

# Radiation exchange in primordial gravitational waves

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Based on arXiv:2310.19071 with Yi Wang and Misao Sasaki

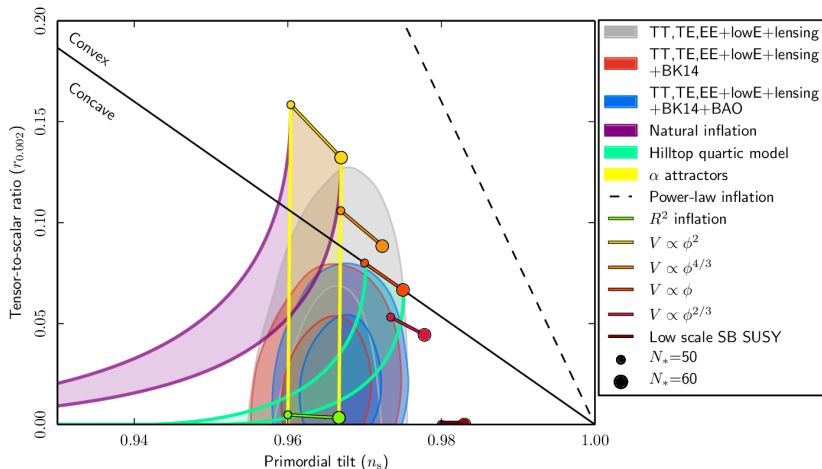
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We computed thermal corrections to primordial gravitational wave (GW) spectrum during **radiation dominant** universe. The **super-horizon** primordial GWs are enhanced.

\*Questions

- 1 Superhorizon conservation of GWs?
- 2 What is the observational implications?

# Cosmological perturbations in the sky



\*Linearized Einstein equation with hydrodynamical approximation:

$$h''_{ij} + 2\frac{a'}{a}h'_{ij} - \partial_k^2 h_{ij} = 0.$$

- 1 Massless.
- 2 Constant superhorizon modes.
- 3 Radiation drives the background dynamics only.
- 4  $h_{ij}$  is a **free field** in the FLRW background.

\*RD universe realized by massless scalar field  $\chi$  with a local thermal state  $\rho_\chi$ .

$$\begin{aligned} S_\chi &= -\frac{1}{2} \int d^4x \sqrt{-g} g^{\mu\nu} \partial_\mu \chi \partial_\nu \chi \\ &= \frac{1}{2} \int d^4x a^2 (\dot{\chi}^2 - \delta^{ij} \partial_i \chi \partial_j \chi) + \dots, \end{aligned}$$

- ① Energy momentum tensor of  $\chi$ .

$$T_{\chi,\mu\nu} = \partial_\mu \chi \partial_\nu \chi - \frac{g_{\mu\nu}}{2} g^{\rho\sigma} \partial_\rho \chi \partial_\sigma \chi$$

- ② Matching condition:

$$T_{\gamma,\mu\nu} = \text{Tr}[\hat{\rho}_\chi \hat{T}_{\chi,\mu\nu}] \rightarrow P = \frac{\rho}{3}$$

\*Individual  $\chi$  feels gravity and vice versa:

$$\begin{aligned} S_\chi &= -\frac{1}{2} \int d^4x \sqrt{-g} g^{\mu\nu} \partial_\mu \chi \partial_\nu \chi \\ &= \dots + \frac{1}{2} \int d^4x a^2 h^i_k \delta^{kj} \partial_j \chi \partial_i \chi \\ &\quad - \frac{1}{4} \int d^4x a^2 h^i_k h^k_l \delta^{lj} \partial_j \chi \partial_i \chi \dots \end{aligned}$$

\*Radiation  $\chi$  and  $h_{ij}$  are coupled.

# Inconsistency?

|                        |                              |
|------------------------|------------------------------|
| Hydrodynamical approx. | (A simple) UV theory         |
| $h_{ij}$ is free       | $h_{ij}$ couples to $\chi$ . |

Something might be missing in the standard cosmological perturbation theory based on hydrodynamical approximation.

\*Hydrodynamical approximation:

$$\text{EoM}(\hat{h}, \hat{\chi}) = 0 \rightarrow \text{Tr} \left[ \hat{\rho} \text{EoM}(\hat{h}, \hat{\chi}) \right]_{\chi} = 0.$$

\* $\hat{h}$  is regarded as external field when averaging  $\chi$ .

→ **Mean field approximation** for  $\chi$ .



# Mean field approximation?

\*Consider effective mass of  $\phi$  in

$$\mathcal{L} = -\frac{1}{2}\partial_\mu\phi\partial^\mu\phi - \frac{1}{2}\partial_\mu\chi\partial^\mu\chi - g\phi^2\chi,$$

\*Wrong answer:

$$m_{\text{eff.},\phi}^2 = 2g\langle\chi\rangle?$$

\*Correct answer: find the pole of the **1-loop propagator**.

\*Hydrodynamics v.s. PT in interaction picture.

\*Interaction Hamiltonian:

$$H_{\text{int}} = \epsilon H_{\hat{h}\hat{\chi}\hat{\chi}} + \epsilon^2 H_{\hat{h}\hat{h}\hat{\chi}\hat{\chi}} + \mathcal{O}(\epsilon^3)$$

\*Operator evolution<sub>[Weinberg 2008]</sub>:

$$h = \hat{h} + \lambda \int_{\tau_R}^{\tau} d\tau_1 [H_{\text{int},1}, \hat{h}] \\ + \lambda^2 \int_{\tau_R}^{\tau} d\tau_1 \int_{\tau_R}^{\tau_1} d\tau_2 [H_{\text{int},2}, [H_{\text{int},1}, \hat{h}]] + \dots$$

\*Inflationary adiabatic vacuum & local equilibrium density operators:

$$\varrho^{\text{tot}} \sim \varrho_{h,\text{adi.}} \otimes \varrho_{\chi,\beta}$$

\*Power spectrum:

$$\text{Tr}[\varrho^{\text{tot}} h_{\mathbf{q}} h_{\mathbf{q}'}] = P_h(q) (2\pi)^3 \delta(\mathbf{q} + \mathbf{q}')$$

\*Feynman graphs in the **in-in formalism**:

$$\Theta(\tau)[\partial\hat{\chi}, \partial\hat{\chi}] = \text{—————}$$

$$\Theta(\tau)[\hat{h}, \hat{h}] = \text{~~~~~}$$

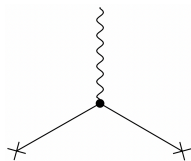
$$\partial\hat{\chi} = \text{—————}\times$$

$$\hat{h} = \text{~~~~~}\times$$

$$\int d^4x = \bullet$$

\*Single scattering with 3-pt interaction:

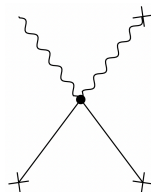
$$\delta \hat{h}_{ij}^{(1,2)}(x) = i \int^{x^0} dy^0 [\hat{H}_{h\chi\chi}(y^0), \hat{h}_{ij}(x)] =$$



\*Gravitational waves are produced by (random motion of) radiation.

\*Single scattering with 4-pt interaction:

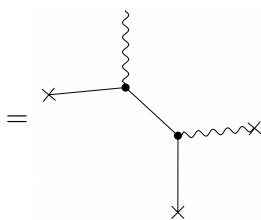
$$\delta \hat{h}_{ij}^{(1,3)}(x) = i \int^{x^0} dy^0 [\hat{H}_{hh\chi\chi}(y^0), \hat{h}_{ij}(x)] =$$



\*The effective mass.

\*Double scattering of 3-pt interaction:

$$\delta \hat{h}_{ij}^{(2,3)}(x) = i \int^{x^0} dy^0 i \int^{y^0} dz^0 \left[ \hat{H}_{h\chi\chi}(z^0), \left[ \hat{H}_{h\chi\chi}(y^0), \hat{h}_{ij}(x) \right] \right]$$

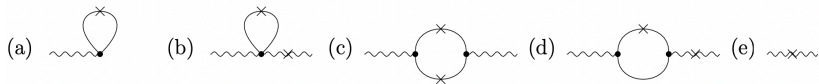


\*Not included in hydrodynamics (beyond MFA).

\*I named this **radiation exchange**.



# GW spectrum up-to 1-loop order



(a) Tadpole

(b) Effective mass

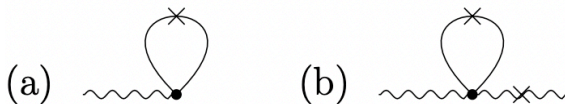
(c) Induced GWs

(d) Radiation exchange

(e) Tree level spectrum

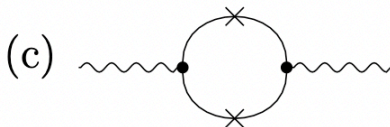
# Cancellation of effective mass

\*General covariance prohibits the mass of graviton:



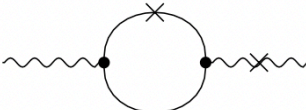
- The effective mass from (b):  $m_{\text{eff}}^2 = 2H^2$ .
- (a) is perturbed by the local thermal state (extra  $h_{ij}$  from coordinate transf.).
- The perturbed tadpole cancels the effective mass.

\*Interactions of plasma generates GWs from nothing.



- Causal production happens only in the subhorizon scale.
- The IR contribution is zero:  $\sim q^3$  (Poissonian)
- IGW does not affect the super horizon primordial GWs.

\*Initial spectrum is enhanced by exchanging radiation.

$$(d) \quad \text{Diagram} = \left( \ln \frac{\tau}{\tau_R} - 1 + \frac{\tau_R}{\tau} \right) P_h^{\text{tree}}$$


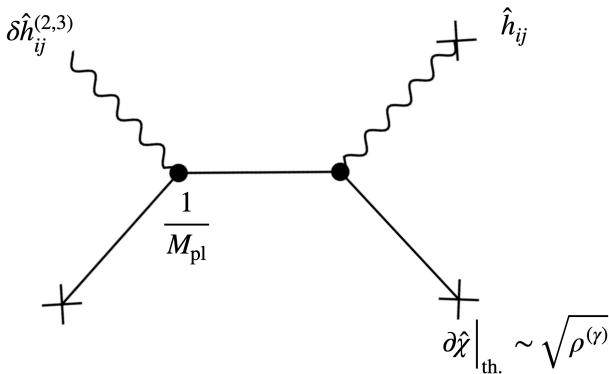
- Reheating time (initial time of RD):  $\tau_R$
- Thermal correction dominates over the tree graph:  
 $\ln \frac{\tau}{\tau_R} = \mathcal{O}(10)$
- IR (super horizon) spectrum varies.

\*Radiation exchange is weird because:

- ① Super horizon spectrum enhances. Violation of causality?
- ② Graviton interactions should be tiny. Why do we have a big effect?
- ③ Is perturbation theory valid? → **Work in progress!**

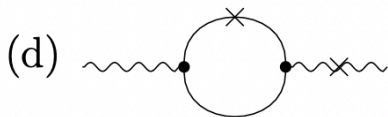
# Why big?

\*Usually, gravitons are optically thin. Why the interaction looks big?



\*BG Friedmann equation:  $3M_{\text{pl}}^2 H^2 = \rho_{\gamma}$ .

# Violation of causality?



\*Diagrams are constructed from causal propagators (based on a local Lagrangian).

- 1 Interactions between GWs and fundamental fields in a radiation fluid are considered.
- 2 These interactions have been missing in the standard framework where the mean-field approximation (MFA) is implicit.
- 3 We went beyond the MFA using the in-in formalism.
- 4 The inflationary GW power spectrum is modified at 1-loop order by the thermal effect, even at super horizon scale.
- 5 1-loop effect is comparable to the tree level since innumerous thermally excited fields contribute.



- 1 Any possibility of cancellation?
- 2 2-loop graphs?
- 3 The same effect on  $\zeta$ ?
- 4 Isocurvature?