UV sensitivity of dark matter production

Oleg Lebedev



University of Helsinki

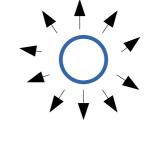
- non-thermal dark matter
- Planck-suppressed operators
- particle production after inflation
- relic abundance

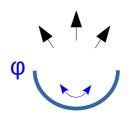
Non-thermal DM has memory !

Production mechanisms (all add up):

- during inflation

- via inflaton oscillations

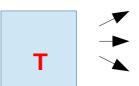




- inflaton decay





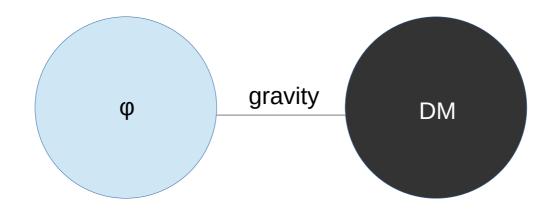


Focus:

inflaton oscillation epoch (preheating)

Assume:

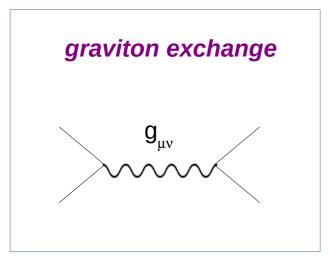
no renormalizable inflaton-dark matter coupling

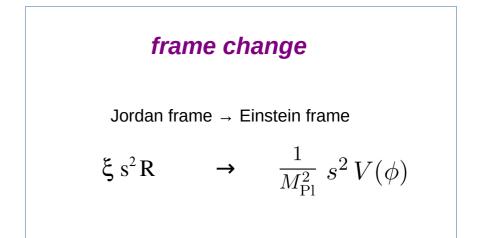


Dim-6 gravity-induced couplings:

$$\Delta \mathcal{L}_6 = \frac{C_1}{M_{\rm Pl}^2} (\partial_\mu \phi)^2 s^2 + \frac{C_2}{M_{\rm Pl}^2} (\phi \partial_\mu \phi) (s \partial^\mu s) + \frac{C_3}{M_{\rm Pl}^2} (\partial_\mu s)^2 \phi^2 - \frac{C_4}{M_{\rm Pl}^2} \phi^4 s^2 - \frac{C_5}{M_{\rm Pl}^2} \phi^2 s^4$$

Simplest examples:





Many more sources: *non-perturbative*, *string states*,...

unknown coefficients On shell, 2 independent dim-6 operators contribute to DM pair-production:

$$\mathcal{O}_3 = rac{1}{M_{
m Pl}^2} \; (\partial_\mu s)^2 \phi^2 \;\;, \;\; \mathcal{O}_4 = rac{1}{M_{
m Pl}^2} \; \phi^4 s^2$$
(supplemented with dim-4 $\mathcal{O}_{
m renorm} = rac{m_\phi^2}{M_{
m Pl}^2} \phi^2 s^2 \;\;$ and 4-DM op $\;\; rac{C_5}{M_{
m Pl}^2} \phi^2 s^4 \;$)

Perturbative DM production rates:

 φ $\phi(t) = \phi_0(t) \cos m_{\phi} t$

$$\begin{array}{l} \mathbf{O_3} \, I \, \mathbf{O_4} & \sim \; \text{energy}^2 \, I \, \phi^2 \, \sim \; \mathbf{m_{\phi}^2} \, I \, \phi^2 \, < < \, 1 \\ \\ \frac{\Gamma \left[\mathcal{O}_3 \right]}{\Gamma \left[\mathcal{O}_4 \right]} \sim \frac{C_3^2}{C_4^2} \, \frac{\omega^4}{\phi_0^4} \quad \sim 10^{-20} \end{array}$$

"graviton exchange" operator much suppressed

E.g.
$$\frac{\tilde{C}_6}{M_{\rm Pl}^4}\phi^6 s^2 + \frac{\tilde{C}_8}{M_{\rm Pl}^6}\phi^8 s^2$$
 more important

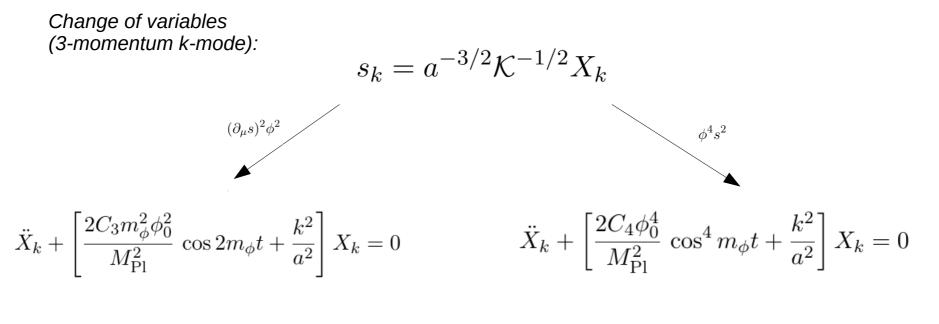
Resonant DM production:

Semiclassical analysis:

$$S = \int d^4x \sqrt{|g|} \left(\frac{1}{2} \mathcal{K}(\phi) g^{\mu\nu} \partial_{\mu} s \,\partial_{\nu} s - \mathcal{V}\right)$$

 $\phi(t) = \phi_0(t) \cos m_{\phi} t$

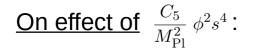
$$\ddot{s} - \frac{1}{a^2} \partial_i \partial_i s + \left(\frac{\dot{\mathcal{K}}}{\mathcal{K}} + 3H\right) \dot{s} + \frac{\mathcal{V}'_s}{\mathcal{K}} = 0$$



resonant production (q >> 1)

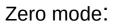
no resonant production $(q \sim 1)$

$$(\partial_{\mu}s)^2\phi^2 \qquad << \qquad \phi^4s^2$$



$$\ddot{s} - \frac{1}{a^2} \partial_i \partial_i s + 3H \,\dot{s} + \frac{4C_5 \phi^2(t)}{M_{\rm Pl}^2} \,s^3 = 0$$

Does not belong to the **Hill type** equations!



25 20

15

0

-5

(z) 5

$$Y'' + \kappa \frac{\cos^2 z}{z^4} Y^3 = 0 \qquad z = m_{\phi}t \quad , \quad \kappa \simeq \frac{14C_5\varphi_0^2}{M_{\rm Pl}^2 a_0^3} \lesssim \mathcal{O}(1)$$

no resonant production for $\kappa \sim 1$

DM abundance via Planck-suppressed operators

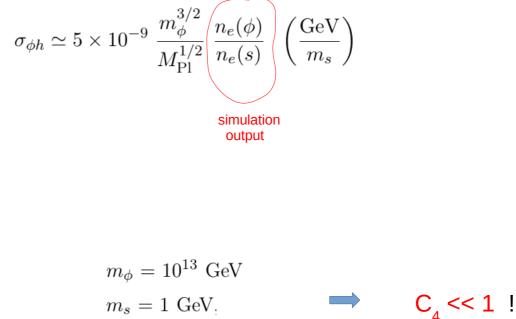
Focus on the most important op

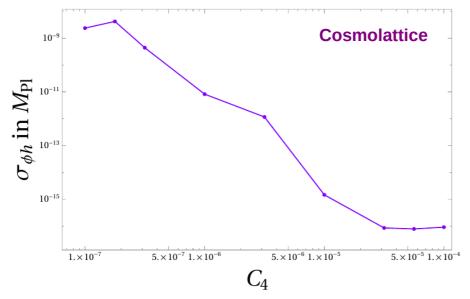
$$rac{C_4}{M_{
m Pl}^2} \phi^4 s^2$$

Reheating:

$$V_{\phi h} = \sigma_{\phi h} \phi H^{\dagger} H \qquad \longrightarrow \qquad \phi \to h h$$

Correct DM abundance requires:





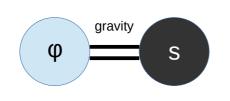
$$\varphi_0 \simeq M_{\rm Pl}$$

Main observation :

Planck—suppressed ("gravity--induced") operators with <u>small</u> Wilson coefficients can account for all of the dark matter !

$$rac{C_4}{M_{
m Pl}^2} \phi^4 s^2$$
 , etc. (dim-8, ...)

Non-thermal DM model building is highly UV sensitive :



- abundance is additive ("memory")
- need to control quantum gravity
- predictivity ?

CONCLUSION

– non-thermal DM production at preheating

- (small) Planck—suppressed operators produce enough DM
- strong UV sensitivity of non-thermal DM (quantum gravity)