

① Compute  $R_{\hat{\mu}\hat{\nu}\hat{\rho}\hat{\sigma}}$  in the free falling astronaut problem when  $r < 2M$  and show that the results are the same as for the  $r > 2M$  case.

② Consider the Killing vector field  $\xi = \partial_t$  of the Schwarzschild geometry. Show that at the horizon  $r = 2M$

$$D_{\xi} \xi^{\mu} = \xi^{\nu} \nabla_{\nu} \xi^{\mu} = \kappa \xi^{\mu},$$

and compute  $\kappa$ .  $\kappa$  is the "surface gravity" of the black hole (Use the Eddington-Finkelstein coordinates for the computation)

( $\Gamma^{\mu}_{\nu\rho}$  for E-F coordinates can be found in the Mathematica notebook p Lecture 10)

A stationary observer is hovering at  $(t, r, \theta, \phi) = (t, R, \frac{\pi}{2}, 0)$ , over a Schwarzschild black hole. Her 4-velocity is  $u^\mu$ , and her acceleration  $a^\mu = u^\nu \nabla_\nu u^\mu = \frac{Du^\mu}{d\tau}$ . If  $\xi = \partial_t$  is the timelike Killing vector field, then:

- Compute  $a^\mu$  and  $a_\mu$  in the  $(t, r, \theta, \phi)$  and the  $(v, r, \theta, \phi)$  coordinate systems (E-F: Eddington-Finkelstein coordinates)

- Compute the scalar  $a^2 = a_\mu a^\mu$  and show that  $a^\mu$  is a spacelike vector for  $r > 2M$

- Compute  $V = (-\xi_\mu \xi^\mu)^{1/2}$ , and show that  $a \cdot V = \kappa = \begin{pmatrix} \text{surface} \\ \text{gravity} \end{pmatrix}$

- Compute  $(\nabla_\mu \xi_\nu)(\nabla^\mu \xi^\nu)$  using E-F coordinates, and show that, at

$$r = 2M, \quad \kappa = -\frac{1}{2} (\nabla_\mu \xi_\nu)(\nabla^\mu \xi^\nu)$$

- Show that  $\eta^\mu = \xi^\nu \nabla_\nu \xi^\mu$  is spacelike for  $r > 2M$ . Is this in contradiction

with  $\xi^\nu \nabla_\nu \xi^\mu = k \xi^\mu$ , since for  $r > 2M$   $\xi^\mu$  is timelike?

③ A particle of rest mass  $m$  is initially at rest relative to a stationary observer  $O_1$  at  $r = \infty$  from a Schwarzschild black hole. The particle is left to fall freely towards the center of the black hole. Another stationary observer  $O_2$  is hovering above the black hole at  $(t, r, \theta, \phi) = (t, R, \frac{\pi}{2}, 0)$ ,  $R = \text{const} > 2M$ .

- What is the speed of the particle as measured by  $O_2$ , when the particle goes through her lab?
- The particle is converted to radiation at  $O_2$ , which is

sent back to  $r = \infty$ . What is the total energy of that radiation, as measured by stationary detectors at  $r = \infty$ ?

- Suppose that when the particle reaches  $O_2$ , it is brought to rest in the  $O_2$ 's lab, and the excess energy is converted to radiation, which is sent back to  $r = \infty$ . What is the energy of the radiation as measured by stationary detectors at  $r = \infty$ ?

④ Rindler Coordinates: Consider the  $(t, r, \theta, \varphi)$  coordinates of the Schwarzschild metric, and make the transformation to the  $(t, \zeta, \theta, \varphi)$  coordinates, where

$$r - 2M = \frac{\zeta^2}{8M}$$

Show that the metric in the new coordinates is given by

$$ds^2 = - \frac{k^2 \zeta^2}{k^2 \zeta^2 + 1} dt^2 + (k^2 \zeta^2 + 1) d\zeta^2 + \frac{1}{4k^2} (k^2 \zeta^2 + 1)^2 d\Omega^2, \text{ where}$$

$$k = \frac{1}{4M}, \quad d\Omega^2 = d\theta^2 + \sin^2 \theta d\varphi^2$$

- Show that for  $r \approx 2M$

$$ds^2 \approx - \kappa^2 \mathcal{I}^2 dt^2 + d\mathcal{I}^2 + \frac{1}{4\kappa^2} d\Omega^2$$

- Consider the coordinate transformation

$$T = \mathcal{I} \sinh(\kappa t)$$

$$X = \mathcal{I} \cosh(\kappa t)$$

and show that the approximate metric above is flat

• Orbital speed: Consider a hovering stationary observer  $O_1$  at  $(t, r, \theta, \phi) = (t, R, \frac{\pi}{2}, 0)$ ,  $R = \text{const} > 2M$ , over a Schwarzschild black hole. A particle falls freely, moving on a  $r = R$  circular trajectory, and goes through the lab of  $O_1$ . Calculate the speed  $v(R)$  of the particle, as measured by  $O_1$ .