Searching for dark matter in extended Higgs sectors

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Corfu 2022 - Workshop on Standard Model and Beyond

Emmy Noether-Programm

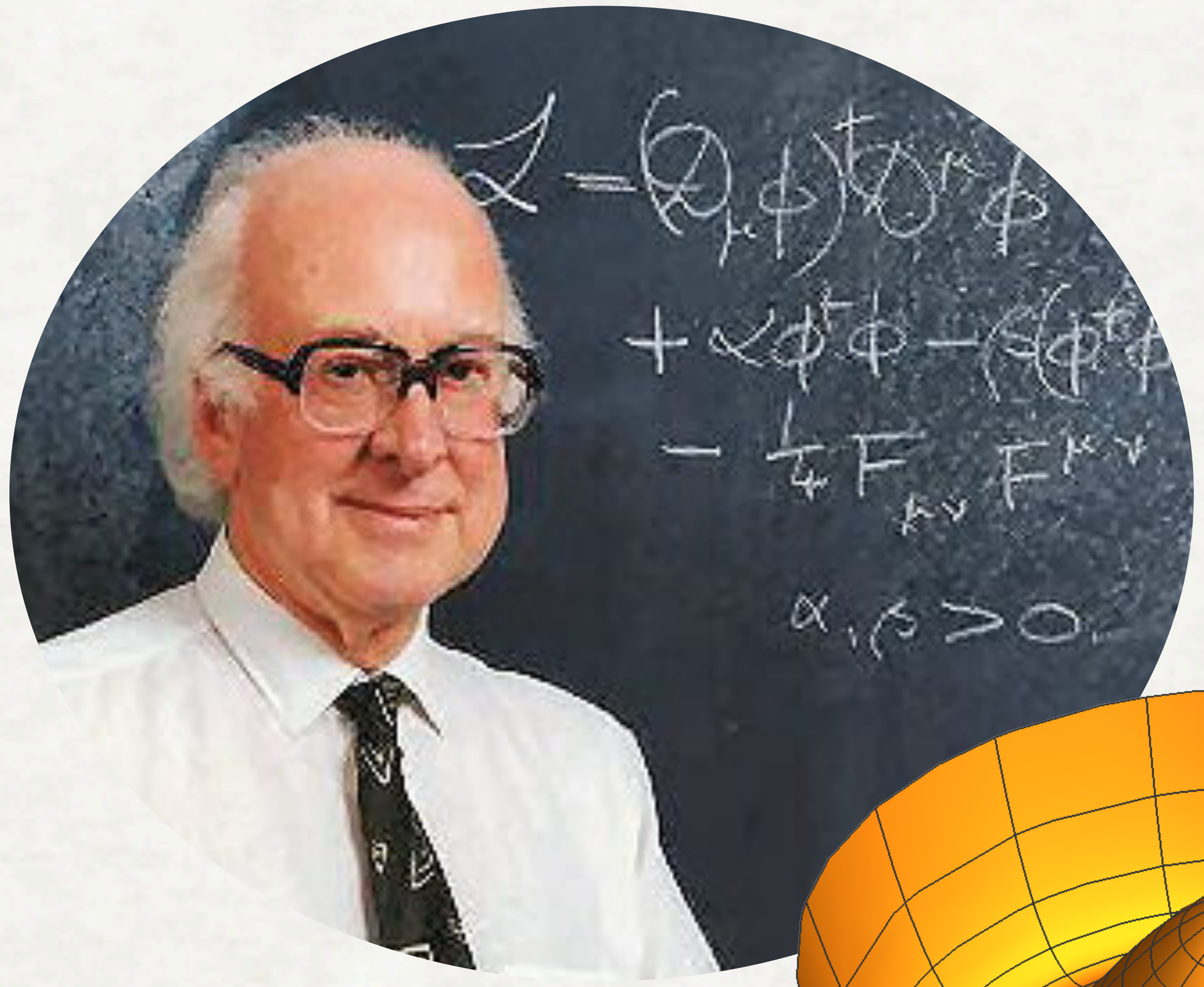
DFG Deutsche Forschungsgemeinschaft

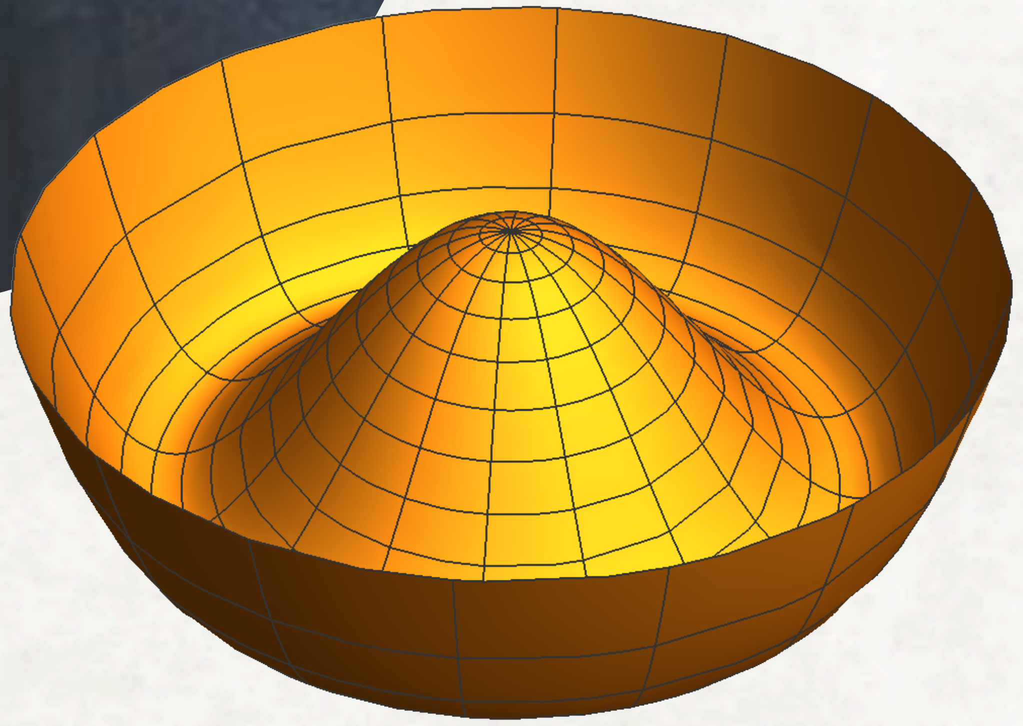
UNI FREIBURG

1964

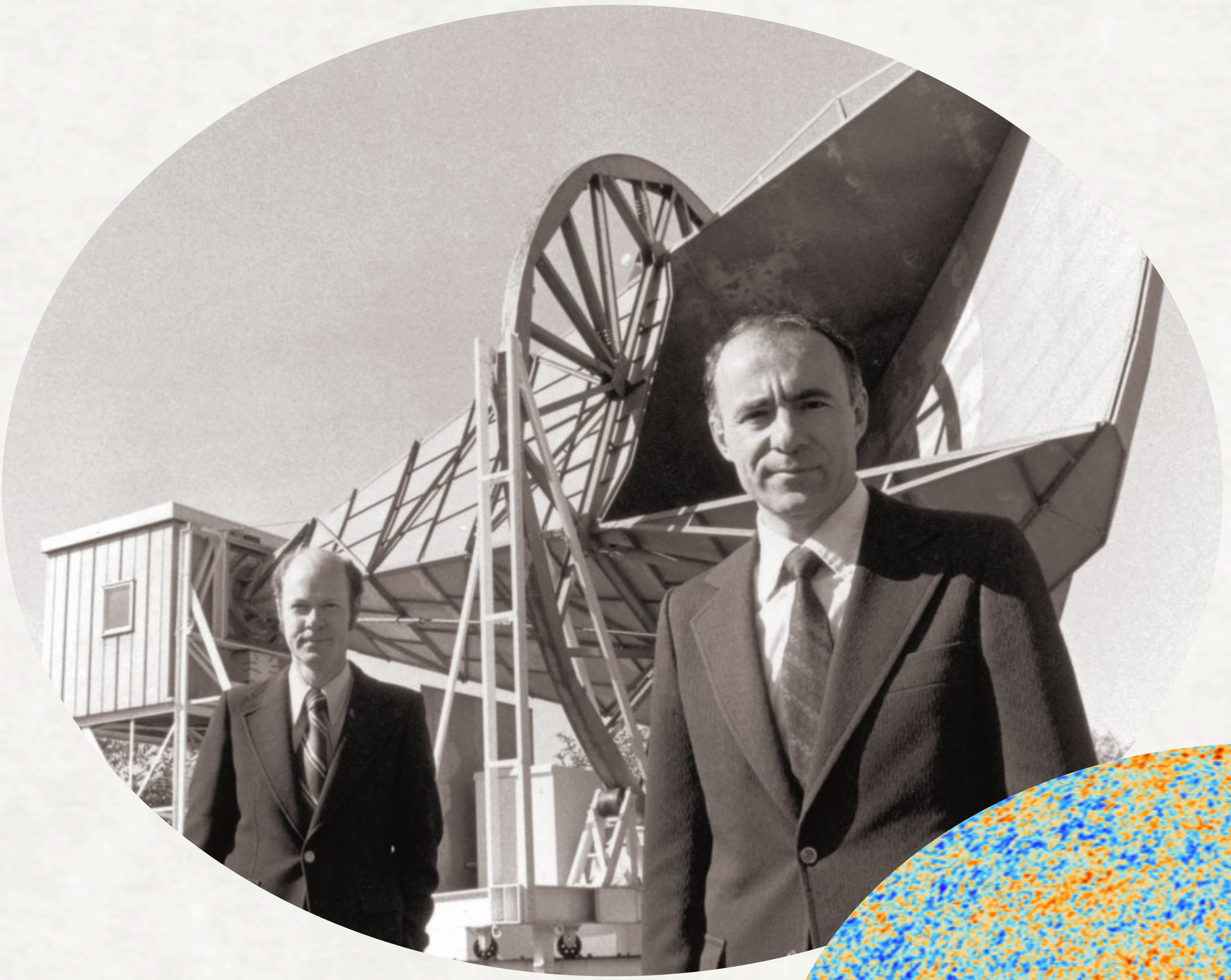
annus mirabilis for particle physics & cosmology

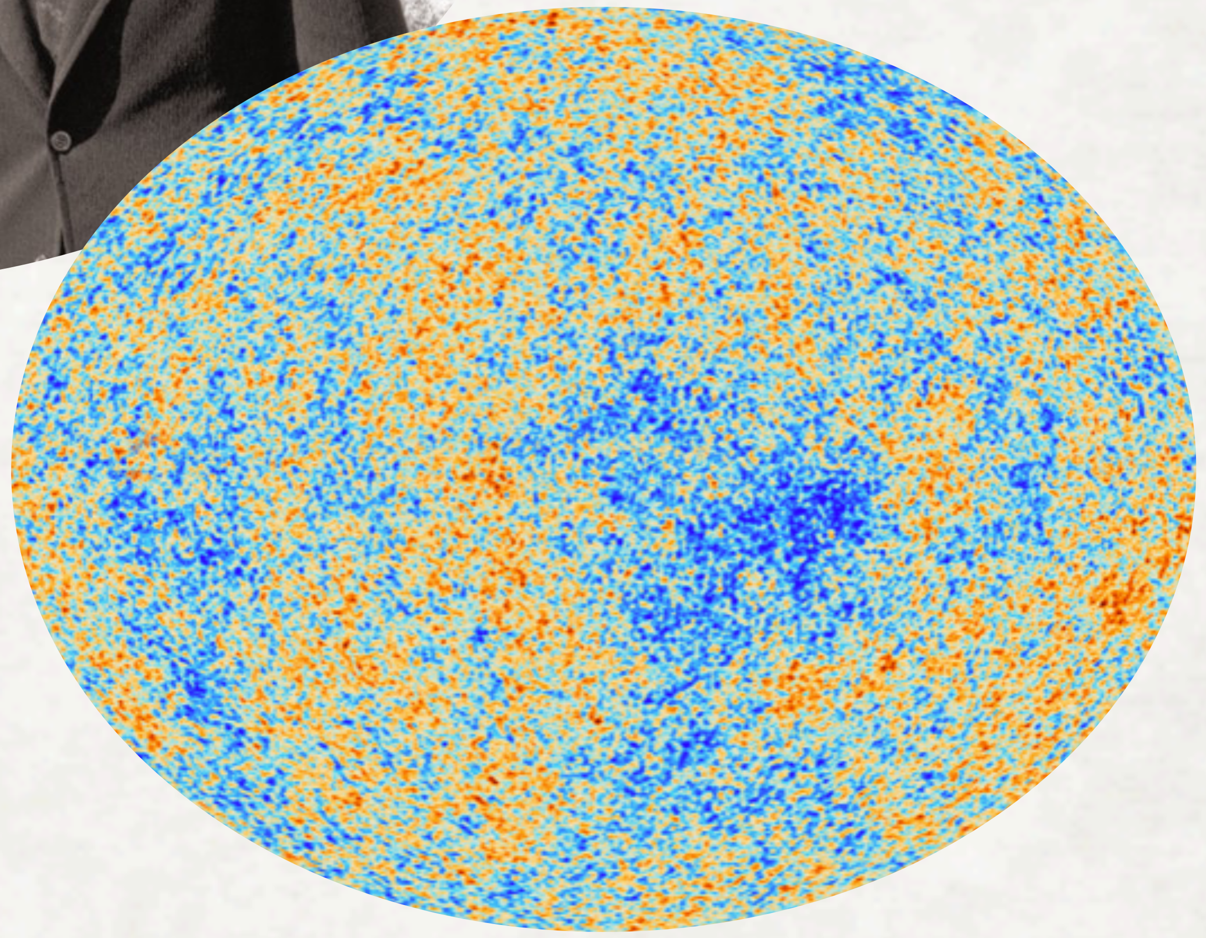
Higgs mechanism





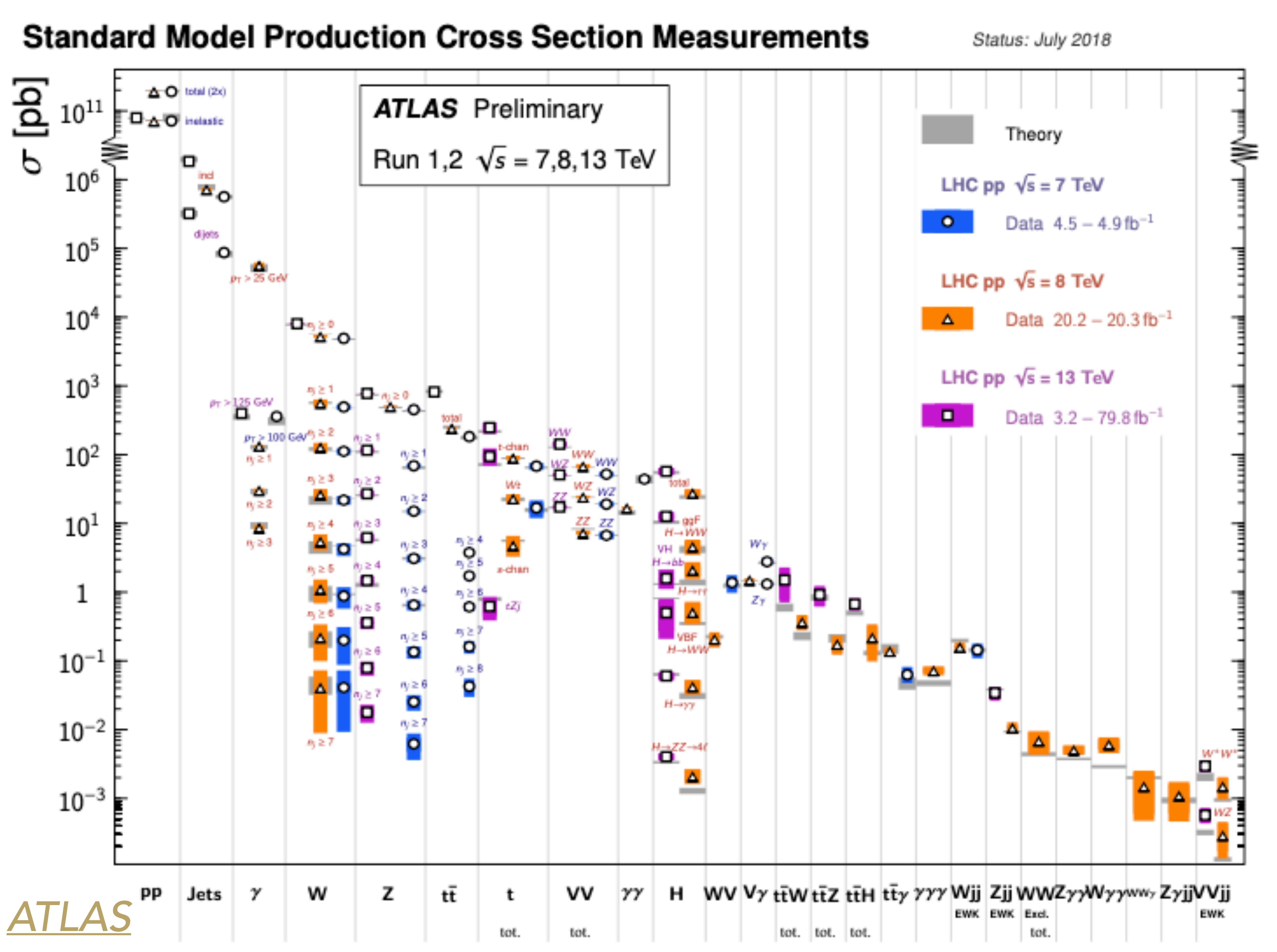
CMB radiation





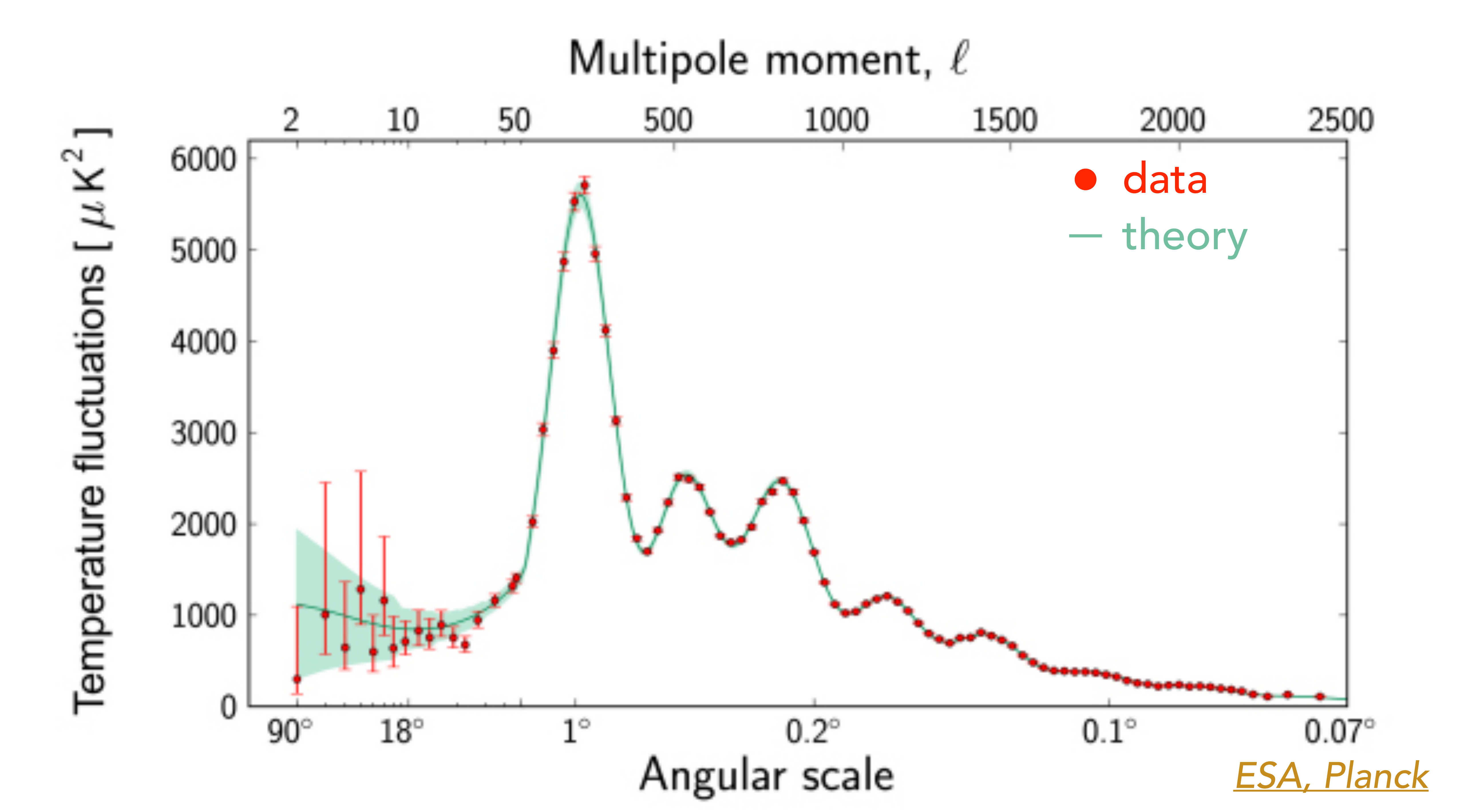
The two Standard Models

particle physics



ATLAS

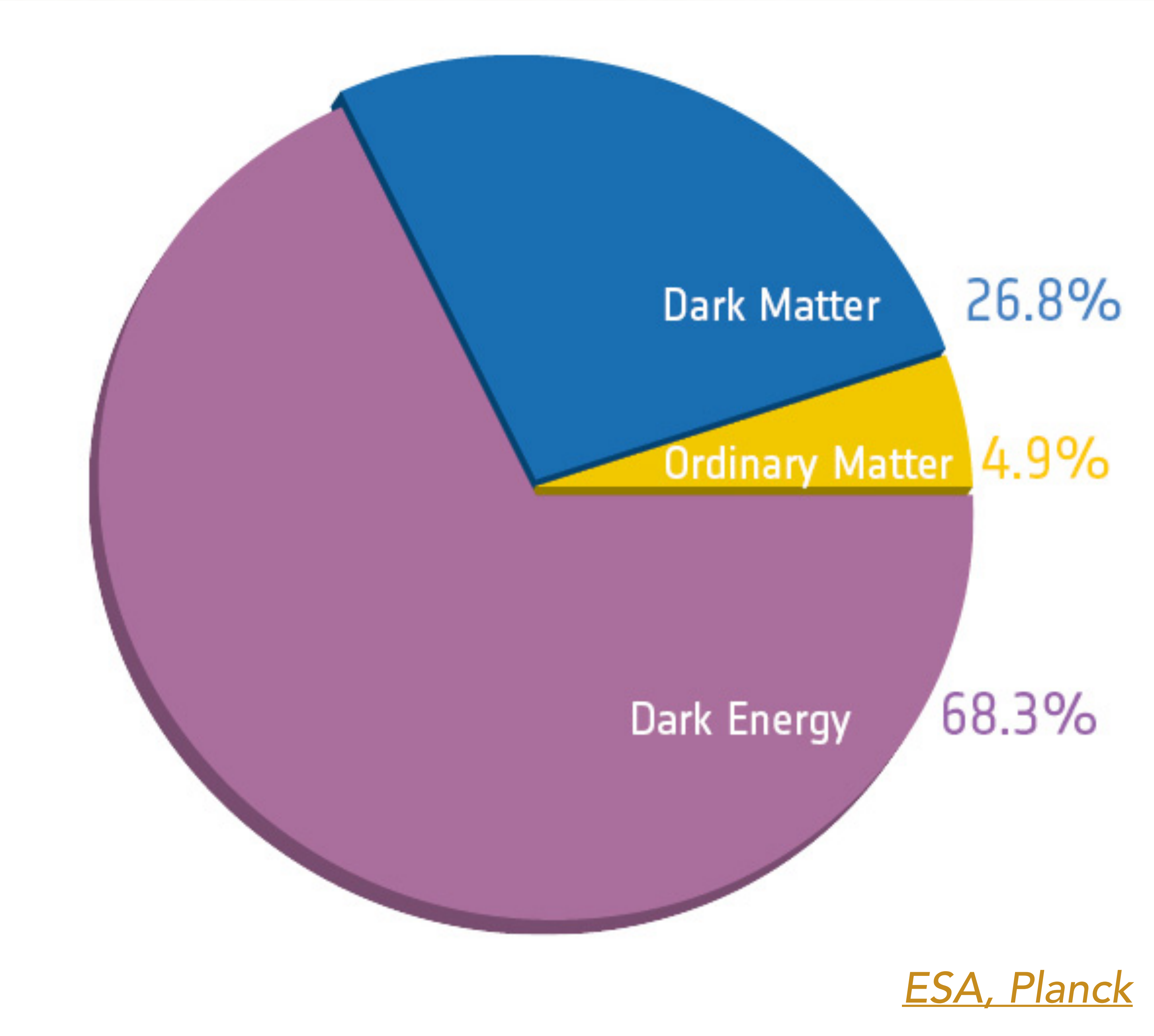
cosmology



ESA, Planck

perfectly describe (almost) everything we measure

Open questions



Dark Energy

What is the nature of DE?

Baryon Asymmetry

why matter dominates over antimatter?

Dark Matter

What is the nature of DM?

+ are these connected with the Higgs sector?

The Higgs - cosmology connection

1. **Scalar bosons ubiquitous in cosmology**

   Brax, RPP 81 (2018) 016902
   Gubitosi et al JCAP 02 (2013) 032
   Bezrukov, Shaposhnikov, PLB 659 (2008) 703
   Germani, Kehagias, PRL 105 (2010) 011302
   Burrage et al, JCAP 11 (2018) 036

2. **EW phase transition can trigger baryogenesis**

   Kuzmin, Rubakov, Shaposhnikov, PLB 155 (1985) 36
   Shaposhnikov, NPB 287 (1987) 757 (1987)
   Nelson, Kaplan, Cohen, NPB 373 (1992) 453

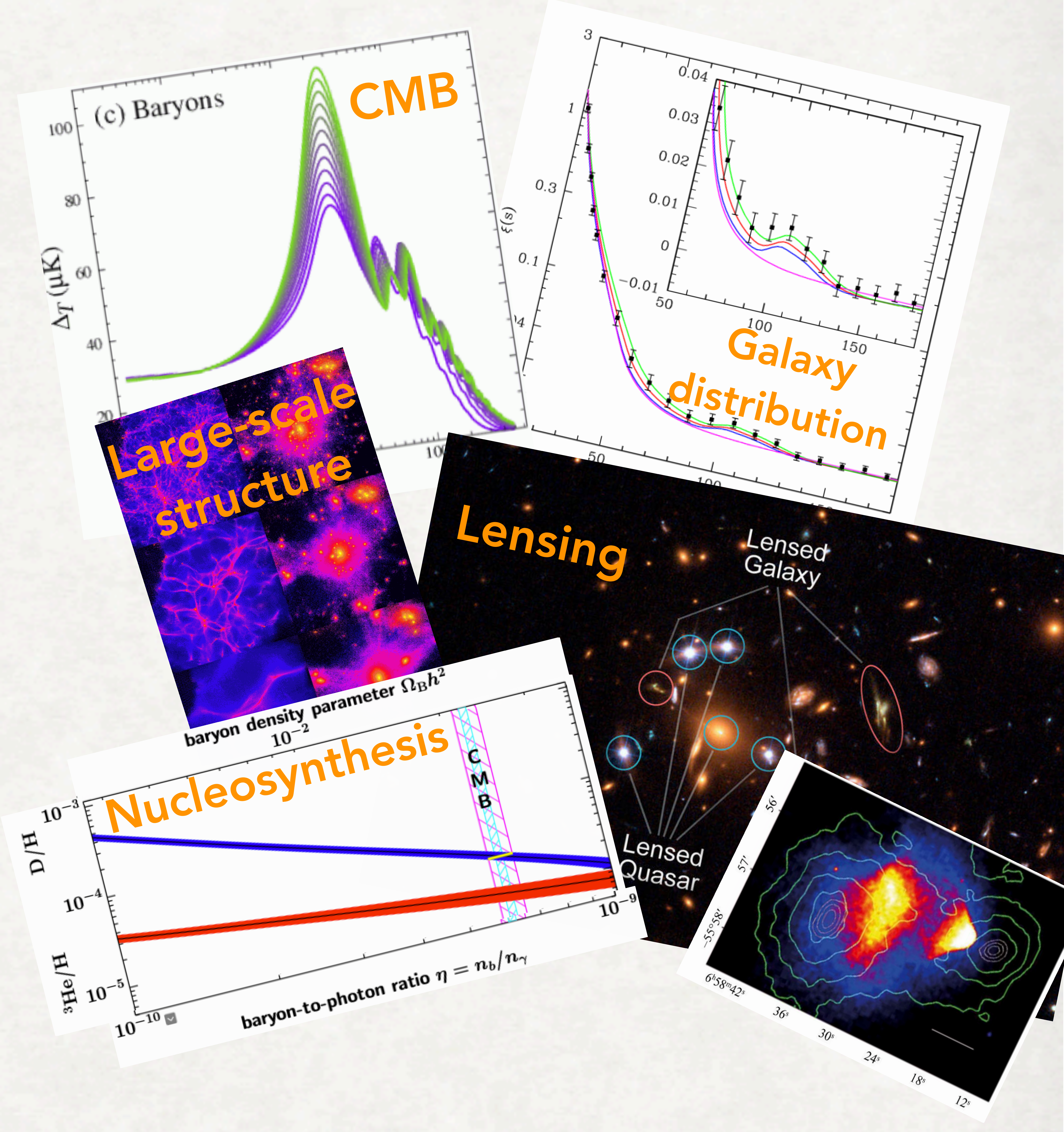
3. **Higgs sector can act as a portal to the dark sector**

   ➡ **this talk**
   based on: SA, Brandt, Haisch 2109.13597 & SA, Haisch 2202.12631

   Silveira, Zee, PLB 161 (1985) 136
   Ipek et al, Phys. Rev. D 90 (2014) 055021
   Bell et al, JCAP 03 (2017) 015
   Berlin et al, JHEP 06 (2014) 078
   Duerr et al, JHEP 09 (2016) 042

DM properties

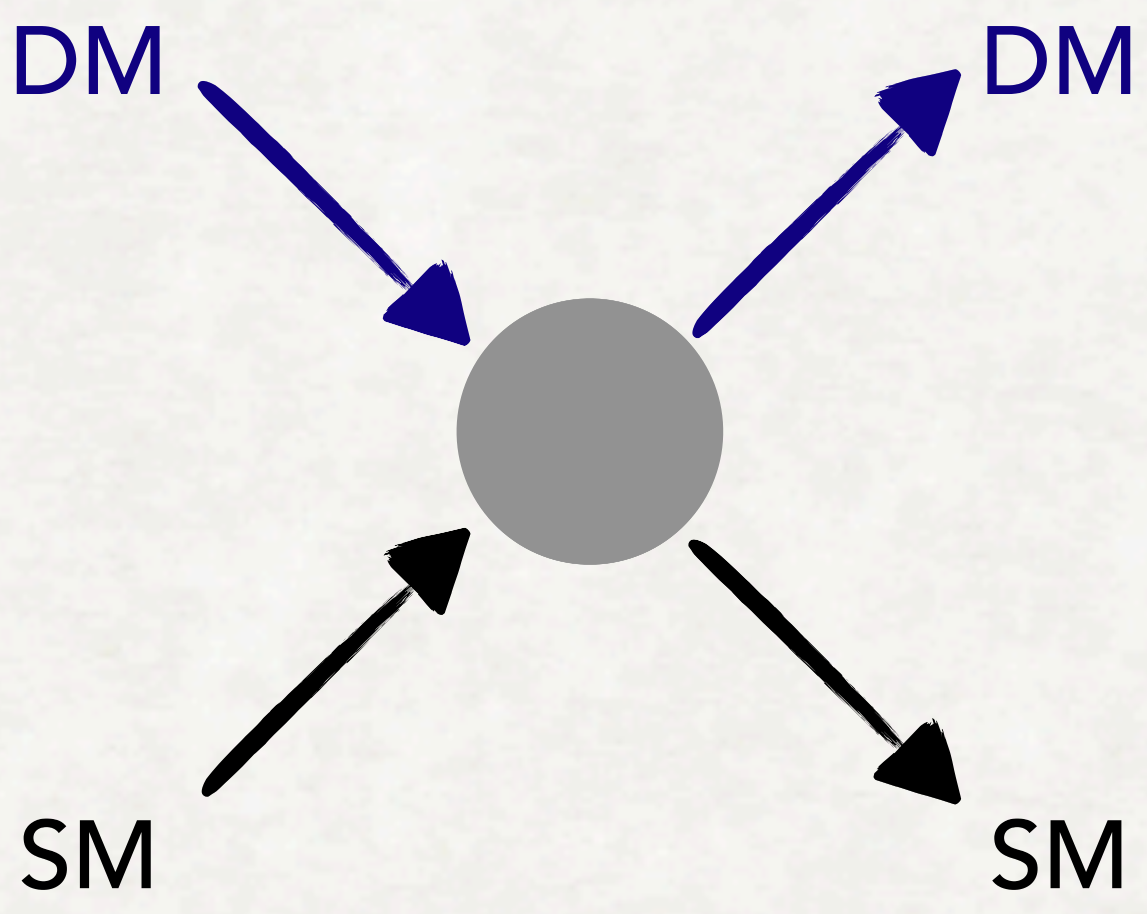
- stable
- weakly interacting
- non-relativistic ("cold")
- non-baryonic
- probably "matter" (not modified gravity)
- can't consist solely of dark astronomical objects (MACHO)



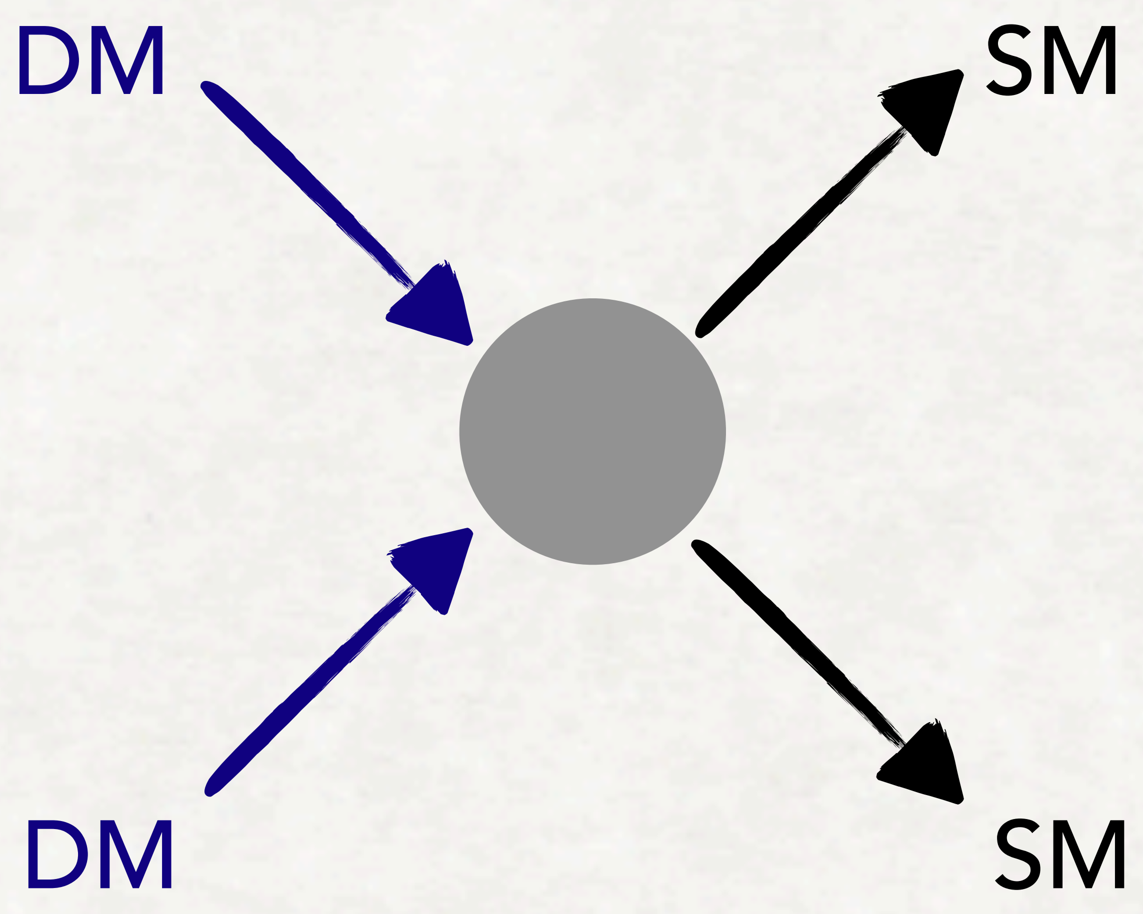
➡ Look for stable weakly interacting massive BSM particles

DM detection

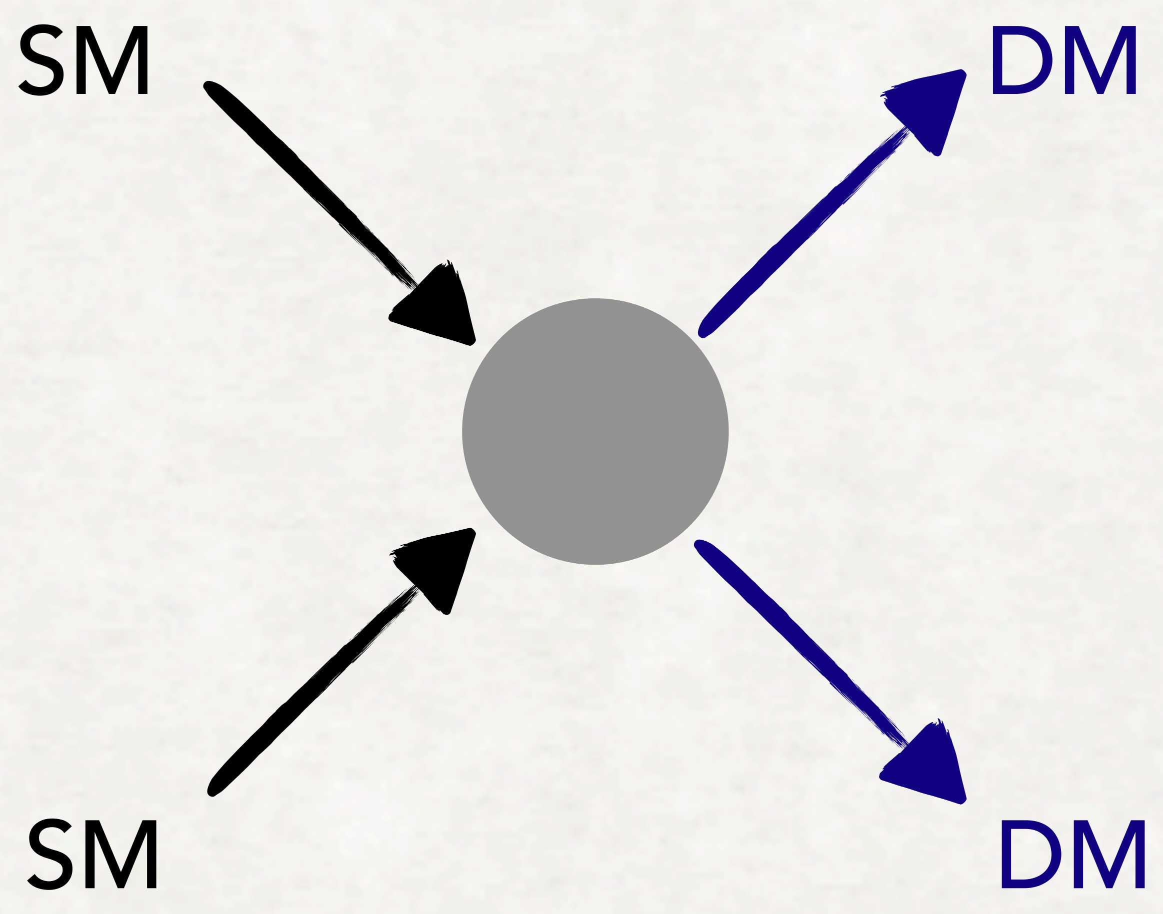
Direct Detection



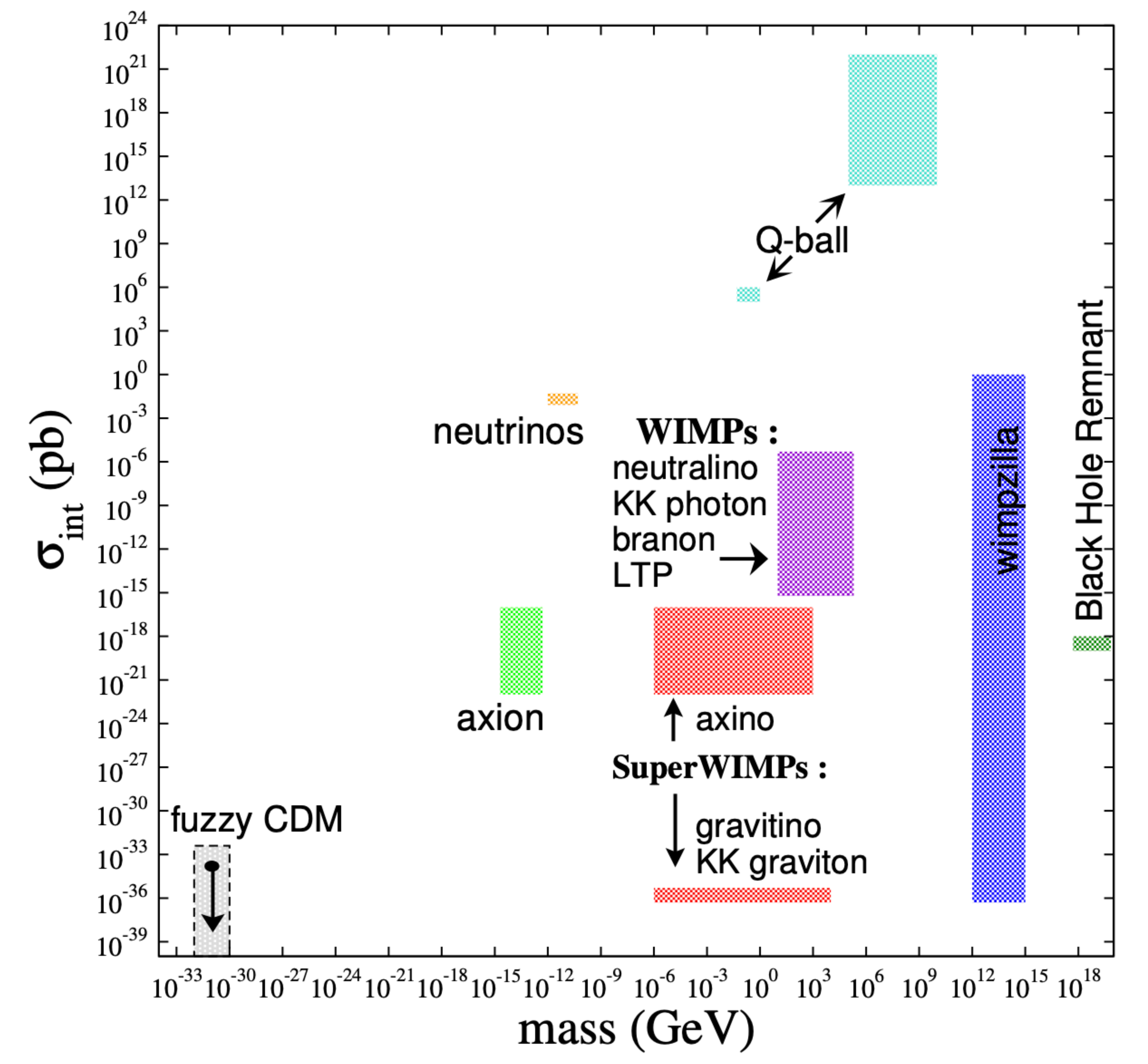
Indirect Detection



Colliders

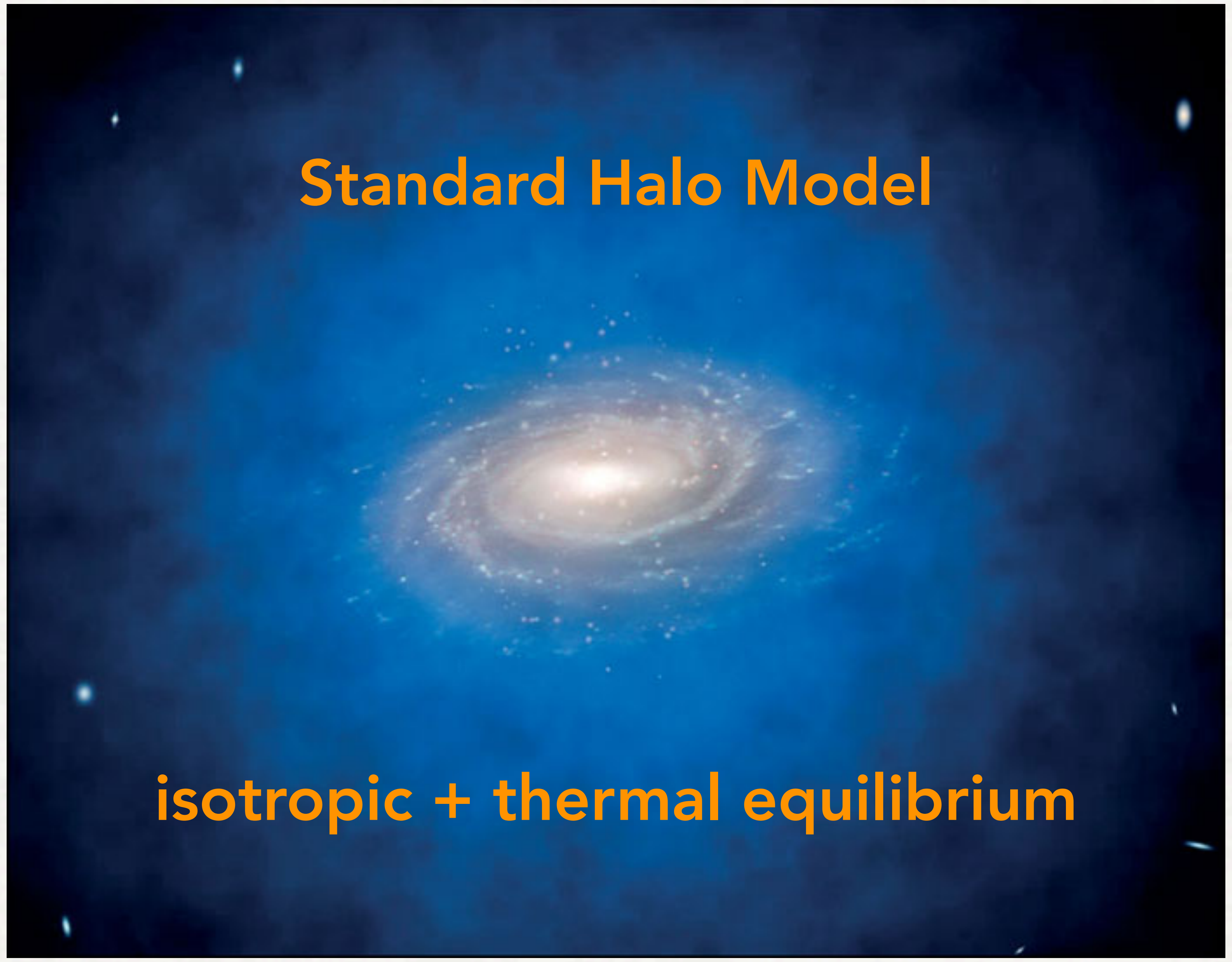


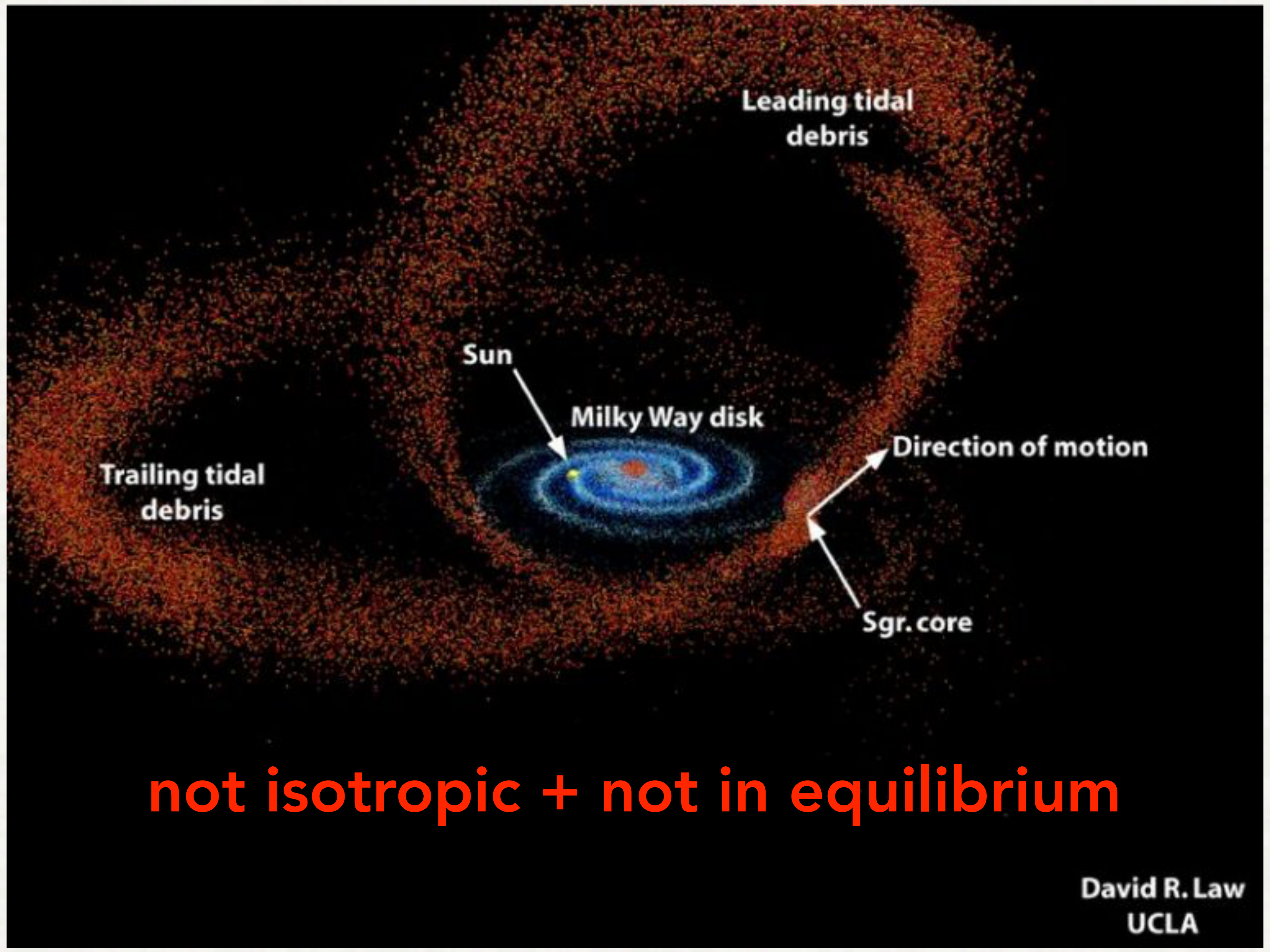
Complementarity is important



DM models predict a vast range of cross-sections/masses

Complementarity is important

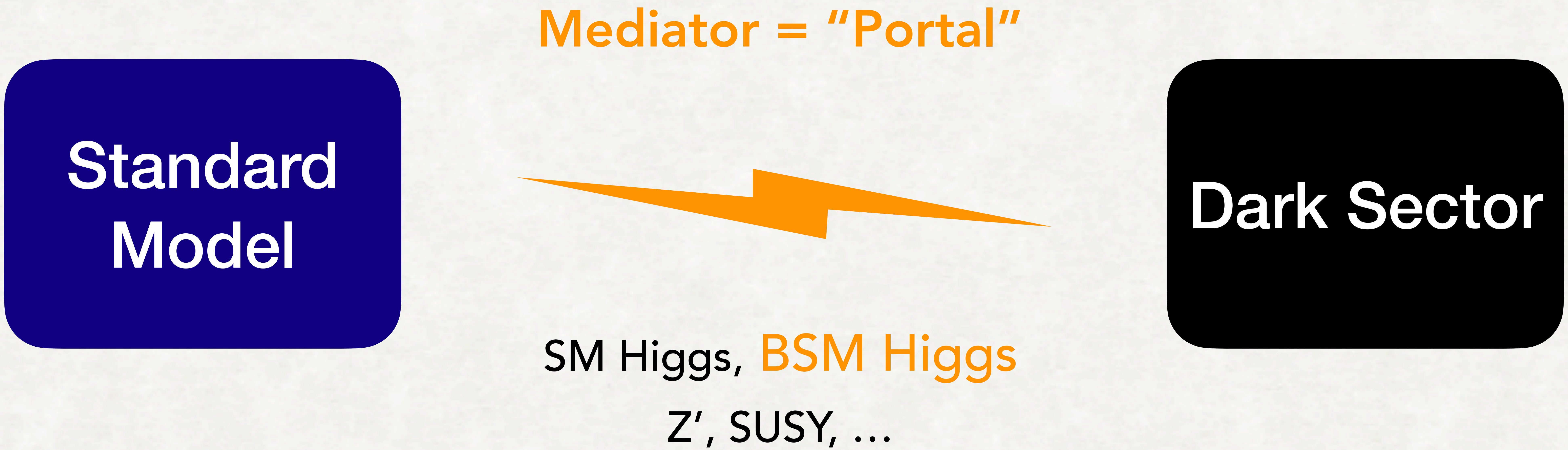




Different experiments sensitive to different assumptions

Purcell et al, JCAP 08(2012) 027 ; Bozorgnia et al, JCAP 07 (2020) 036 ; Necib et al, Nature Astron. 4 (2020) 11

Dark Matter interactions



In the following

- DM is assumed to be a SM singlet
- we concentrate on BSM Higgs portals

Classes of models

EFT

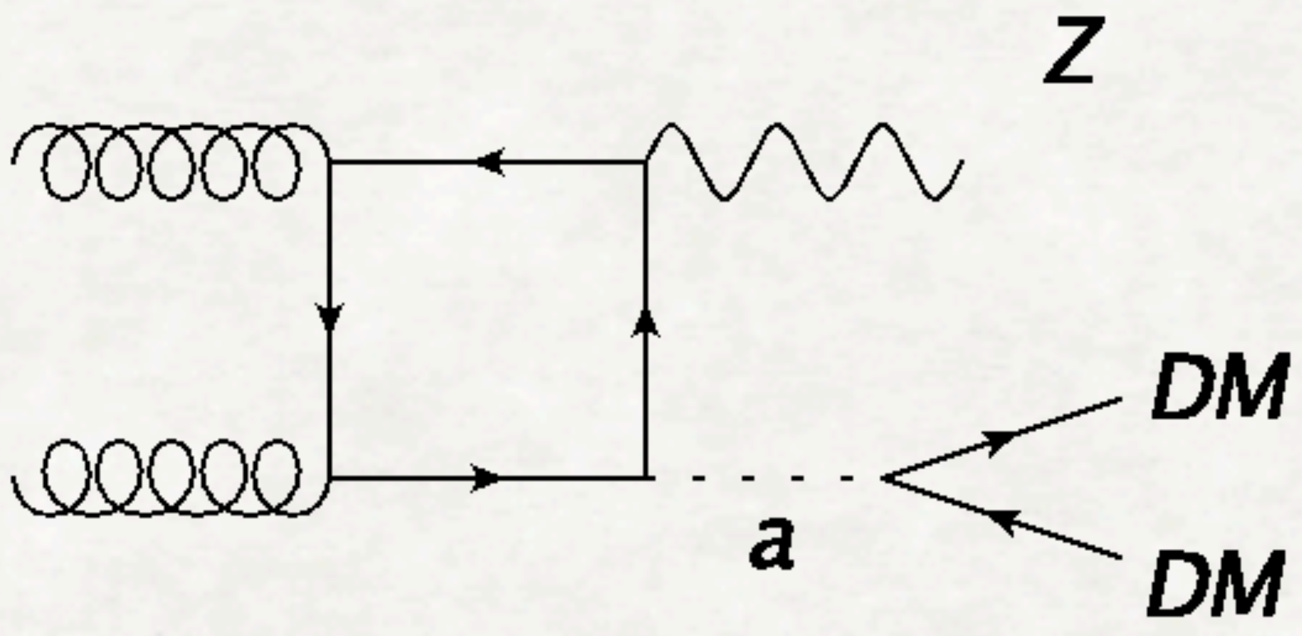
$$\mathcal{L} = \mathcal{L}_{\rm SM} + \frac{c}{\Lambda^{d-4}}\mathcal{O}^{(d)}_{\rm DM}$$

- agnostic to microscopic (UV) theory

✓ **only 1 parameter (Λ)**

◉ **breaks down at high energies**

Simplified



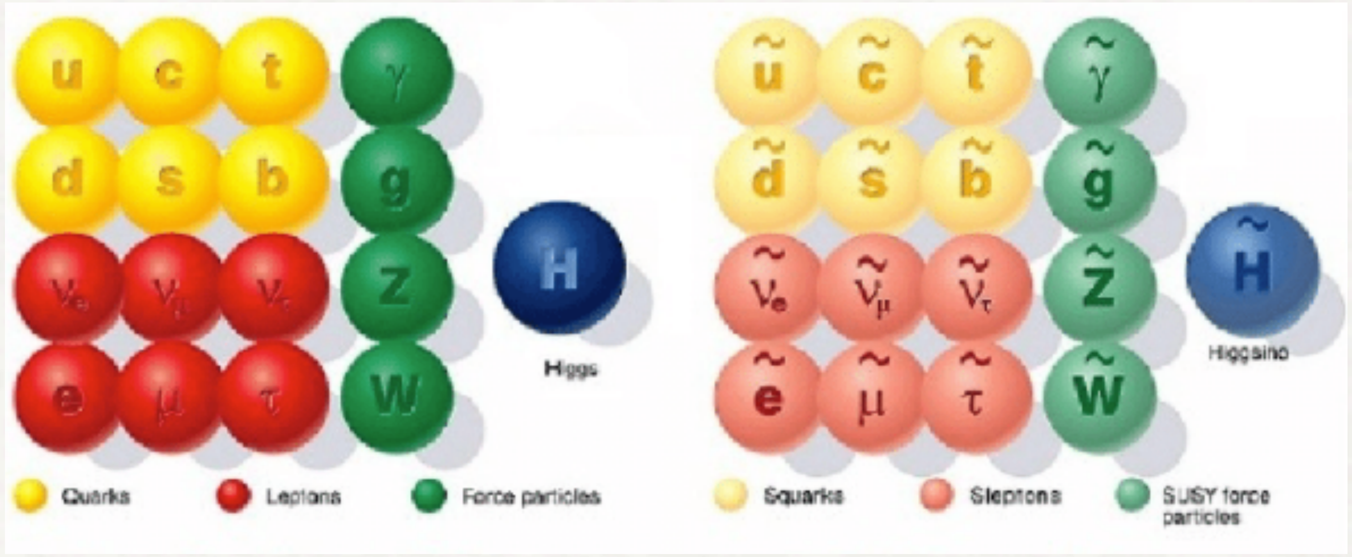
- add **only 1/few mediator** coupling DM to SM

✓ **very few parameters** (g_{SM}, g_{DM}, m_a)

✓ **easier to constrain**

➡ workhorse for DM searches

Complete



- **e.g. MSSM**
- **several new particles**

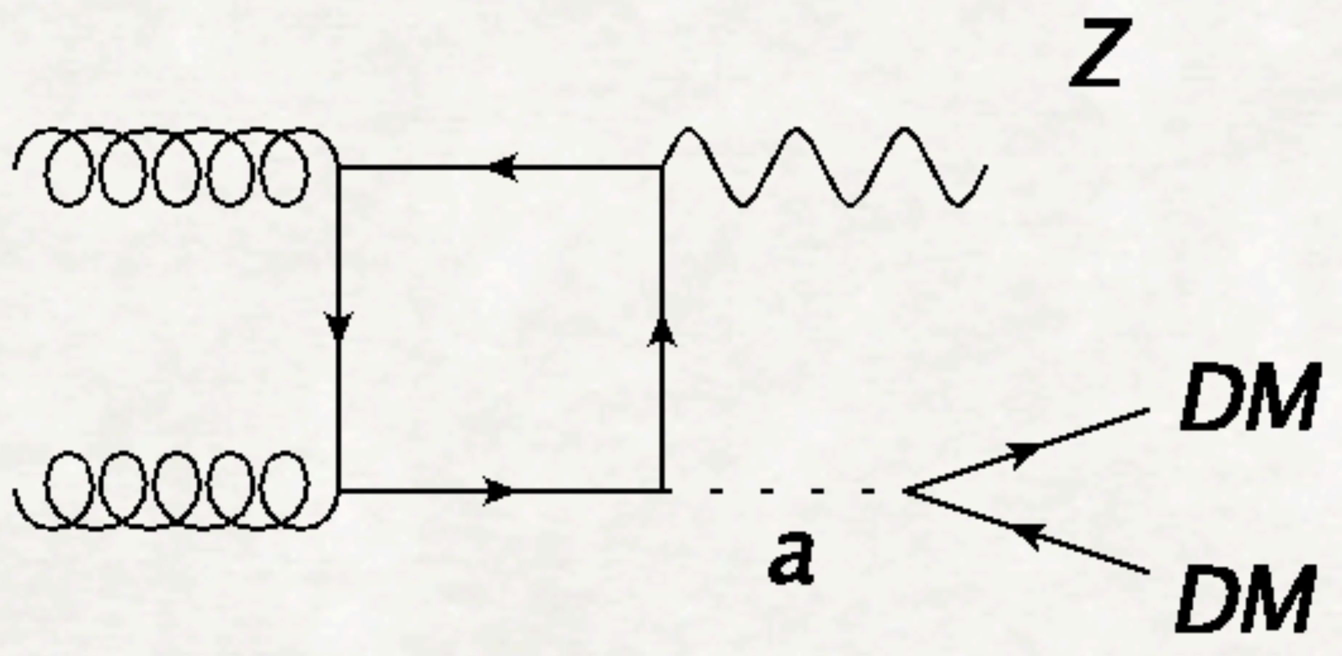
✓ **valid at high energy scales**

✓ **predictions for everything**

◉ **many parameters**

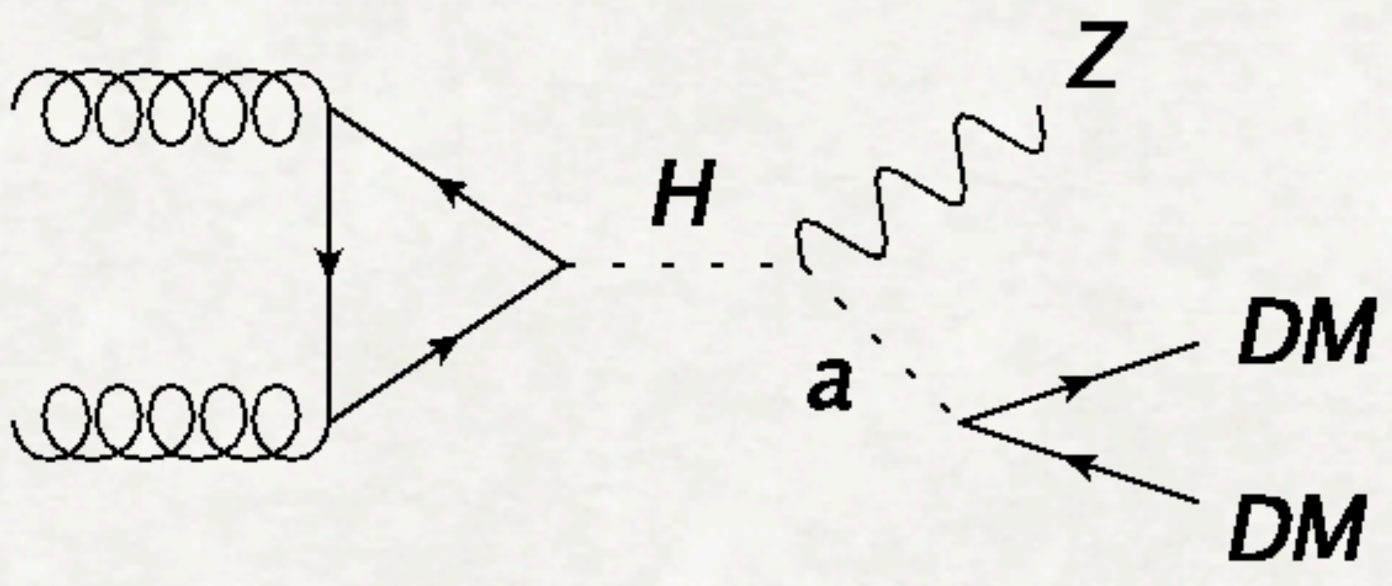
◉ **very hard to constrain**

The need for extended Higgs sectors



If a is singlet ⇒ unitarity violation

$$\mathcal{M} \sim \ln^2 s$$



Extended Higgs sector:

- fixes unitarity violation

★ bonus: resonant signatures

[Englert et al, 1604.07975](); [Haisch, Polesselo, 1812.00694](); [Bell et al, 1503.07874]()

The 2 Higgs Doublet Model

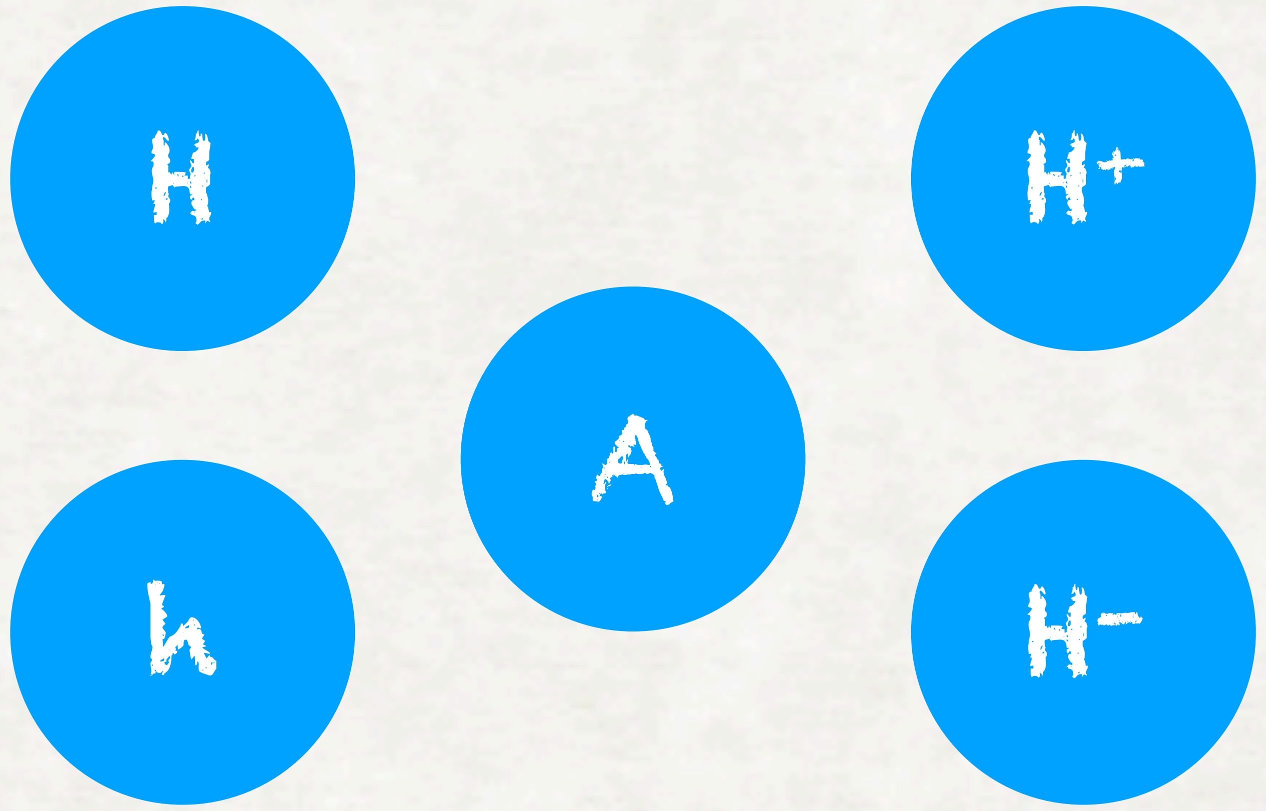
- 5 Higgs bosons

- 5 parameters considered:
  m_A, m_H, $m_{H\pm}$
  α: mixing between H, h
  tan**β**: ratio of vacuum expectation values

- **Different Yukawa structures**
  suppressed/enhanced couplings to fermions

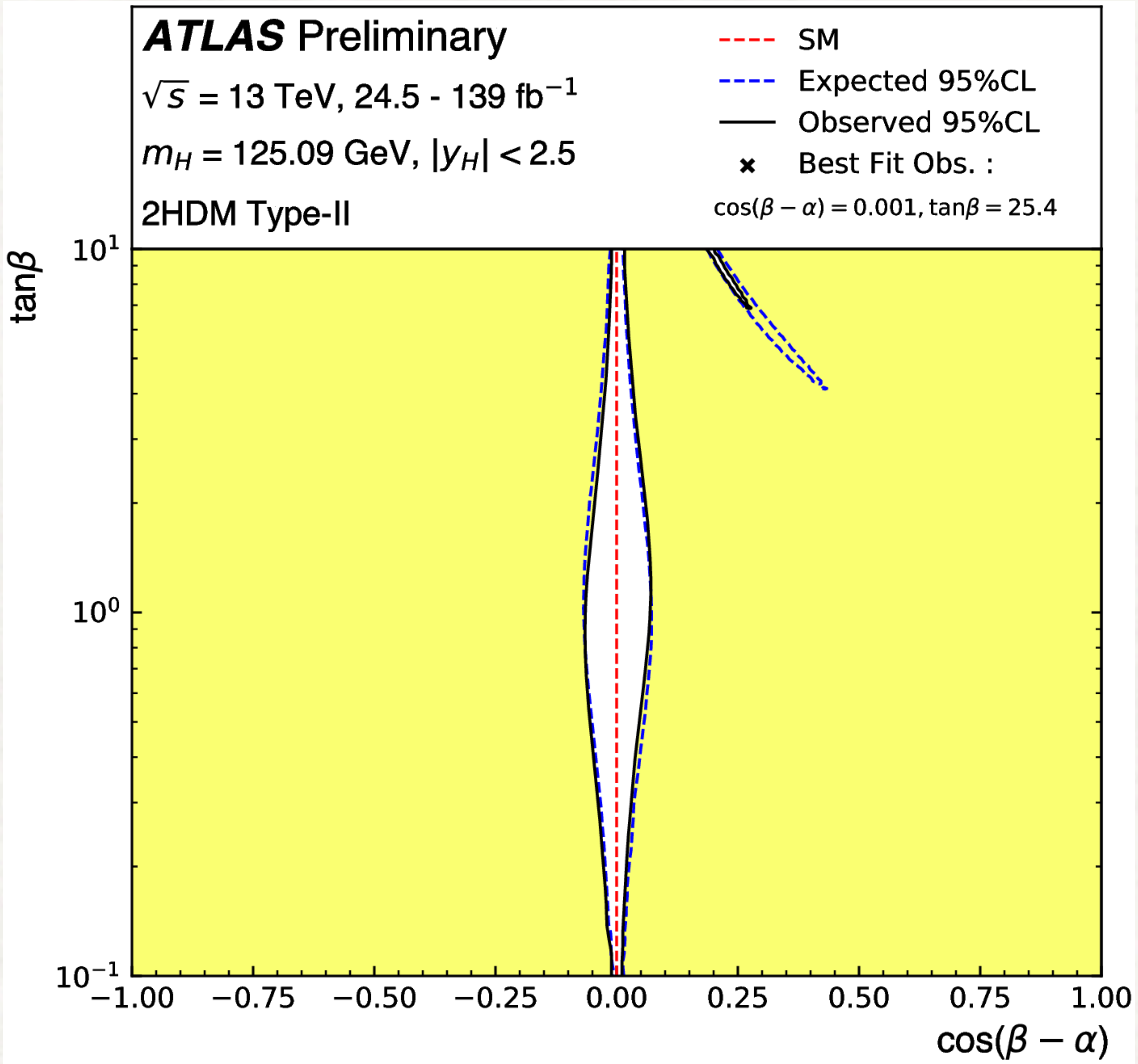
- Alignment limit: cos(β-α)=0
  h has the same couplings as the SM Higgs

- related to other models (e.g. axion, MSSM, …)



Branco et al, Phys.Rept. 516 (2012) 1
Rompotis, Ferrari, Symmetry 2021,13,2144

Constraints



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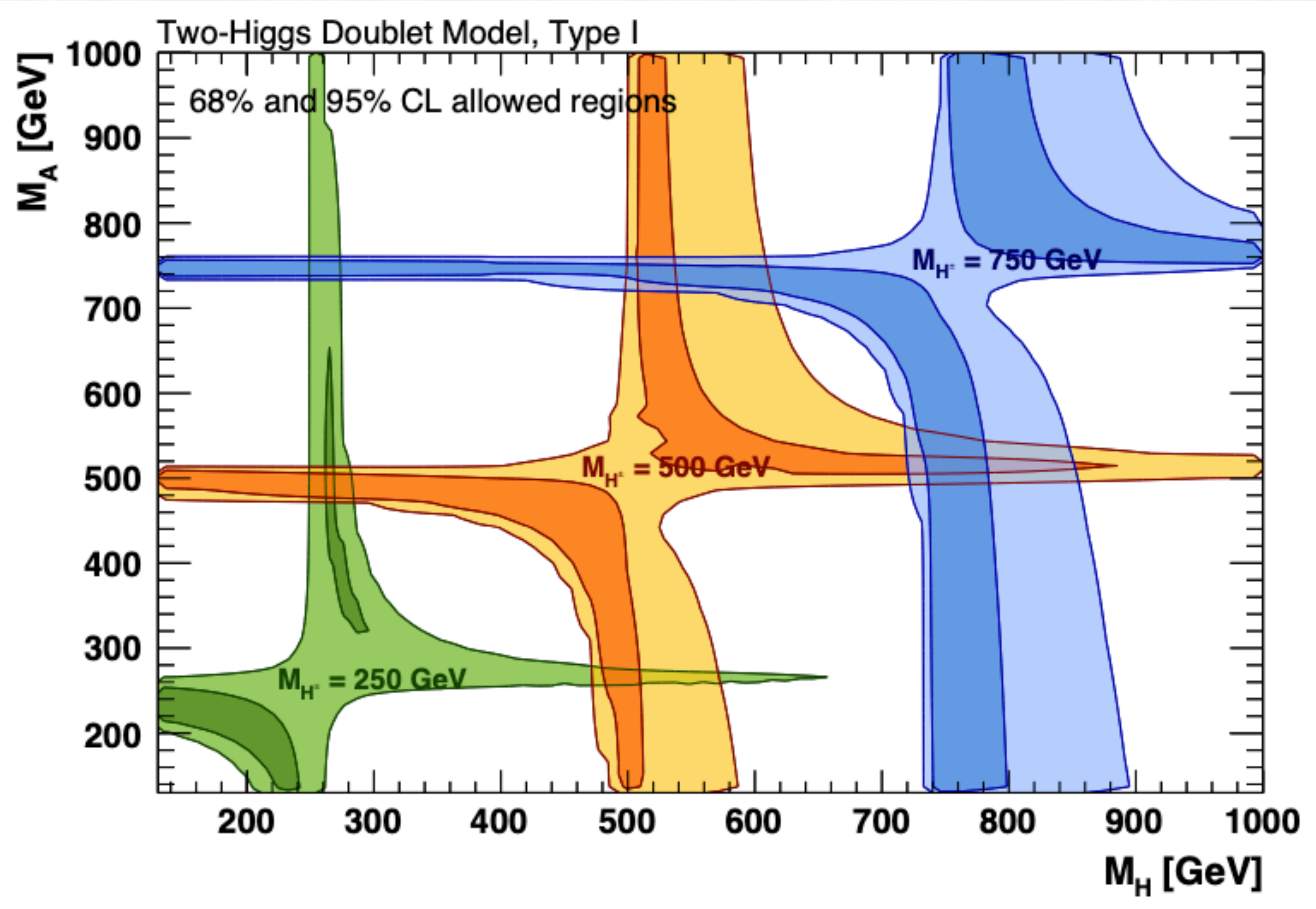


Gfitter, EPJC 78 (2018) 675

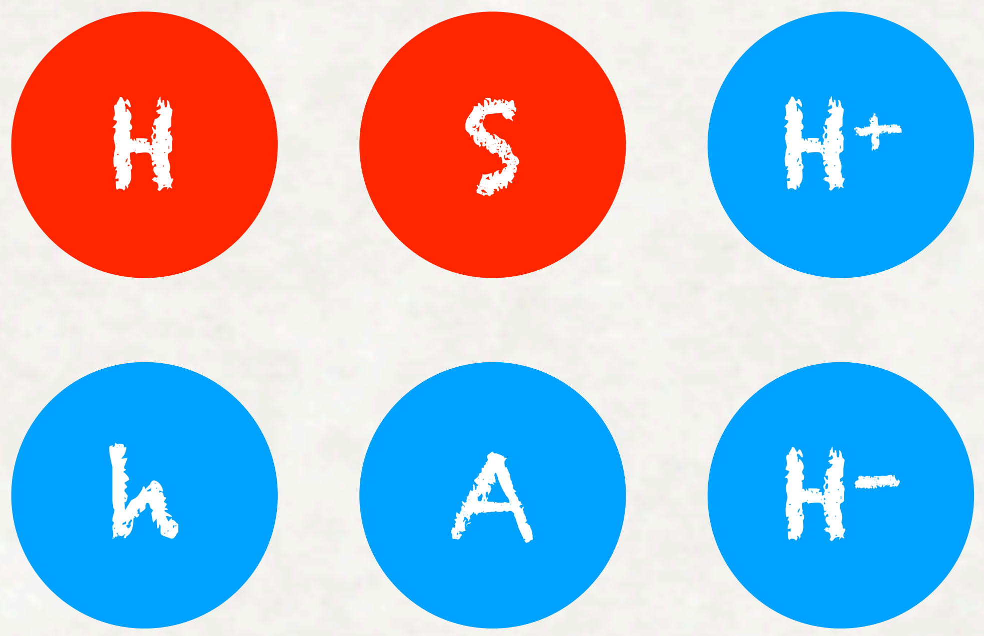
- Higgs coupling measurements: we are close to alignment limit
- $H^\pm$ must be degenerate with A or H
  - ➡ in the following we only consider $\cos(\boldsymbol{\beta}-\boldsymbol{\alpha})=0$ & $m_A=m_H=m_{H\pm}$

Model #1: 2HDM + scalar

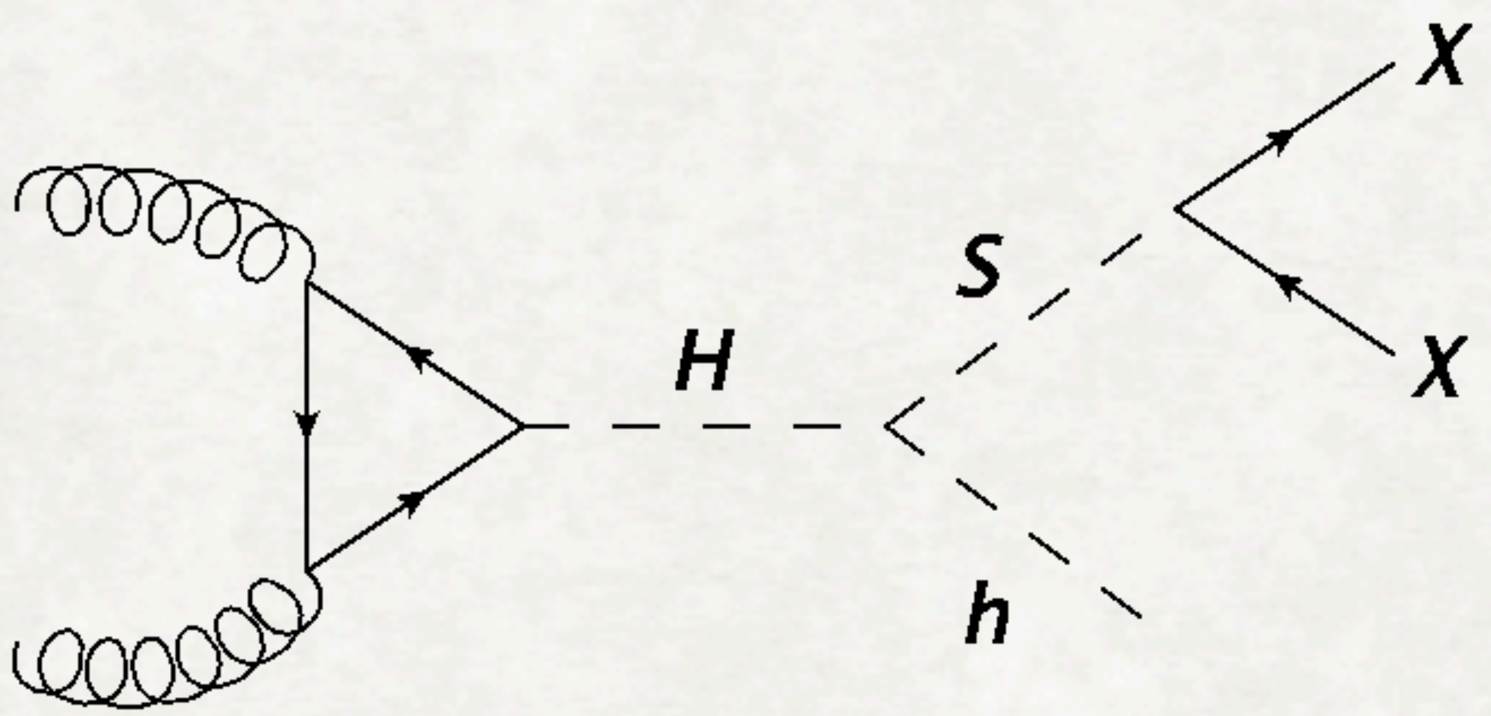
- 2HDM Type-II
- Extra **scalar mediator S** that couples to DM
- Mixing between CP-even scalars
- 6 Higgs bosons
- Resonant signatures

$$H = \cos\theta \tilde{H} + \sin\theta \tilde{S}$$

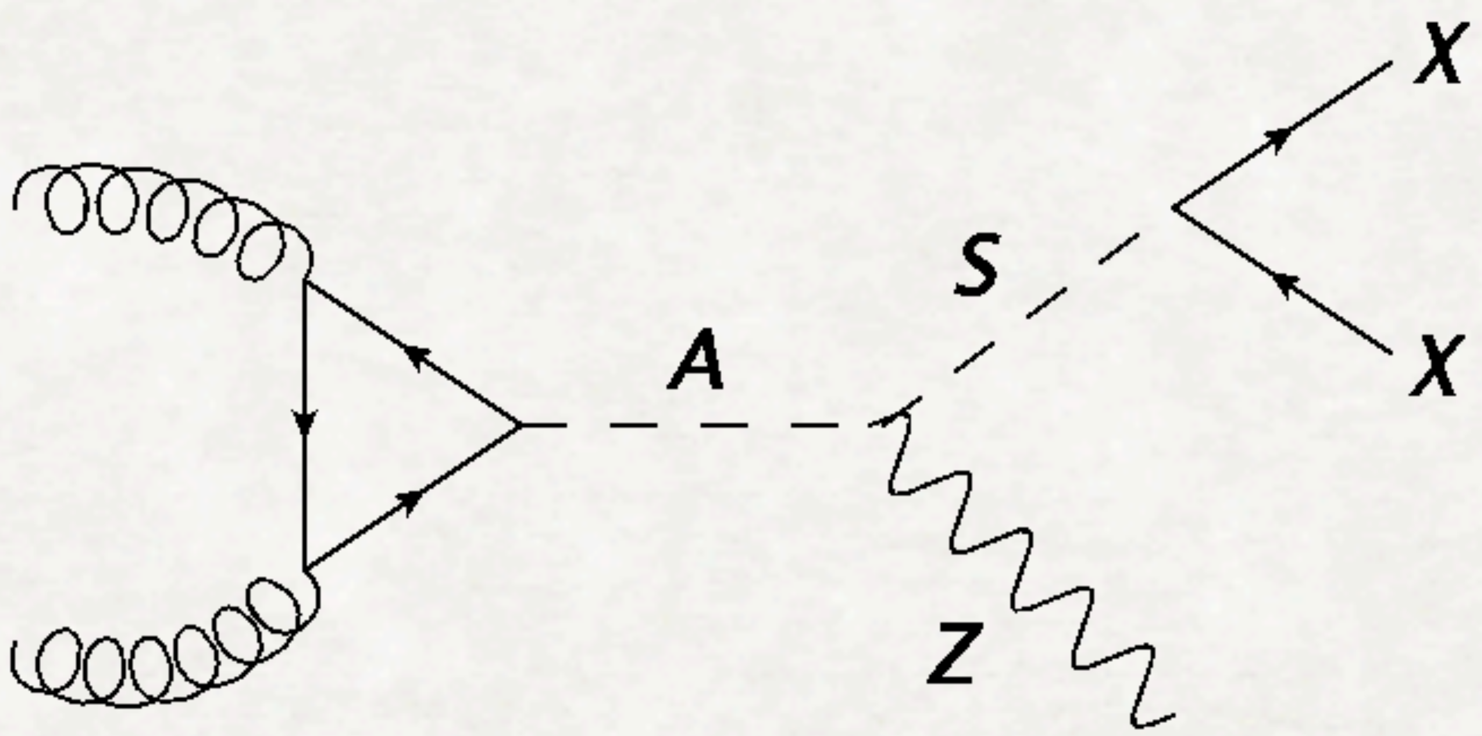
$$S = -\sin\theta \tilde{H} + \cos\theta \tilde{S}$$



$h+E_T^{miss}$



$Z+E_T^{miss}$

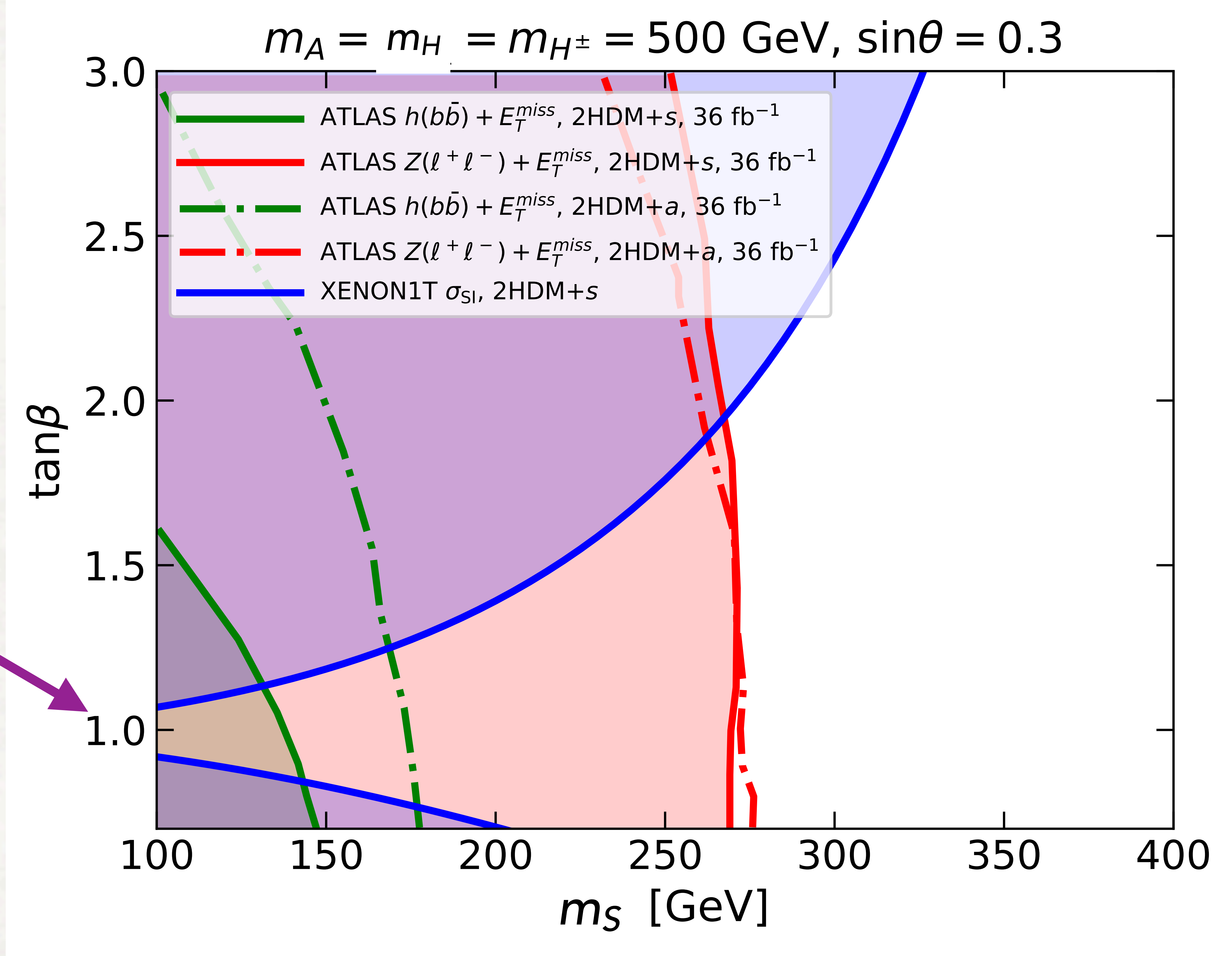


Bell et al, 1612.03475; Arcadi et al, 2001.10540

2HDM + scalar: constraints

- Not very much explored @ LHC
- Scalar mediator ⇒ dominant constraints from direct detection

★ DD experiments blind in certain regions
- scalars are degenerate (m_S=m_H)
- tan$\beta \cong 1$
- even for models that are considered DD territory, LHC can provide complementary constraints



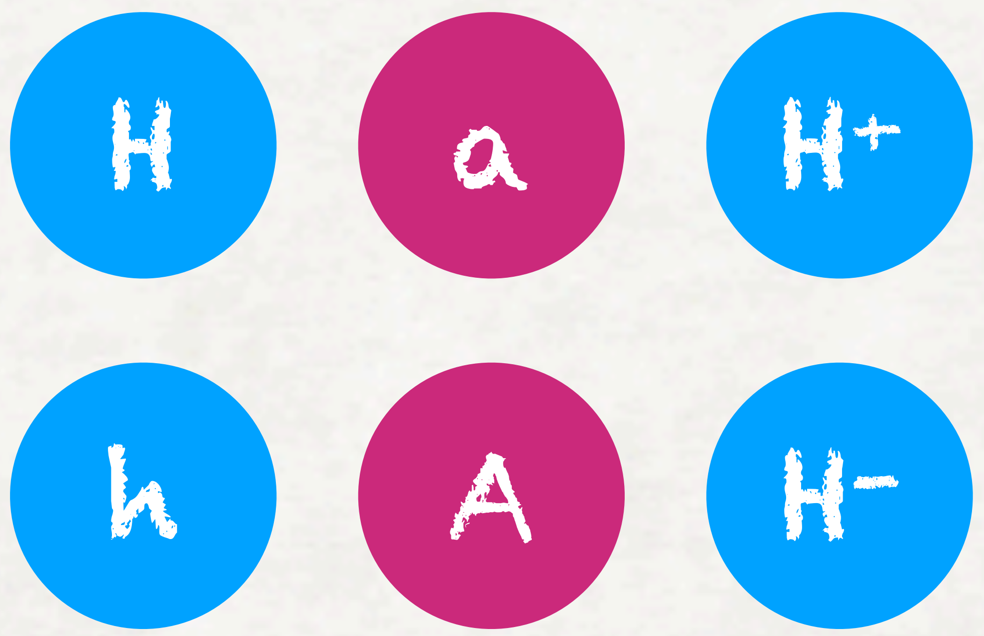
SA, Brandt, Haisch, Symmetry 13 (2021) 2406

Model #2: 2HDM + pseudoscalar

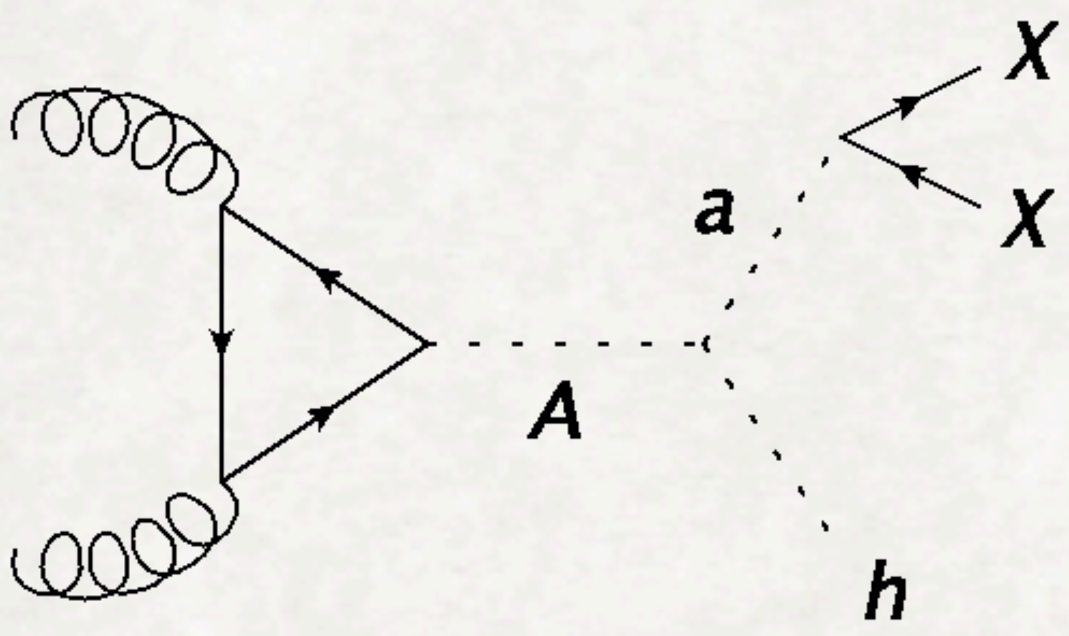
- 2HDM Type-II
- Mixing between CP-odd Higgses
- 6 Higgs bosons
- Extra **pseudoscalar mediator a** that couples to DM
  - suppressed DD constraints
  - originally proposed to explain Fermi-LAT excess
- Very rich phenomenology: colliders + ID + DD

$$A = \cos\theta \tilde{A} + \sin\theta \tilde{a}$$

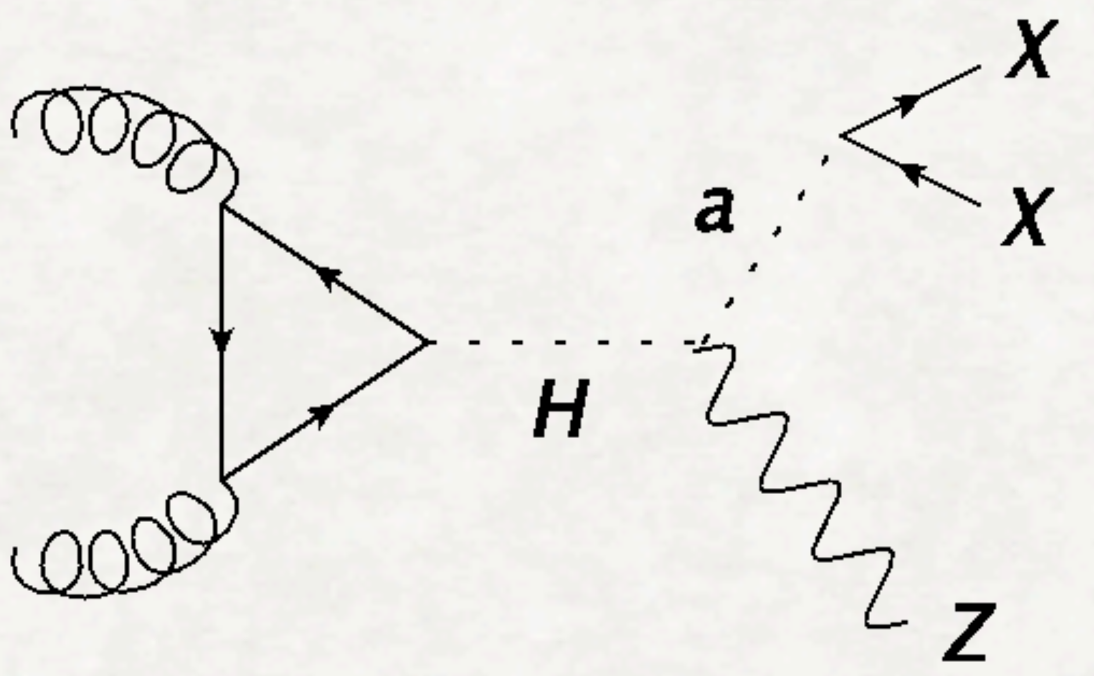
$$a = -\sin\theta \tilde{A} + \cos\theta \tilde{a}$$



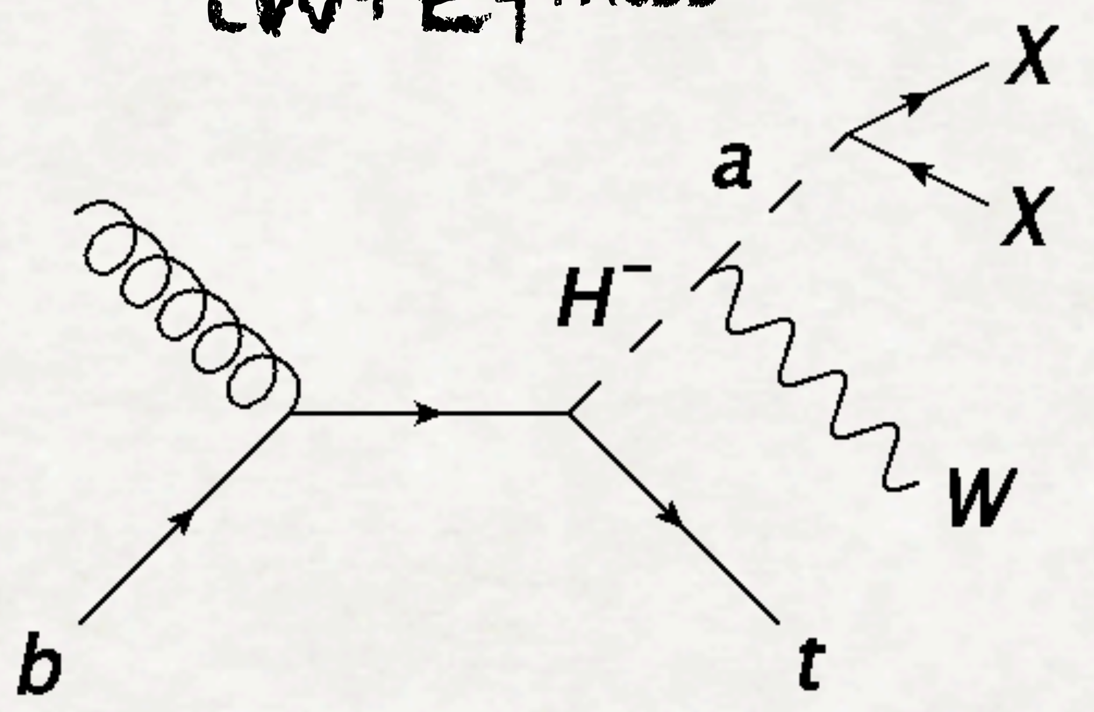
$h + E_T^{miss}$



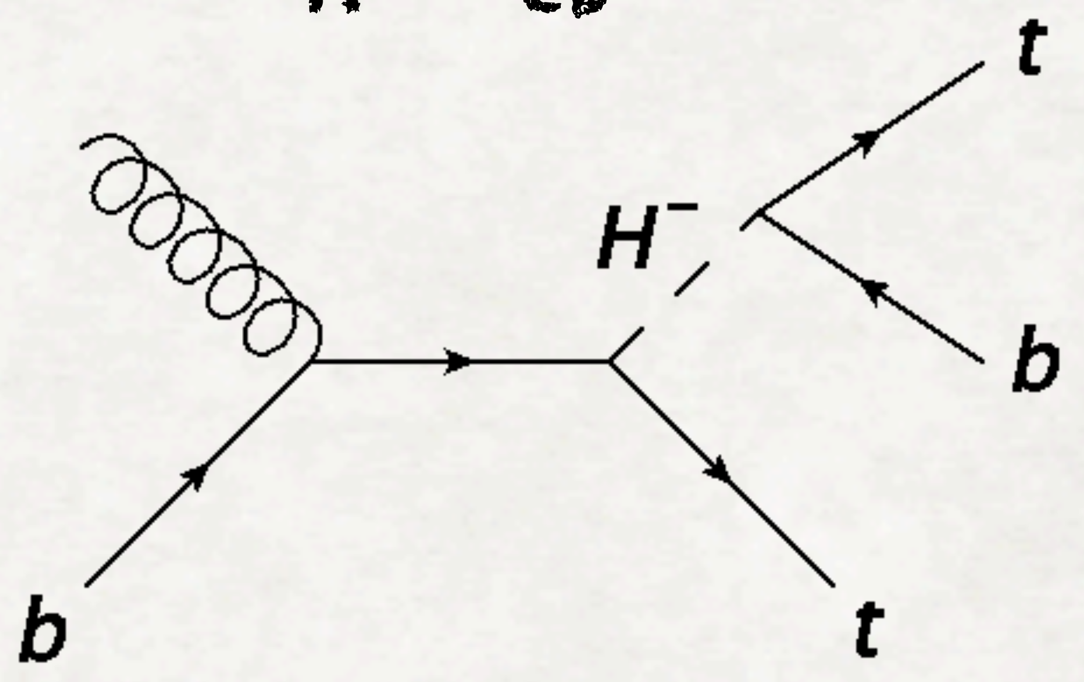
$Z + E_T^{miss}$



$tW + E_T^{miss}$

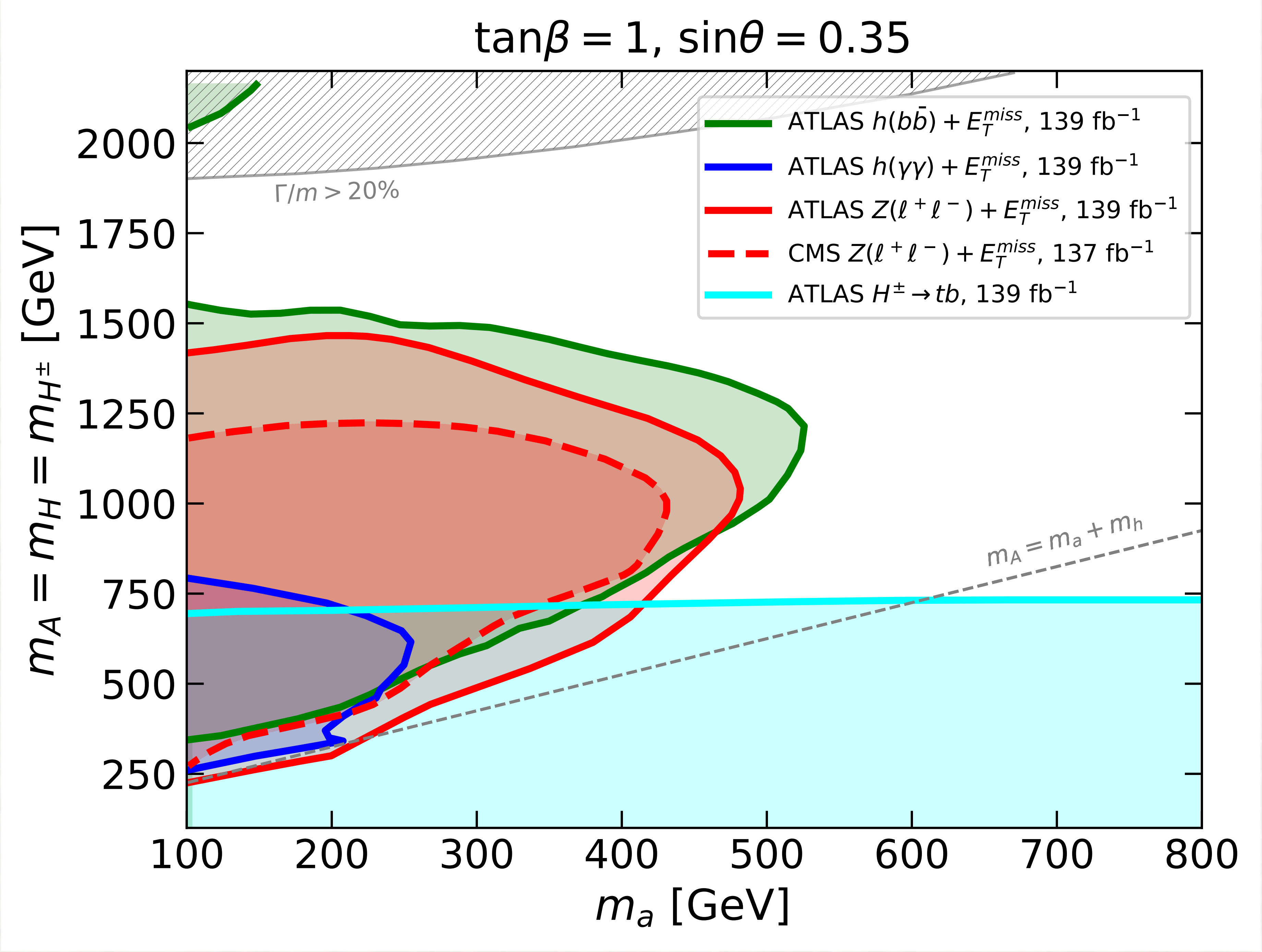


$H^{\pm} \rightarrow tb$

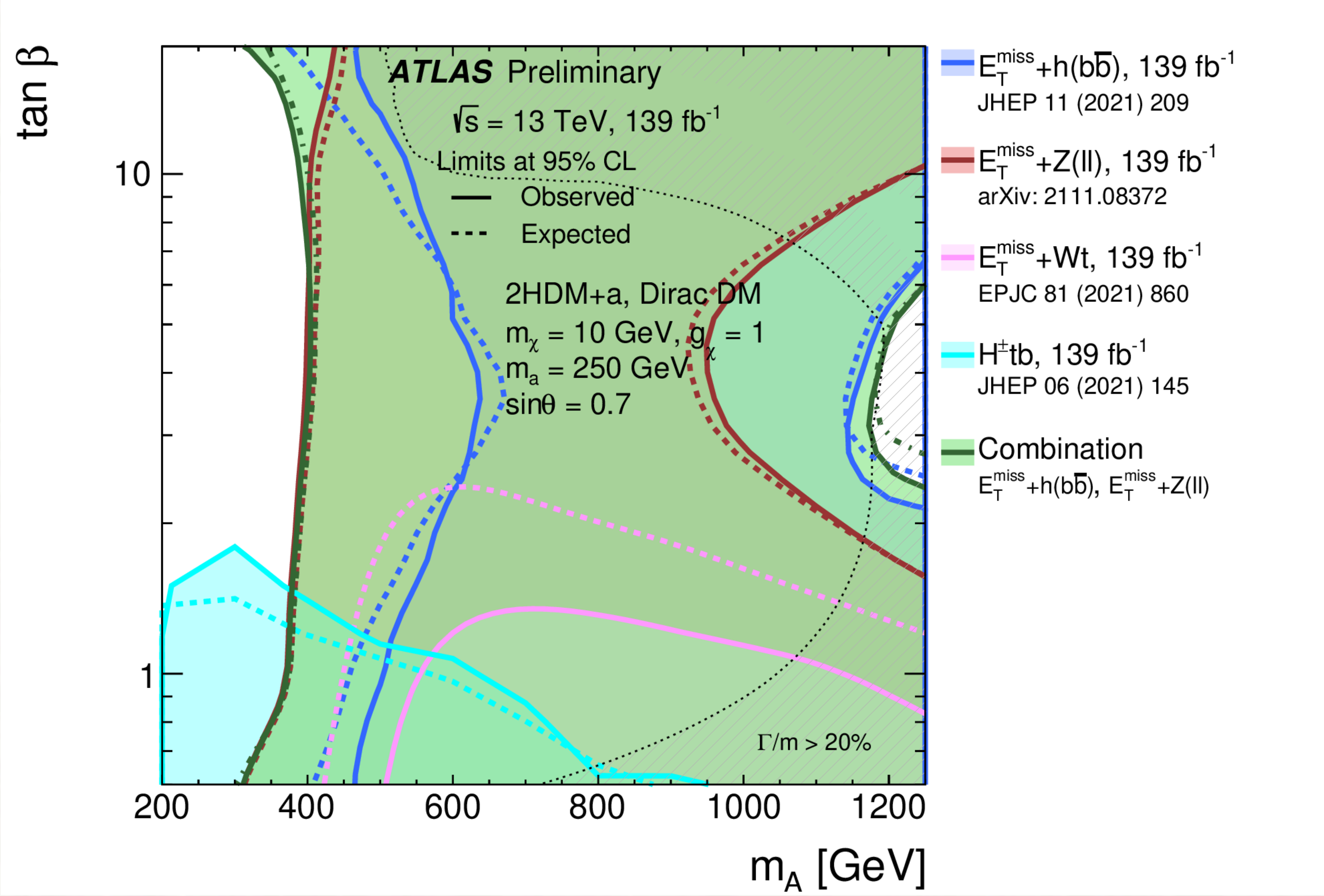


Ipek, et al, 1404.3716 ; No, 1509.01110 ; Goncalves, et al, 1611.04593 ; Bauer, et al, 1701.07427 ; Abe et al, 1810.09420

2HDM + pseudoscalar: constraints



[SA, Brandt, Haisch, Symmetry 13 (2021) 2406]



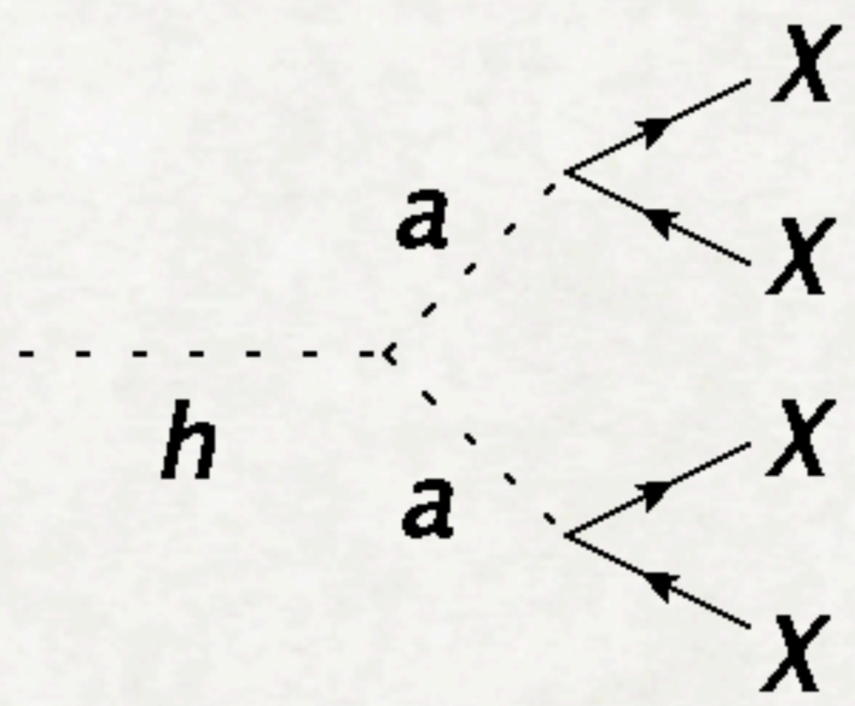
[ATL-PHYS-PUB-2021-045]

- A lot of parameter space excluded, $m_a \gtrsim 500$ GeV, $m_A \gtrsim 1$TeV for a range of mixing angles
- Goal: close sensitivity gaps (e.g. low m_A,m_a at intermediate tan**β**)

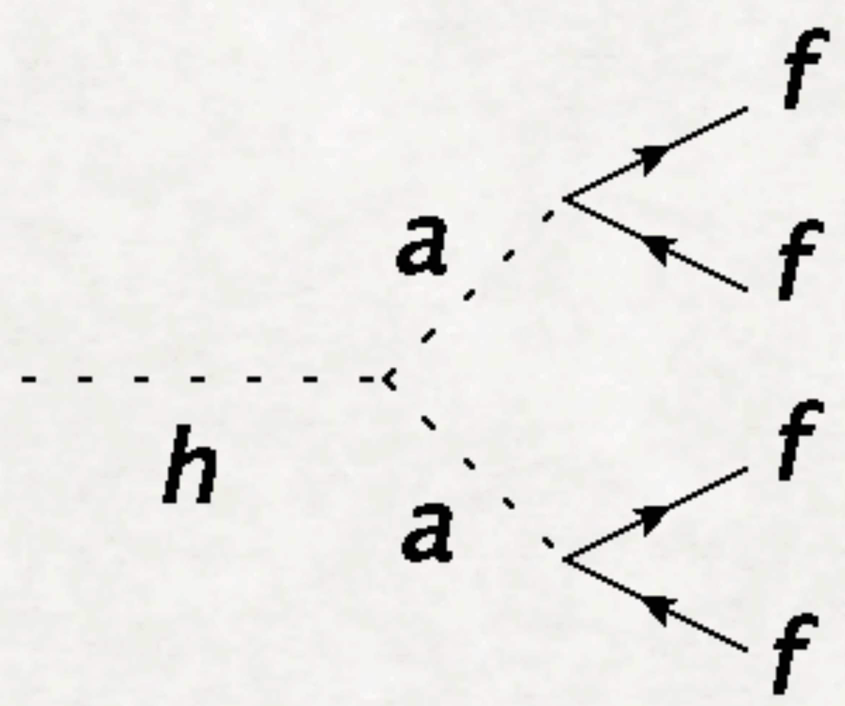
2HDM + pseudoscalar: complementary searches

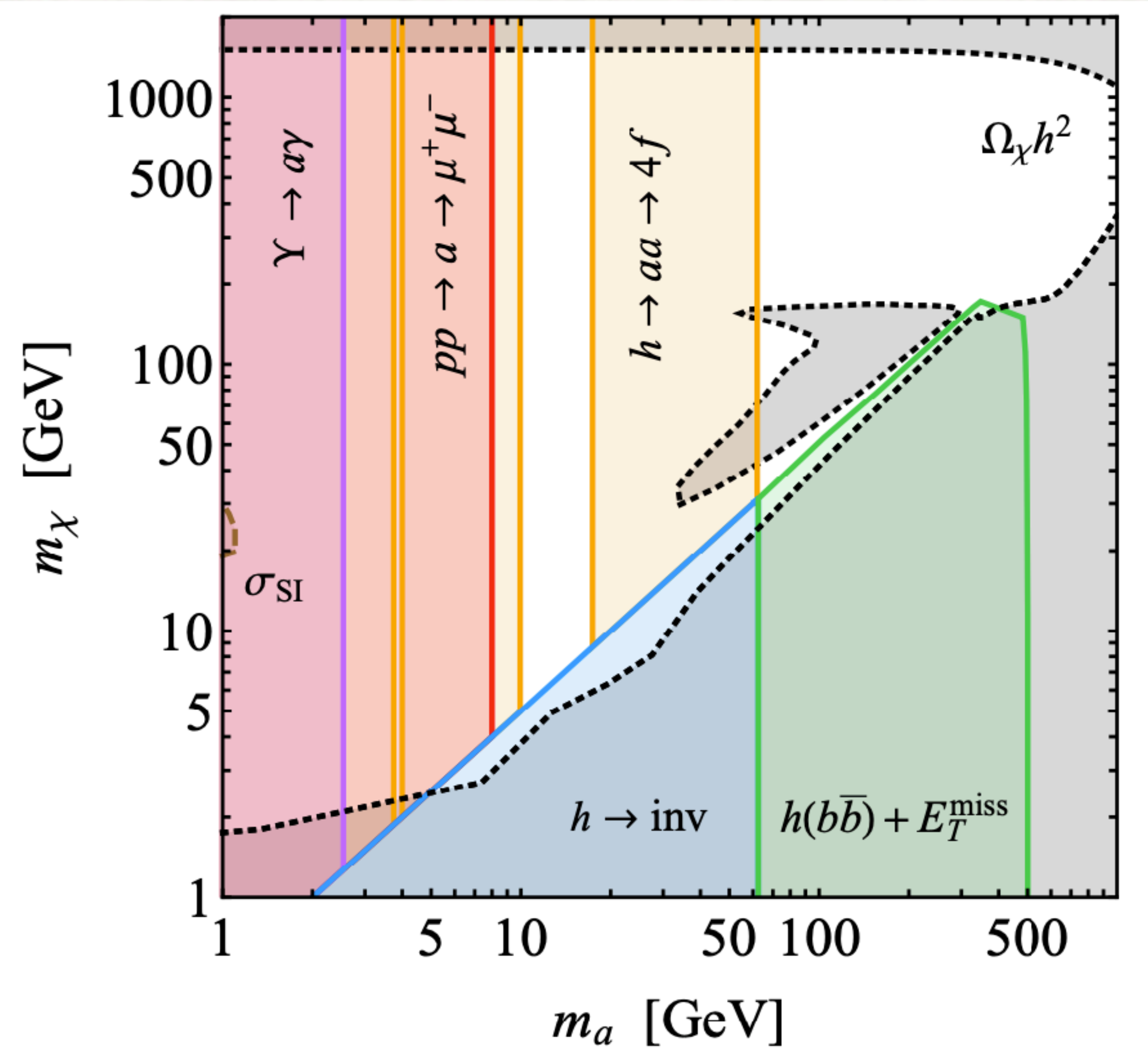
- $m_a > m_h/2$ & low m_χ: X+E_T^{miss}
- $m_a < m_h/2$ & $m_\chi < m_a/2$: h→invisible
- $m_a < m_h/2$ & $m_\chi > m_a/2$: h→4 fermions
  - generally when h→aa is open the model is tightly constrained from Higgs width unless finely tuned

h→inv



h→4f



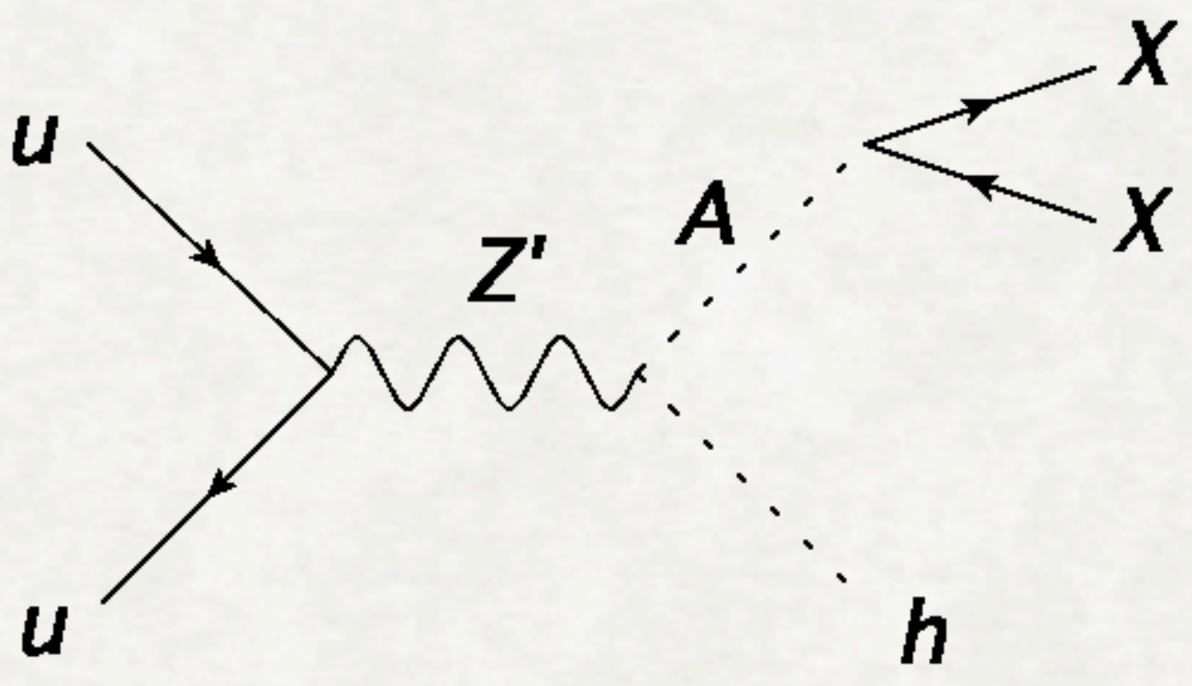


SA, Haisch, 2202.12631

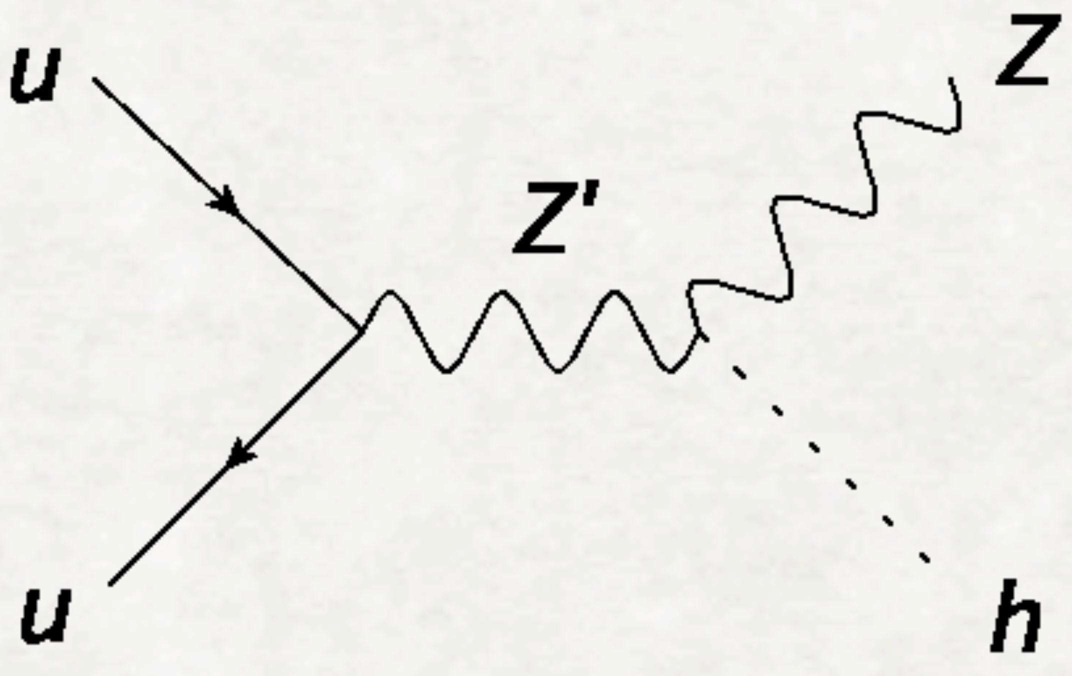
Model #3: 2HDM + vector

- 2HDM Type-II + Z′ only coupling to up-type quark ⇒ evades dilepton constraints
- CP-odd Higgs A couples to DM particles
- Large h+E_T^{miss} signal (highly boosted Higgs in contrast to 2HDM+a)
- Also constraints from EW measurements and dijets

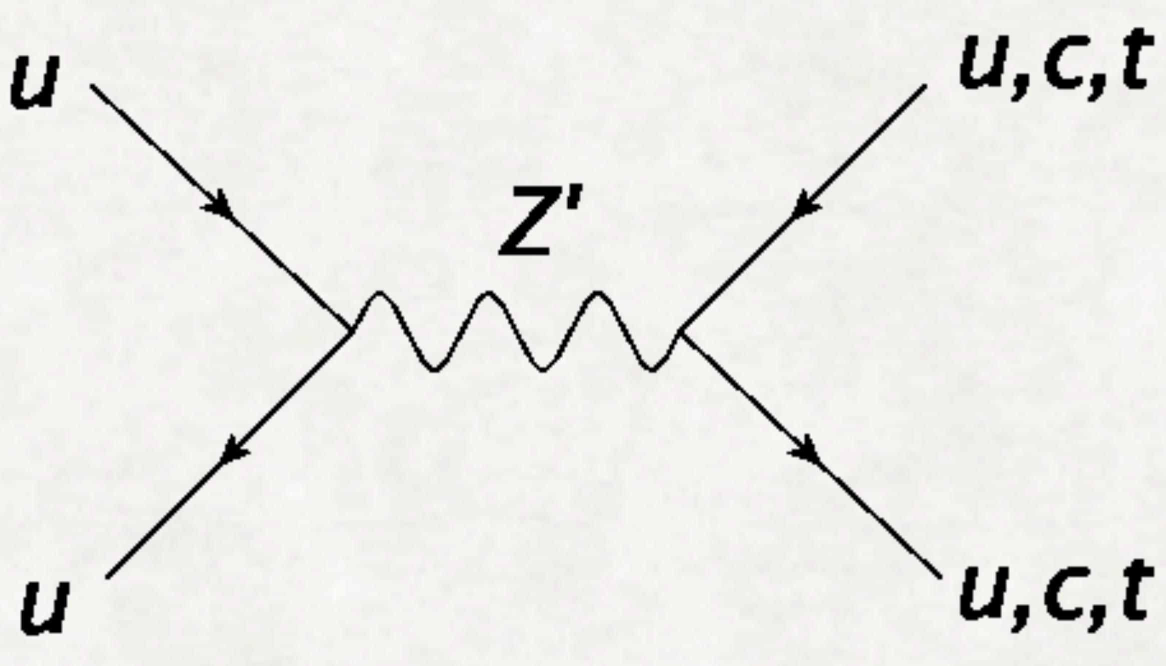
h+E_T^{miss}



Zh resonance

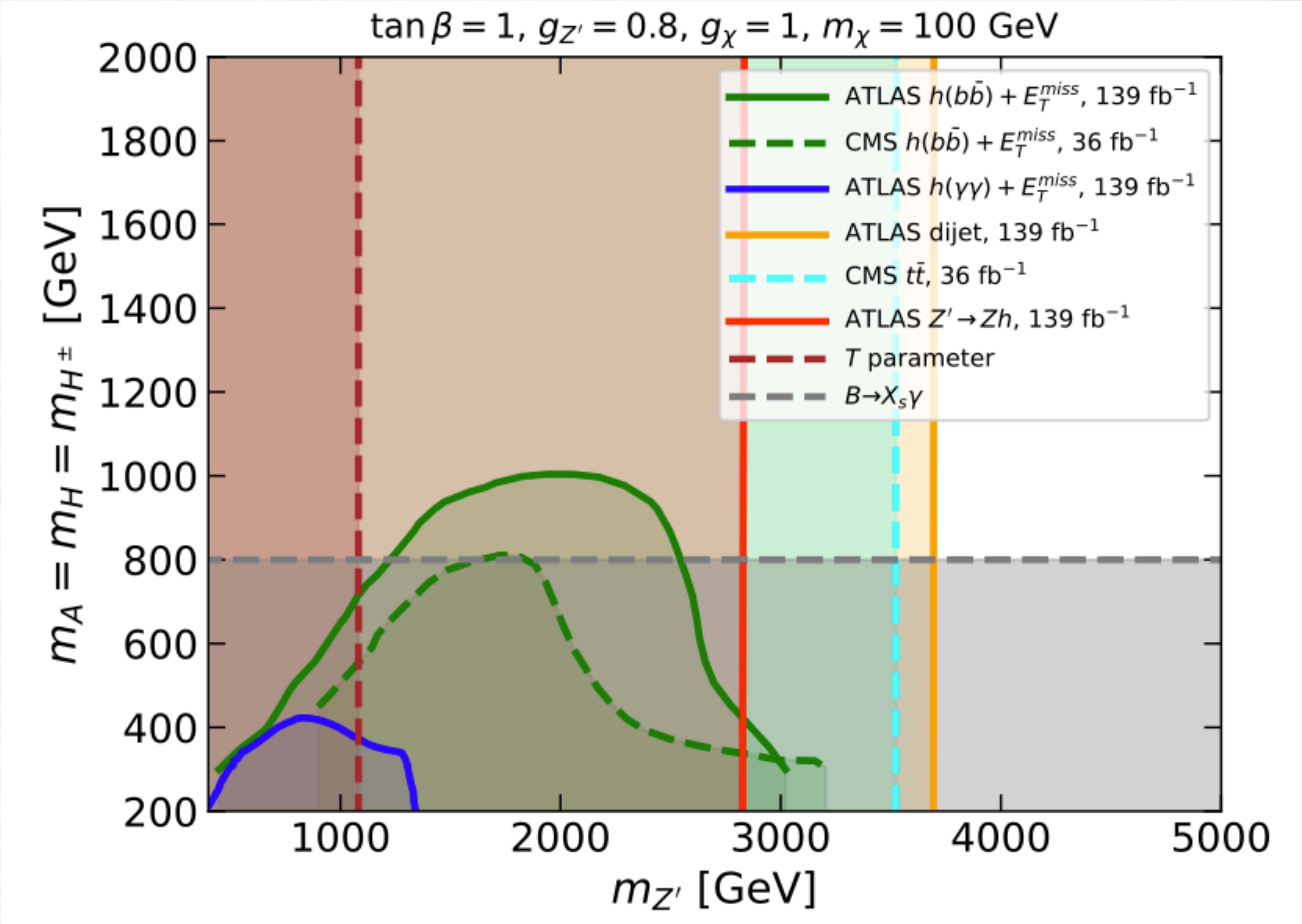


dijet/$t\bar{t}$ resonance



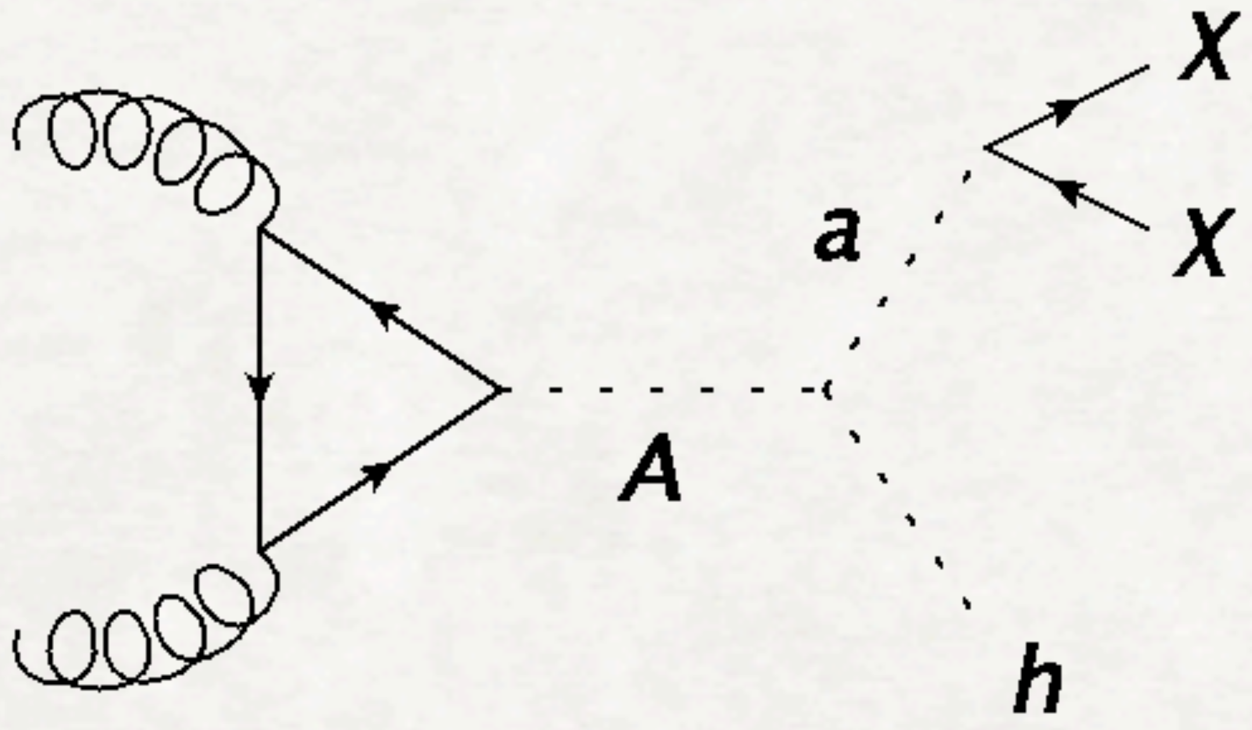
2HDM + vector: constraints

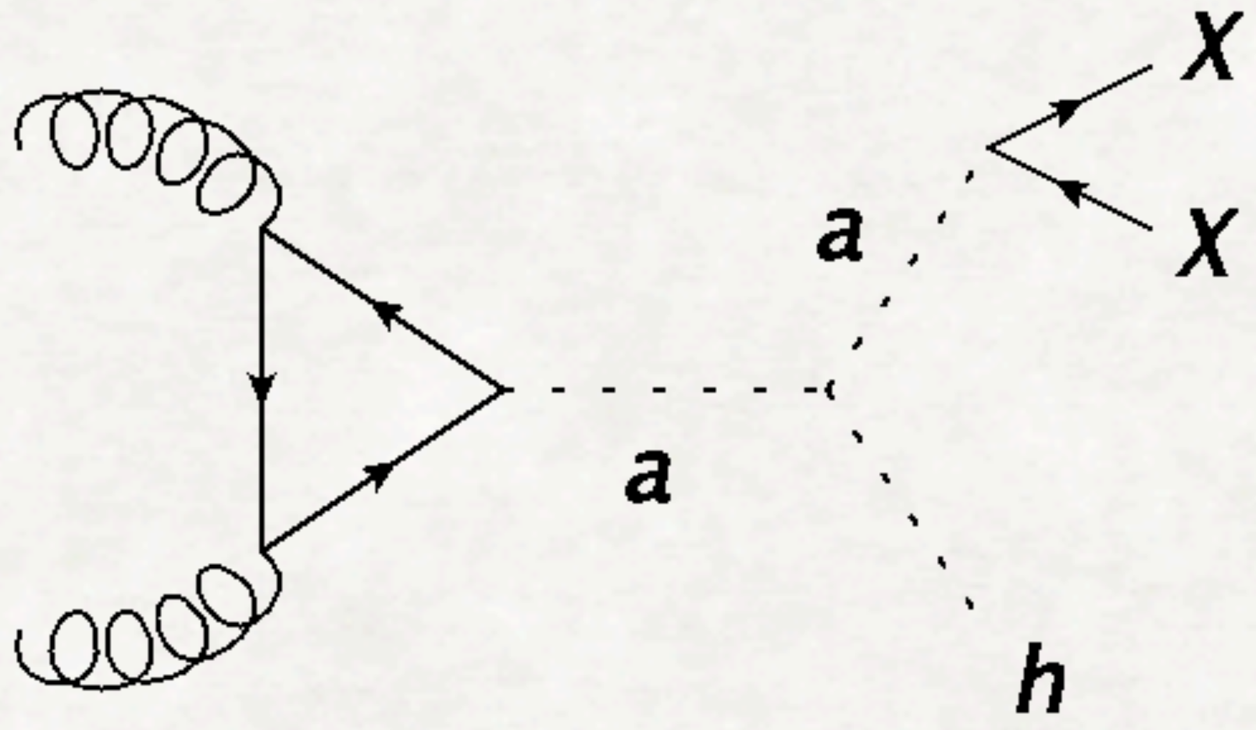
- $m_{Z'}$ excluded up to 2-3 TeV for $m_A \lesssim$ 1 TeV
- EW and flavour measurements provide significant constraints
- Zh and dijet resonances provide better constraints (and this seems hard to avoid)
   - ➡ "DM searches" don't always provide the best constraints to DM models

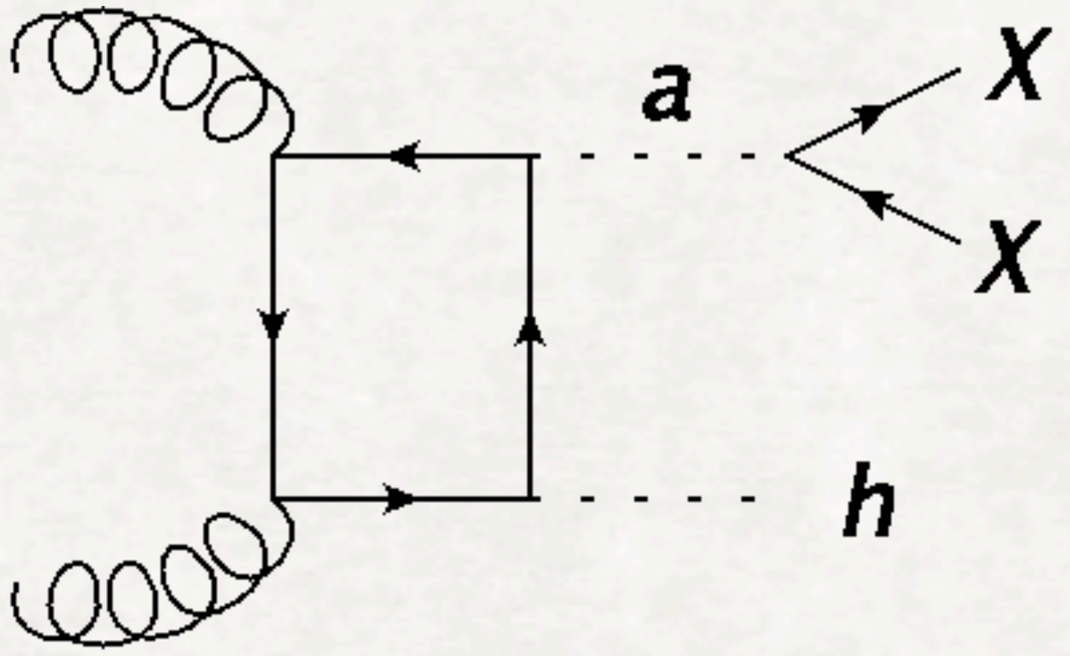


SA, Brandt, Haisch, Symmetry 13 (2021) 2406

“Model-independent” Limits





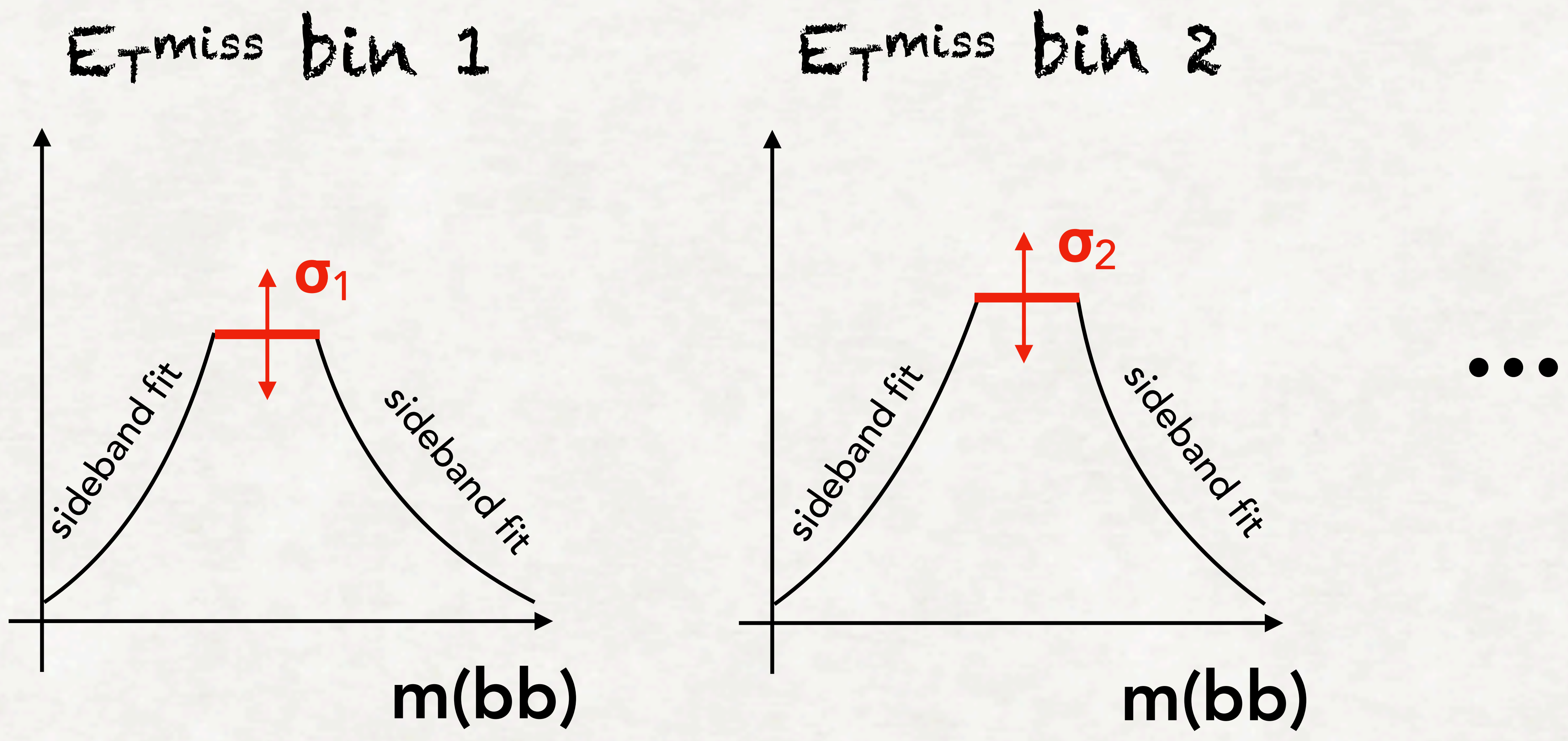


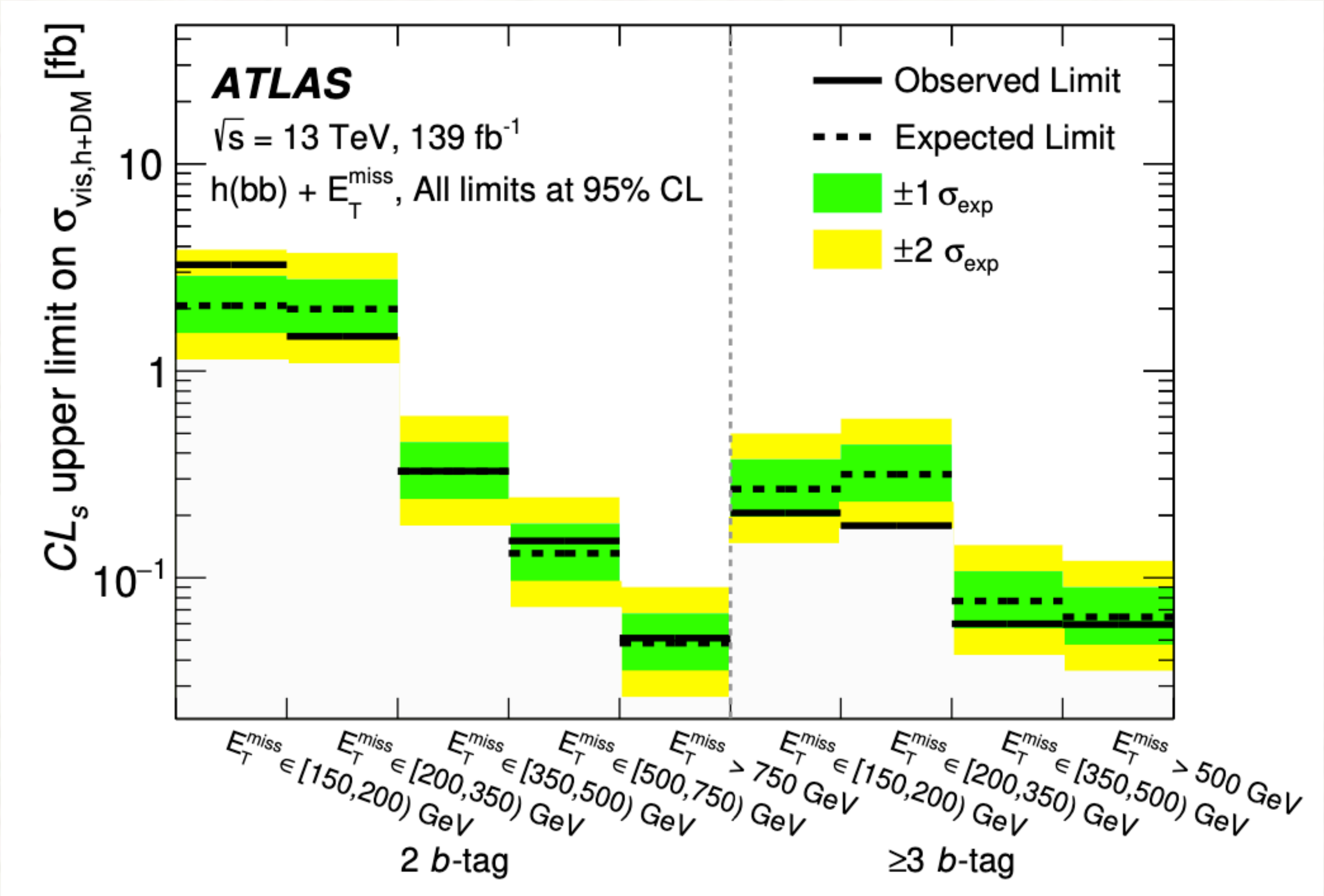
- Parameter choices can affect signal characteristics (e.g. softer E_T^{miss})

➡ **How to produce limits that are easy to re-interpret?**

"Model-independent" Limits

- Present constraints in terms of $\sigma^{\text{vis}}_{h+\text{DM}} \equiv \sigma_{h+\text{DM}} \times \text{BR}(h \to b\bar{b}) \times (\mathcal{A} \times \epsilon)$





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- Maximum cross-section of a signal-like resonance that the data can accommodate in bins of a given variable (e.g. E_T^{miss})
- Folding with Axϵ theorists can re-interpret results in any model with SM-like Higgs

Conclusions

Dark Matter: among the few evidence for new physics

Multifaceted approach necessary

➡ different experiments + different analyses

Higgs sector(s) can provide a portal to DM

➡ studying SM and BSM Higgs sectors crucial

Simple models increasingly ruled out - we need:

➡ systematic approach: combinations + re-interpretations

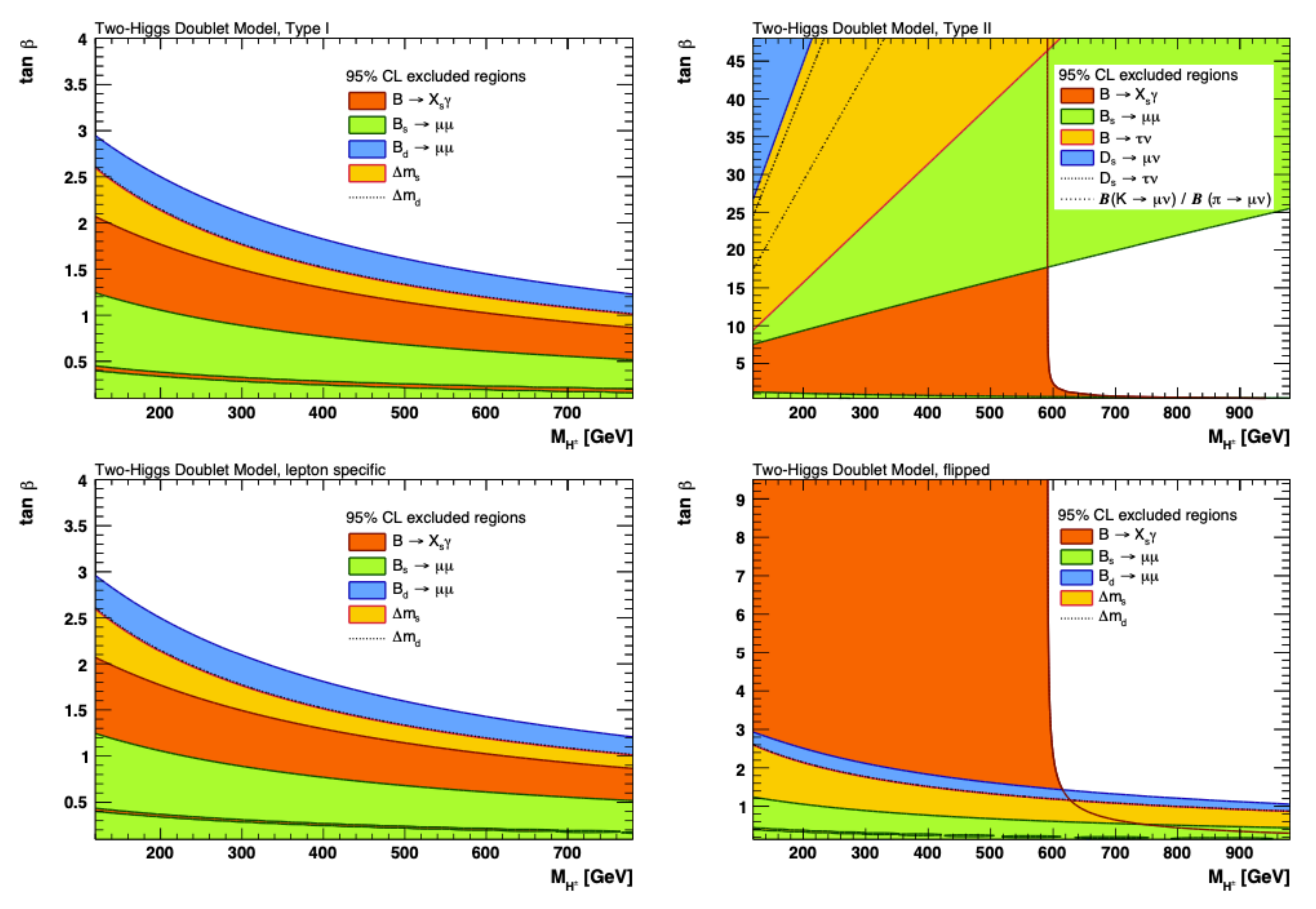
➡ exploration of all possible final states

Backup

2HDM couplings

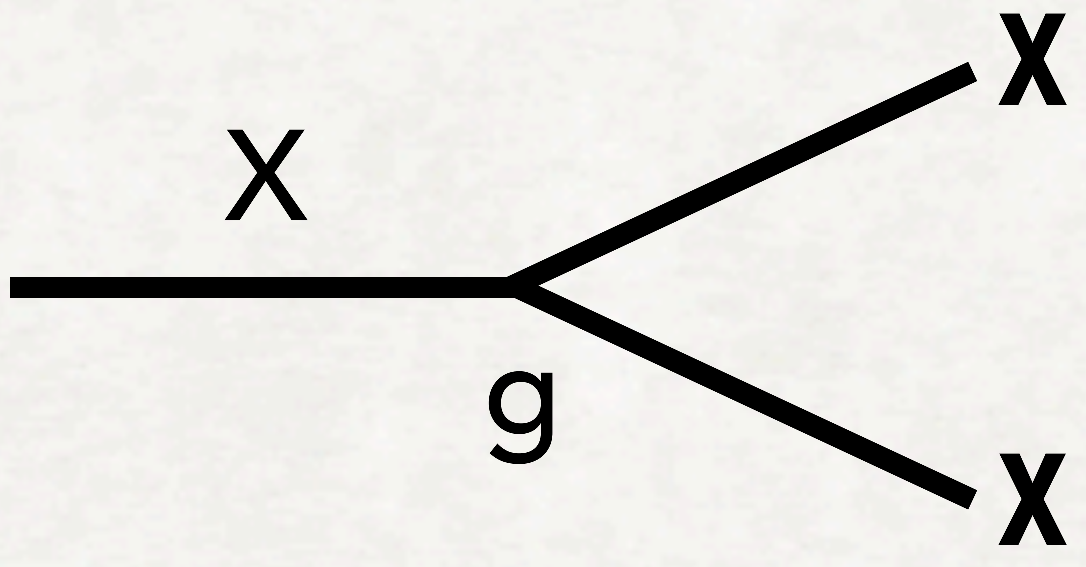
Coupling modifier	Type I	Type II
$\xi(h,u)$	$s_{\beta-\alpha}+c_{\beta-\alpha}/t_\beta$	$s_{\beta-\alpha}+c_{\beta-\alpha}/t_\beta$
$\xi(h,d)$, $\xi(h,l)$		$s_{\beta-\alpha}-c_{\beta-\alpha}t_\beta$
$\xi(H,u)$	$c_{\beta-\alpha}-s_{\beta-\alpha}/t_\beta$	$c_{\beta-\alpha}-s_{\beta-\alpha}/t_\beta$
$\xi(H,d)$, $\xi(H,l)$		$c_{\beta-\alpha}+s_{\beta-\alpha}t_\beta$
$\xi(A,u)$	$1/t_\beta$	$1/t_\beta$
$\xi(A,d)$, $\xi(A,l)$		t_β
$\xi(h,VV)$	$s_{\beta-\alpha}$	
$\xi(H,VV)$	$c_{\beta-\alpha}$	
$\xi(A,VV)$	0	

2HDM constraints



Higgs vs Z portal

from U.Haisch



$$\Gamma(X \to \chi\bar{\chi}) \sim g^2 m_X$$

$$\text{BR}(X \to \chi\bar{\chi}) = \frac{\Gamma(X \to \chi\bar{\chi})}{\Gamma_X}$$

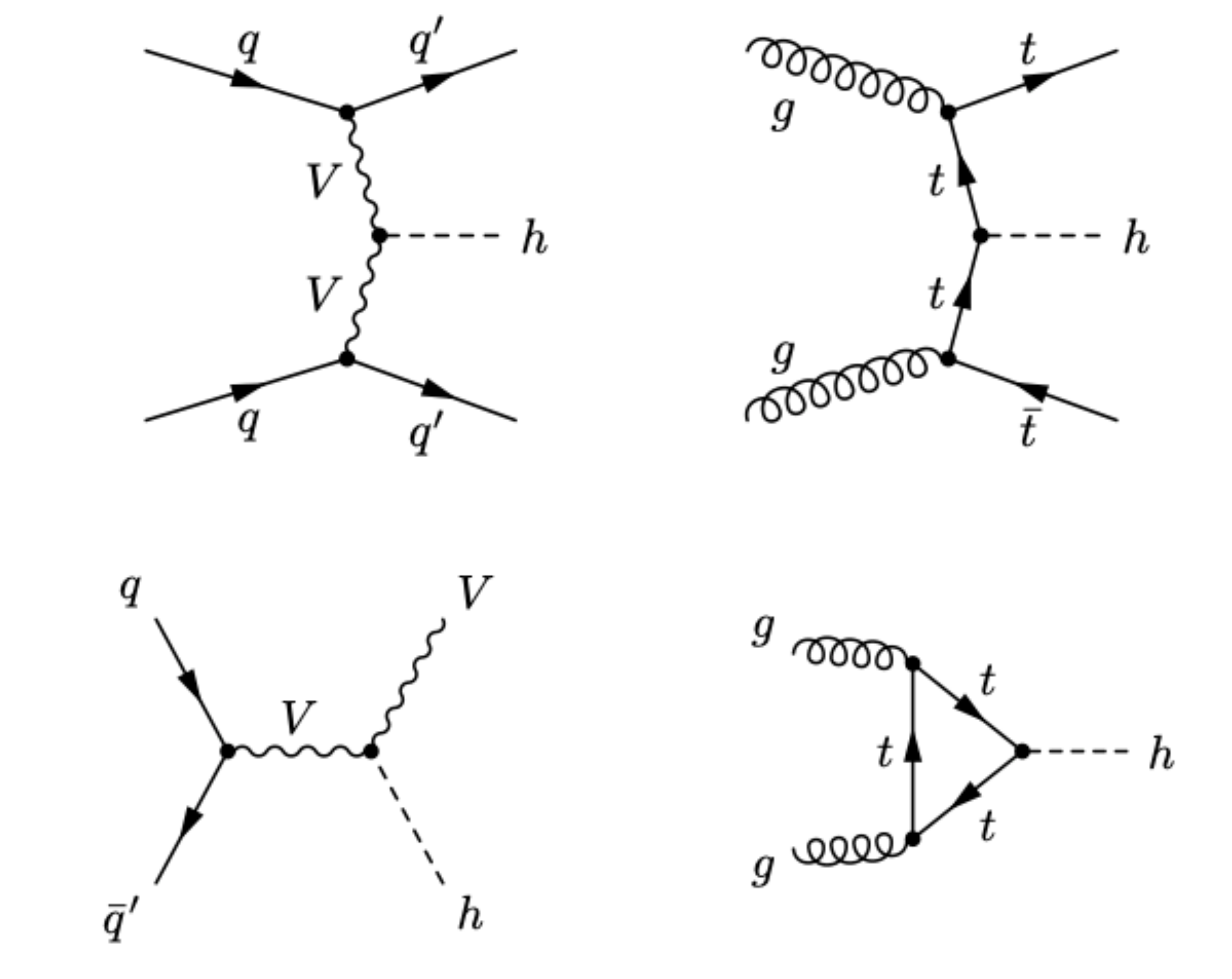
$$g \sim \frac{\Gamma_X \cdot \text{BR}(X \to \chi\bar{\chi})}{m_X}$$

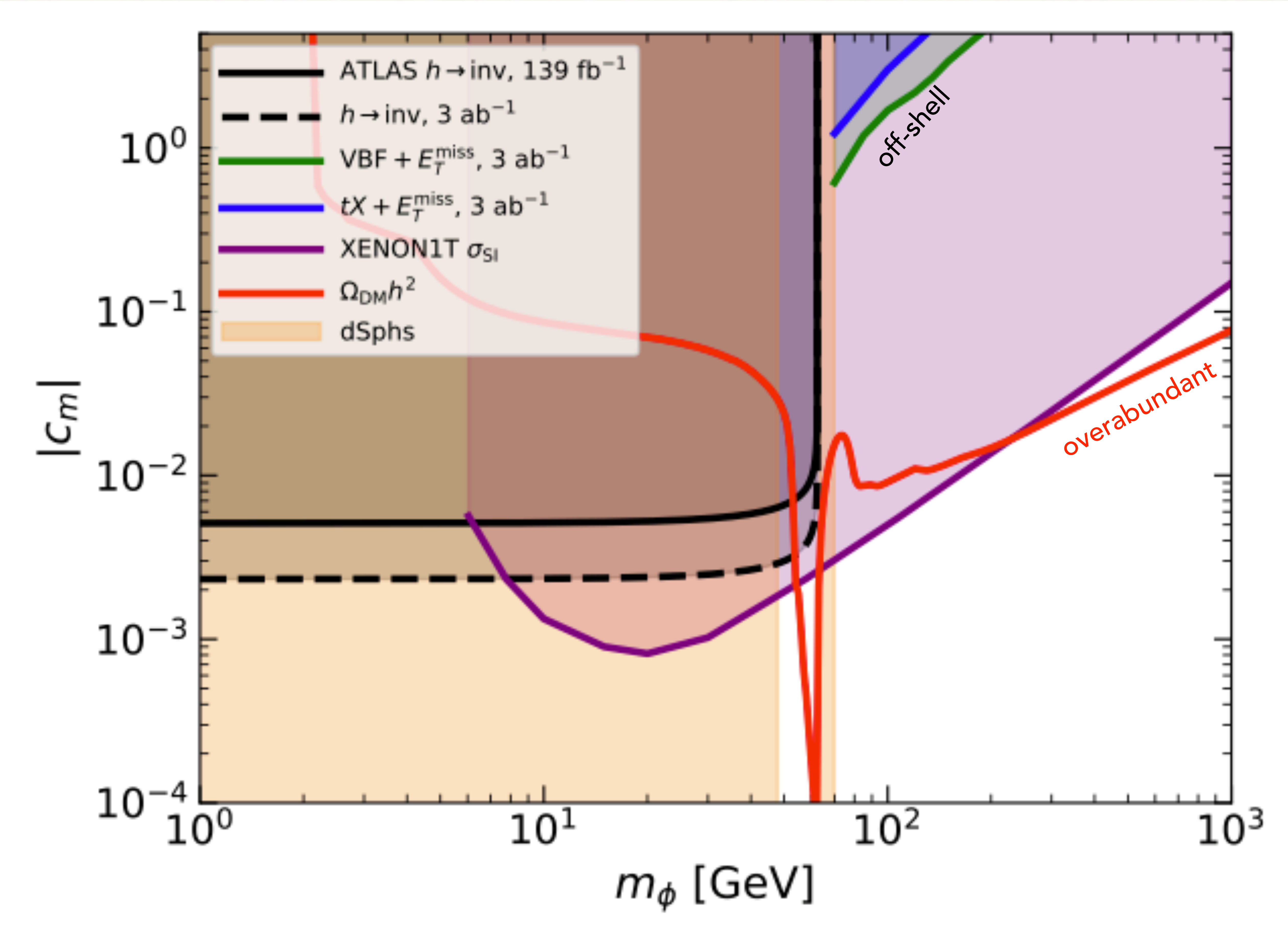
$$\Gamma_Z \simeq 2.5\ \text{GeV}, \Gamma(Z \to \chi\bar{\chi}) \lesssim 2\ \text{MeV} \Rightarrow g \lesssim 0.03$$

$$\Gamma_h \simeq 4.1\ \text{MeV}, \Gamma(Z \to \chi\bar{\chi}) \lesssim 0.5\ \text{MeV} \Rightarrow g \lesssim 0.01$$

SM Higgs portal

$c_m \phi^2 (H^\dagger H)$



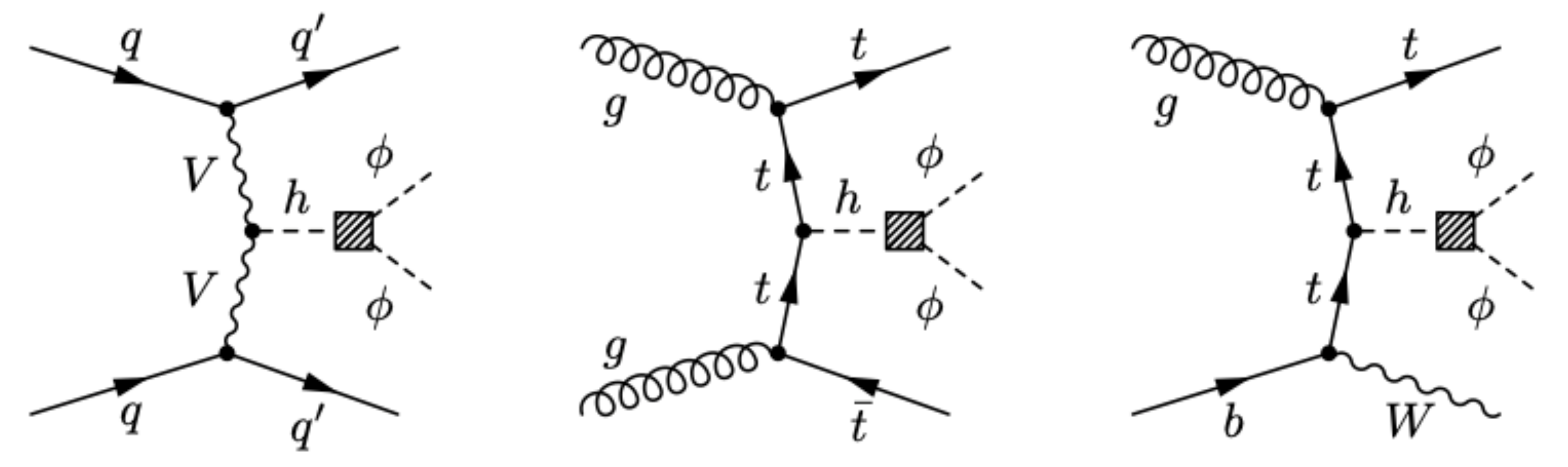


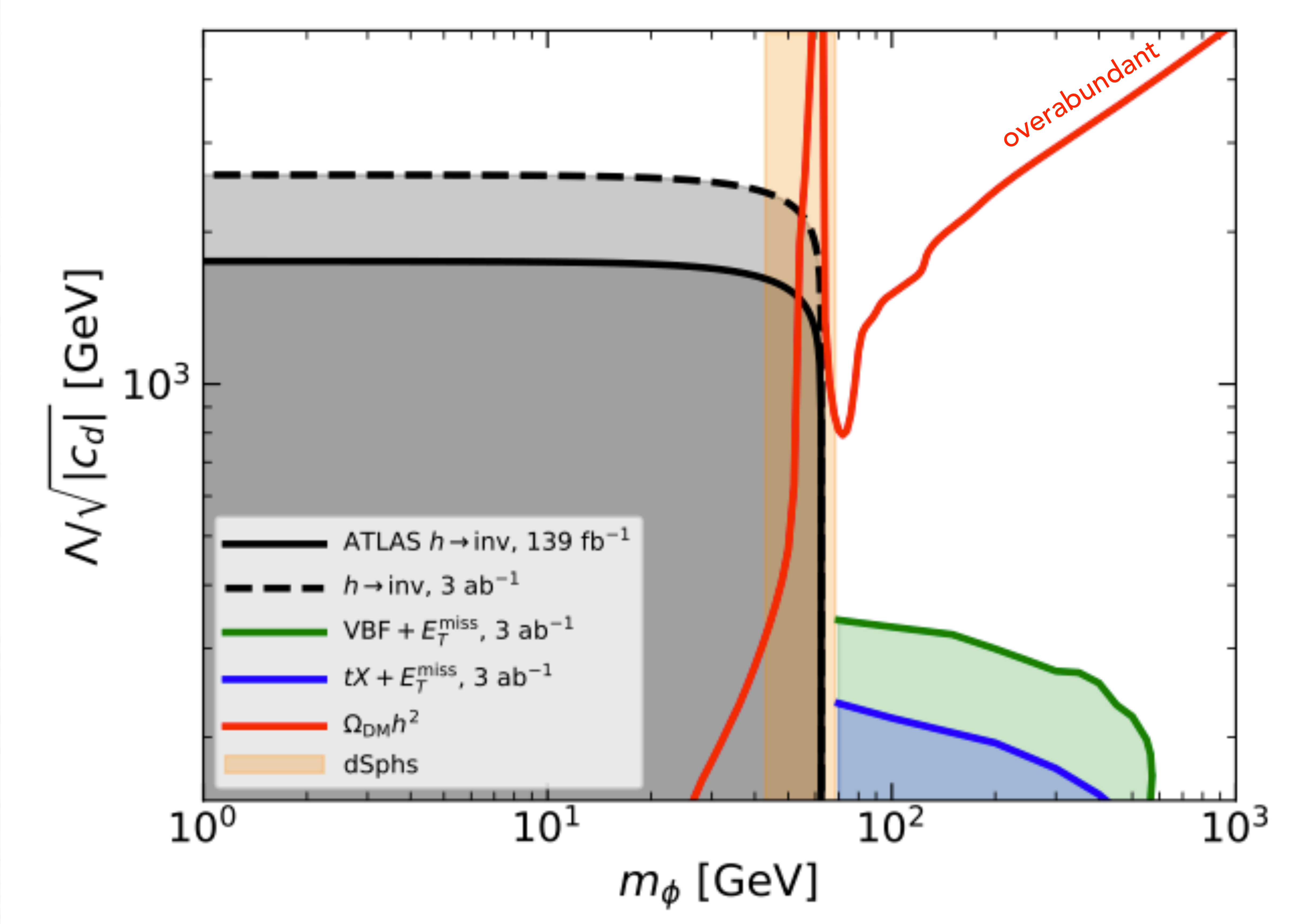
- DM scalar - dim 4 operator
- LHC constraints relevant for m < 5 GeV
- ID constraints from Fermi-LAT assume **ϕϕ**→bb and Ωh^2=0.12, so model dependent

Rudhofer et al, 1910.04170; Haisch et al, 2107.12389; SA, Brandt, Haisch 2109.13597

Derivative Higgs portal

$$\frac{c_d}{\Lambda^2}\left(\partial_\mu\phi^2\right)\left(\partial^\mu(H^\dagger H)\right)$$





- dim 6 EFT operator
- arise in models with global symmetry breaking, where DM is a pNGB - **Λ** ~ scale of symmetry breaking
- kinetic dependence of interaction suppresses DD constraints

Rudhofer et al, 1910.04170; Haisch et al, 2107.12389; SA, Brandt, Haisch 2109.13597

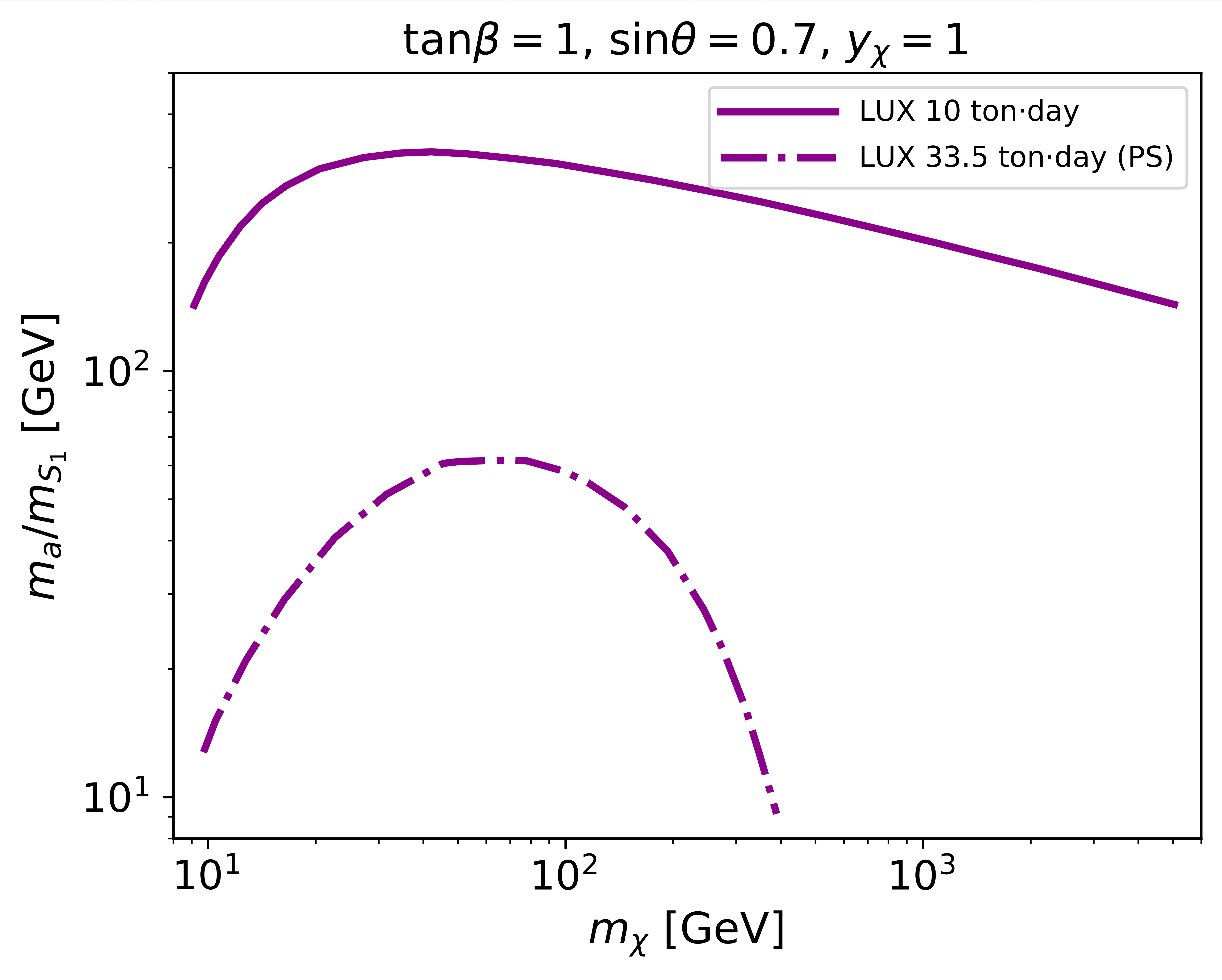
2HDM + scalar: WIMP-nucleon cross-section

Wilson coefficient of $\chi\bar{\chi}N\bar{N}$

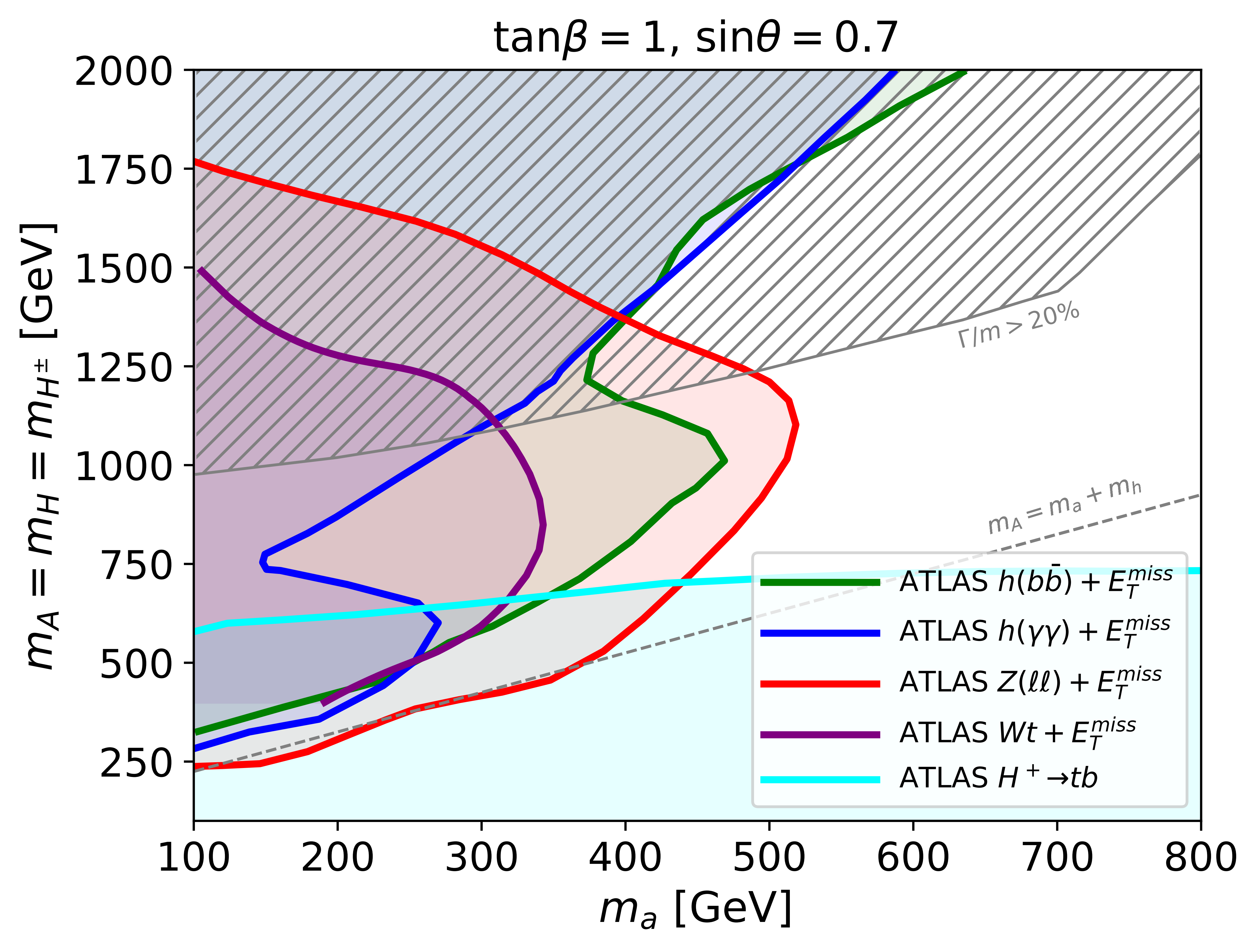
$$c_N = \frac{m_N}{v}\frac{y_\chi \sin(2\theta)}{2}\left(\frac{1}{m_{S_1}^2} - \frac{1}{m_{S_2}^2}\right) \times \left[\cot\beta\, f_{T_u}^N - \tan\beta \sum_{q=d,s} f_{T_q}^N + \frac{4\cot\beta - 2\tan\beta}{27} f_{T_G}^N\right]$$

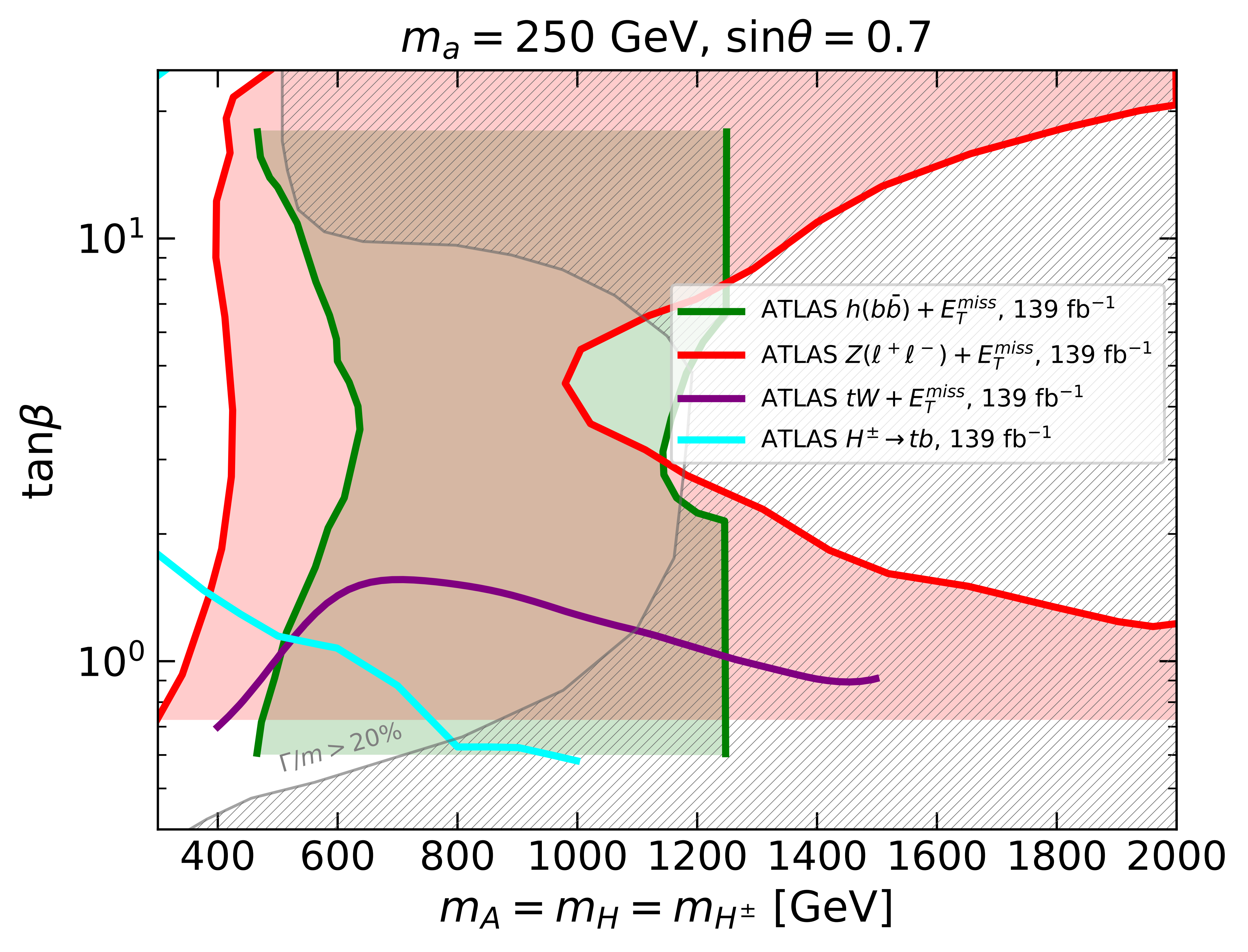
- Up and down-quark contributions interfere destructively in Type-II
- Numerically close to 0 for tan**β** ≅ 1

DD scalar vs pseudoscalar

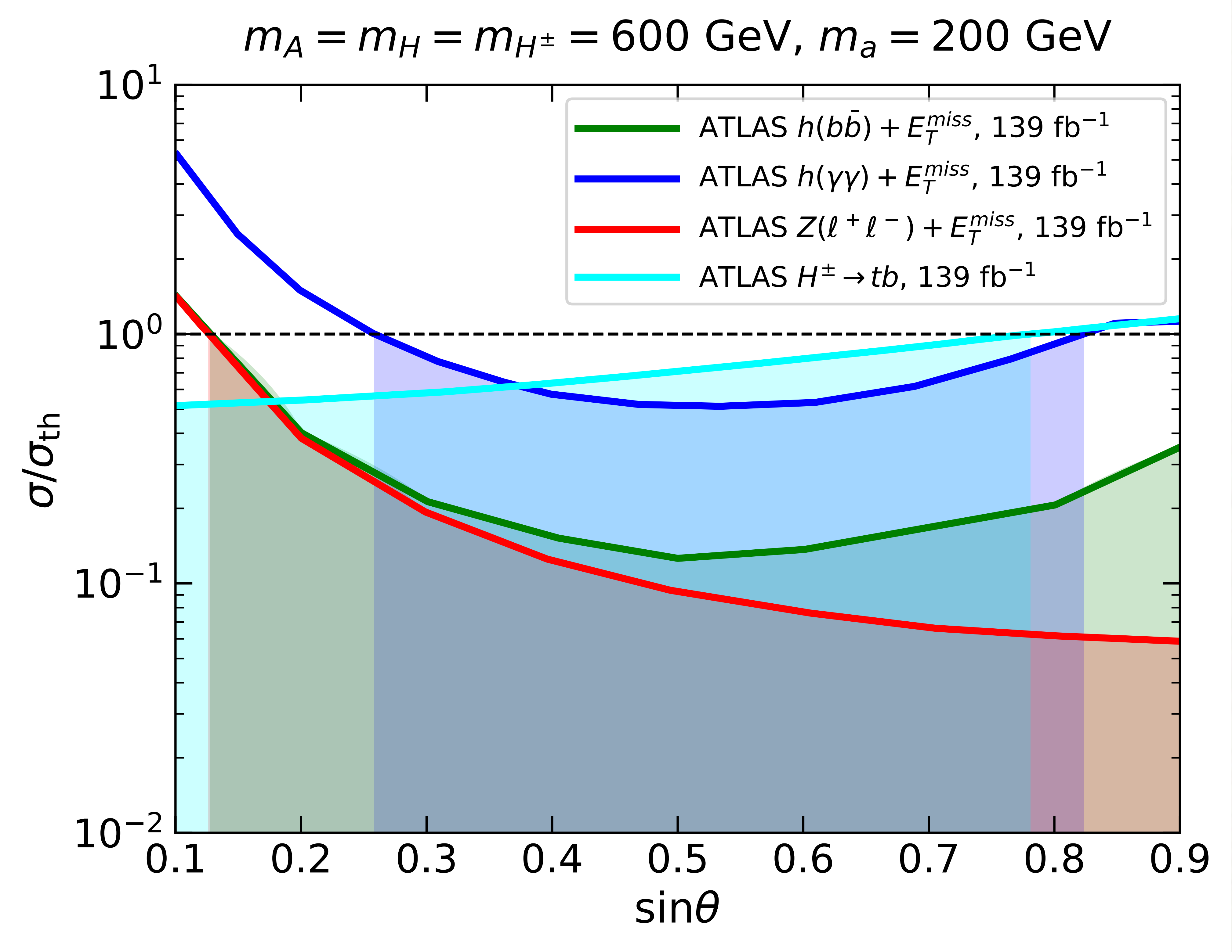


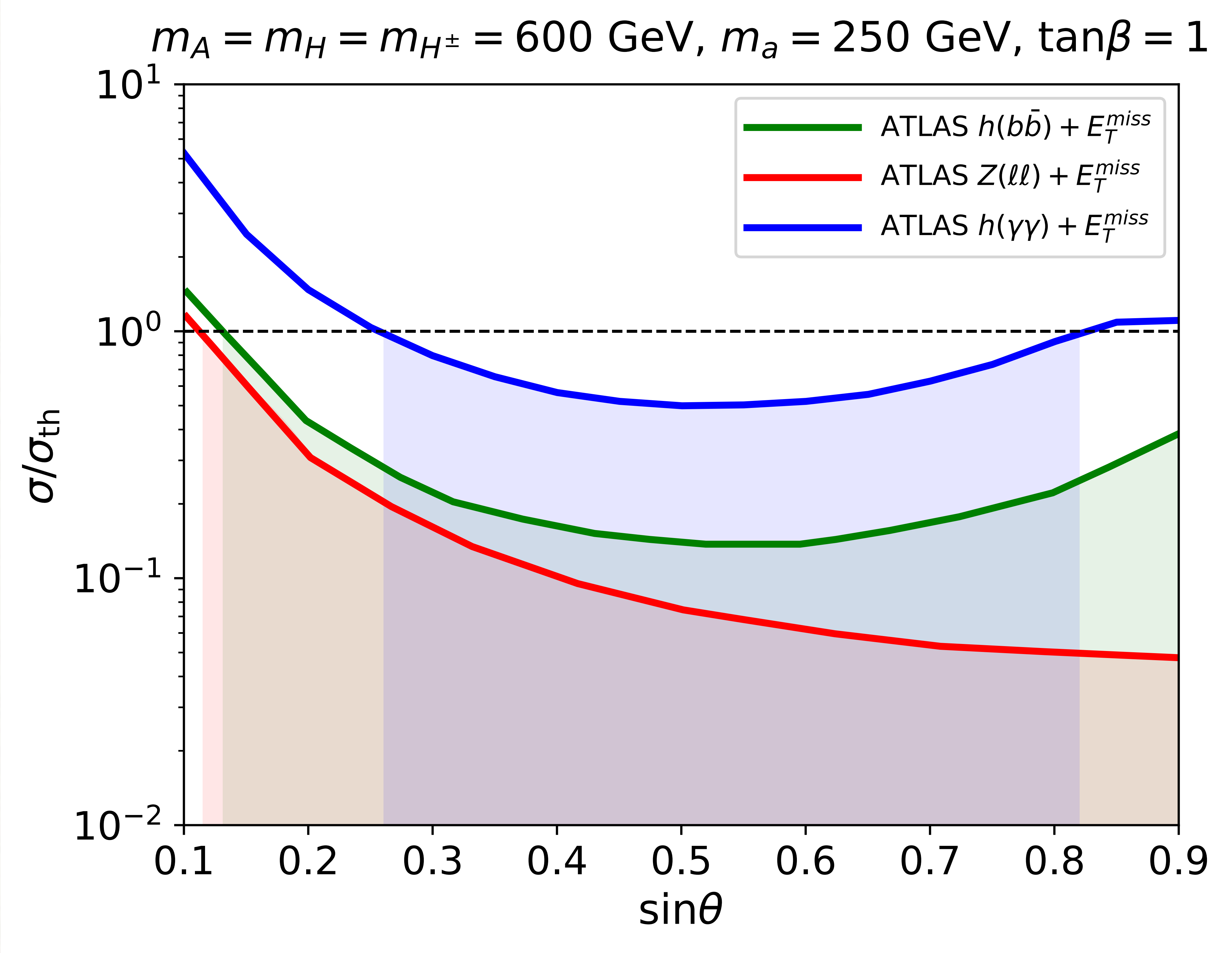
2HDM+a - large mixing



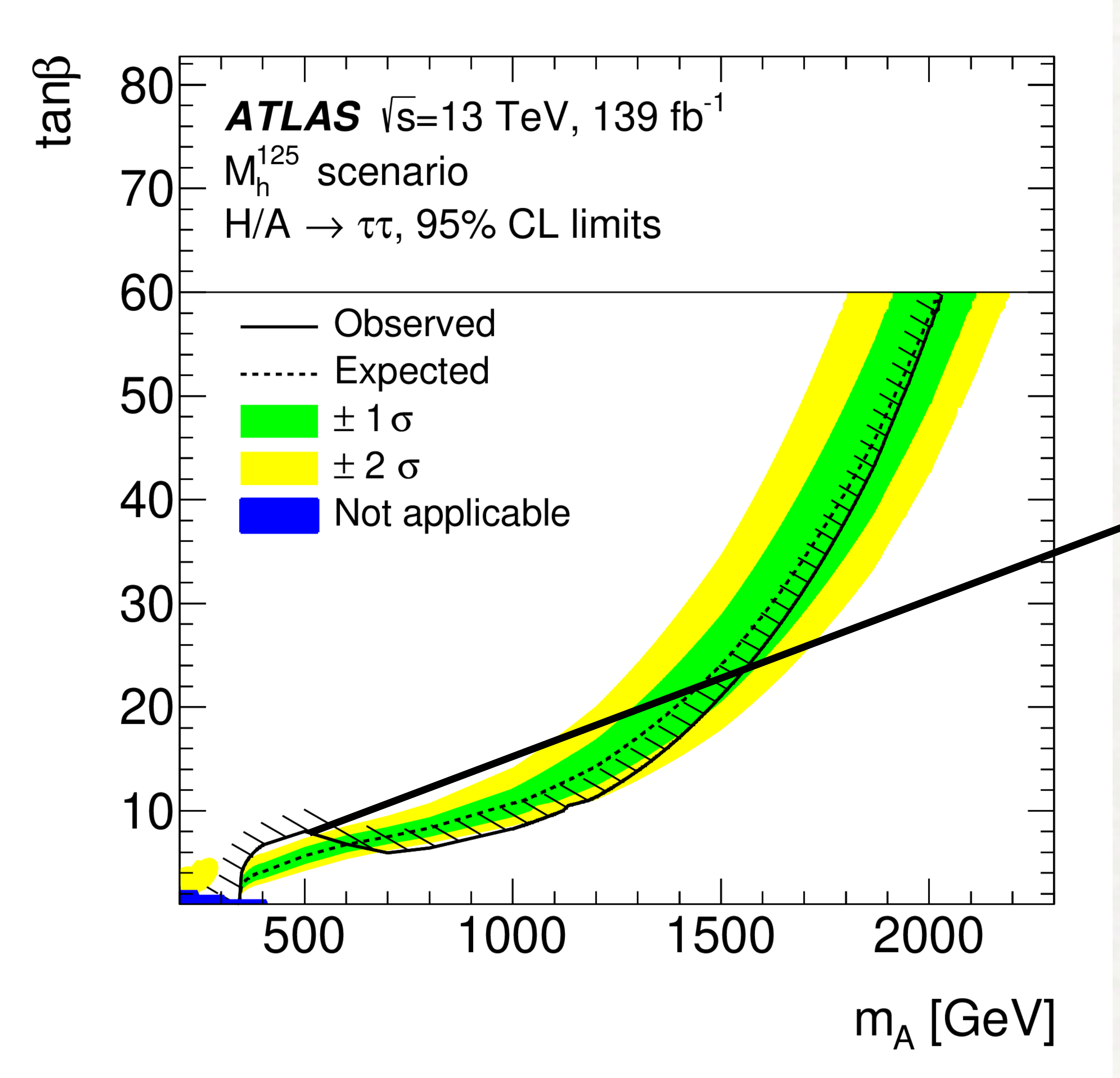


2HDM+a - mixing angle scan

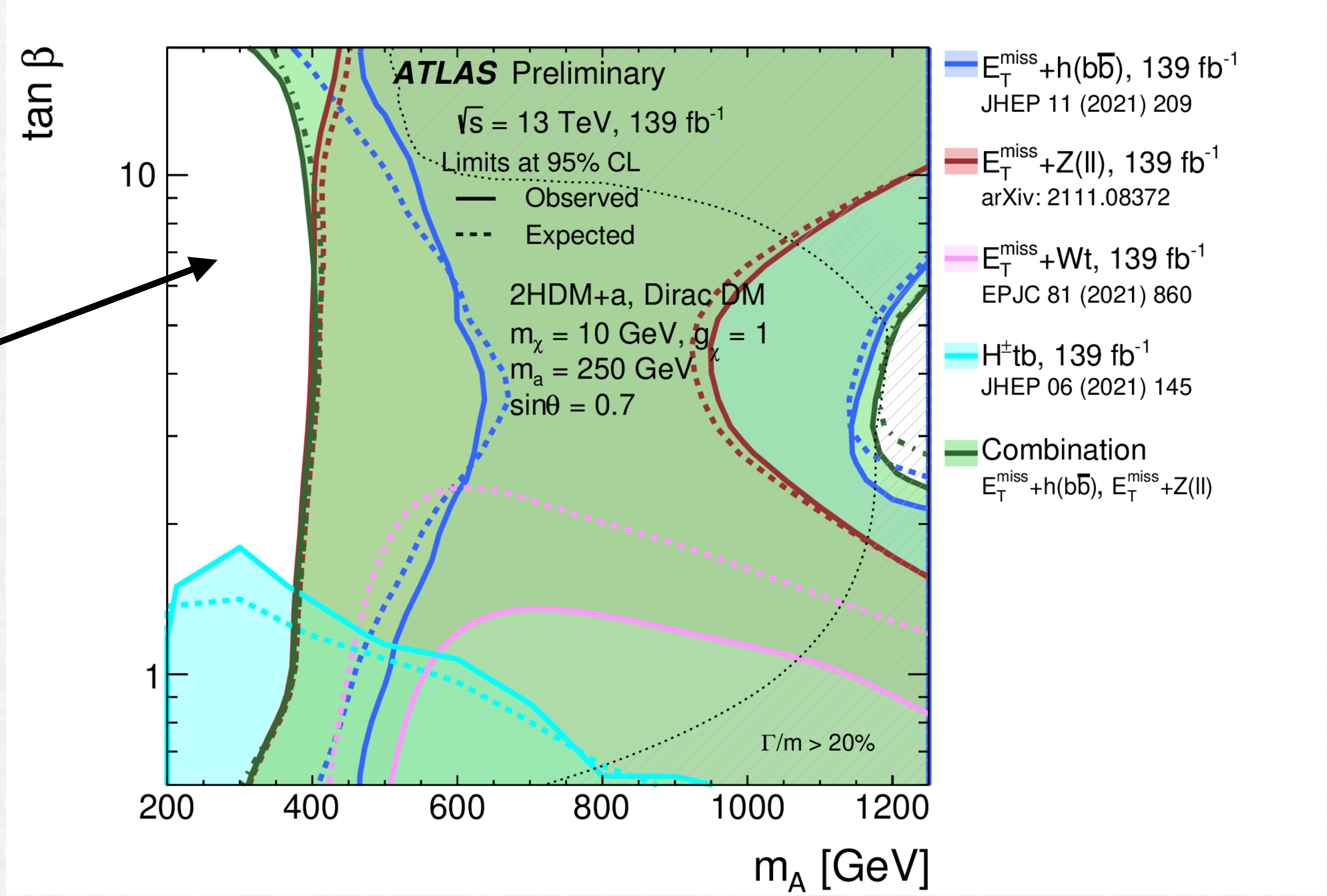




Constraints from taus



[Phys. Rev. Lett. 125 (2020) 051801]



Higgs width constraints

$$\Gamma(h \to aa) = \frac{g_{haa}^2 m_h}{32\pi}\sqrt{1-\frac{4m_a^2}{m_h^2}} \quad > \Gamma_h^{u.l.} = 1.1 \text{ GeV}$$

$$g_{haa} = \frac{1}{m_h v}\left[2\left(m_A^2 - m_a^2 + \frac{m_h^2}{2} - \lambda_3 v^2\right)\sin^2\theta - 2(\lambda_{P1}\cos^2\beta + \lambda_{P2}\sin^2\beta)v^2\cos^2\theta\right]$$

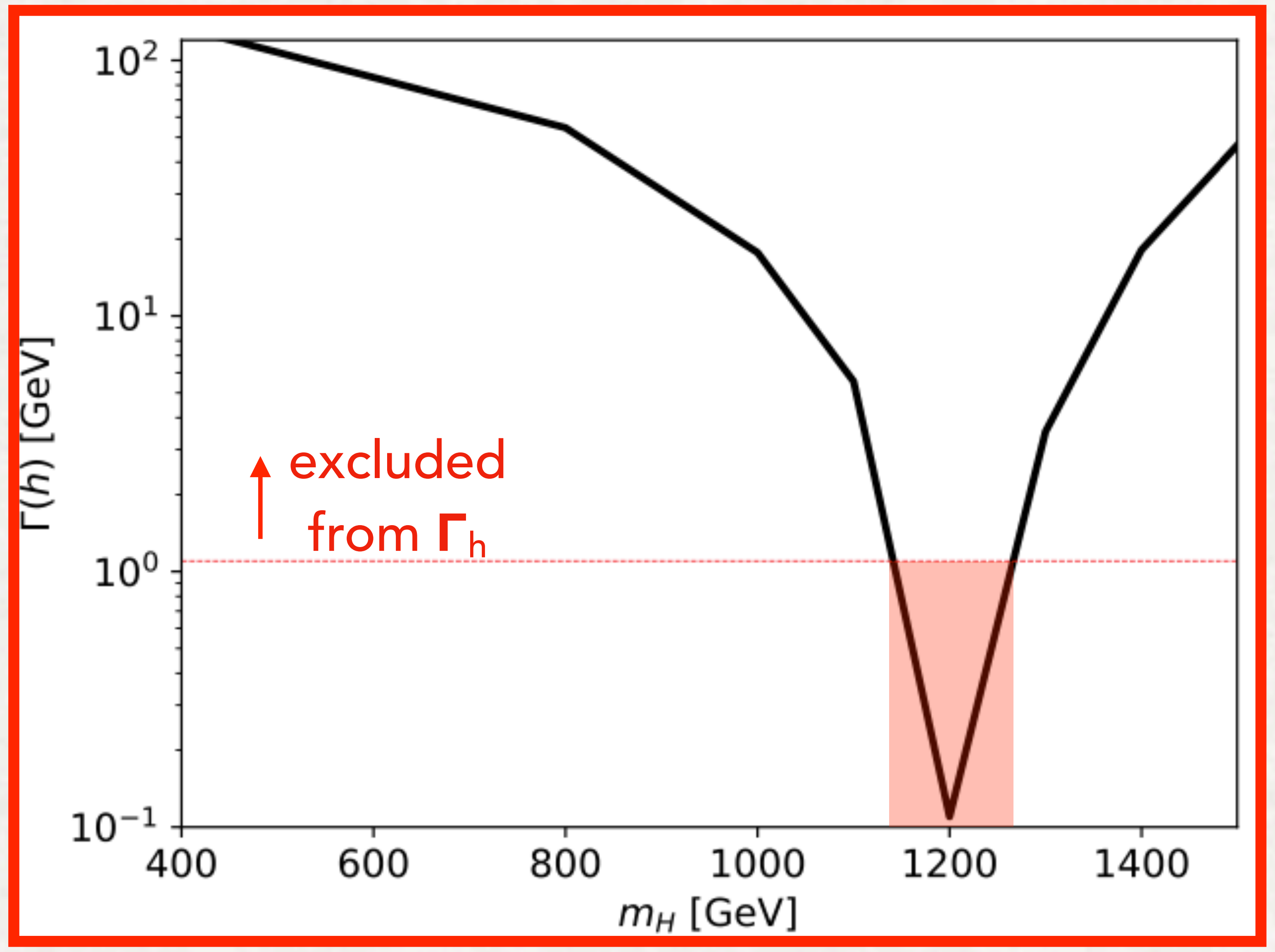
$$= \frac{1}{m_h v}\left[2\left({m_A}^2 - m_a^2 + \frac{m_h^2}{2}\right)\sin^2\theta - 2\lambda_3 v^2\right]$$

$\cos(\beta-\alpha)=0$

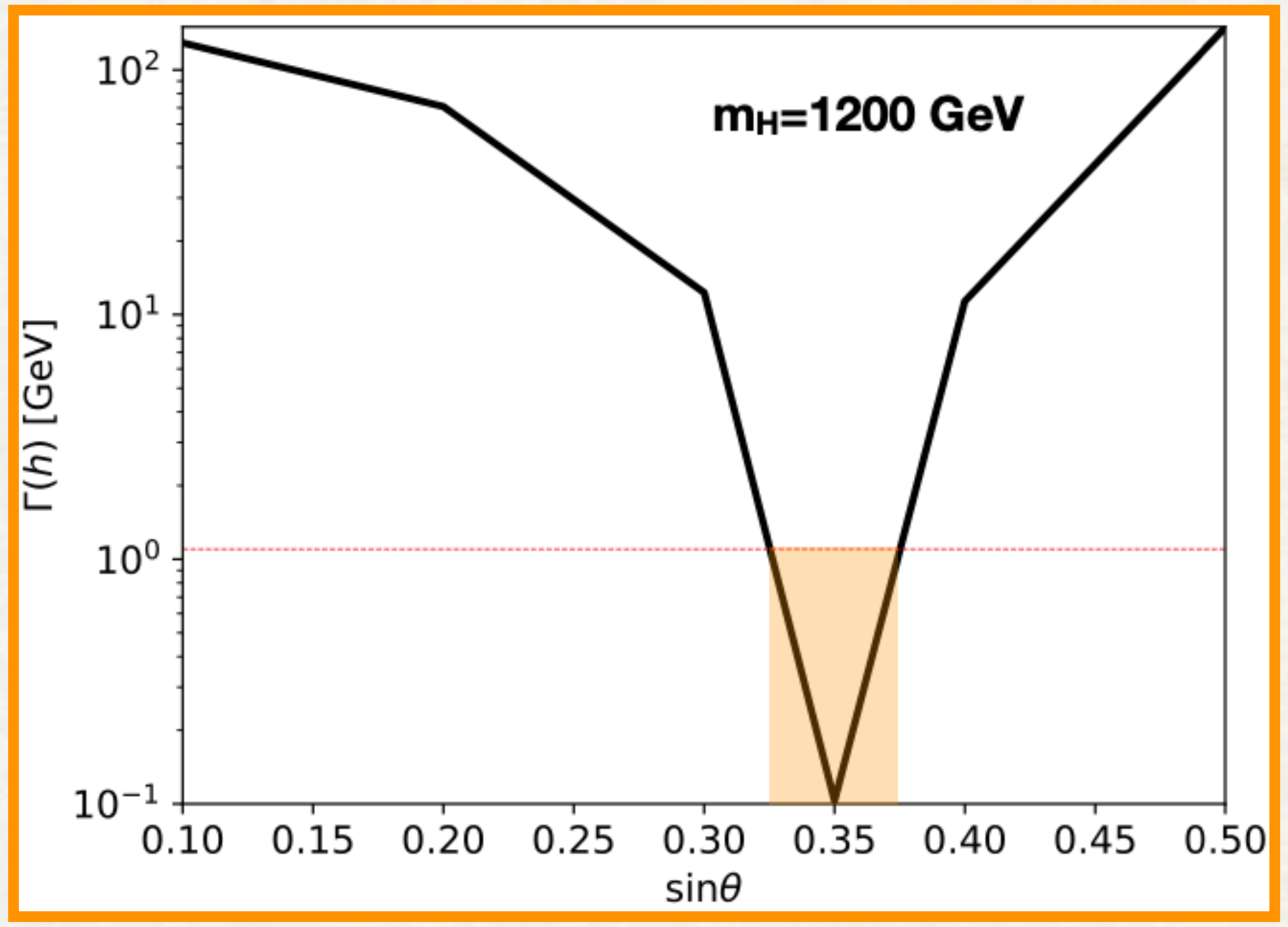
$m_A=m_H=m_{H+}$

$\lambda_3=\lambda_{P1}=\lambda_{P2}$

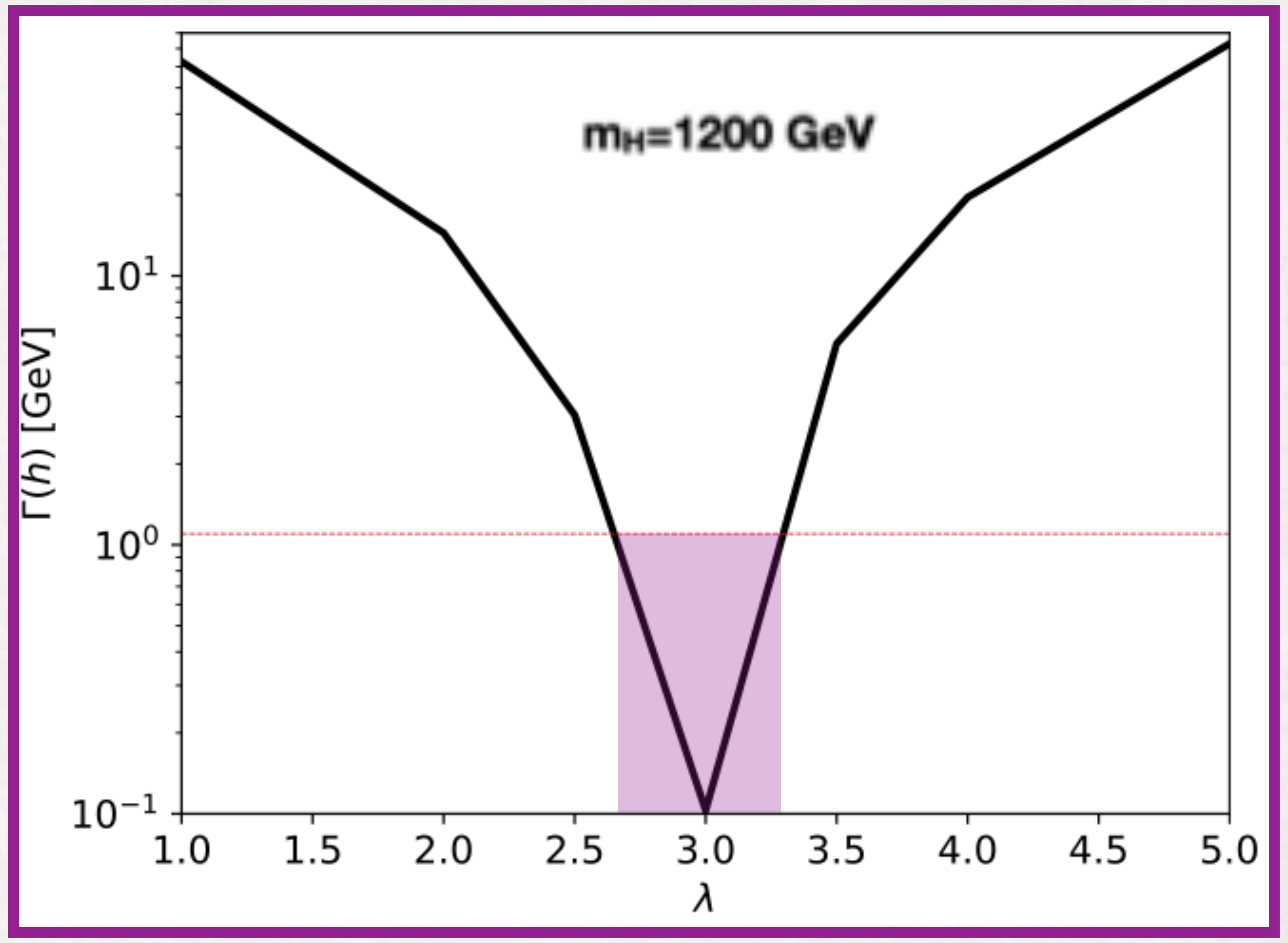
m_A



$\sin\theta$



λ_3



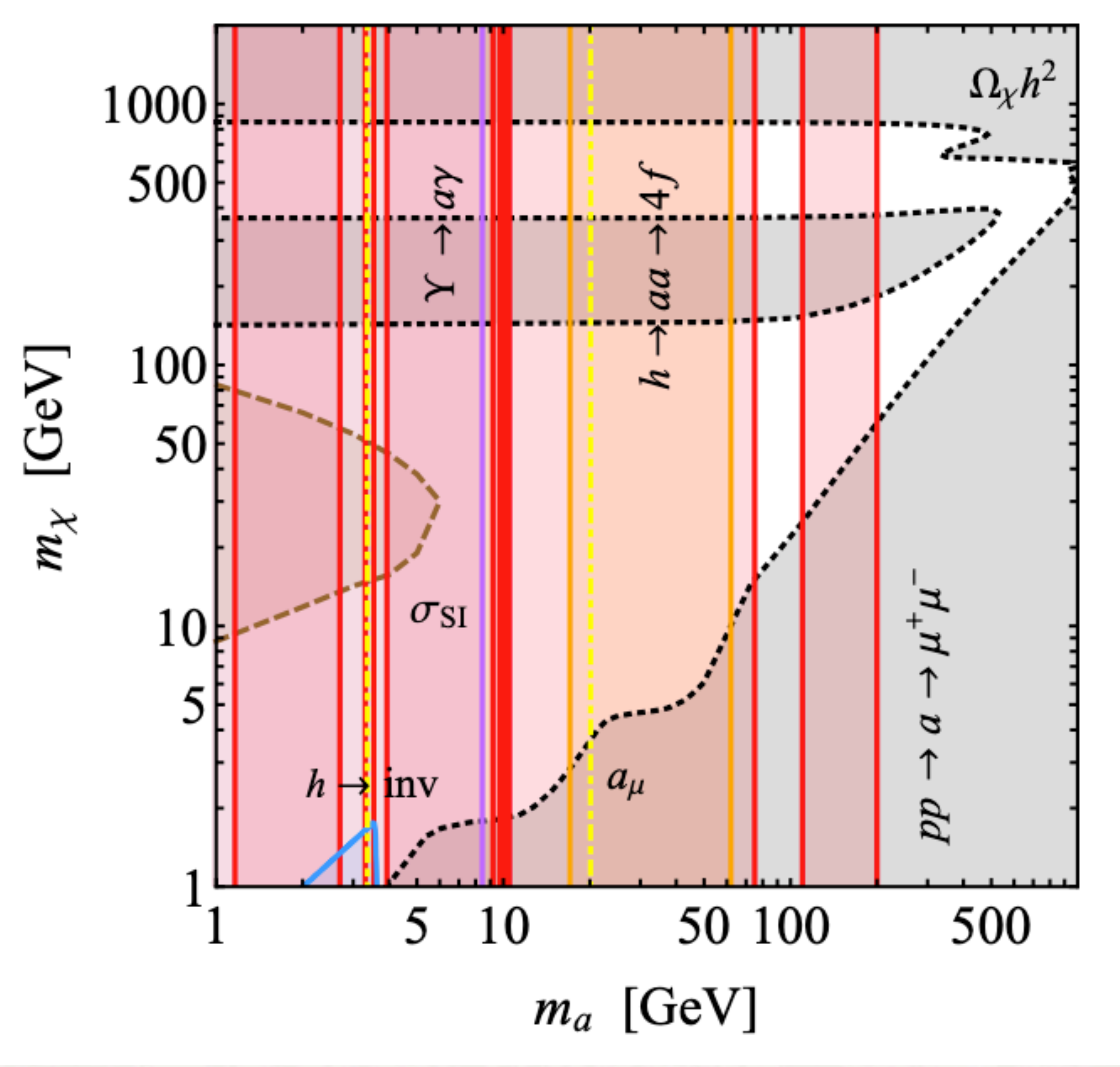
LHC vs g-2

- Original idea from Arcadi, Djouadi & Queiroz (2112.11902) to simultaneously explain DM & muon g-2
- Can also evade constraints from $\boldsymbol{\Gamma}_h$
- Large tan$\boldsymbol{\beta}$ and small m_a needed to get the correct sign for $\boldsymbol{\delta a}_\mu$
- h→4f extend down to very low m_χ because $\boldsymbol{\Gamma}(a\to\chi\chi) \sim y_\chi^2 \cos^2\boldsymbol{\theta} = 0.005$
- h→inv has small BR and MET spectrum very soft so mono-h(bb) has no sensitivity

<u>However</u>

1. g-2 motivated region already ruled out
2. Non-perturbative Haa coupling (g_{Haa} ~ 40) leading to $\Gamma_H > m_H$ over the whole m_a-m_χ plane

$$\{m_A, \tan\beta, \sin\theta, \lambda_3, y_\chi\} = \{1.0\,\text{TeV}, 40, 0.7, 8, 0.1\}$$



SA, Haisch, 2202.12631

2HDM+Z′: coupling scan