBRA, TOPOLOGY, ANALYSIS, GEOMETRY

ALGEBRAIC APPROACH

- algebras (generators and relations): nice, but?
- depends on the choice of the (polynomial) basis (!)

TOPOLOGICAL APPROACH

- C* algebras (normed, complete algebras): nice, but?
- topology does not see any smoothness (!)

FUNCTIONAL ANALYSIS

- analysis (mostly) on operator algebras
- the core of (spectral) noncommutative approach
- **•** unifies algebraic and topological approach

GEOMETRY AND SMOOTHNESS

Classical geometry is differential

- an orientable manifold *M*, smooth functions, *C*∞(*M*),
- differential algebra Ω(*M*), metric *g* µν, Laplace operator ∆,
- spin*^c* structure(s), real spin structure, Dirac operator

GEOMETRY AND SMOOTHNESS

Classical geometry is differential

- an orientable manifold *M*, smooth functions, *C*∞(*M*),
- differential algebra Ω(*M*), metric *g* µν, Laplace operator ∆,
- spin*^c* structure(s), real spin structure, Dirac operator
- All definitions use (sometimes in a hidden way)
	- **o** differentiation (derivatives)
	- **differential operators** and their properties

GEOMETRY AND SMOOTHNESS

Classical geometry is differential

- an orientable manifold *M*, smooth functions, *C*∞(*M*),
- differential algebra Ω(*M*), metric *g* µν, Laplace operator ∆,
- spin*^c* structure(s), real spin structure, Dirac operator
- All definitions use (sometimes in a hidden way)
	- **o** differentiation (derivatives)
	- **differential operators** and their properties

Differential operators - as **operators on a Hilbert space**

- **•** came with the dawn of quantum mechanics
- the core of noncommutative (spectral) approach

The **significance** of (differential) operators

Most of classical geometry can be encoded in terms of bounded and unbounded operators acting on a separable Hilbert space.

The **significance** of (differential) operators

Most of classical geometry can be encoded in terms of bounded and unbounded operators acting on a separable Hilbert space.

How do we reconstruct geometry ?

The **significance** of (differential) operators

Most of classical geometry can be encoded in terms of bounded and unbounded operators acting on a separable Hilbert space.

How do we reconstruct geometry ?

1 differential calculus: $da = [D, a]$ and FODC $Ω¹(M)$

The **significance** of (differential) operators

Most of classical geometry can be encoded in terms of bounded and unbounded operators acting on a separable Hilbert space.

How do we reconstruct geometry ?

- **1** differential calculus: $da = [D, a]$ and FODC $Ω¹(M)$
- ² metric: *d*(*x*, *y*) = sup||[*D*,*f*]||≤¹ |*f*(*x*) − *f*(*y*)|

The **significance** of (differential) operators

Most of classical geometry can be encoded in terms of bounded and unbounded operators acting on a separable Hilbert space.

How do we reconstruct geometry ?

- **1** differential calculus: $da = [D, a]$ and FODC $Ω¹(M)$
- ² metric: *d*(*x*, *y*) = sup||[*D*,*f*]||≤¹ |*f*(*x*) − *f*(*y*)|
- ³ additional connection (if spinors twisted by a vector bundle)

The **significance** of (differential) operators

Most of classical geometry can be encoded in terms of bounded and unbounded operators acting on a separable Hilbert space.

How do we reconstruct geometry ?

- **1** differential calculus: $da = [D, a]$ and FODC $Ω¹(M)$
- ² metric: *d*(*x*, *y*) = sup||[*D*,*f*]||≤¹ |*f*(*x*) − *f*(*y*)|
- ³ additional connection (if spinors twisted by a vector bundle)

 2990

⁴ dimension (growth of eigeinvalues: $N_D(Λ) \sim Λ^d$),

The **significance** of (differential) operators

Most of classical geometry can be encoded in terms of bounded and unbounded operators acting on a separable Hilbert space.

How do we reconstruct geometry ?

- **1** differential calculus: $da = [D, a]$ and FODC $Ω¹(M)$
- ² metric: *d*(*x*, *y*) = sup||[*D*,*f*]||≤¹ |*f*(*x*) − *f*(*y*)|
- ³ additional connection (if spinors twisted by a vector bundle)
- ⁴ dimension (growth of eigeinvalues: $N_D(Λ) \sim Λ^d$),
- **6** integral (exotic traces, Wodzicki residue, NC residue)

The **significance** of (differential) operators

Most of classical geometry can be encoded in terms of bounded and unbounded operators acting on a separable Hilbert space.

How do we reconstruct geometry ?

- **1** differential calculus: $da = [D, a]$ and FODC $Ω¹(M)$
- ² metric: *d*(*x*, *y*) = sup||[*D*,*f*]||≤¹ |*f*(*x*) − *f*(*y*)|
- ³ additional connection (if spinors twisted by a vector bundle)
- ⁴ dimension (growth of eigeinvalues: $N_D(Λ) \sim Λ^d$),
- ⁵ integral (exotic traces, Wodzicki residue, NC residue)
- ⁶ scalar of curvature, metric tensor, Einstein tensor via spectral computations

Spectral Metric and Einstein Functionals, L.Dabrowski, A.Sitarz, P.Zalecki, Advances in Mathematics, Volume 427, 2[02](#page-13-0)[3,](#page-15-0) [1](#page-6-0)[0](#page-7-0)[9](#page-14-0)[1](#page-15-0)[28](#page-0-0)

Connes' spectral triple approach

Algebra \mathcal{A} , its faithful representation π on a Hilbert space \mathcal{H} , a selfadjoint, unbounded operator *D*, satisfying several conditions:

¹ ∀*a* ∈ A [*D*, π(*a*)] ∈ *B*(H), *D* [−]¹ has compact resolvant

Connes' spectral triple approach

Algebra \mathcal{A} , its faithful representation π on a Hilbert space \mathcal{H} , a selfadjoint, unbounded operator *D*, satisfying several conditions:

- ¹ ∀*a* ∈ A [*D*, π(*a*)] ∈ *B*(H), *D* [−]¹ has compact resolvant
- \bullet even ST: $\exists \gamma \in \mathscr{A}' : \gamma^2 = 1, \gamma = \gamma^\dagger, \gamma D + D \gamma = 0,$

Connes' spectral triple approach

Algebra \mathcal{A} , its faithful representation π on a Hilbert space \mathcal{H} , a selfadjoint, unbounded operator *D*, satisfying several conditions:

- ¹ ∀*a* ∈ A [*D*, π(*a*)] ∈ *B*(H), *D* [−]¹ has compact resolvant
- \bullet even ST: $\exists \gamma \in \mathscr{A}' : \gamma^2 = 1, \gamma = \gamma^\dagger, \gamma D + D \gamma = 0,$
- **3** more structure like reality condition, Hochschild cycle etc.

Connes' spectral triple approach

Algebra \mathcal{A} , its faithful representation π on a Hilbert space \mathcal{H} , a selfadjoint, unbounded operator *D*, satisfying several conditions:

- ¹ ∀*a* ∈ A [*D*, π(*a*)] ∈ *B*(H), *D* [−]¹ has compact resolvant
- \bullet even ST: $\exists \gamma \in \mathscr{A}' : \gamma^2 = 1, \gamma = \gamma^\dagger, \gamma D + D \gamma = 0,$
- ³ more structure like reality condition, Hochschild cycle etc.
- ⁴ assure that *D* is a first order differential operator

Connes' spectral triple approach

Algebra $\mathcal A$, its faithful representation π on a Hilbert space $\mathcal X$, a selfadjoint, unbounded operator *D*, satisfying several conditions:

- ¹ ∀*a* ∈ A [*D*, π(*a*)] ∈ *B*(H), *D* [−]¹ has compact resolvant
- \bullet even ST: $\exists \gamma \in \mathscr{A}' : \gamma^2 = 1, \gamma = \gamma^\dagger, \gamma D + D \gamma = 0,$
- **3** more structure like reality condition, Hochschild cycle etc.
- ⁴ assure that *D* is a first order differential operator
- ⁵ ...+ a lot of conditions assuring smoothness

Theorem [Connes]

 $\mathscr{A}=C^\infty(M),\, M$ spin Riemannian compact manifold, $\mathscr{H}=L^2(S),$ (sections of spinor bundle) and *D* the Dirac operator on *M* then $(\mathcal{A}, \mathcal{H}, D)$ is a spectral triple (with a real structure).

Accomplishments and failures

The Standard Model

A finite spectral triple describes the Standard Model based on an algebra $\mathbb{C} \oplus \mathbb{H} \oplus M_3(\mathbb{C})$. A **nonproduct** geometry removes fermion doubling and explains CP-violation in geometrical terms

Accomplishments and failures

The Standard Model

A finite spectral triple describes the Standard Model based on an algebra $\mathbb{C} \oplus \mathbb{H} \oplus M_3(\mathbb{C})$. A **nonproduct** geometry removes fermion doubling and explains CP-violation in geometrical terms

Accomplishments and failures

The Standard Model

A finite spectral triple describes the Standard Model based on an algebra C ⊕ H ⊕ *M*3(C). A **nonproduct** geometry removes fermion doubling and explains CP-violation in geometrical terms A. Bochniak, A. Sitarz, A spectral geometry for the Standard Model without the fermion doubling Phys. Rev. D 101, 075038

Quantum spaces

The spectral triple approach allows to study geometry of quantum spheres (including Podles sphere and *SUq*(2).)

But..

The spectral approach is **limited** to Euclidean metrics. A version of a spectral triple formalism for indefinite metric is possible (using the Krein space formalism) - but there is no generic spectral approach (noncommutative residue).

Instead: a model of $N + 1$ NC geometries? The **classical** simplest $N + 1$ geometry.

$$
ds^2 = \eta dt^2 + \gamma_{ij} dx^i dx^j,
$$

where η determines the signature of the metric, $\eta = \pm 1$. The extrinsic curvature reduces to:

$$
\mathcal{K}_{ij}=-\frac{1}{2}\frac{d\gamma_{ij}(t)}{dt},\qquad \frac{d\gamma^{ij}(t)}{dt}=2\gamma^{im}\gamma^{jk}\mathcal{K}_{mk}
$$

Ricci tensor

$$
R_{ik}=r_{ik}+\frac{1}{\eta}\dot{K}_{ik}+\frac{2}{\eta}\gamma^{pm}K_{mk}K_{ip}-\frac{1}{\eta}KK_{ik},
$$

and

$$
R_{00}=\partial_t(\gamma^{mi}K_{mi})-\gamma^{pm}\gamma^{rt}K_{mr}K_{pt}.
$$

A model of $N + 1$ NC geometries.

The curvature of the **classical** simplest $N + 1$ geometry.

$$
R=r-\frac{1}{\eta}\big(K^2+K_{ij}K^{ij}-2\dot{K}\big)
$$

FRLW geometries

Assuming $\gamma_{ij}(x, t) = a(t)\zeta_{ij}(x)$ we obtain:

$$
R_{ik} = r_{ik} + \frac{1}{\eta} \left(-\frac{1}{2}\ddot{a} + \frac{2-N}{4}\frac{\dot{a}^2}{a} \right) \zeta_{ik},
$$

\n
$$
R_{00} = N \left(\frac{d}{dt} \left(-\frac{1}{2}\frac{\dot{a}}{a} \right) - \frac{1}{4}\frac{\dot{a}^2}{a^2} \right),
$$

\n
$$
R = \frac{1}{a}r - \frac{1}{\eta} \left(\frac{N^2 + N}{4} \left(\frac{\dot{a}^2}{a^2} \right) + N \frac{d}{dt} \left(\frac{\dot{a}}{a} \right) \right).
$$

,

Can the curvature be recovered spectrally? The question: If we have an *N*-dimensional Euclidean space and a family of Laplace operators ∆(*t*) that are given by the metrics $\gamma_{ii}(t)$, can we compute the scalar of curvature using the spectral properties of ∆(*t*) ?

Can the curvature be recovered spectrally? The question: If we have an *N*-dimensional Euclidean space

and a family of Laplace operators ∆(*t*) that are given by the metrics $\gamma_{ii}(t)$, can we compute the scalar of curvature using the spectral properties of ∆(*t*) ?

The $3 + 1$ case.

We start with a family $\Delta_3(x, t)$ of 3-dimensional Laplace-type operators set by the metric $g_{ij}(x)$:

$$
\Delta_3(x,t)=-\frac{1}{\sqrt{g(t)}}\partial_j(\sqrt{g(t)}g^{jk}(t)\partial_k).
$$

and we extend it to a 4-dimensional operator acting on $M_3 \times \mathbb{R}$:

$$
\Delta_4(x,t)=-\partial_t^2+c(t)\partial_t+\Delta_3(t),
$$

where

$$
c(t)=-\frac{1}{\sqrt{g}}\frac{d\sqrt{g}}{dt}.
$$

Can the curvature be recovered spectrally? The question: If we have an *N*-dimensional Euclidean space

and a family of Laplace operators ∆(*t*) that are given by the metrics $\gamma_{ii}(t)$, can we compute the scalar of curvature using the spectral properties of ∆(*t*) ?

The $3 + 1$ case.

We start with a family $\Delta_3(x, t)$ of 3-dimensional Laplace-type operators set by the metric $g_{ij}(x)$:

$$
\Delta_3(x,t)=-\frac{1}{\sqrt{g(t)}}\partial_j(\sqrt{g(t)}g^{jk}(t)\partial_k).
$$

and we extend it to a 4-dimensional operator acting on $M_3 \times \mathbb{R}$:

$$
\Delta_4(x,t)=-\partial_t^2+c(t)\partial_t+\Delta_3(t),
$$

where

$$
c(t)=-\frac{1}{\sqrt{g}}\frac{d\sqrt{g}}{dt}.
$$

Curvature for $3 + 1$ geometry.

The scalar of curvature.

We compute the scalar of curvature using purely spectral methods:

> $R \sim \mathsf{Wres}(\Delta_4^{-1})$ $\begin{pmatrix} -1 \\ 4 \end{pmatrix}$

where the Wodzicki residue depends solely on the spectral properties of the operator Δ .

How to compute it?

The Wodzicki residue is easily computable using the calculus of pseudodifferential operators, if a ΨDO *T* has an expansion in homogeneous symbols:

$$
T = T_p(\xi, x) + T_{p-1}(\xi, x) + T_{p-2}(\xi, x) + \cdots
$$

then

$$
\mathsf{Wres}(\mathcal{T}) = \int_M \int_{|\xi|=1} \mathcal{T}_{-n}(\xi, x).
$$

 2990

Curvature for $3 + 1$ geometry.

The scalar of curvature.

So using the symbol of homogeneity -4 of $\Delta(t)^{-1}$ we obtain:

$$
\mathsf{Wres}(f(x,t)\Delta_4^{-1}(t)) = \int dt \int_M \sqrt{g(t)} f(x,t) \Big[r(x,t) + \\ + \Big(-\frac{\pi^2}{6} (2K_{ij}K^{ij} + K^2) - \frac{\pi^2}{2} K^2 + \pi^2 K \Big) \Big].
$$

Curvature for $3 + 1$ geometry.

The scalar of curvature.

So using the symbol of homogeneity -4 of $\Delta(t)^{-1}$ we obtain:

$$
\mathsf{Wres}(f(x,t)\Delta_4^{-1}(t)) = \int dt \int_M \sqrt{g(t)} f(x,t) \Big[r(x,t) + \\ + \Big(-\frac{\pi^2}{6} (2K_{ij}K^{ij} + K^2) - \frac{\pi^2}{2} K^2 + \pi^2 K \Big) \Big].
$$

Can it be rewritten?

Next, can we rewrite it using only Δ₃(*t*)? The 3-dimensional scalar of curvature is

$$
\mathsf{Wres}\big(f(x,t)\Delta_3^{-\frac{1}{2}}(t)\big)\sim \int dt \int_M \sqrt{g(t)}f(x,t)r(x,t)
$$

No go theorem

There exist no functional depending on

$$
\mathsf{Wres}(\dot{\Delta}_3(t)^2 \Delta_3^{-\frac{7}{2}}(t))
$$

and
$$
\mathsf{Wres}(\ddot{\Delta}_3(t) \Delta_3^{-\frac{5}{2}}(t))
$$

that recovers the 4-dimensional Eistein-Hilbert functional

No go theorem

There exist no functional depending on

$$
\mathsf{Wres}(\dot{\Delta}_3(t)^2 \Delta_3^{-\frac{7}{2}}(t))
$$

and
$$
\mathsf{Wres}(\ddot{\Delta}_3(t) \Delta_3^{-\frac{5}{2}}(t))
$$

that recovers the 4-dimensional Eistein-Hilbert functional

Covariant derivative

We introduce a covariant derivative (acting on operators):

$$
\nabla_t = \partial_t + \mathbf{c},
$$

where *c* is an operator (!) and require that the volume functional is covariantly constant:

$$
\mathsf{Wres}\bigg(f(x,t)(\nabla_t)\Delta_3(t)^{-\frac{3}{2}}\bigg)=0
$$

Solution

There exist a functional depending on

$$
\text{Wres}\bigg(\big(\nabla_t\Delta_3(t)\big)^2\Delta_3^{-\frac{7}{2}}(t)\bigg)
$$
\n
$$
\text{and} \qquad \text{Wres}\bigg(\big(\nabla_t^2\Delta_3(t)\big)\Delta_3^{-\frac{5}{2}}(t)\bigg)
$$

that recovers the 4-dimensional Eistein-Hilbert functional

Solution

There exist a functional depending on

$$
\mathsf{Wres}\bigg(\big(\nabla_t\Delta_3(t)\big)^2\Delta_3^{-\frac{7}{2}}(t)\bigg)
$$
\n
$$
\mathsf{wres}\bigg(\big(\nabla_t^2\Delta_3(t)\big)\Delta_3^{-\frac{5}{2}}(t)\bigg)
$$

that recovers the 4-dimensional Eistein-Hilbert functional

But.

- no proof that such covariant derivative always exists
- assumptions on covariantly constant volume (?)

Examples (apart from the classical ones).

 Q^{α}

- Examples (apart from the classical ones).
- The spectral triple picture (so $D(t)$, $\dot{D}(t)$, $\ddot{D}(t)$)

- Examples (apart from the classical ones).
- The spectral triple picture (so $D(t)$, $\dot{D}(t)$, $\ddot{D}(t)$)
- Torsion: see Spectral Torsion **Spectral Torsion L.Dabrowski, A.Sitarz, P. Zalecki Commun. Math. Phys. 405, 130 (2024).**

- Examples (apart from the classical ones).
- The spectral triple picture (so $D(t)$, $\dot{D}(t)$, $\ddot{D}(t)$)
- Torsion: see Spectral Torsion **Spectral Torsion L.Dabrowski, A.Sitarz, P. Zalecki Commun. Math. Phys. 405, 130 (2024).**
- Almost noncommutative geometries ?

- Examples (apart from the classical ones).
- The spectral triple picture (so $D(t)$, $\dot{D}(t)$, $\ddot{D}(t)$)
- **Torsion: see Spectral Torsion Spectral Torsion L.Dabrowski, A.Sitarz, P. Zalecki Commun. Math. Phys. 405, 130 (2024).**
- Almost noncommutative geometries ?
- Evolution of fuzzy spaces ?

THANK YOU !