

Θέμα 1^ο - Καμπύλες

$$ds^2 = -x^2 dt^2 + dx^2$$

$$x = 1 + \frac{1}{2} \cos \lambda$$

$$t = 1 + \frac{1}{2} \sin \lambda \quad 0 \leq \lambda < 2\pi$$

$$dx = -\frac{1}{2} \sin \lambda d\lambda$$

$$dt = \frac{1}{2} \cos \lambda d\lambda$$

$$ds^2 = - \left(1 + \frac{1}{2} \cos \lambda\right)^2 \left(\frac{1}{2} \cos \lambda d\lambda\right)^2 + \left(-\frac{1}{2} \sin \lambda d\lambda\right)^2$$

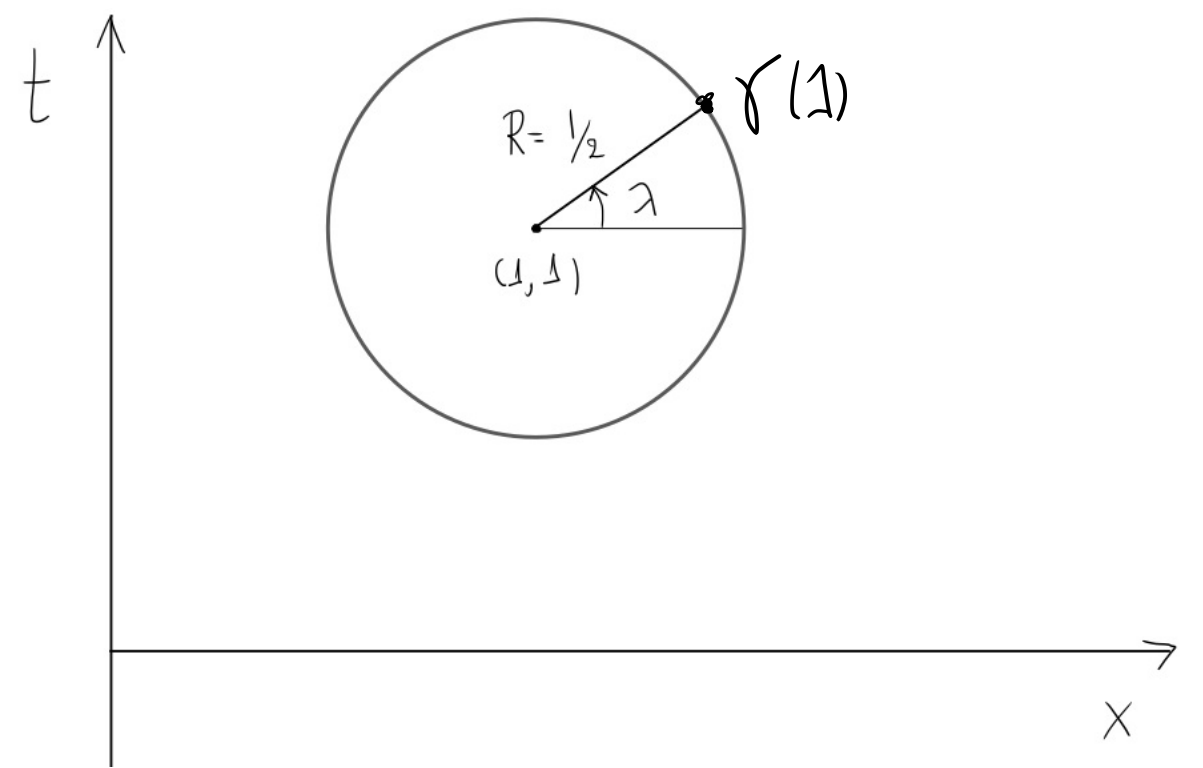
$$= \frac{1}{16} (4 - 8 \cos^2 \lambda - 4 \cos^3 \lambda - \cos^4 \lambda) d\lambda^2 = f(\lambda) d\lambda^2$$

$$S = \int |ds^2|^{1/2} = \int_0^{2\pi} \sqrt{|f(\lambda)|} d\lambda$$

$$\lambda = 0, f(0) = \frac{1}{16} (4 - 8 - 4 - 1) = -\frac{9}{16} < 0 \text{ timelike}$$

$$f\left(\frac{\pi}{3}\right) = \frac{23}{256} > 0 \text{ spacelike}$$

$$f\left(\frac{\pi}{2}\right) = \frac{1}{4} > 0 \text{ spacelike}$$



Θέμα 2 - Το Υπεβολικό Σημείο

$$ds^2 = \frac{1}{y^2} (dx^2 + dy^2) \quad y > 0$$

$$(g_{\mu\nu}) = \begin{pmatrix} \frac{1}{y^2} & 0 \\ 0 & \frac{1}{y^2} \end{pmatrix} \quad (g^{\mu\nu}) = \begin{pmatrix} y^2 & 0 \\ 0 & y^2 \end{pmatrix}$$

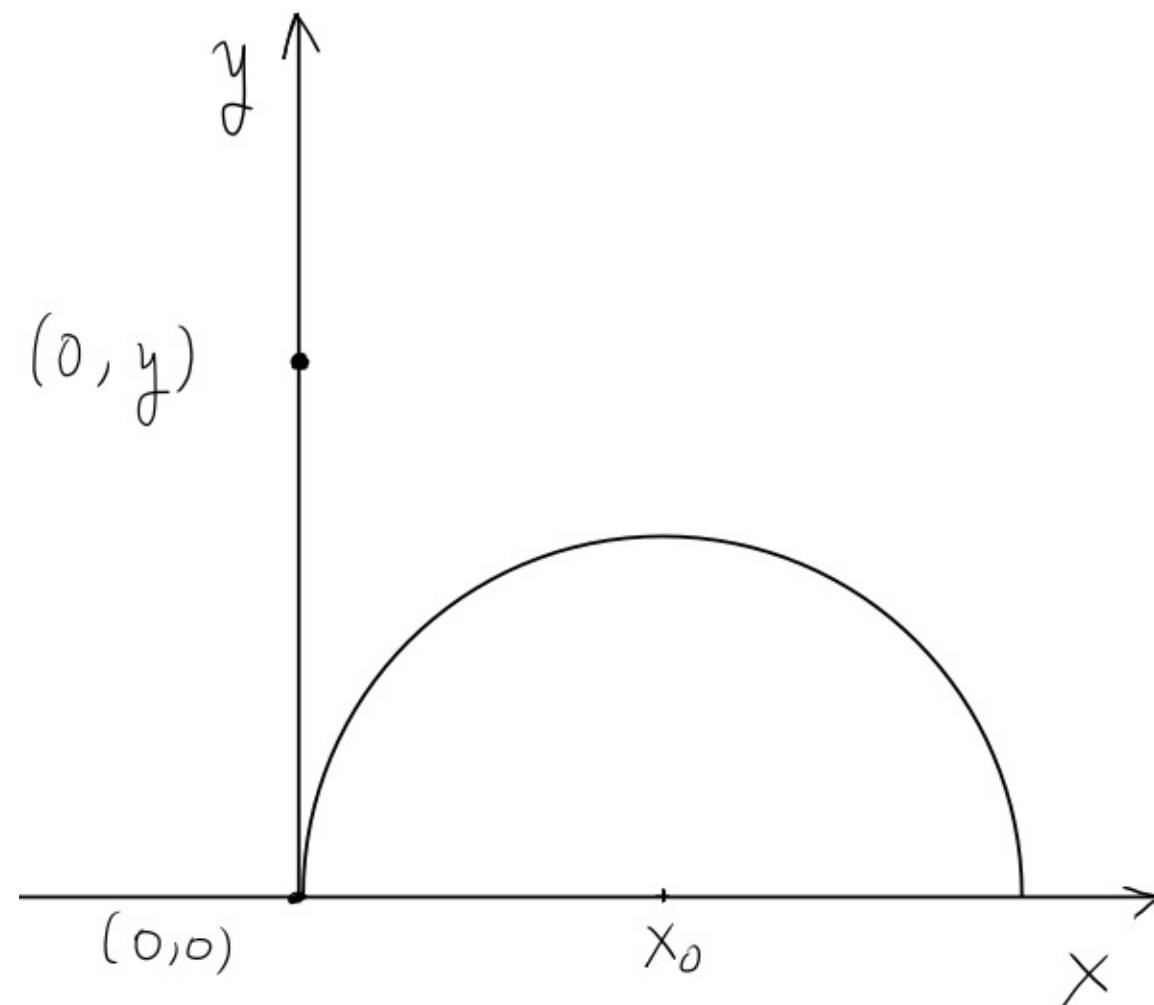
Η μετρική είναι διαγώνια, άρα χρησιμοποιούμε τις

$$g_{\mu\alpha} \Gamma^{\alpha}_{\nu\beta} = \Gamma_{\mu\nu\beta} = \frac{1}{2} \left(\partial_{\nu} g_{\mu\beta} + \partial_{\beta} g_{\mu\nu} - \partial_{\mu} g_{\nu\beta} \right)$$

$$g_{11} \Gamma^1_{21} = \Gamma_{121} = \frac{1}{2} \left(\partial_2 g_{11} + \cancel{\partial_1 g_{12}} - \cancel{\partial_1 g_{21}} \right) \Rightarrow \frac{1}{y^2} \Gamma^1_{21} = \frac{1}{2} \partial_y \left(\frac{1}{y^2} \right) = \frac{1}{2} \left(-\frac{2}{y^3} \right) \Rightarrow \Gamma^1_{21} = -\frac{1}{y}$$

$$g_{22} \Gamma^2_{11} = \Gamma_{211} = \frac{1}{2} \left(\cancel{\partial_1 g_{21}} + \cancel{\partial_1 g_{12}} - \partial_2 g_{11} \right) \Rightarrow \frac{1}{y^2} \Gamma^2_{11} = -\frac{1}{2} \partial_y \left(\frac{1}{y^2} \right) = -\frac{1}{2} \left(-\frac{2}{y^3} \right) \Rightarrow \Gamma^2_{11} = \frac{1}{y}$$

$$g_{22} \Gamma^2_{22} = \Gamma_{222} = \frac{1}{2} \left(\partial_2 g_{22} + \cancel{\partial_2 g_{22}} - \cancel{\partial_2 g_{22}} \right) \Rightarrow \frac{1}{y^2} \Gamma^2_{22} = \frac{1}{2} \partial_y \left(\frac{1}{y^2} \right) = \frac{1}{2} \left(-\frac{2}{y^3} \right) \Rightarrow \Gamma^2_{22} = -\frac{1}{y}$$



Εξισώσεις γεωδαισιολογικών: $\ddot{x}^\mu + \Gamma^\mu_{\nu\rho} \dot{x}^\nu \dot{x}^\rho = 0$ $\dot{x}^\mu = \frac{dx^\mu}{d\tau}$

$\mu=1$ $\ddot{x} + 2\Gamma^1_{21} \dot{y} \dot{x} = 0 \Rightarrow \ddot{x} - 2\dot{x} \frac{\dot{y}}{y} = 0$

$\mu=2$ $\ddot{y} + \Gamma^2_{11} \dot{x} \dot{x} + \Gamma^2_{22} \dot{y} \dot{y} = 0 \Rightarrow \ddot{y} + \frac{1}{y} (\dot{x})^2 - \frac{1}{y} (\dot{y})^2 = 0$

Killing Vector Field (KVF)

Επειδή $\partial_x g_{\mu\nu} = 0$ το $\xi = \partial_x$ είναι KVF

Αρα η ποσότητα $\xi^\mu u_\mu = g_{\mu\nu} \xi^\mu u^\nu$ διατηρείται λόγω της

γεωδαισιολογίας, δηλ. $\frac{d}{d\tau} (\xi^\mu u_\mu) = u^\nu \nabla_\nu (\xi^\mu u_\mu) = 0$. Αλλά

$$\xi^\mu u_\mu = g_{\mu\nu} \xi^\mu u^\nu = g_{11} \xi^1 u^1 = \frac{1}{y^2} \cdot 1 \cdot \dot{x}$$

Ορίζουμε $\mathcal{P}^\mu u_\mu \equiv \frac{1}{R}$ πάνω στην ημισφαίριση, οπότε

$$\frac{1}{f^2} \dot{x} = \frac{1}{R} \Rightarrow \dot{x} = \frac{f^2}{R} \quad (1) \quad \boxed{\dot{x} \neq 0}$$

Αλλά, αν τ είναι affine parameter, τότε

$$u_\mu u^\mu = 1 \Rightarrow g_{\mu\nu} u^\mu u^\nu = 1 \Rightarrow g_{11} (\dot{x})^2 + g_{22} (\dot{y})^2 = 1 \Rightarrow \frac{1}{f^2} (\dot{x})^2 + \frac{1}{f^2} (\dot{y})^2 = 1$$

$$\Rightarrow (\dot{x})^2 + (\dot{y})^2 = f^2 \quad (2)$$

$$(2) \stackrel{(1)}{\Rightarrow} \frac{1}{R^2} f^4 + \dot{y}^2 = f^2 \Rightarrow \dot{y}^2 = f^2 \left(1 - \frac{f^2}{R^2} \right) \Rightarrow \dot{y} = \pm f \left(1 - \frac{f^2}{R^2} \right)^{1/2}$$

$$\text{Τότε: } \frac{dx}{dy} = \frac{dx/d\tau}{dy/d\tau} = \frac{f^2/R}{\pm f \left(1 - \frac{f^2}{R^2} \right)^{1/2}} = \pm \frac{f}{R} \left(1 - \frac{f^2}{R^2} \right)^{-1/2}$$

$$\Rightarrow x = \pm \int dy \frac{y}{R} \left(1 - \frac{y^2}{R^2}\right)^{-1/2} = \pm R \left(1 - \frac{y^2}{R^2}\right)^{1/2} + x_0$$

$$\Rightarrow (x - x_0)^2 + y^2 = R^2 \quad \text{ημικύκλιο κέντρου } (x_0, 0), \text{ ακτίνας } R$$

Αν $\dot{x} = 0 \Rightarrow R = \infty$, και από την $\Delta^{\hat{=}}$ βγαίνει ότι $x = x_0$

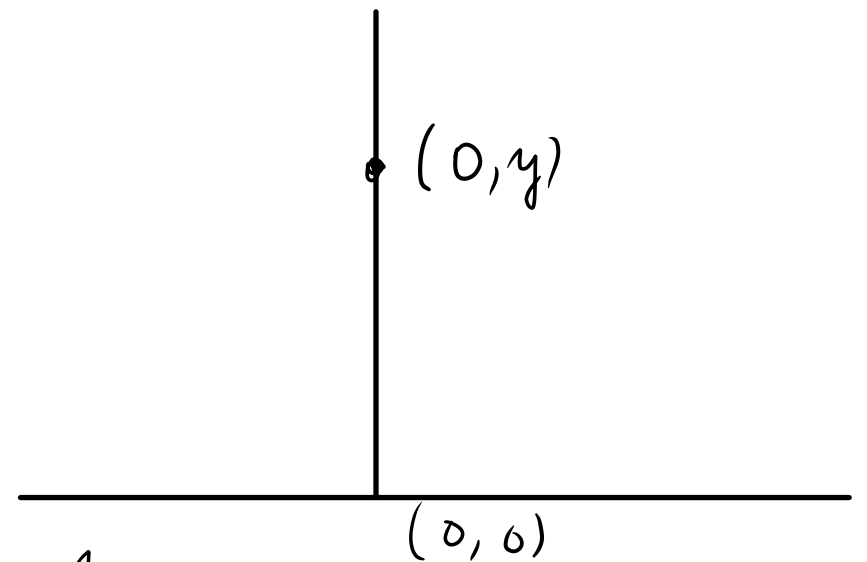
(ημικυκλίες παράλληλες στον άξονα y)

Στο $\Delta \equiv$ εσωτερικά $\hat{=}$ υπάρχει το μήκος της

$$\text{καμπύλης} \quad y = z \quad x = 0$$

$$dy = dz \quad dx = 0$$

$$ds^2 = g_{\mu\nu} dx^\mu dx^\nu = g_{zz} dy dy = \frac{1}{z^2} dz^2 \Rightarrow s = \int_0^y \frac{dz}{z} = \ln z \Big|_0^y = \infty$$



Παρατήρηση: Ολτες οι γεωδαισιαιες που περνουν απο το (x, y) εχουν ανεπει ο μικρος κειχει να "φταει" στα αξονα x : Θεωρετε το ηφικουλο

$$\begin{aligned} x &= x_0 + R \cos z \\ y &= R \sin z \end{aligned} \left. \begin{array}{l} \\ \\ \end{array} \right\} \Rightarrow \begin{array}{l} dx = -R \sin z dz \\ dy = R \cos z dz \end{array} \Rightarrow ds^2 = \frac{1}{y^2} (dx^2 + dy^2) = \frac{1}{R^2 \sin^2 z} (R^2 \sin^2 z dz^2 + R^2 \cos^2 z dz^2) \\ 0 < z < \pi \qquad \qquad \qquad = \frac{1}{\sin^2 z} dz^2$$

Για ορετα μικρο z :

$$s = \int_0^{z_0} \frac{dz}{\sin z} = \ln \tan \frac{z}{2} \Big|_0^{z_0} = +\infty$$

Θέμα 3 - Τανυστική Ενέργεια - Ορμή

$$\left. \begin{aligned} \nabla_\mu A_\nu &= \partial_\mu A_\nu - \Gamma^\rho_{\mu\nu} A_\rho \\ \nabla_\nu A_\mu &= \partial_\nu A_\mu - \Gamma^\rho_{\nu\mu} A_\rho \end{aligned} \right\} \Rightarrow \nabla_\mu A_\nu - \nabla_\nu A_\mu = \partial_\mu A_\nu - \partial_\nu A_\mu - (\Gamma^\rho_{\mu\nu} - \Gamma^\rho_{\nu\mu}) A_\rho = \partial_\mu A_\nu - \partial_\nu A_\mu$$

Αρα, κάτω από μεταβολή $\delta g^{\mu\nu}$ έχω $\delta F_{\mu\nu} = 0$

$$\delta S = \delta \left(\int \sqrt{-g} \left(-\frac{1}{4} g^{\mu\rho} g^{\nu\sigma} F_{\mu\nu} F_{\rho\sigma} \right) \right)$$

$$= -\frac{1}{4} \int \delta \sqrt{-g} F^2 - \frac{1}{4} \int \sqrt{-g} \delta g^{\mu\rho} g^{\nu\sigma} F_{\mu\nu} F_{\rho\sigma} - \frac{1}{4} \int \sqrt{-g} g^{\mu\rho} \delta g^{\nu\sigma} F_{\mu\nu} F_{\rho\sigma}$$

$$= -\frac{1}{4} \int \sqrt{-g} \left(-\frac{1}{2} g_{\mu\nu} \delta g^{\mu\nu} \right) F^2 - \frac{1}{4} \int \sqrt{-g} \delta g^{\mu\nu} g^{\rho\sigma} F_{\mu\rho} F_{\nu\sigma} - \frac{1}{4} \int \sqrt{-g} g^{\sigma\rho} \delta g^{\nu\mu} F_{\sigma\nu} F_{\rho\mu}$$

$\rho \rightarrow \nu \quad \nu \rightarrow \rho$
 $\sigma \rightarrow \mu \quad \mu \rightarrow \sigma$

$\begin{matrix} F_{\rho\mu} \\ -F_{\mu\rho} \end{matrix}$
 $\begin{matrix} F_{\rho\mu} \\ -F_{\mu\rho} \end{matrix}$

$$= -\frac{1}{2} \int \sqrt{-g} \left[g^{\rho\sigma} F_{\mu\rho} F_{\nu\sigma} - \frac{1}{4} g_{\mu\nu} F^2 \right] \delta g^{\mu\nu}$$

$$\Rightarrow T_{\mu\nu} = -\frac{2}{\sqrt{-g}} \frac{\delta S}{\delta g^{\mu\nu}} = F_{\mu\rho} F_{\nu}{}^{\rho} - \frac{1}{4} g_{\mu\nu} F^2$$

Thienvote tvea $g_{\mu\nu} = \eta_{\mu\nu} \Rightarrow F^{0i} = -F_{0i} = \Sigma^i \quad F^{ij} = \epsilon^{ijk} B^k = F_{ij}$

$$T_{00} = F_{0\rho} F_0{}^\rho - \frac{1}{4} \eta_{00} F^2 = F_{0i} F_{0i} + \frac{1}{4} F^2$$

$$F^2 = F_{\mu\nu} F^{\mu\nu} = F_{0i} F^{0i} + F_{i0} F^{i0} + F_{ij} F^{ij} = 2 F_{0i} F^{0i} + F_{ij} F^{ij}$$

$$= 2(-\Sigma_i) \Sigma_i + \epsilon^{ijk} B_k \epsilon_{ijl} B_l = -2 \Sigma_i \Sigma_i + (\delta_{ij} \delta_{kl} - \delta_{jl} \delta_{ki}) B_k B_l$$

$$= -2 \Sigma^2 + (3 \delta_{kl} - \delta_{kl}) B_k B_l = -2 \Sigma^2 + 2 \delta_{kl} B_k B_l = -2(\Sigma^2 - B^2)$$

$$T^{00} = T_{00} = \mathcal{E}^2 - \frac{1}{4} 2(\mathcal{E}^2 - \mathcal{B}^2) = \frac{\mathcal{E}^2 + \mathcal{B}^2}{2}$$

$$T_{0i} = F_{0\rho} F_i{}^\rho - \frac{1}{4} \eta_{0i} F^2 = F_{0j} F_i{}^j = (-\mathcal{E}_j) \epsilon_{ijk} \mathcal{B}_k = -\epsilon_{ijk} \mathcal{E}_j \mathcal{B}_k = -(\vec{\mathcal{E}} \times \vec{\mathcal{B}})_i$$

$$T^{0i} = -T_{0i} = +(\vec{\mathcal{E}} \times \vec{\mathcal{B}})_i$$

$$T_{ij} = F_{i\rho} F_j{}^\rho - \frac{1}{4} \eta_{ij} F^2 = F_{i0} F_j{}^0 + F_{ik} F_j{}^k - \frac{1}{4} \delta_{ij} F^2$$

$$= \mathcal{E}_i (-\mathcal{E}_j) + \epsilon_{ikl} \mathcal{B}_l \epsilon_{jkm} \mathcal{B}_m - \frac{1}{4} \delta_{ij} F^2$$

$$= -\mathcal{E}_i \mathcal{E}_j + (\delta_{lm} \delta_{ij} - \delta_{lj} \delta_{im}) \mathcal{B}_l \mathcal{B}_m - \frac{1}{4} \delta_{ij} F^2$$

$$= -\mathcal{E}_i \mathcal{E}_j + \delta_{ij} (\delta_{lm} \mathcal{B}_l \mathcal{B}_m) - \mathcal{B}_i \mathcal{B}_j + \frac{1}{4} \delta_{ij} 2(\mathcal{E}^2 - \mathcal{B}^2)$$

$$= (-\mathcal{E}_i \mathcal{E}_j + \frac{1}{2} \delta_{ij} \mathcal{E}^2) + (-\mathcal{B}_i \mathcal{B}_j + \frac{1}{2} \mathcal{B}^2) = T^{ij}$$

Θέμα 4: Καμπυλότητα Από δευ έχω $\Gamma^{\alpha}_{\mu\nu} = [\nabla_{\mu}, \nabla_{\nu}](T^{\alpha}_{\beta\gamma} \omega_{\alpha} V^{\beta} W^{\gamma}) = 0$

$$\nabla_{\mu} \nabla_{\nu} T^{\alpha}_{\beta\gamma} \omega_{\alpha} V^{\beta} W^{\gamma} =$$

$$\begin{aligned} & \nabla_{\mu} \nabla_{\nu} T \omega V W + \cancel{\nabla_{\nu} T \nabla_{\mu} \omega} V W + \cancel{\nabla_{\nu} T \omega} \cancel{\nabla_{\mu} V} W + \cancel{\nabla_{\nu} T \omega} \cancel{V} \nabla_{\mu} W \\ & \cancel{\nabla_{\mu} T \nabla_{\nu} \omega} V W + T \cancel{\nabla_{\mu} \nabla_{\nu} \omega} V W + T \cancel{\nabla_{\nu} \omega} \cancel{\nabla_{\mu} V} W + T \cancel{\nabla_{\nu} \omega} \cancel{V} \cancel{\nabla_{\mu} W} \\ & \cancel{\nabla_{\mu} T \omega} \cancel{\nabla_{\nu} V} W + T \cancel{\nabla_{\mu} \omega} \cancel{\nabla_{\nu} V} W + T \omega \cancel{\nabla_{\mu} \nabla_{\nu} V} W + T \omega \cancel{\nabla_{\nu} V} \cancel{\nabla_{\mu} W} \\ & \cancel{\nabla_{\mu} T \omega} \cancel{V} \nabla_{\nu} W + T \cancel{\nabla_{\mu} \omega} \cancel{V} \nabla_{\nu} W + T \omega \cancel{\nabla_{\mu} V} \cancel{\nabla_{\nu} W} + T \omega \cancel{V} \cancel{\nabla_{\mu} \nabla_{\nu} W} \end{aligned}$$

$$\nabla_{\nu} \nabla_{\mu} T^{\alpha}_{\beta\gamma} \omega_{\alpha} V^{\beta} W^{\gamma} =$$

$$\begin{aligned} & \nabla_{\nu} \nabla_{\mu} T \omega V W + \cancel{\nabla_{\mu} T \nabla_{\nu} \omega} V W + \cancel{\nabla_{\mu} T \omega} \cancel{\nabla_{\nu} V} W + \cancel{\nabla_{\mu} T \omega} \cancel{V} \nabla_{\nu} W \\ & \cancel{\nabla_{\nu} T \nabla_{\mu} \omega} V W + T \cancel{\nabla_{\nu} \nabla_{\mu} \omega} V W + T \cancel{\nabla_{\mu} \omega} \cancel{\nabla_{\nu} V} W + T \cancel{\nabla_{\mu} \omega} \cancel{V} \nabla_{\nu} W \\ & \cancel{\nabla_{\nu} T \omega} \cancel{\nabla_{\mu} V} W + T \cancel{\nabla_{\nu} \omega} \cancel{\nabla_{\mu} V} W + T \omega \cancel{\nabla_{\nu} \nabla_{\mu} V} W + T \omega \cancel{\nabla_{\mu} V} \cancel{\nabla_{\nu} W} \\ & \cancel{\nabla_{\nu} T \omega} \cancel{V} \nabla_{\mu} W + T \cancel{\nabla_{\nu} \omega} \cancel{V} \nabla_{\mu} W + T \omega \cancel{\nabla_{\nu} V} \cancel{\nabla_{\mu} W} + T \omega \cancel{V} \cancel{\nabla_{\nu} \nabla_{\mu} W} \end{aligned}$$

$$\Rightarrow [\nabla_\mu, \nabla_\nu](T\omega \vee W) = ([\nabla_\mu, \nabla_\nu]T)\omega \vee W + T([\nabla_\mu, \nabla_\nu]\omega) \vee W + T\omega([\nabla_\mu, \nabla_\nu]V) \vee W + T\omega \vee ([\nabla_\mu, \nabla_\nu]W)$$

$$= 0 \quad \Rightarrow$$

$$([\nabla_\mu, \nabla_\nu]T^\alpha{}_{\beta\gamma})\omega_\alpha V^\beta W^\gamma = +T^\alpha{}_{\beta\gamma} R^\lambda{}_{\alpha\mu\nu}\omega_\lambda V^\beta W^\gamma - T^\alpha{}_{\beta\gamma}\omega_\alpha R^\beta{}_{\lambda\mu\nu}V^\lambda W^\gamma - T^\alpha{}_{\beta\gamma}\omega_\alpha V^\beta R^\gamma{}_{\lambda\mu\nu}W^\lambda$$

$$= (R^\alpha{}_{\lambda\mu\nu}T^\lambda{}_{\beta\gamma} - R^\lambda{}_{\beta\mu\nu}T^\alpha{}_{\lambda\gamma} - R^\lambda{}_{\gamma\mu\nu}T^\alpha{}_{\beta\lambda})\omega_\alpha V^\beta W^\gamma$$

$$\Rightarrow [\nabla_\mu, \nabla_\nu]T^\alpha{}_{\beta\gamma} = R^\alpha{}_{\lambda\mu\nu}T^\lambda{}_{\beta\gamma} - R^\lambda{}_{\beta\mu\nu}T^\alpha{}_{\lambda\gamma} - R^\lambda{}_{\gamma\mu\nu}T^\alpha{}_{\beta\lambda}$$