



Aditya Batra

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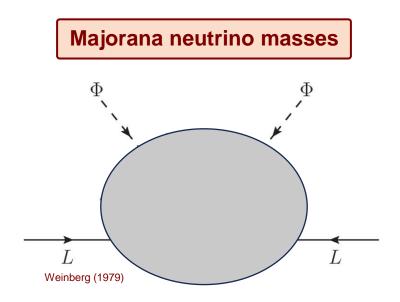




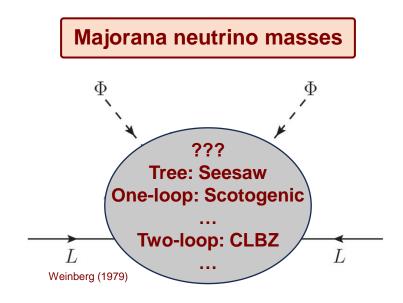


- Neutrino flavour oscillations which imply massive neutrinos and lepton mixing;
- Observed dark matter abundance;
- Strong CP problem: Lack of a theoretical explanation for the non-observation of the neutron electric dipole moment which indicates that strong interactions preserve CP symmetry.

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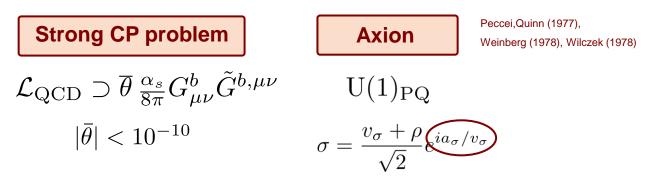
Strong CP problem

$$\mathcal{L}_{\text{QCD}} \supset \overline{\theta} \, \frac{\alpha_s}{8\pi} G^b_{\mu\nu} \tilde{G}^{b,\mu\nu}$$
$$|\bar{\theta}| < 10^{-10}$$

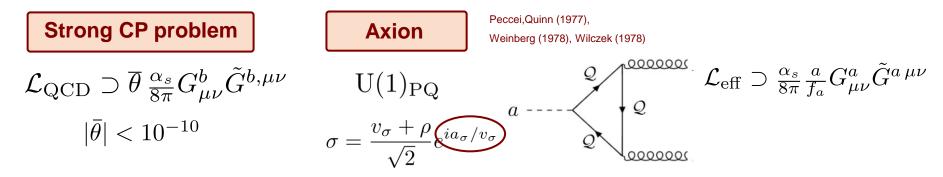
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Strong CP problem	Axion	Peccei,Quinn (1977), Weinberg (1978), Wilczek (1978)
$\mathcal{L}_{ m QCD} \supset \overline{ heta} rac{lpha_s}{8\pi} G^b_{\mu u} ilde{G}^{b,\mu u}$	${ m U}(1)_{ m PQ}$	
$ \bar{\theta} < 10^{-10}$	$\sigma = \frac{v_{\sigma} + \rho}{\sqrt{2}} e^{ia_{\sigma}}$	v/v_{σ}

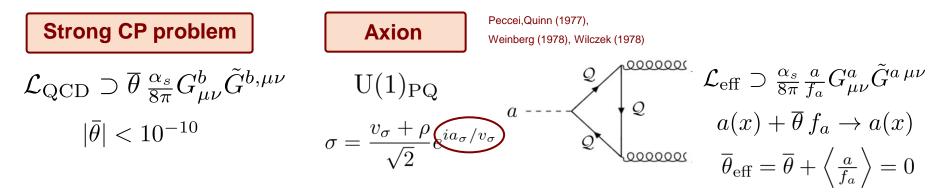
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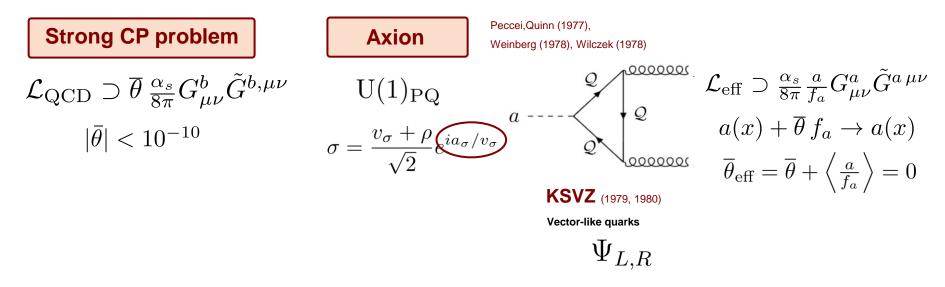
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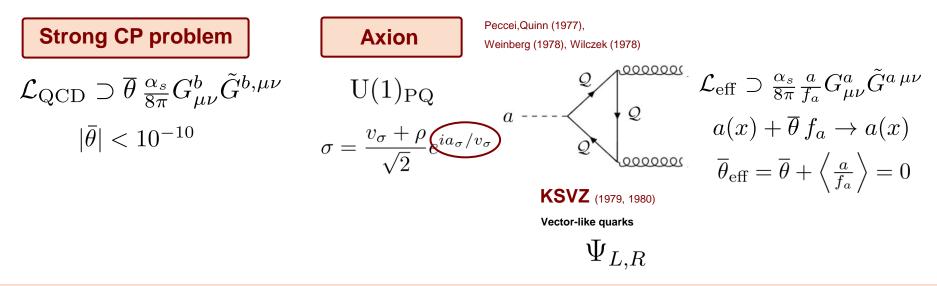


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Our approach:

New class of models where **neutrino masses** are **radiatively generated by coloured particles** which **simultaneously** solve through the PQ mechanism the **strong CP problem.** The predicted **axion** particle accounts for **dark matter**.

Fields	$SU(3)_c \otimes SU(2)_L \otimes U(1)_Y$	$U(1)_{\rm PQ}$	Multiplicity
$\overline{\Psi_L}$	$[(p,q), 2n \pm 1, 0]$	ω	n_{Ψ}
Ψ_R	$[(p,q), 2n \pm 1, 0]$	0	n_{Ψ}
σ	(1, 1, 0)	ω	1
η	[(p,q), 2n, 1/2]	0	n_{η}
χ	$[(p,q),2n\pm 1,0]$	0	n_{χ}

	Fields	$SU(3)_c \otimes SU(2)_L \otimes U(1)_Y$	$U(1)_{\rm PQ}$	Multiplicity
Vector-like quarks	$\overline{\Psi_L} \Psi_R$	$[(p,q), 2n \pm 1, 0] \ [(p,q), 2n \pm 1, 0]$	ω_0	n_{Ψ} n_{Ψ}
Complex scalar singlet		(1, 1, 0) $[(p, q), 2n, 1/2]$	ω 0	1
Coloured scalars	η χ	[(p,q), 2n, 1/2] $[(p,q), 2n \pm 1, 0]$	0	n_{η} n_{χ}

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		(1 , 1 , 0) [(p, q), 2n, 1/2]		

 $\sigma = \frac{v_{\sigma} + \rho}{\sqrt{2}} e^{ia_{\sigma}/v_{\sigma}}$

Complex

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Complex

Yukawa Lagrangian

$$-\mathcal{L}_{\mathrm{Yuk.}} \supset \mathbf{Y}_{\Psi} \overline{\Psi_L} \Psi_R \sigma + \frac{1}{2} \mathbf{Y}_{\chi_j} \Psi_R^T C \chi_j \Psi_R + \mathbf{Y}_i \bar{L} \eta_i^* \Psi_R + \mathrm{H.c.}$$

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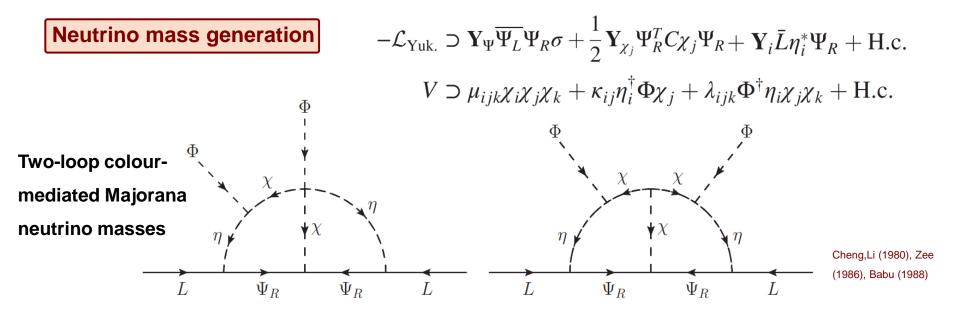
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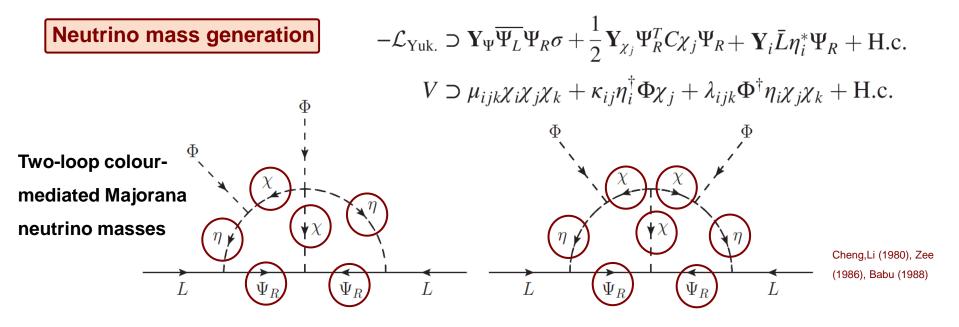
Scalar Potential

$$V \supset \mu_{ijk} \chi_i \chi_j \chi_k + \kappa_{ij} \eta_i^{\dagger} \Phi \chi_j + \lambda_{ijk} \Phi^{\dagger} \eta_i \chi_j \chi_k + \text{H.c.}$$

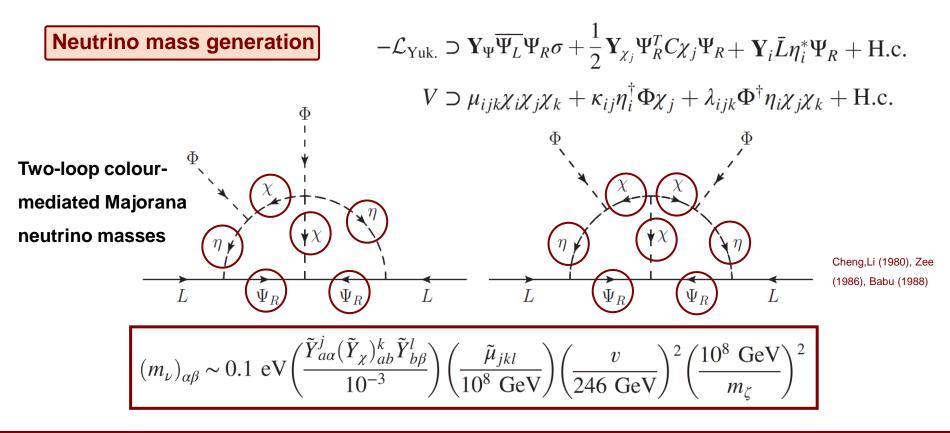
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$$V \supset \mu_{ijk} \chi_{i} \chi_{j} \chi_{k} + \kappa_{ij} \eta_{i}^{\dagger} \Phi \chi_{j} + \lambda_{ijk} \Phi^{\dagger} \eta_{i} \chi_{j} \chi_{k} + \text{H.c.}$$

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Axion decay constant

$$f_a = \frac{f_{\rm PQ}}{N} = \frac{v_\sigma}{\sqrt{2}N}$$

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QCD axion mass relation

$$m_a = 5.70(7) \left(\frac{10^{12} \text{ GeV}}{f_a}\right) \ \mu\text{eV}_{\text{Cortona et al.(2016)}}$$

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$$N = 2n_{\Psi}\omega(2n\pm 1)T(p,q)$$

Axion-to-photon coupling

$$g_{a\gamma\gamma} = \frac{\alpha_e}{2\pi f_a} \left[\frac{E}{N} - 1.92(4) \right]_{\text{Cortona et al.(2016)}}$$

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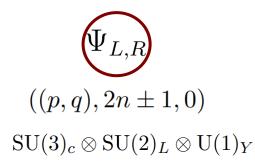
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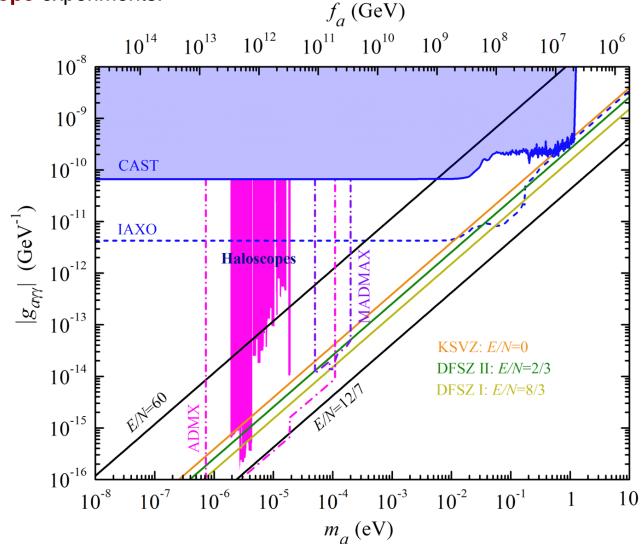
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					$SU(2)_L$		
$\Psi_{L,R}$	E/N	r	3	5	7	9	11
$((p,q),2n\pm 1,0)$		3	4	12	24	40	60
$\mathrm{SU}(3)_c \otimes \mathrm{SU}(2)_L \otimes \mathrm{U}(1)_Y$		6	8/5	24/5	48/5	16	24
	$SU(3)_c$	10	8/9	8/3	16/3	80/9	40/3
		15	1	3	6	10	15
		15/	4/7	12/7	24/7	40/7	60/7

The different models can be probed through the **axion-to-photon coupling** at **helioscope** and **haloscope** experiments.



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Coloured scalars

Vector-like quarks

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Lead to potentially dangerous stable coloured/baryonic and electrically charged relics

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Axion dark matter via the misalignment mechanism in the pre-inflationary scenario

Callan et al. (1978); Gross et al. (1981); Dimopoulos et al. (2008)

$$\Omega_a h^2 \simeq \Omega_{\rm CDM} h^2 \frac{\theta_0^2}{2.15^2} \left(\frac{f_a}{2 \times 10^{11} \text{ GeV}} \right)^{\frac{7}{6}}$$

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Isocurvature fluctuations are constrained by CMB data setting a bound on the inflationary scale

$$H_I \lesssim \frac{0.9 \times 10^7}{\Omega_a h^2 / \Omega_{\rm CDM} h^2} \left(\frac{\theta_0}{\pi} \frac{f_a}{10^{11} \text{ GeV}} \right) \text{ GeV}$$

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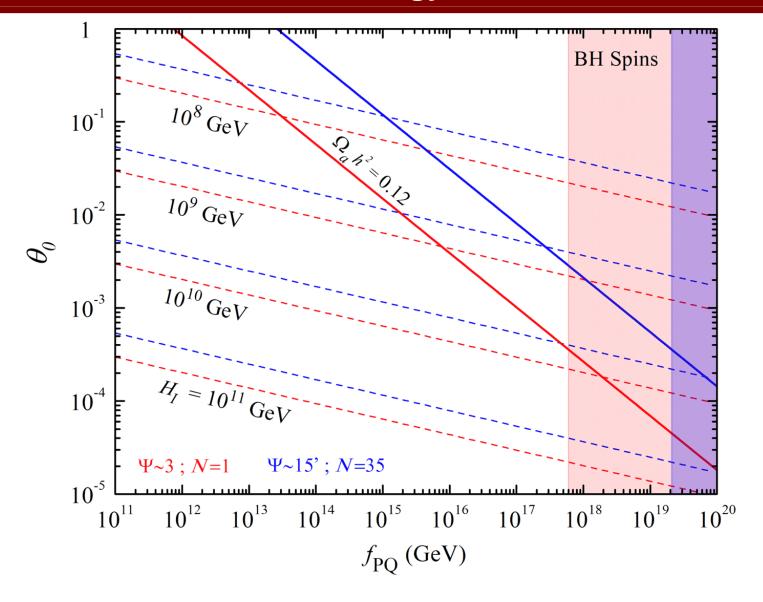
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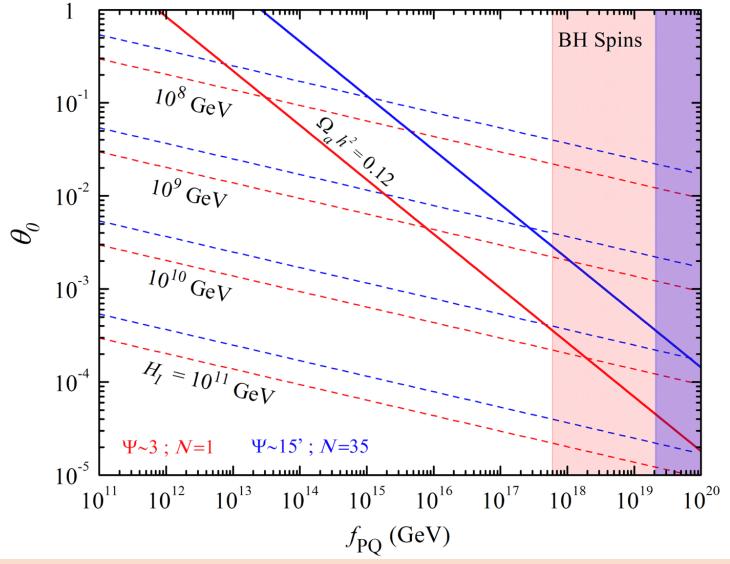
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For $\theta_0 \sim O(1)$ axions can account for the **full CDM budget**, provided $\mathbf{f_a} \sim \mathbf{5} \times \mathbf{10^{11}}$ GeV, a region currently under scrutiny at **haloscopes**. Aditya Batra – Corfu2024 – August 31, 2024

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Thank you!



Post-inflationary scenario

	Fields	${ m SU}(3)_{ m c}\otimes { m SU}(2)_{ m L}\otimes { m U}(1)_{ m Y}$	$\rm U(1)_{PQ}$	Multiplicity
Fermions	$ \Psi_L$	$(3,1,-1/3\ (2/3))$	0	n_{Ψ}
	$ \Psi_R$	$(3,1,-1/3\ (2/3))$	$-1/n_{\Psi}$	n_{Ψ}
	σ	(1, 1, 0)	$1/n_{\Psi}$	1
Scalars	$\mid \eta$	$(3, 2, 1/6 \ (7/6))$	$-1/n_{\Psi}$	n_{η}
	χ_1	(6 , 1 ,2/3~(-4/3))	$2/n_{\Psi}$	n_{χ_1}
	χ_2	$(3,1,-1/3\ (2/3))$	$-1/n_{\Psi}$	n_{χ_2}

 $-\mathcal{L}_{\text{Yuk.}} \supset \mathbf{Y}_{\Psi} \overline{\Psi_L} \Psi_R \sigma + \frac{1}{2} \mathbf{Y}_{\chi_{1j}} \Psi_R^T C \ \chi_{1j} \Psi_R + \mathbf{Y}_i \overline{L} \ \eta_i^* \Psi_R + \mathbf{m} \ \overline{\Psi_L} d_R \ (+\mathbf{m} \ \overline{\Psi_L} u_R) + \text{H.c.} ,$

 $V \supset \mu_{ijk} \chi_{1i} \chi_{2j} \chi_{2k} + \kappa_{ij} \eta_i^{\dagger} \Phi \chi_{2j} + \lambda_{ijk} \Phi^{\dagger} \eta_i \chi_{1j} \chi_{2k} + \text{H.c.}$

