

# Double Beta Decay Experiments – Focusing on AMoRE project

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Institute for Basic Science

2021. 9. 7.

Corfu, “Workshop on the Standard Model and Beyond”

# I will talk about

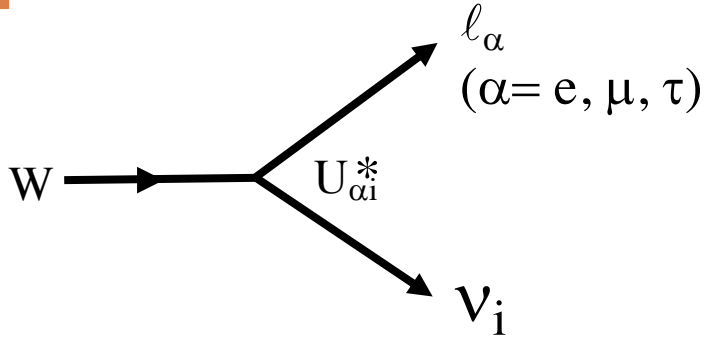
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1. Motivation of Neutrinoless DBD
2. Current status & Comparisons
3. Status of AMoRE experiment

# Status of Neutrino mixing & oscillation

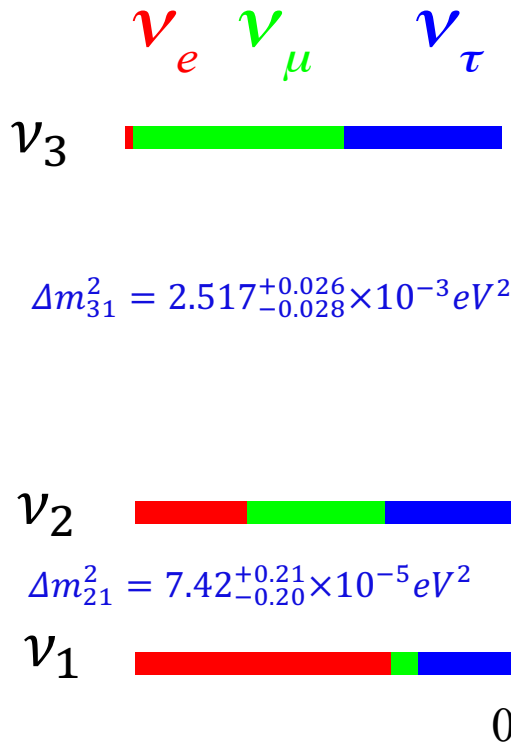
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## Parameters

- 3 mixing angles.
- 2 mass differences
- 1 Dirac CP phase
- + 2 Majorana phase (if Majorana neutrinos)

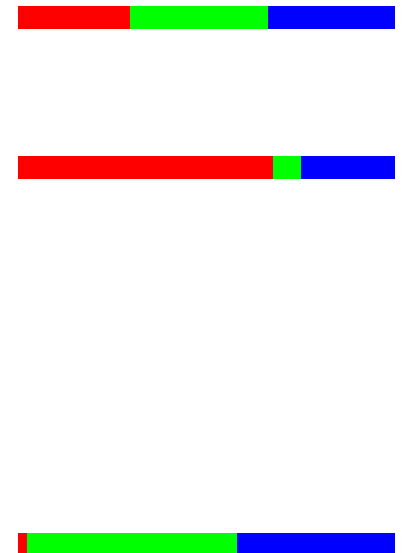
## Normal Ordering



$m$

	NO
$\sin^2 \theta_{12}$	$0.304_{-0.012}^{+0.012}$
$\sin^2 \theta_{23}$	$0.573_{-0.020}^{+0.016}$
$\sin^2 \theta_{13}$	$0.02219_{-0.00063}^{+0.00062}$
$\delta^{CP}$	$197_{-24}^{+27}$

## Inverted Ordering



$$P(\nu_\alpha \rightarrow \nu_\alpha) = 1 - \sin^2 2\theta \cdot \sin^2 \left( \frac{\Delta m^2 L}{E_\nu} \right)$$

# Experiments for neutrino mass

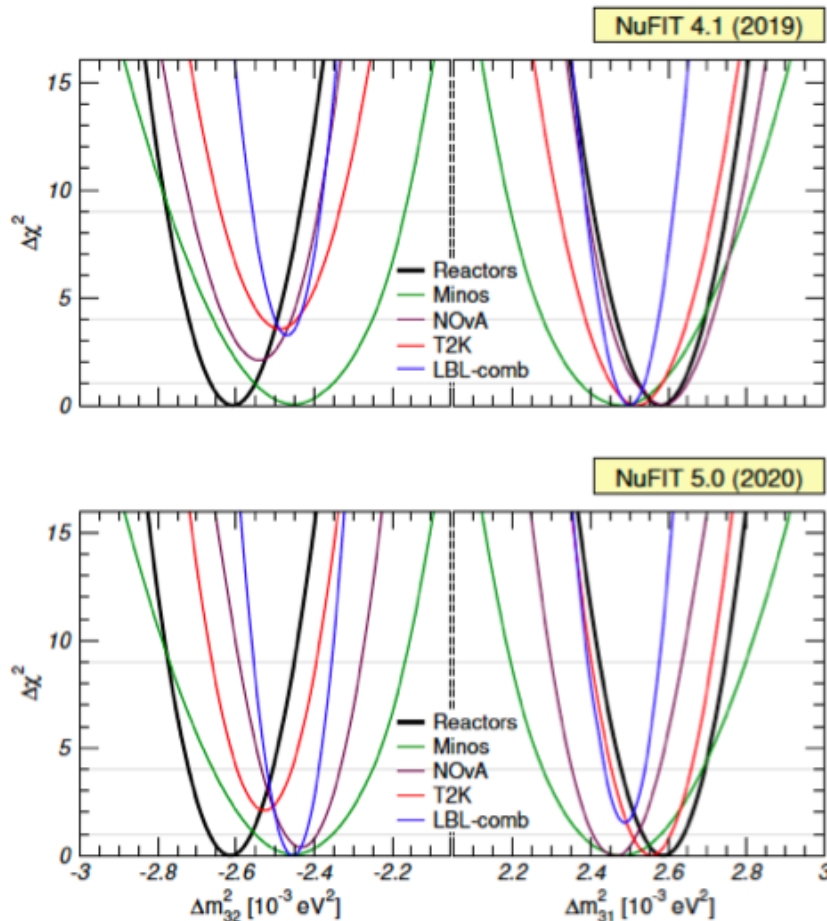
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- Mass Ordering ?
- Mass limits from single beta decay
- Are neutrinos and antineutrinos the same particles ?
- Mass limits from  $0\nu\beta\beta$  experiments.
- Neutrino mass from cosmology observation
- Are there sterile neutrinos ?



# Updates for mass ordering in 2020

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Esteban et al., JHEP09 2020, 178

- The best fit remains for the normal mass ordering, however, with reduced significance.
- Without SK-atm, inverted ordering is disfavored only with  $\Delta\chi^2 = 2.7(1.6\sigma)$  to be compared with  $\Delta\chi^2 = 6.2(2.5\sigma)$  in 2019.
- This change is driven by the new LBL results from T2K and NOvA.
- For experimentalists, we need to assume Normal Ordering.

Eligio Lisi @ Neutel 2021

Cosmo data can add from  $\sim 0$  to  $\sim 0.7\sigma$  to the  $\sim 2.7\sigma$  oscillation preference for NO

→ overall  $\sim 2.7\text{--}3.4\sigma$  hint for NO vs IO

# Motivation of $0\nu\beta\beta$ - Most promising BSM Physics !

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SM has only left(right)-handed massless (anti)neutrinos.

## Dirac Neutrino Masses

$$L_D = -m_D(\overline{\nu}_R\nu_L + \overline{\nu}_L\nu_R)$$

- Lepton # is conserved.  $y^\nu$ : Yukawa Coupling  $\sim 10^{-12}$
- Higgs mechanism needs right-handed neutrinos,  $\nu_R$ .

## Majorana Masses

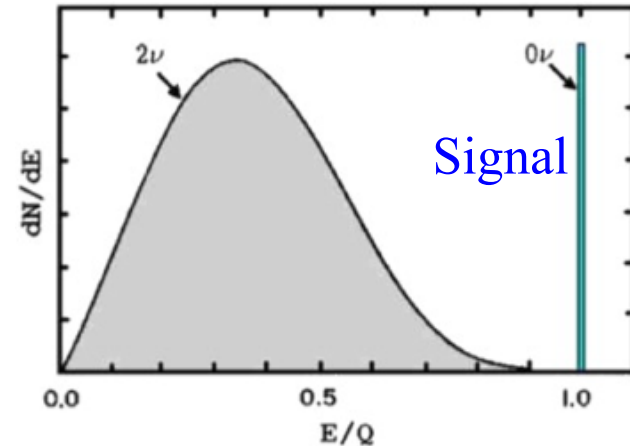
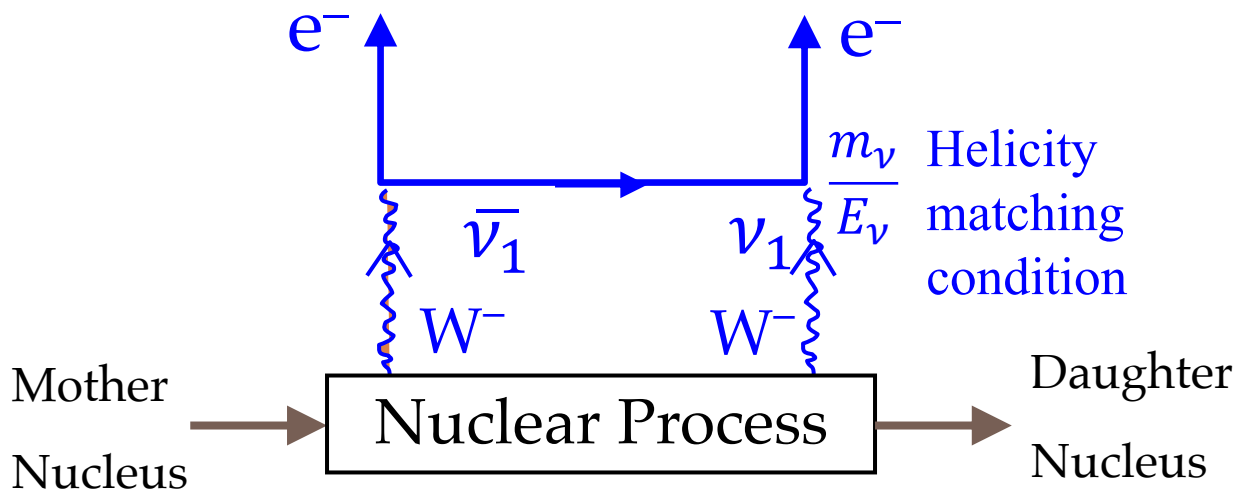
- “right-handed Majorana mass term” can be ;  $L_R = -m_R/2[(\overline{\nu}_R)^c\nu_R + \overline{\nu}_R(\nu_R)^c]$
- $(\nu)^c = \nu \rightarrow$  Majorana particle (No L# is needed)
- See-Saw Mechanism gives two Majorana mass eigenstates,

$$m_1 \simeq \frac{m_D^2}{m_R}$$
$$m_2 \simeq m_R$$

# How to test if neutrinos are Majorana particles ?

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- Seek Neutrinoless Double Beta decay ( $0\nu\beta\beta$ )



- 1939, Furry already suggested to search  $0\nu\beta\beta$  to check Majorana's theory. Furry PR56, 1184(1939)

## More quantitatively...

- $T_{1/2}$  of  $0\nu\beta\beta$  decay are inversely proportional to effective  $0\nu\beta\beta$  neutrino mass,  $m_{\beta\beta}$ , which is a function of neutrino masses, mixing angles, and Majorana phases.

for light neutrino exchange model.

$$\begin{array}{c}
 \text{Phase} \\
 \text{factor} \\
 \left[ T_{1/2}^{0\nu} \right]^{-1} = G_{0\nu} \underbrace{|M_{0\nu}|^2}_{\text{Nuclear Matrix Element}} \underbrace{\left( \frac{m_{\beta\beta}}{m_e} \right)^2}_{\text{Neutrino Mass}}
 \end{array}$$

**Half-life Measured**
**Nuclear Matrix Element**
**Neutrino Mass**

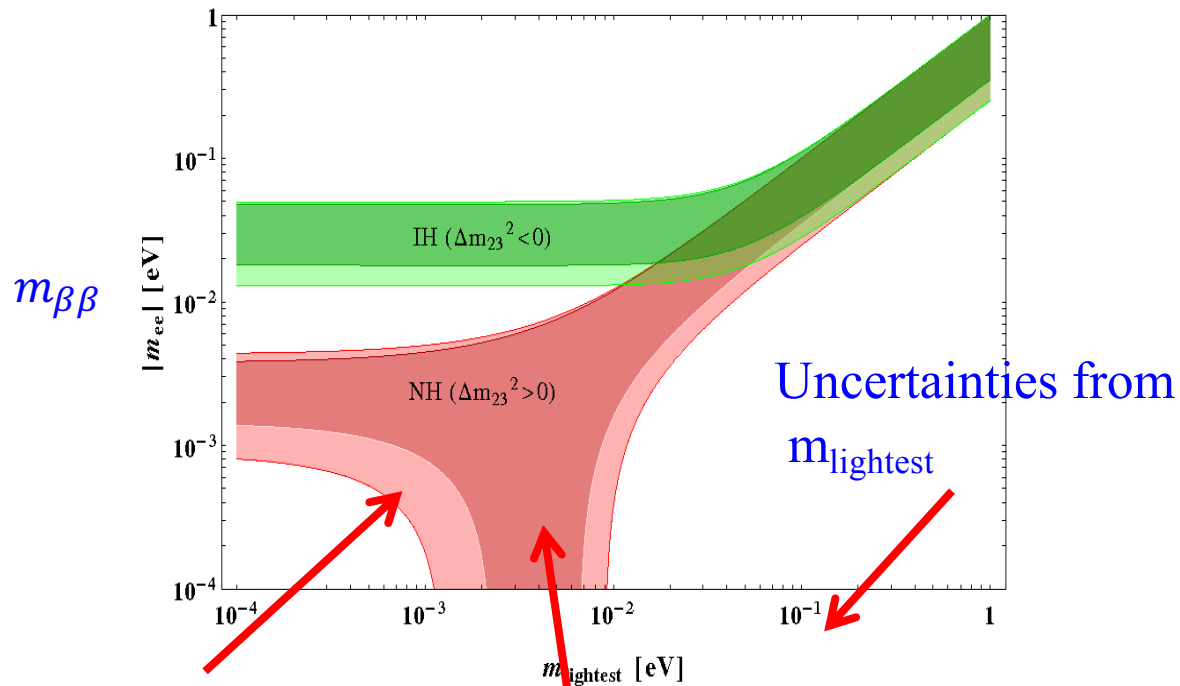
effective  $0\nu\beta\beta$  neutrino mass is ;

$$m_{\beta\beta} = \left| \sum_{k=1}^3 U_{ek}^2 m_k \right| = \left| c_{13}^2 c_{12}^2 e^{2i\eta_1} m_1 + c_{13}^2 s_{12}^2 e^{2i\eta_2} m_2 + s_{13}^2 e^{-2i\delta} m_3 \right|$$

$$T_{1/2}^{0\nu} \rightarrow m_{\beta\beta}$$

# Expectation on effective $0\nu\beta\beta$ mass

for light neutrino exchange model.

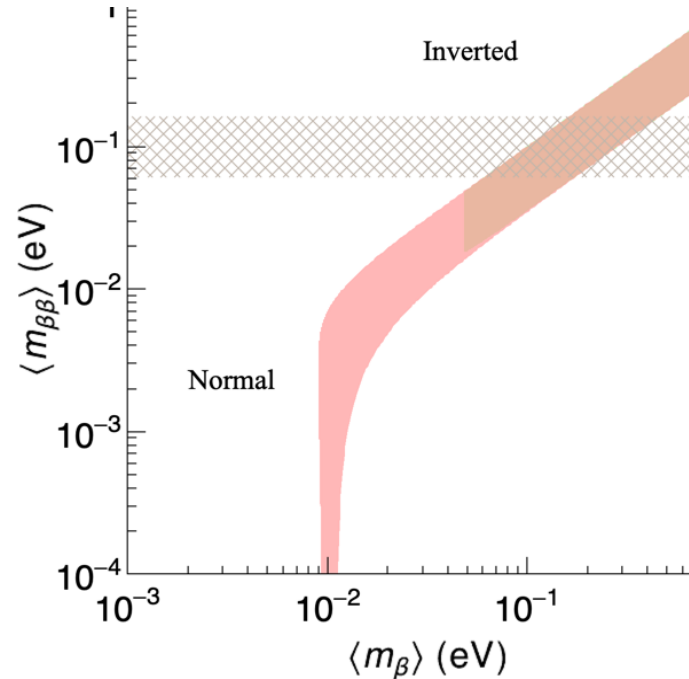
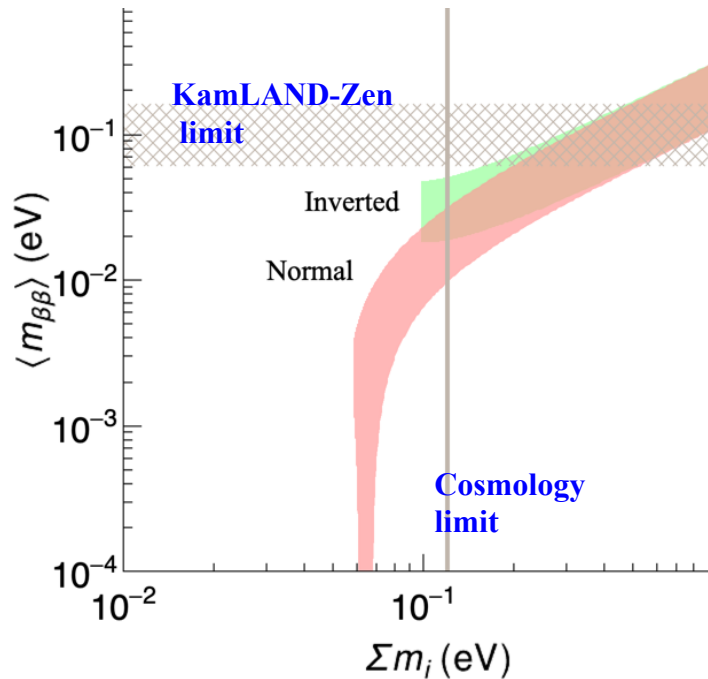
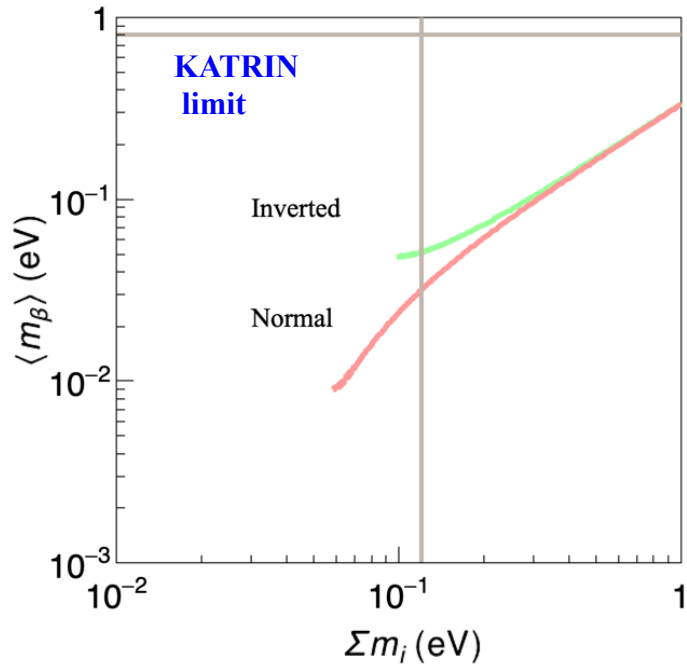


Uncertainties from mixing angles

Uncertainties from phases

# Current Mass Limits

- Neutrino mass is ultra small, and we don't understand its origin. It is related to if neutrinos are Majorana particles.
- Neutrino mass is constrained by beta decays and cosmology.

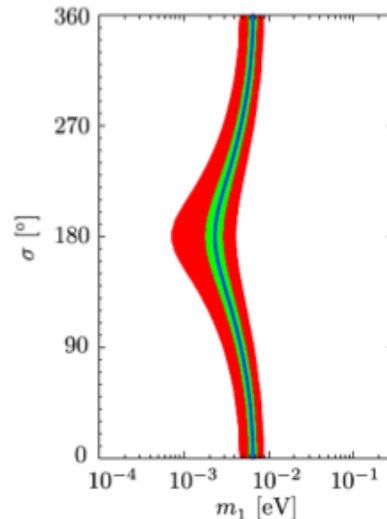
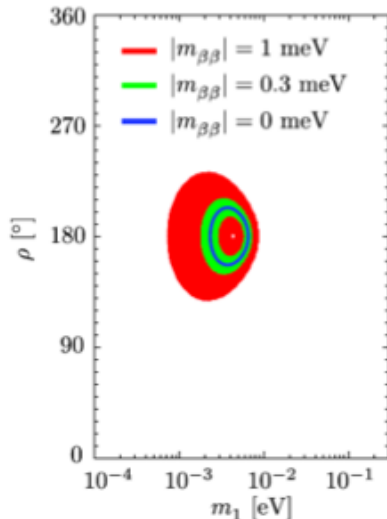
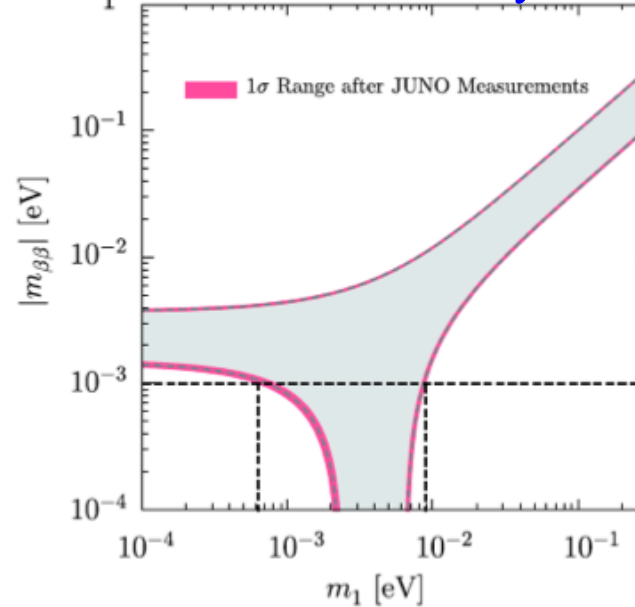
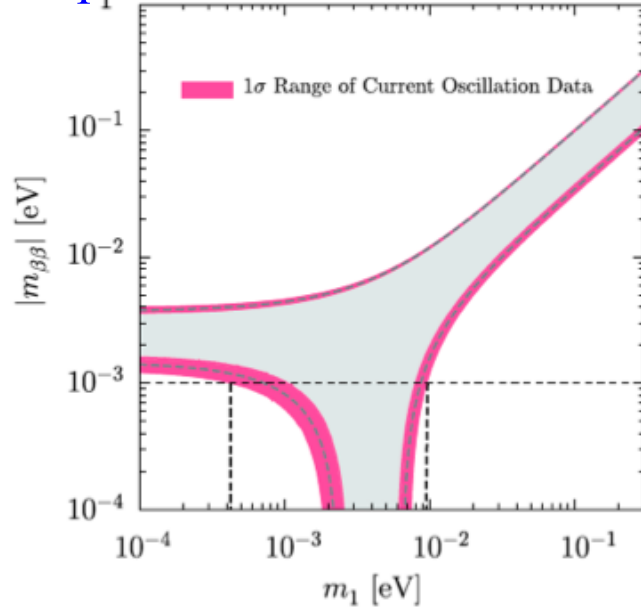




# Towards Majorana phases

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Majorana phase can be studied when effective mass sensitivity is below  $\sim 1\text{meV}$



$$m_{\beta\beta} = \left| \sum_{k=1}^3 U_{ek}^2 m_k \right|$$

$$= \left| c_{13}^2 c_{12}^2 e^{i\rho} m_1 + c_{13}^2 s_{12}^2 m_2 + s_{13}^2 e^{i\sigma} m_3 \right|$$

But, what about the uncertainty in nuclear matrix elements ?

Cao et al., CPC44, 031001 (2020)



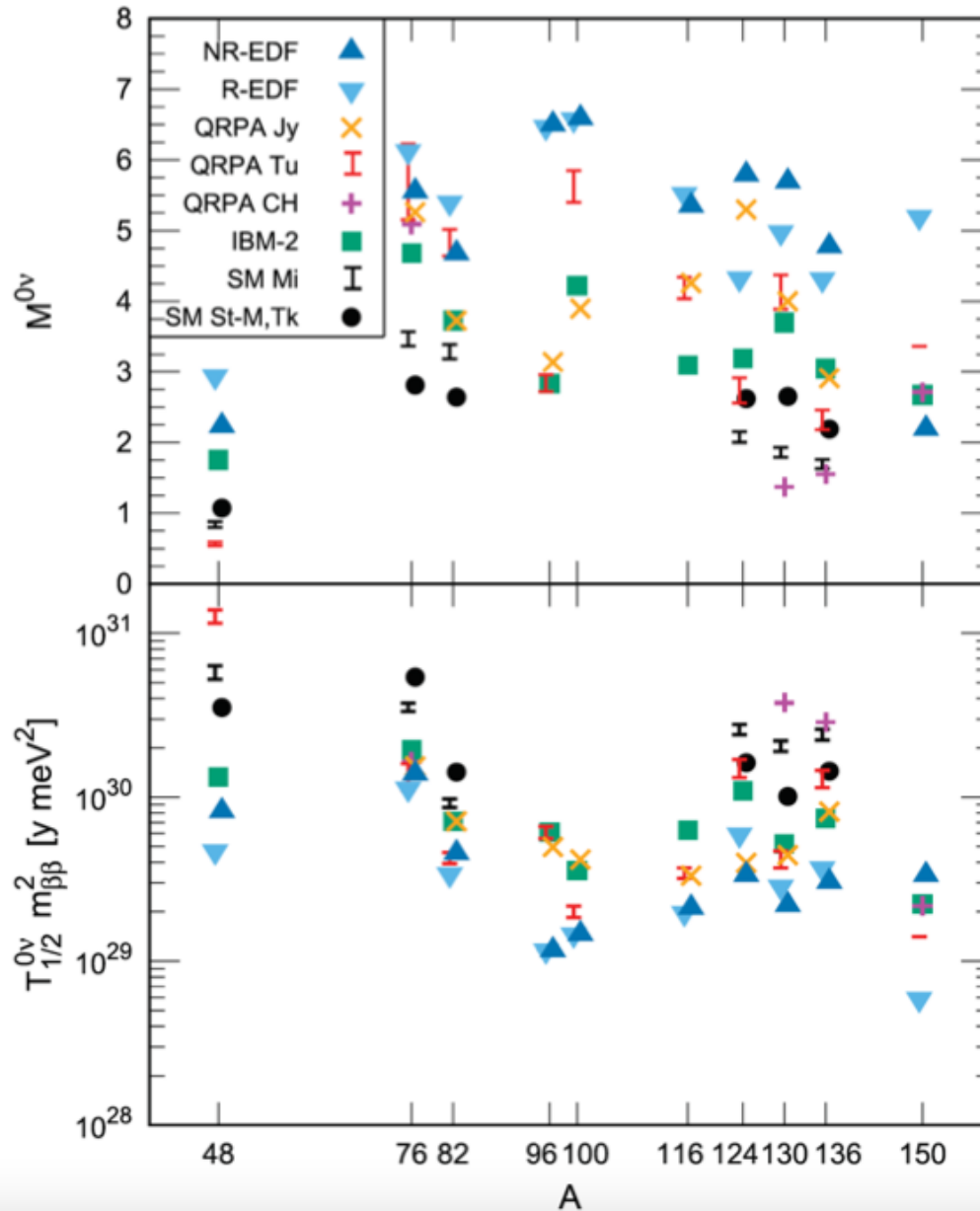
# In summary, $0\nu\beta\beta$ will

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- confirm
  - Neutrinos are Majorana particles and have Majorana masses.
  - Lepton number non-conservation.
- support on
  - See-Saw model of the neutrino mass.
  - Leptogenesis to account for the baryon asymmetry of the universe.

# Best isotope for experiment ?

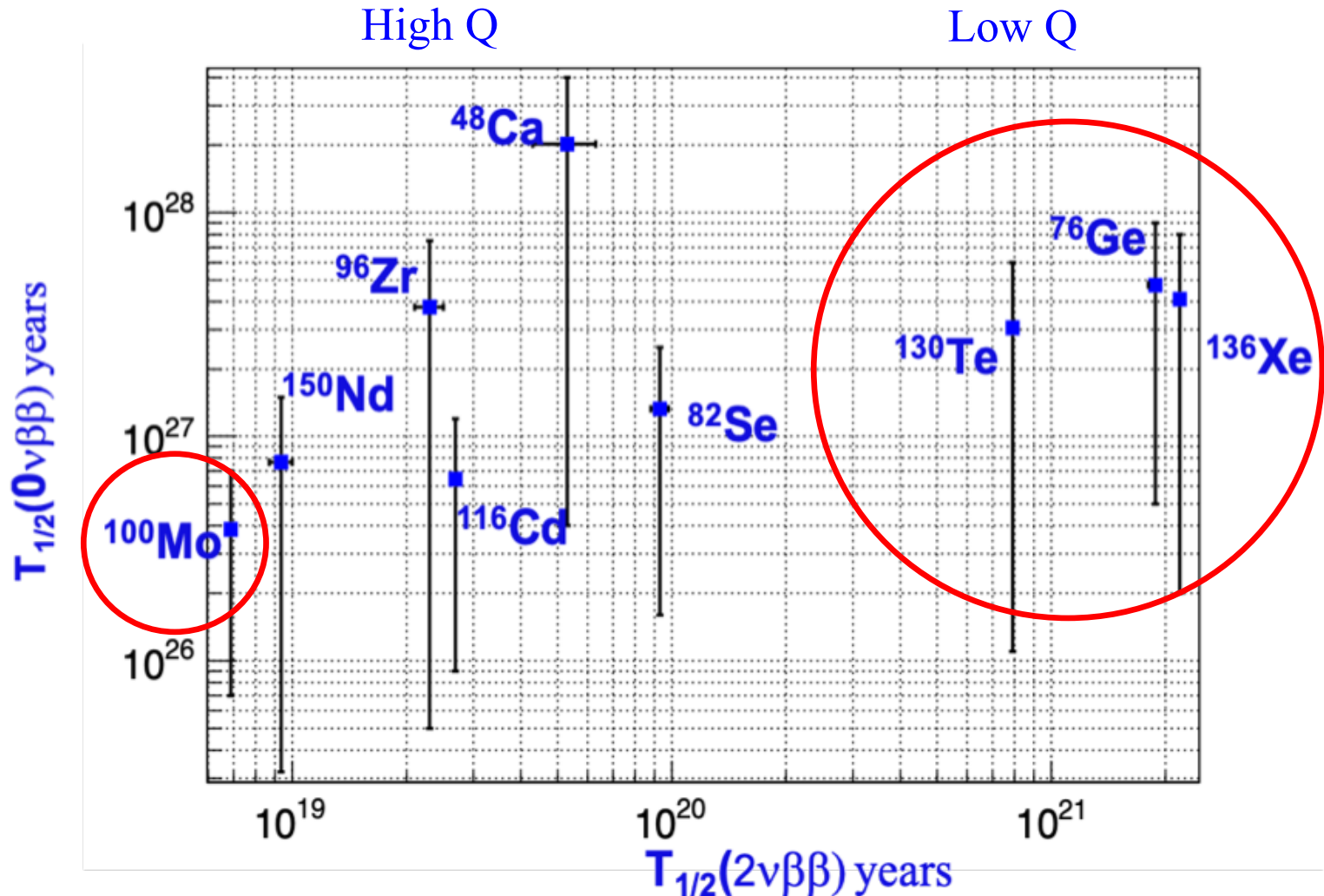
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Engel & Menendez  
RPP 80(2017)

# $0\nu\beta\beta$ vs $2\nu\beta\beta$ $T(1/2)$

- Many isotopes are used for  $0\nu\beta\beta$  searches.
- A correlation between  $2\nu\beta\beta$  half-life(measured) vs  $0\nu\beta\beta$  half-life calculated with various models for inverted mass ordering(20-49 meV).

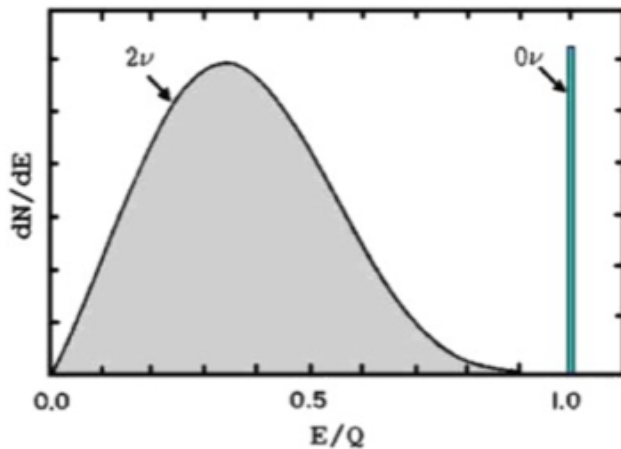


# Now, how sensitive are the $0\nu\beta\beta$ experiments ?

- $0\nu\beta\beta$  needs a good energy resolution and extremely low backgrounds.
- Sensitivities on the half-life depends on background and exposure

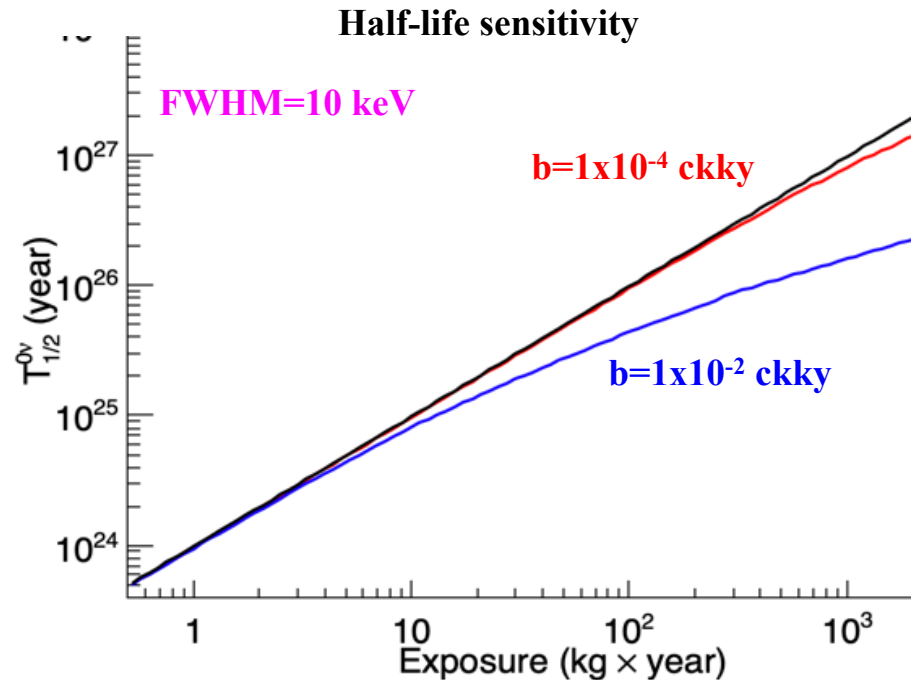
$$T_{1/2}^{0\nu} \propto \sqrt{\frac{MT}{b\Delta E}} \text{ (for finite backgrounds)} \quad T_{1/2}^{0\nu} \propto MT \text{ (for "0" backgrounds)}$$

Signal : sharp peak @ Q-value

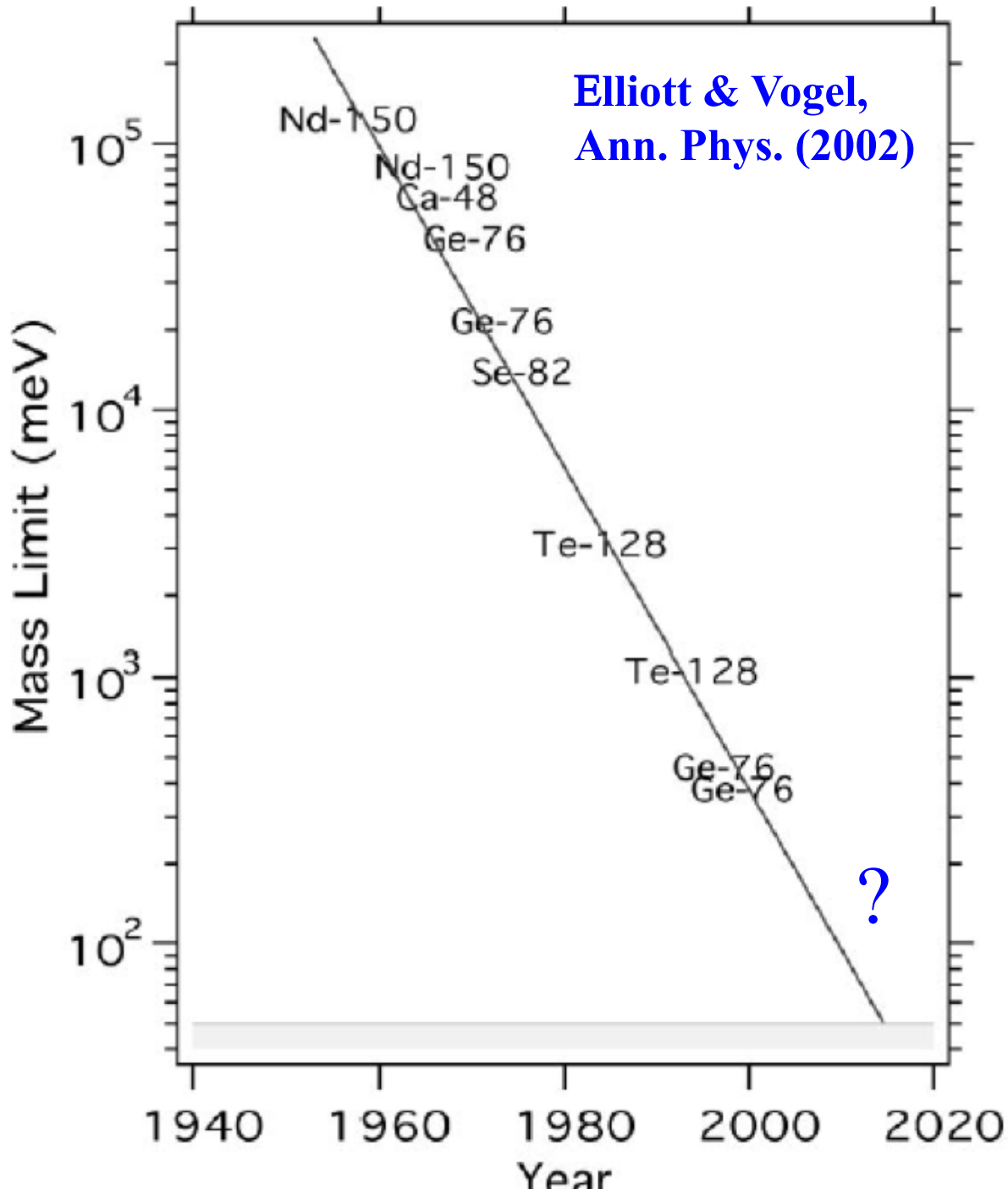


Background Unit :

ckky=counts/(keV kg year)

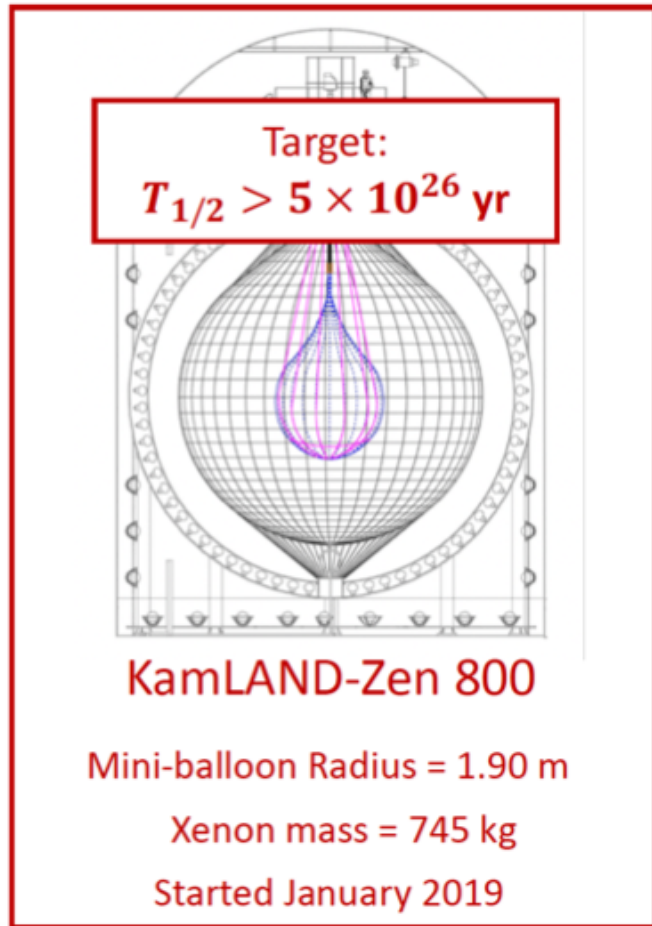






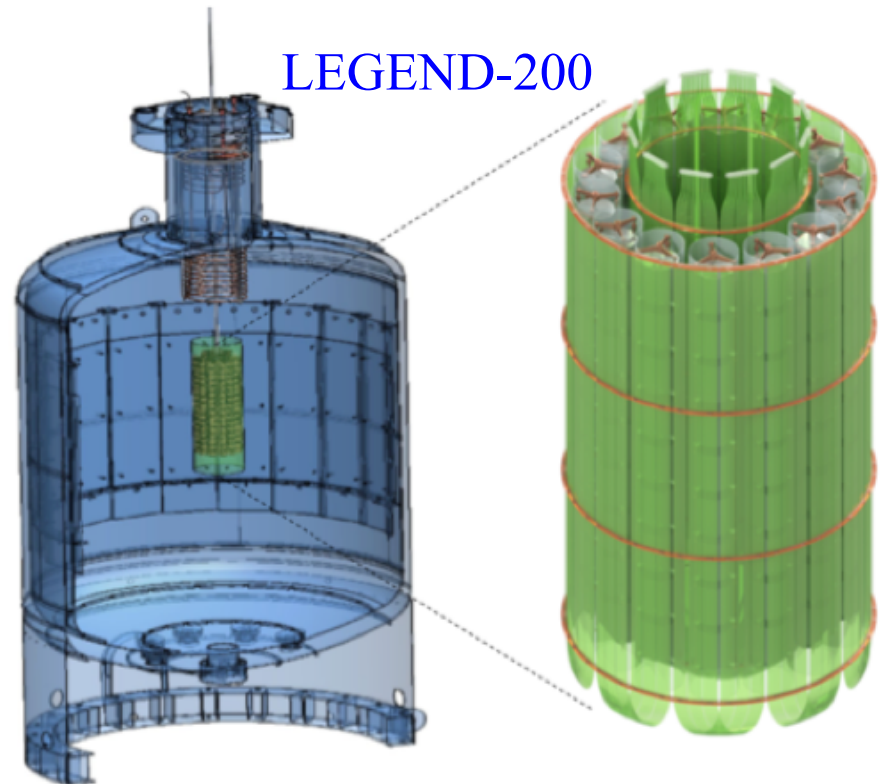
# Future experiments

Current



SnowMass 2021

KamLAND-Zen



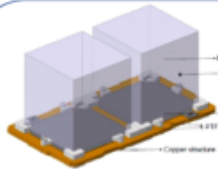
+ LEGEND-1000

+ KamLAND2-Zen

# CUPID

CUPID pre-CDR [arXiv:1907.09376](https://arxiv.org/abs/1907.09376) upgrade to CDR ongoing

- Single module:  $\text{Li}_2^{100}\text{MoO}_4$ , 45x45x45 mm – ~ 280 g
- 57 towers of 14 floors with 2 crystals each - 1596 crystals
- ~240 kg of  $^{100}\text{Mo}$  with >95% enrichment
- ~ $1.6 \times 10^{27}$   $^{100}\text{Mo}$  atoms
- No reflecting foil
- Ge light detector as in CUPID-Mo, CUPID-0



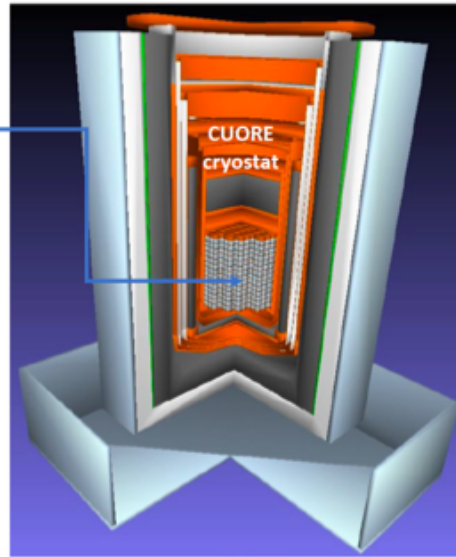
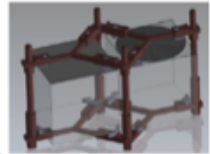
## Baseline design

Gravity stacked structure  
Crystals thermally interconnected

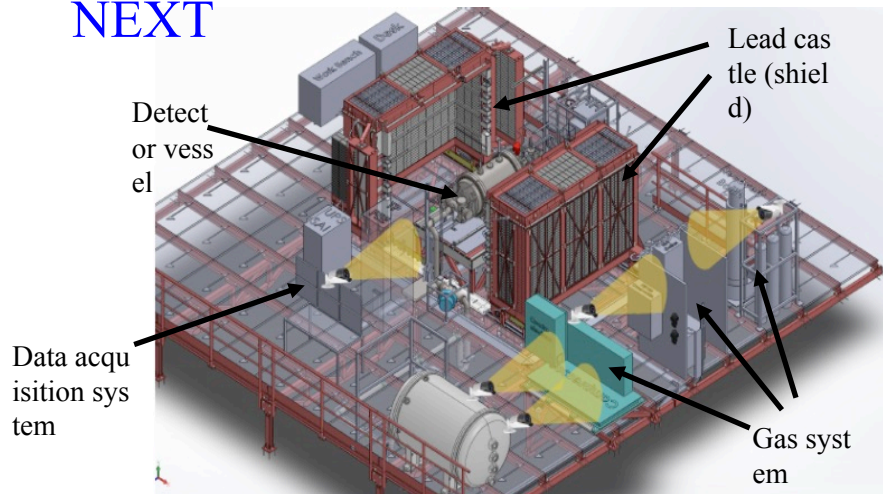
Tests ongoing

## Alternative design

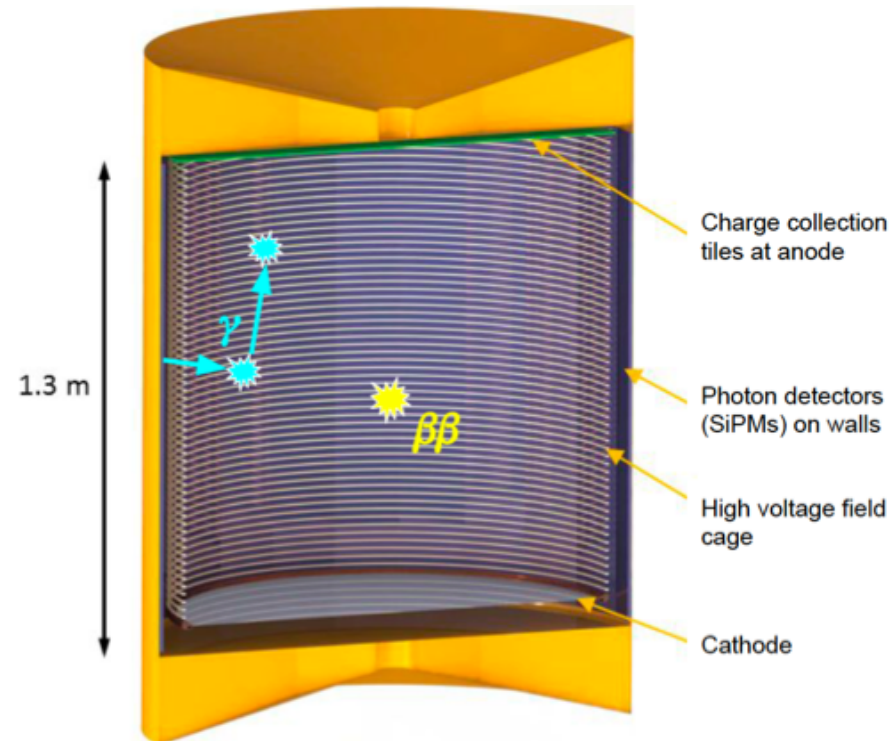
Crystals thermally independent  
No Cu holder for light detectors



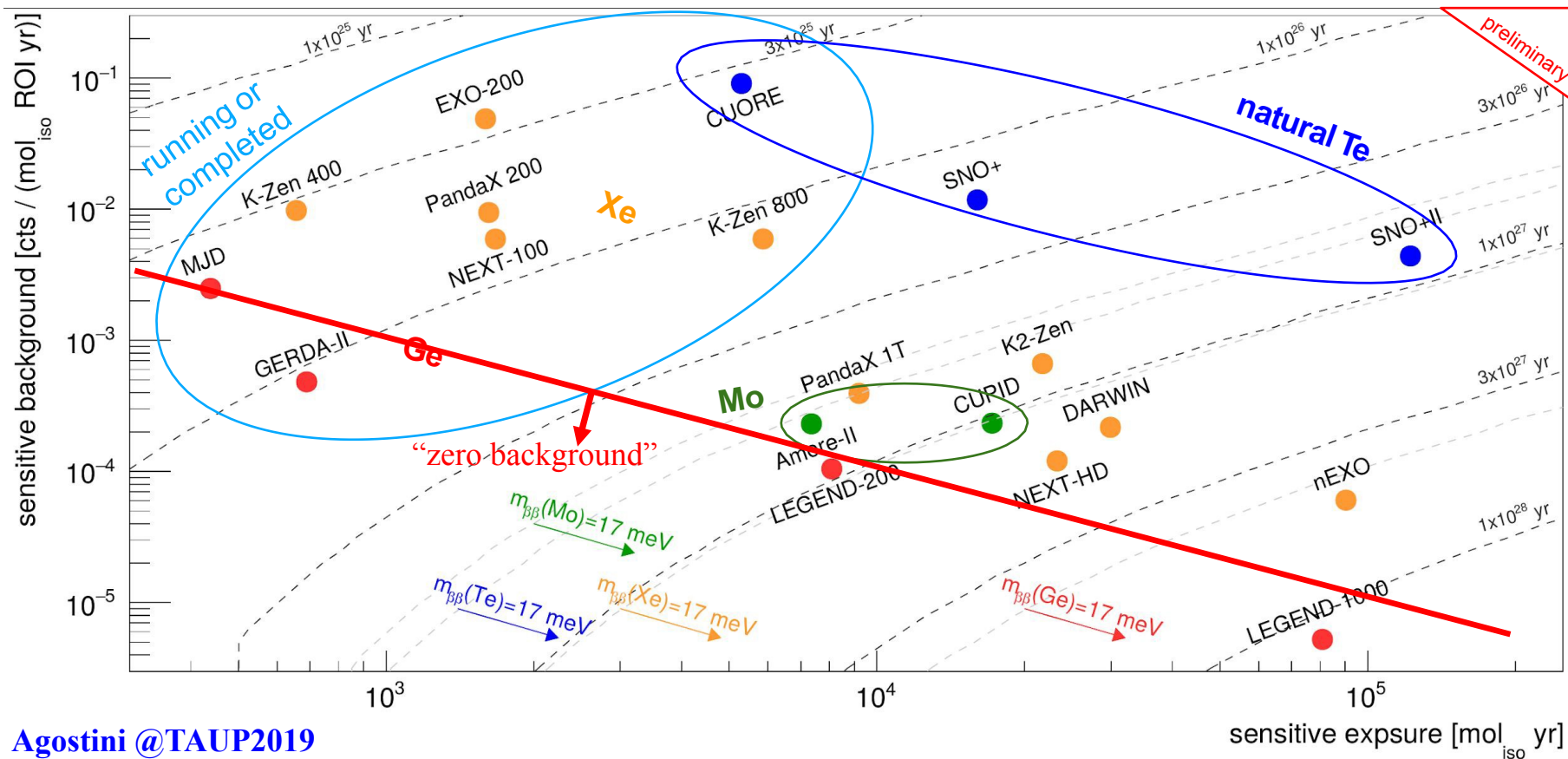
# NEXT



# nEXO



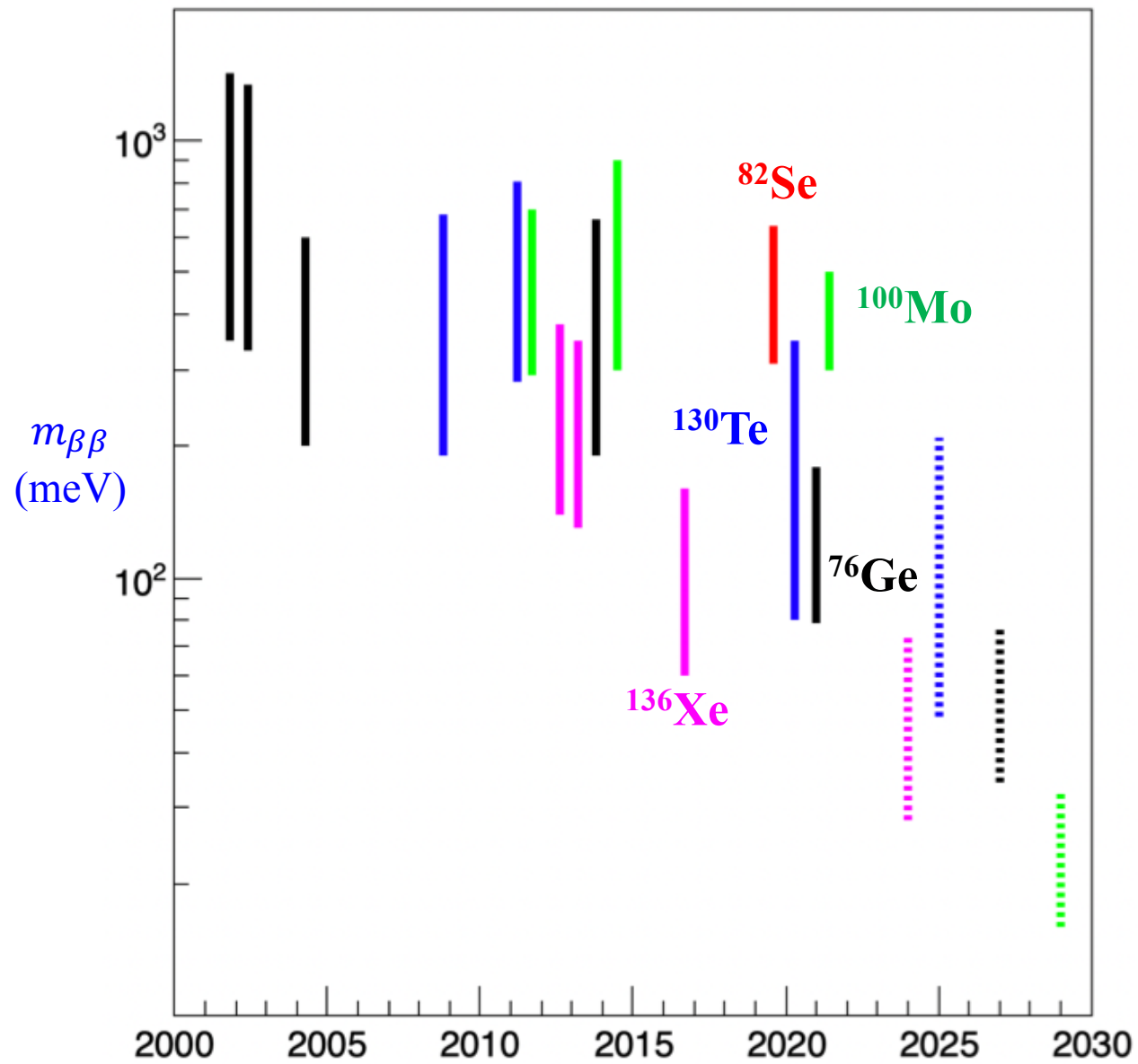
# Comparison with other experiments.



Agostini @TAUP2019

- AMoRE-II is comparable to CUPID, LEGEND-200, KamLAND2-ZEN.
- IBS(CUP) has a MOU with INFN(Gran Sasso) to collaborate between AMoRE and CUPID.

# Recent Limits & Perspectives

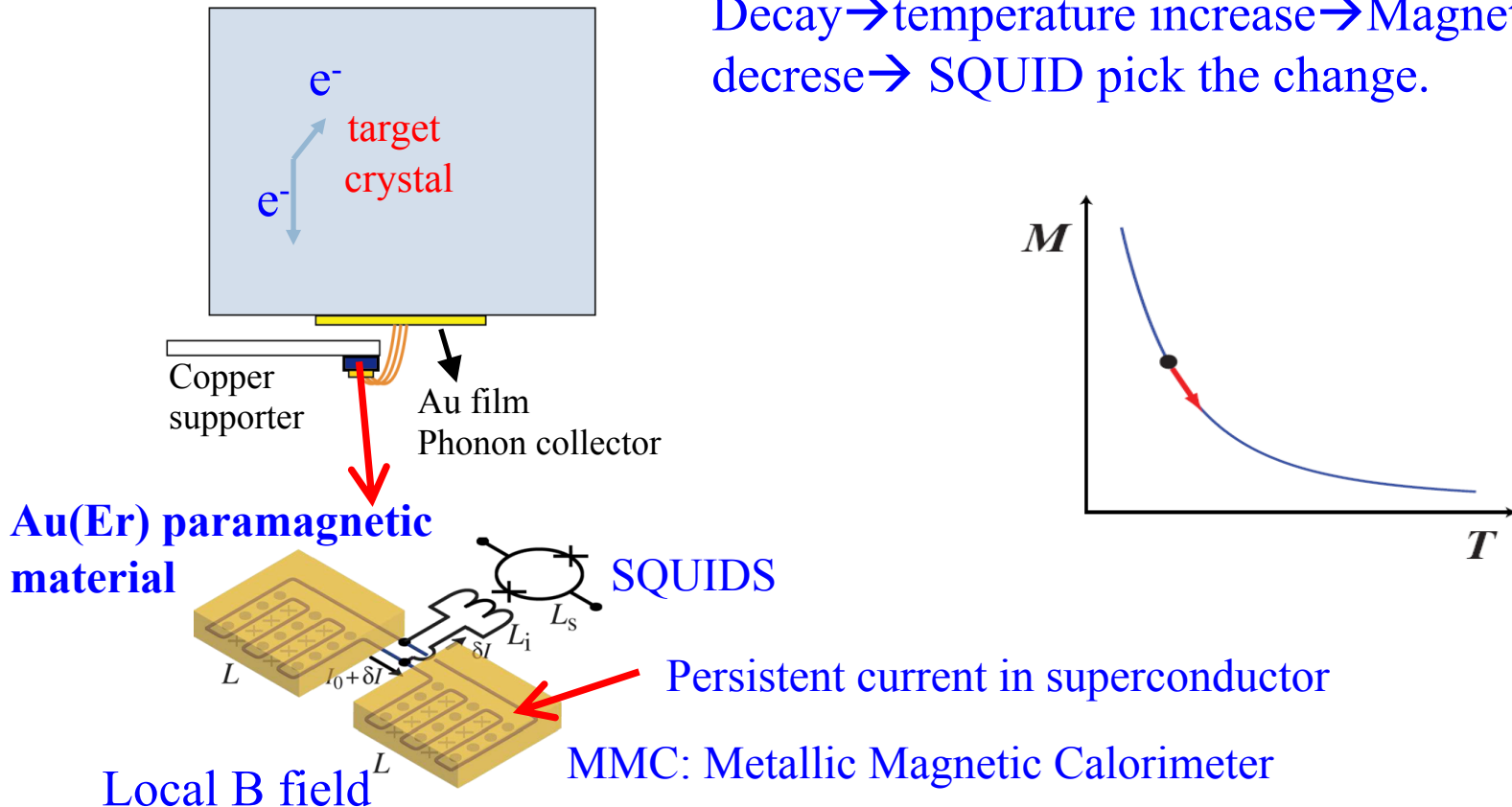


# Principle of AMoRE Detector

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- Use Mo containing Scintillating Bolometer :  $(^{40}\text{Ca},\text{X})^{100}\text{MoO}_4 + \text{MMC}$
- For Each crystal, phonon and photon sensors made of MMCs+SQUIDS to separate alphas (background) and betas (signal). **Highly Technical !**

Decay  $\rightarrow$  temperature increase  $\rightarrow$  Magnetization decrease  $\rightarrow$  SQUID pick the change.

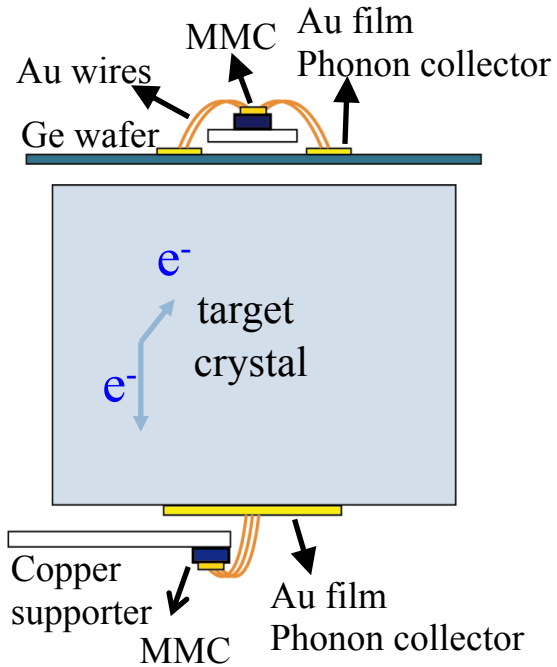




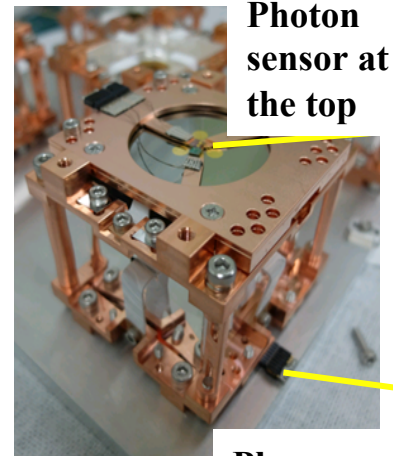
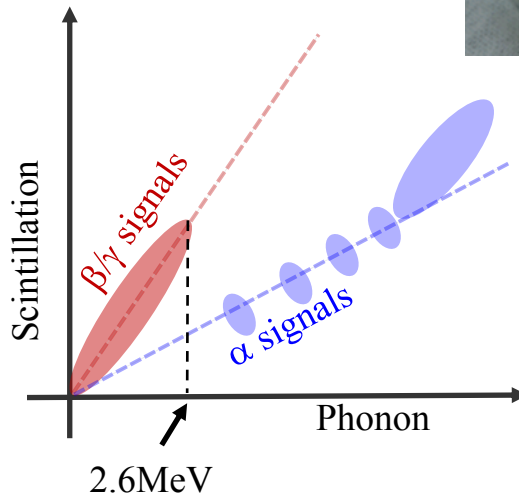
# Real AMoRE Detector

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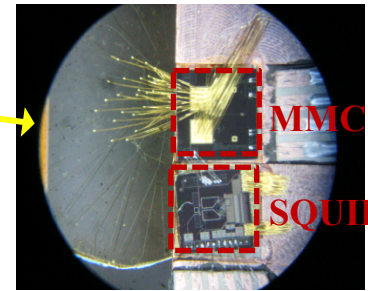
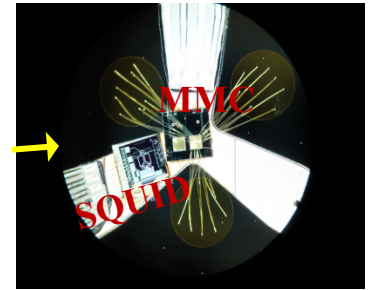
- Additionnally detect lights from scintillating crystal → can remove continuous alpha backgrounds.



Light detector added

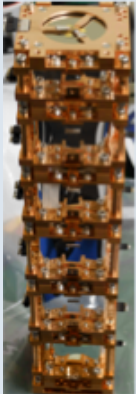
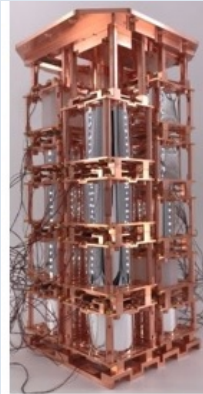
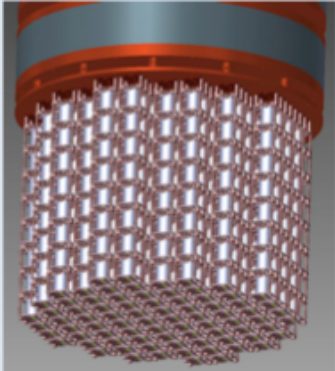


Photon sensor at the top



Phonon sensor at the bottom

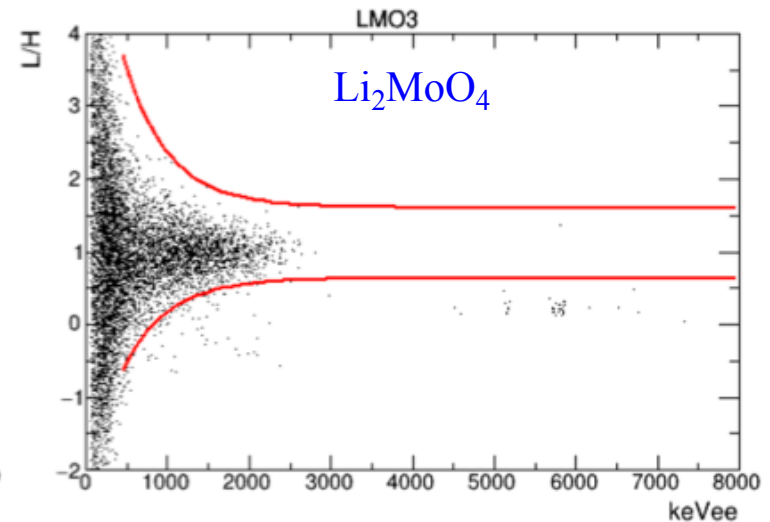
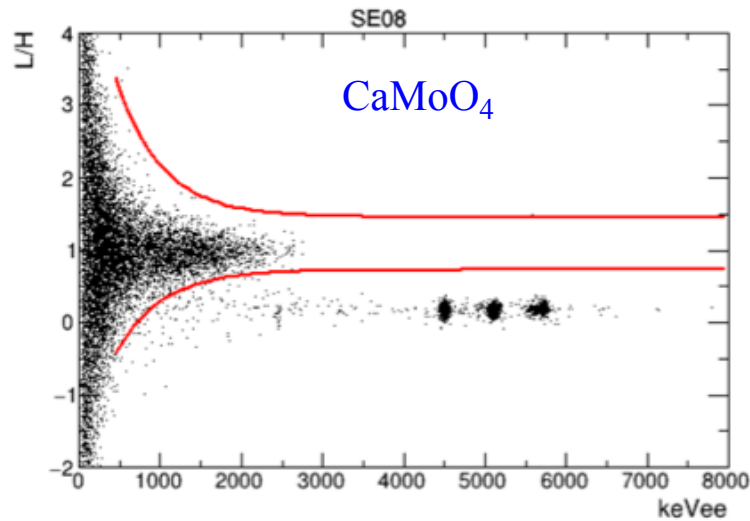
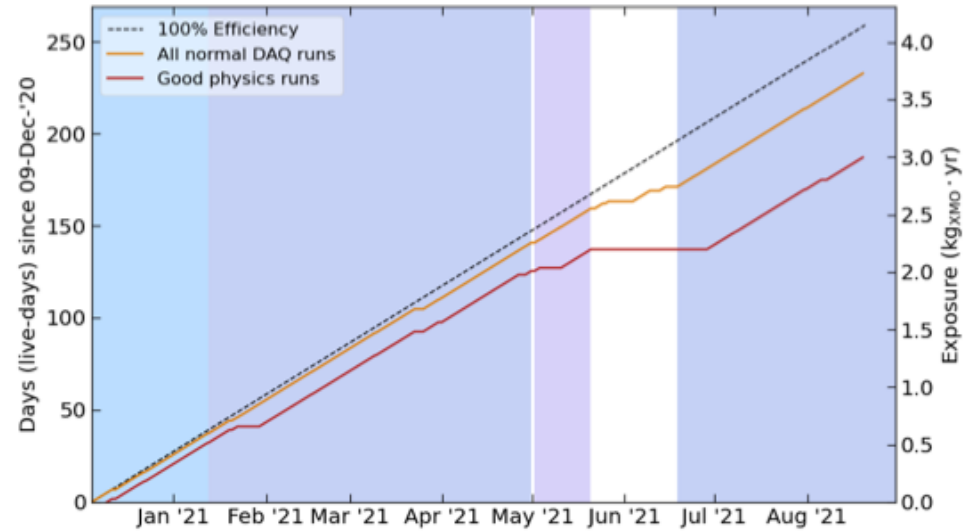
# Plan of AMoRE Project

Phases	AMoRE-Pilot	AMoRE-I	AMoRE-II
Detector Setup (Not in scale)			
Crystals	$^{40}\text{Ca}^{100}\text{MoO}_4$ (CMO)	$(^{40}\text{Ca},\text{Li}_2)^{100}\text{MoO}_4$	$\text{Li}_2^{100}\text{MoO}_4$ (LMO)
Crystal # & Mass	6, 1.9kg	18, 6.2kg	596, 178kg
Background Goal(ckky)	$10^{-1}$	$<10^{-2}$	$<10^{-4}$
$T_{1/2}$ (year)	$1.0 \times 10^{23}$	$7.0 \times 10^{24}$	$8.0 \times 10^{26}$
$m_{\beta\beta}$ (meV)	1200-2100	140-270	13-25
Location/Schedule	Y2L / 2015-2018	Y2L / 2020-2022	Yemilab / 2022-2027

# AMoRE-I

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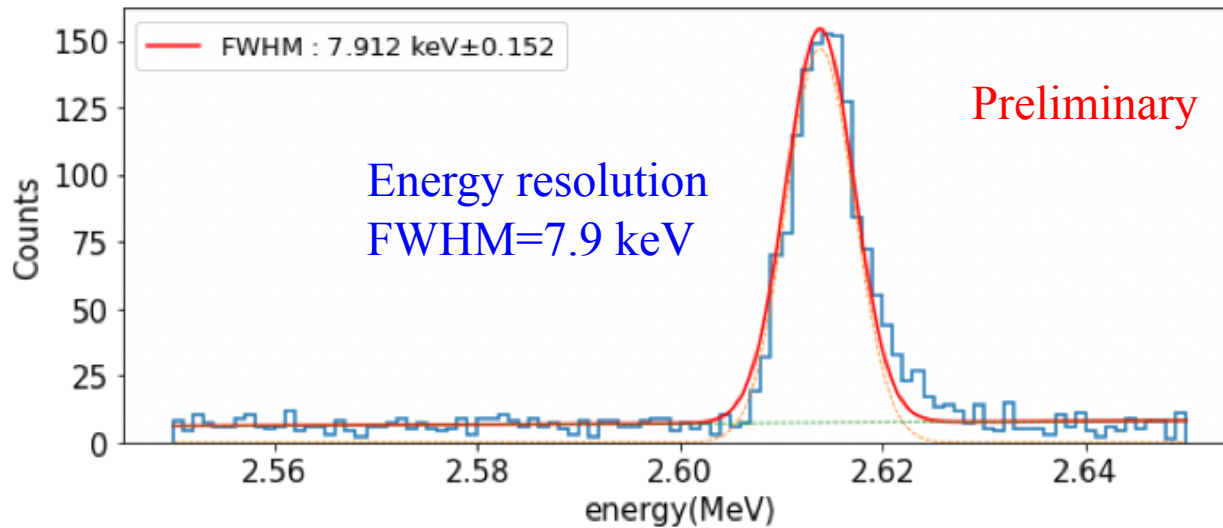
- AMoRE-I run began Aug. 2020 @ Y2L
- Purpose :
  - Check detector performance (LMO)
  - Understand background better.
- Same cryostat as pilot. 18 crystals.



# AMoRE-II status

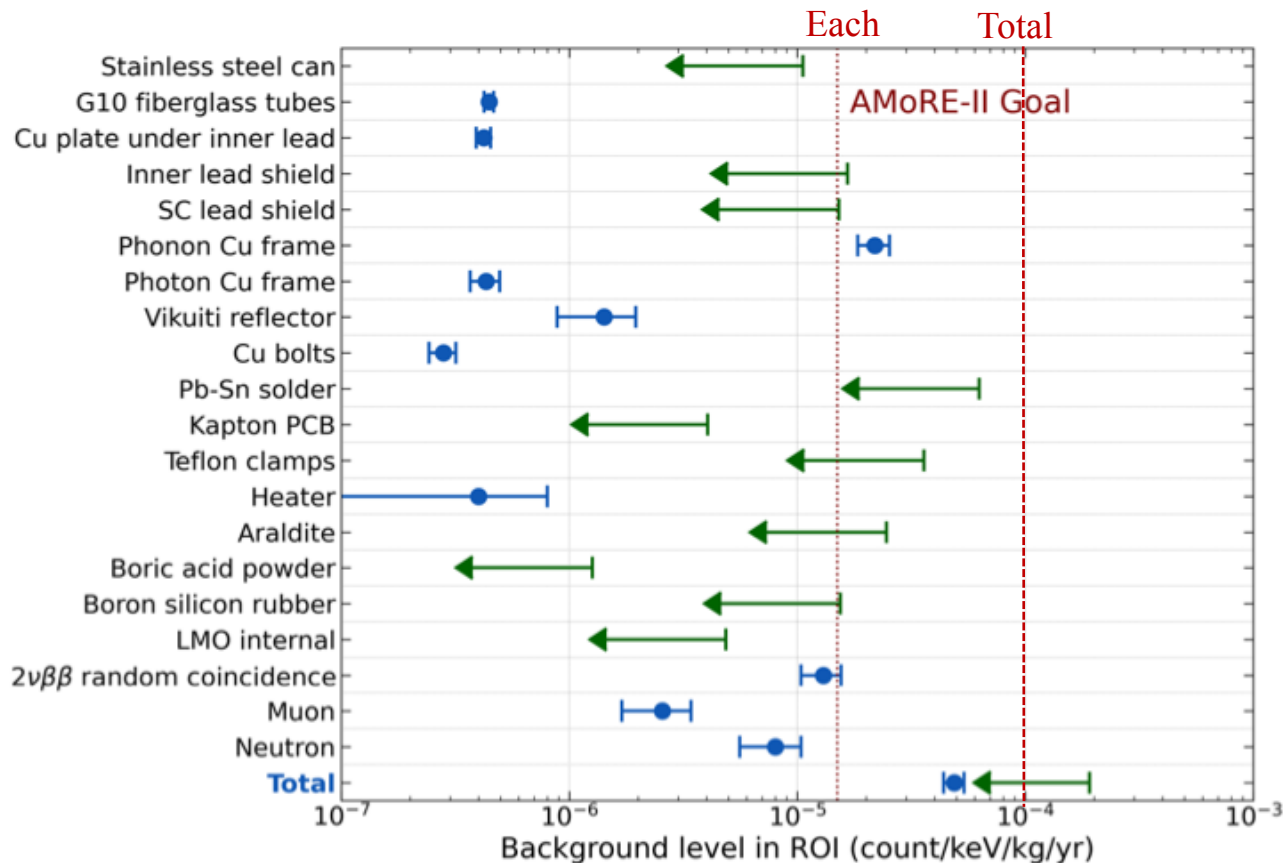
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- Tested larger crystal to reduce channel numbers for bigger experiment.
- Preliminary 6cm(D)x6cm(H) crystal result is promising.



# Expectation of AMoRE-II Backgrounds

1. All materials inside Pb shielding are measured by ICP-MS and(or) HPGe.
2. All outside sources are estimated by measurements and Geant4 simulation.
3. Muon induced backgrounds are simulated with the expected muon flux.
4. Projected total background level is  $4.9(5) \times 10^{-5}$  ckky and  $< 2 \times 10^{-4}$  ckky for a limiting values.



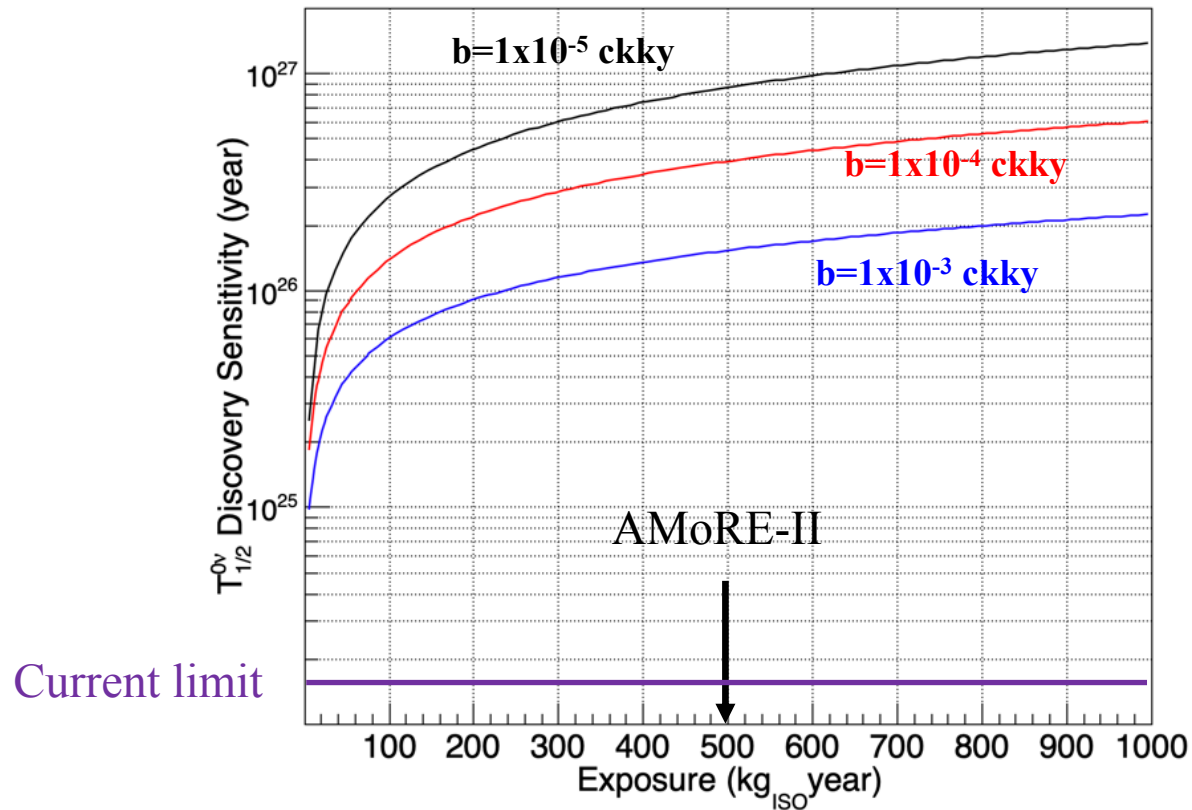
# Discovery Sensitivities

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Discovery sensitivity :

The half-life for which an experiment has a 50% chance to measure a signal above background with a significance of at least 3 sigma (99.7%).  $4 \times 10^{26}$  years

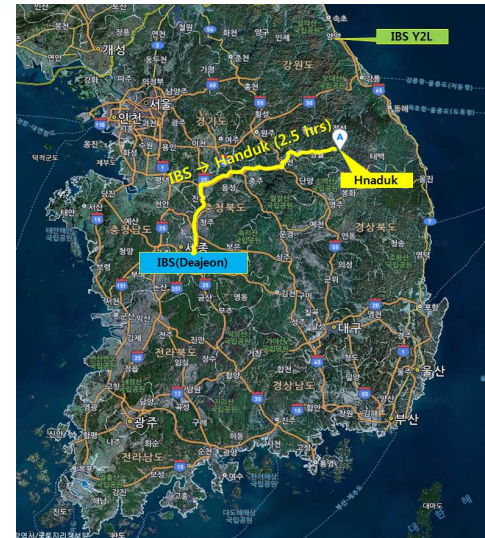
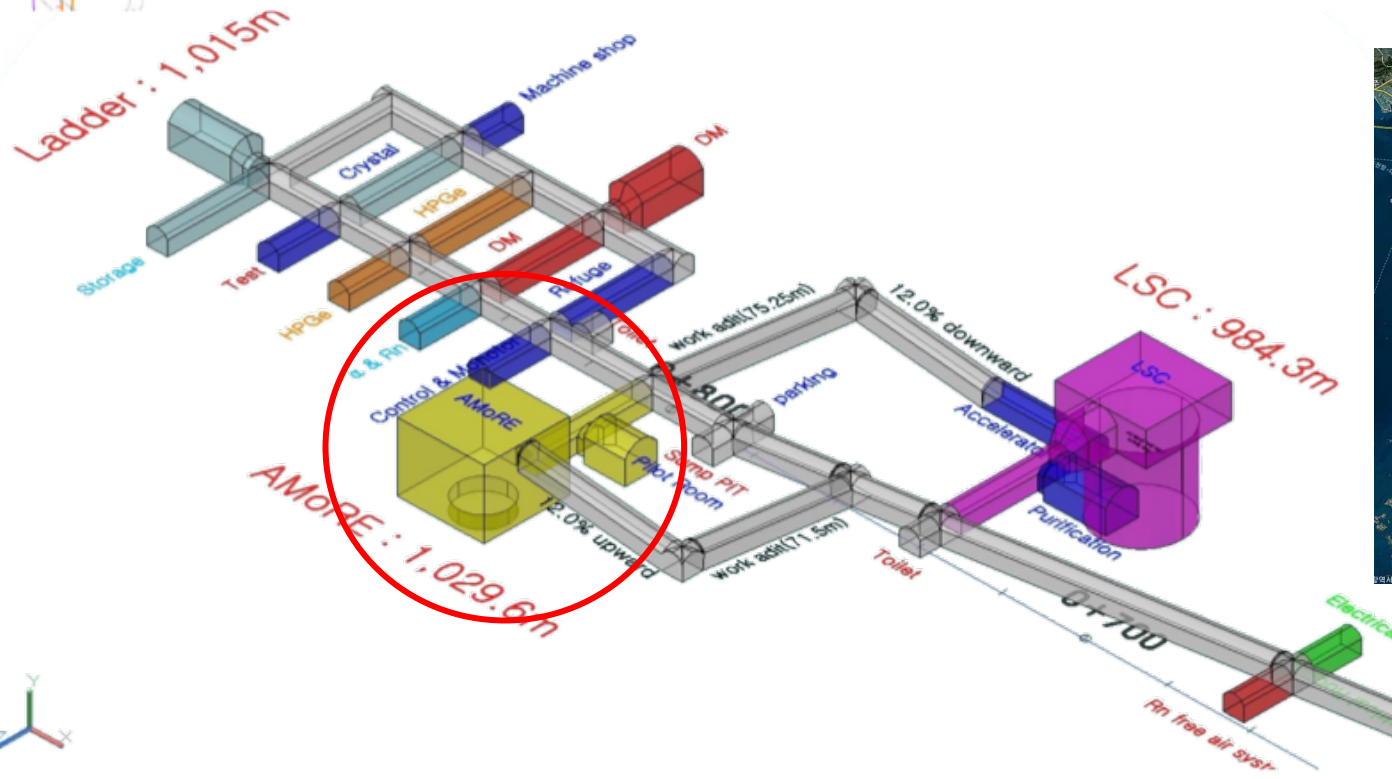
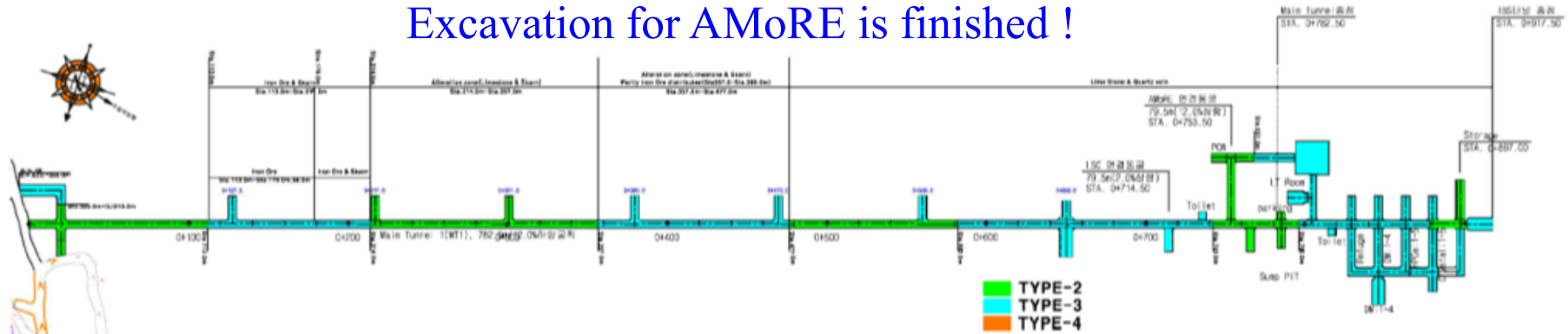
Detection efficiency : 0.7, Background normalized to isotope mass



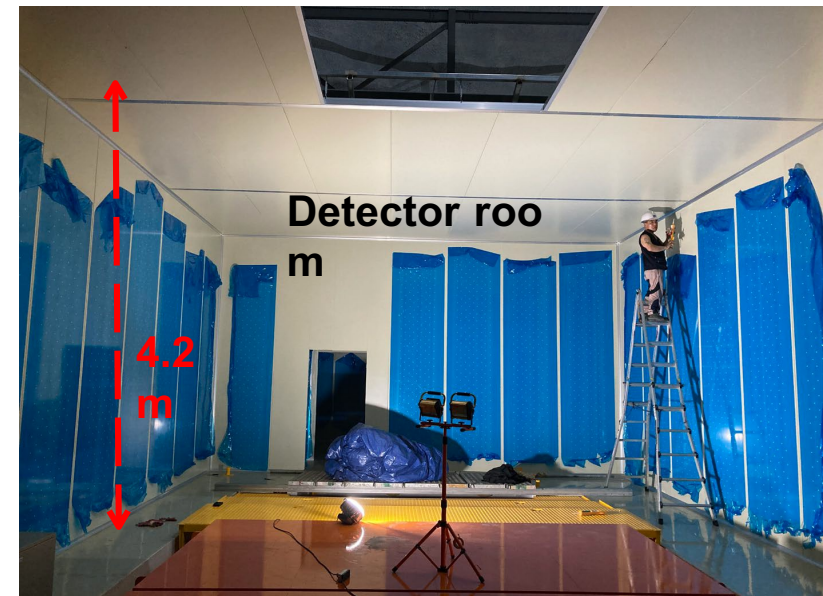
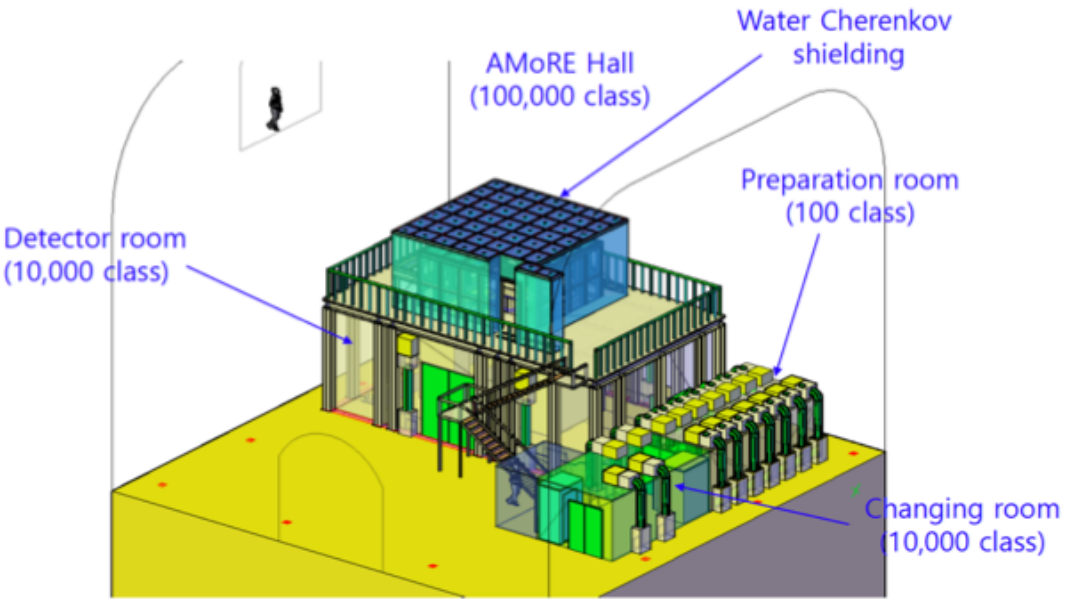


# AMoRE-II @ Yemilab

Excavation for AMoRE is finished !



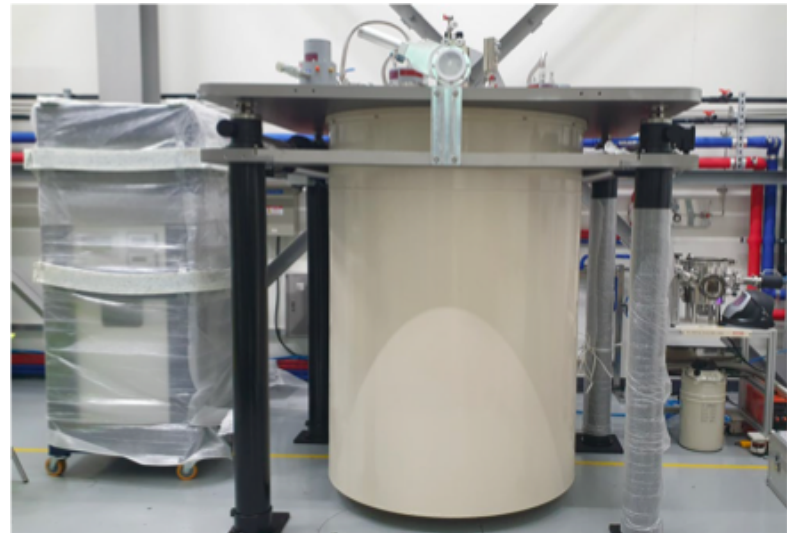
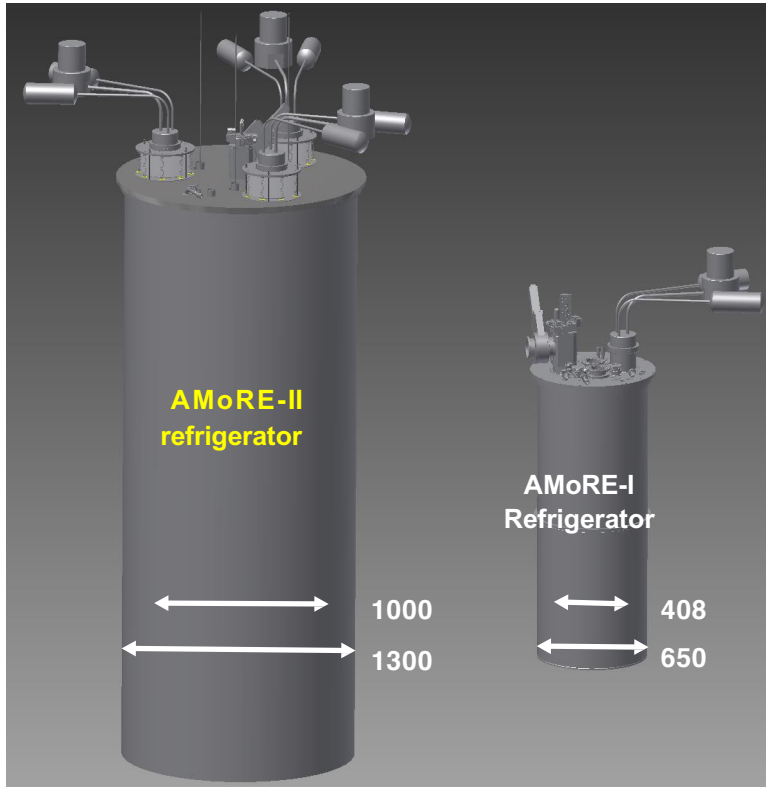
# Construction is going on..



# AMoRE-II refrigerator

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- Big and powerful dilution Refrigerator
- Three PTR (PT420 RM)
- 2.4 mW @ 120 mK,
- $> 5 \mu\text{W}$  @ 10 mK
- Delivered to IBS in Aug. 2020.



## Future

- It should be more clear when we have first data of AMoRE-II with 90 crystals.
- Modular expansion is possible, increasing # of detectors.
- After AMoRE-II, ton scale experiment can be considered. ~ CUPID 1ton.
- Can we try to do research about possibility to enrich Mo-100 ?
- CUPID & AMoRE discuss to collaborate for future combination.



# Summary

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- $0\nu\beta\beta$  is one of the best probe for BSM physics.
- Experimentally will reach IO region  $\sim 2030$ , and will continue to lower the sensitivities.
- AMoRE experiment aims to be sensitive close to  $10^{27}$  year range for  $^{100}\text{Mo}$  isotope and will be installed by end of 2022.
- Other experiments, LEGEND, CUPID, nEXO are actively pursued.
- $0\nu\beta\beta$  should be pursued for multiple isotopes.
- $0\nu\beta\beta$  can be discovered at anytime with new sensitivities.

# Neutrino 2022

## SEOUL

2022. 5. 30 ~ 6. 4 , COEX

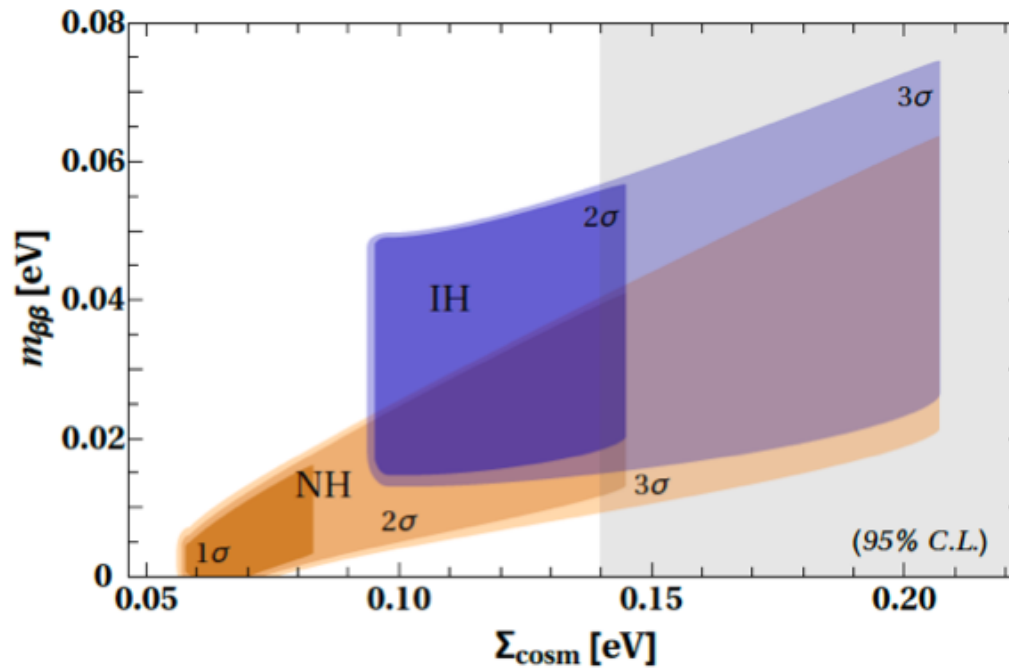
Korea



# Cosmological constraints

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Dell’Oro et al., *JCAP* 12 (2015) 023



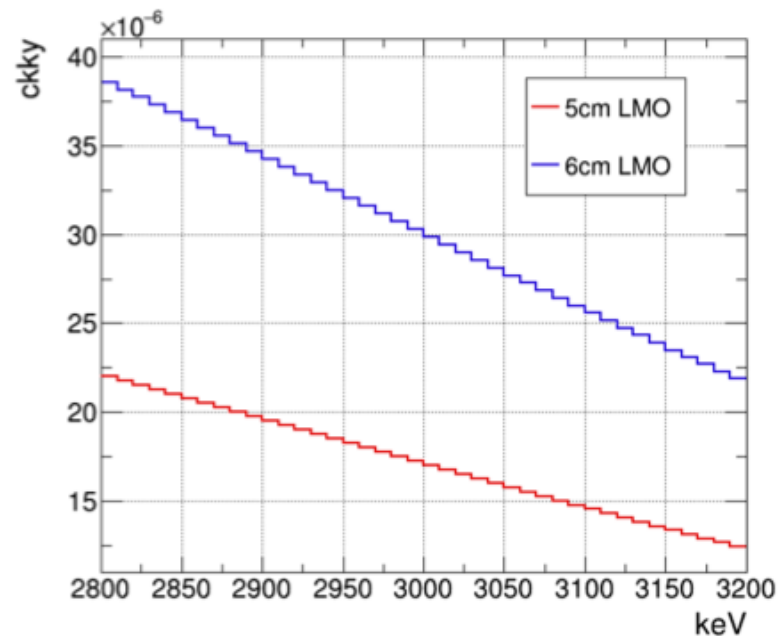
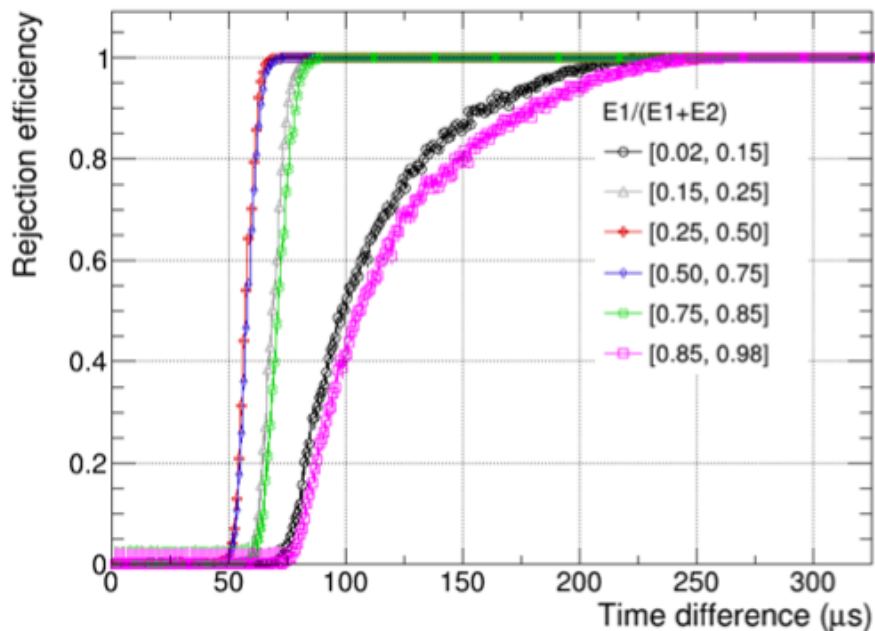
$$\Sigma < 84 \text{ meV} \quad (1\sigma \text{ C.L.})$$

$$\Sigma < 146 \text{ meV} \quad (2\sigma \text{ C.L.})$$

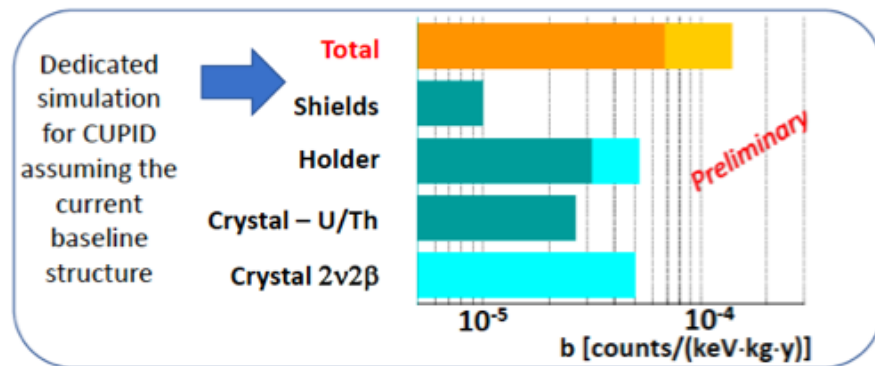
$$\Sigma < 208 \text{ meV} \quad (3\sigma \text{ C.L.})$$

# Pileup background estimation

- A realistic estimation assuming real spectra and noise data from AMoRE-pilot
- Crystal size is important – pile up event rate is proportional to square of single rates.
- 6cm crystal is acceptable.



Compare with CUPID-Mo

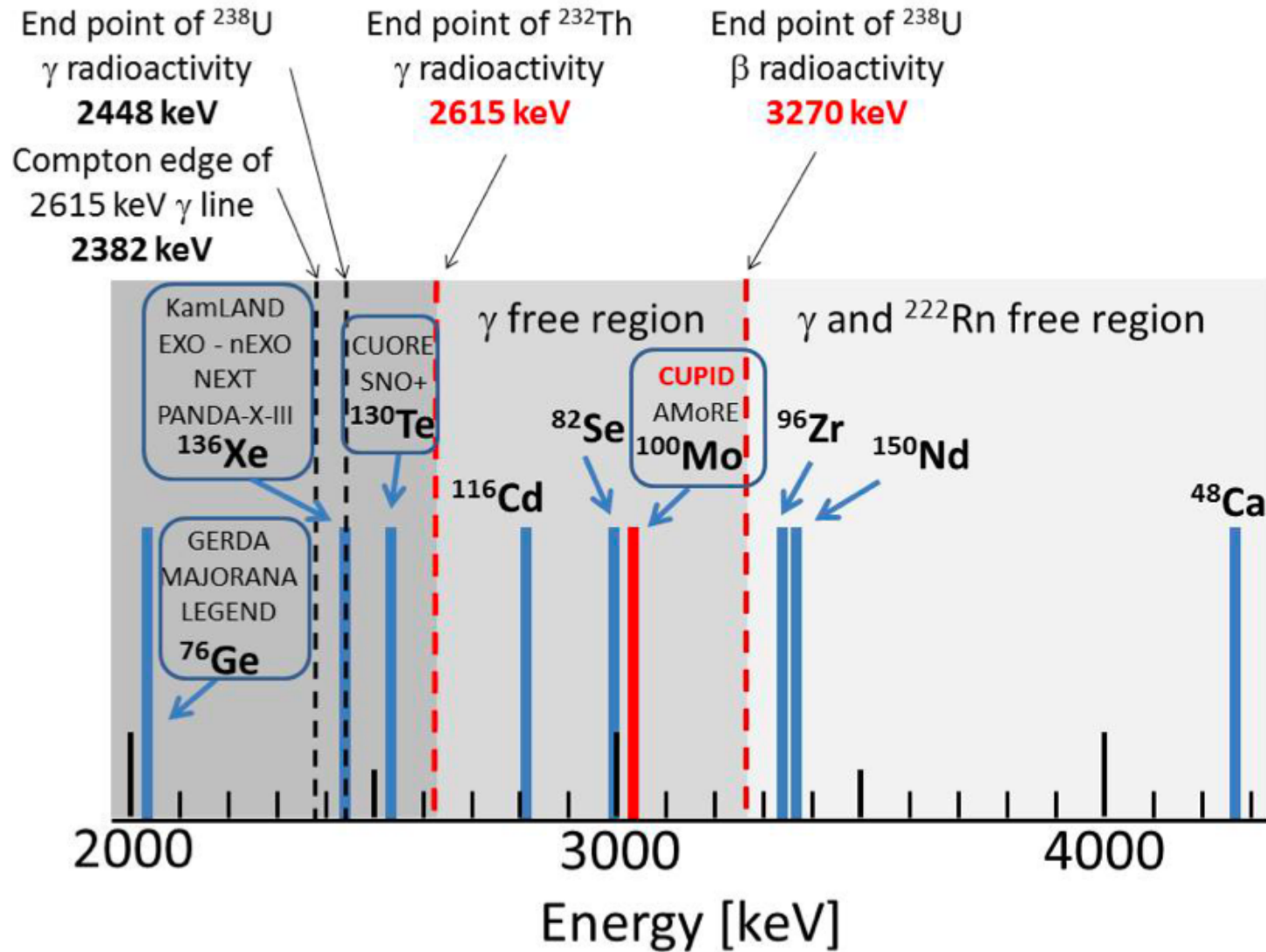






# Q value vs backgrounds

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# Current best results for $0\nu\beta\beta$

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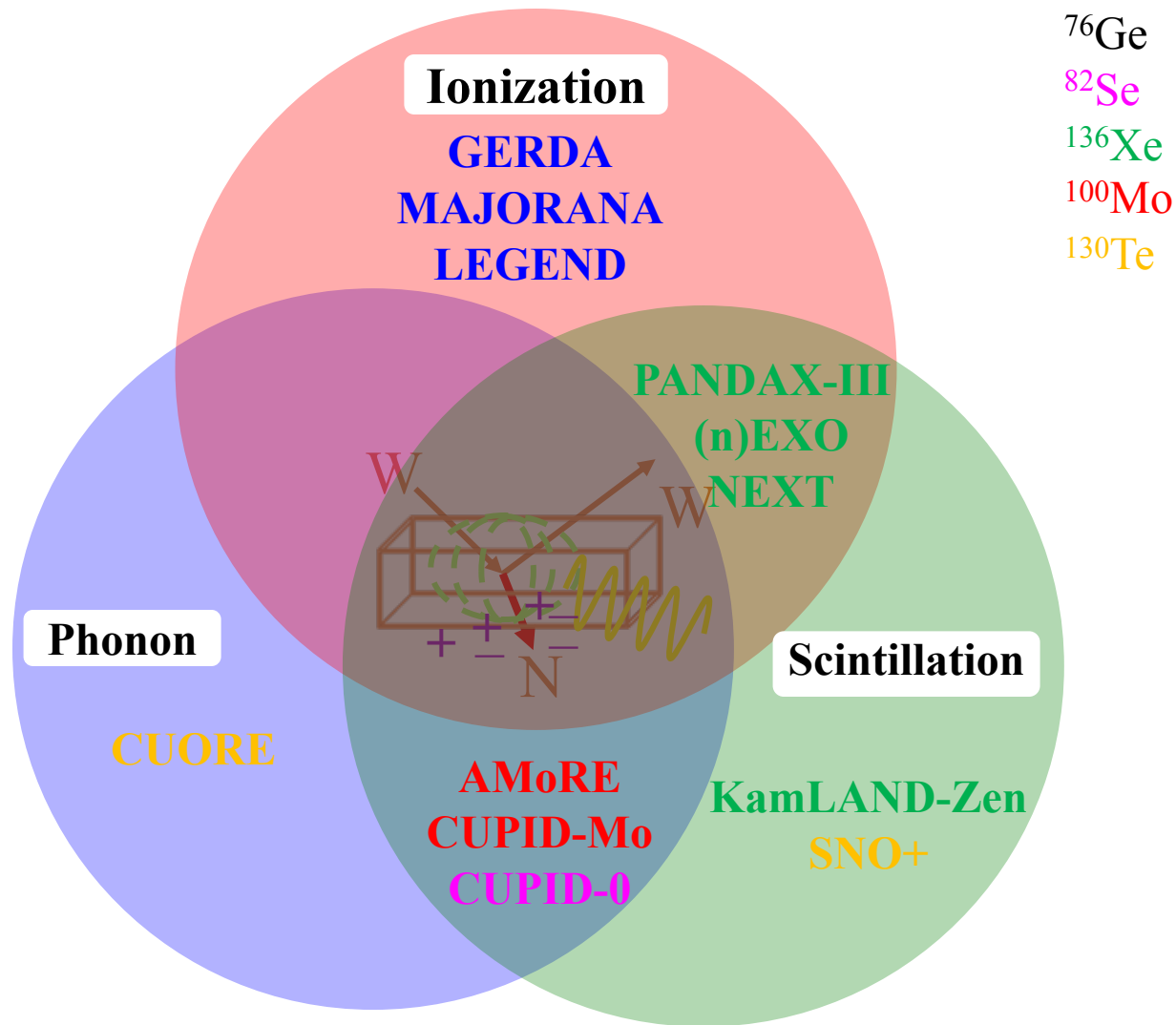
2021. 4. 10

Nucl.	Q (keV)	Abun. (%)	$T_{1/2}^{2\nu}$ ( $10^{20}$ Y)	Exp	$T_{1/2}^{0\nu}$ ( $10^{24}$ Y)	M (eV)	Ref.
$^{48}\text{Ca}$	4270.0	0.187	0.53(0.1)	CANDLES	> 0.058	<3.1-15.4	PRC 78 058501 (2008)
$^{76}\text{Ge}$	2039.1	7.8	18.8(0.8)	GERDA-II	>180	<0.079-0.18	PRL125, 252502 (2020)
$^{82}\text{Se}$	2997.9	9.2	0.93(0.05)	CUPID-0	> 2.4	<0.38-0.77	PRL120, 232502 (2018)
$^{100}\text{Mo}$	3034.4	9.6	0.0688(0.0025)	CUPID-Mo	>1.5	<0.3-0.5	arXiv:2011.13243(2020)
$^{116}\text{Cd}$	2813.4	7.6	0.269(0.009)	AURORA	> 0.19	<1-1.8	nulc-ex/1601.05578.
$^{130}\text{Te}$	2527.5	34.5	7.91(0.21)	CUORE	> 32	<0.08-0.35	PRL124, 122501 (2020)
$^{136}\text{Xe}$	2458.0	8.9	21.8(0.5)	KamLAND-Zen	> 107	<0.06-0.16	PRL117, 082503 (2016)
$^{150}\text{Nd}$	3371.4	5.6	0.0934(0.0065)	NEMO-3	> 0.02	<1.6-5.3	PRD 94 072003 (2016)

Bolometer, Scintillation, Ionization

# Detector Techniques for $0\nu\beta\beta$

Similar techniques are used as direct dark matter experiments

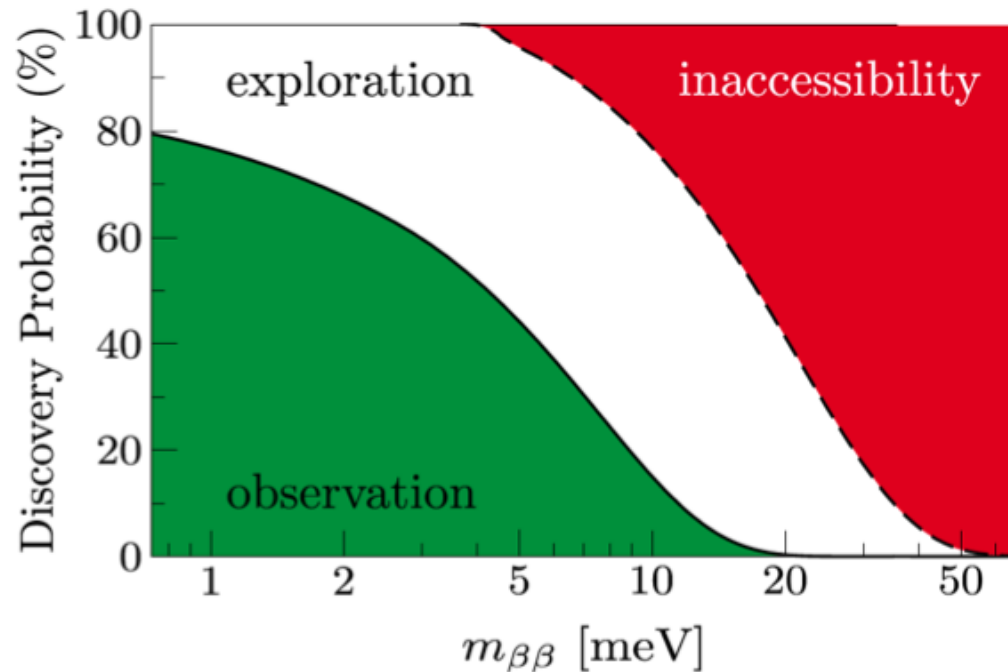


# Discovery potential for $0\nu\beta\beta$

“Discovery probabilities of Majorana neutrinos based on cosmological data”,

M. Agostini et al. *PRD* 103 (2021) 3, 033008

- Normal Ordering assumed.
- Planck cosmological limit imposed.



Experiments with sensitivities of the order of 10 meV  $\rightarrow$  discovery power between 20 and 80%.

## A prediction for the decade of the 20'

1. We'll learn that  $m_3$  is the heaviest and  $\theta_{23} > 45^\circ$ .
2. Cosmological measurements of the sum of neutrino masses will have interesting upper and non-zero lower limits.
3. The value of  $\delta_{CP}$  will not be zero or  $\pi$ .
4. The MiniBooNE low E excess will be explained.
5. The LSND excess will not be explained.
6. DUNE will mostly be on time.
7.  $0\nu\beta\beta$  limits will improve but there will not be a signal.
8. At least one of these predictions will be wrong.

- Murray Goodman –

In “Long-Baseline Neutrino Oscillation Newsletters”

# Review

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- If total background in ROI  $< 1$ , the sensitivity is

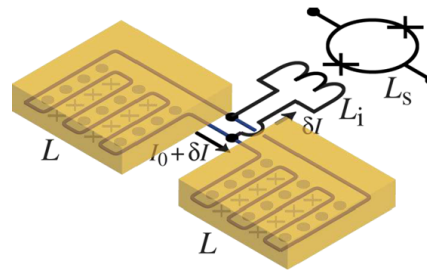
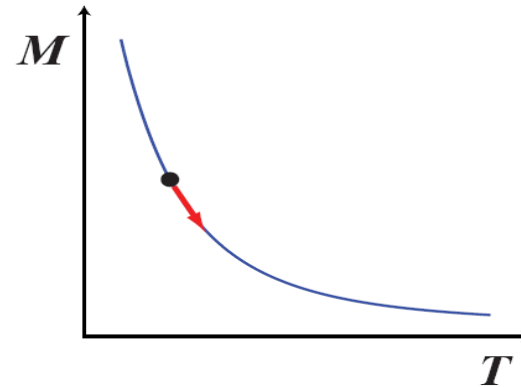
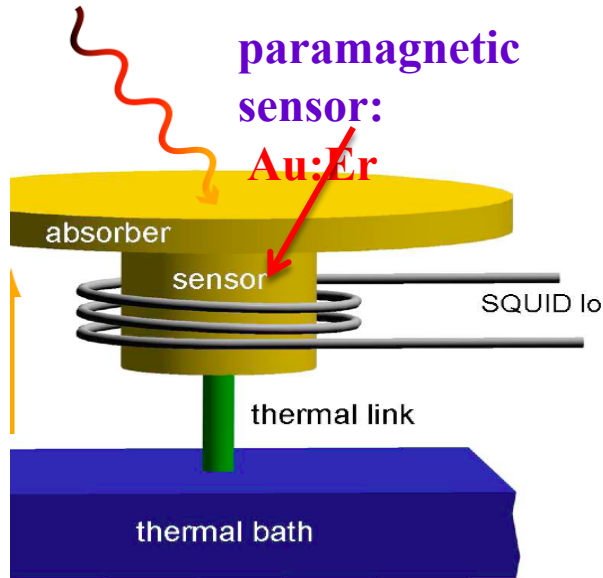
$$S_{0\nu} \propto MT,$$

Background =  $MT(\Delta E)B < 1 \rightarrow MT < \frac{1}{(\Delta E)B}$  : Exposure for “0” background.

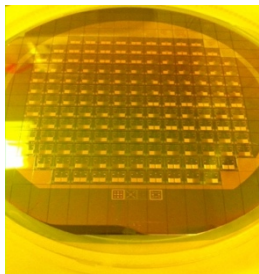
Exp	$(\Delta E)$ (keV)	B (ckty)	MT (t y)	T1/2(year)
LEGEND-200	3	$2 \times 10^{-1}$	1.7	$\sim 5 \times 10^{27}$
KAMLamd2-Zen	200	0.01	0.5	$\sim 2 \times 10^{27}$
CUPID	5	0.1	2	$\sim 10^{27}$
nEXO	25	$1.4 \times 10^{-2}$	3	$\sim 10^{28}$
AMoRE	10	0.1	1	$\sim 8 \times 10^{26}$

# Low temperature MMC sensor

MMC (Metallic Magnetic Calorimeter)



MMC: Metallic Magnetic Calorimeter



- All fabrication can be done at CUP, IBS

Total = 157

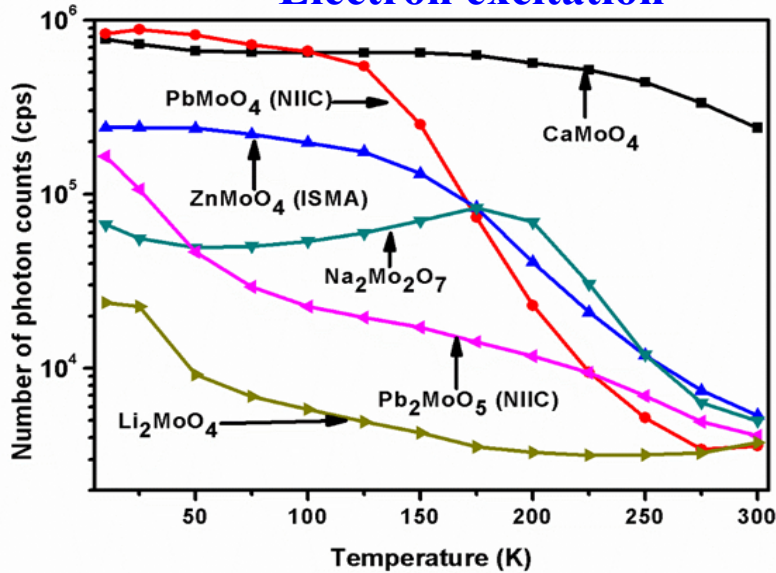
(Au:Er one sided + Au:Er both sided + test pattern(33))



# AMoRE-II Crystal Decision

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## Electron excitation

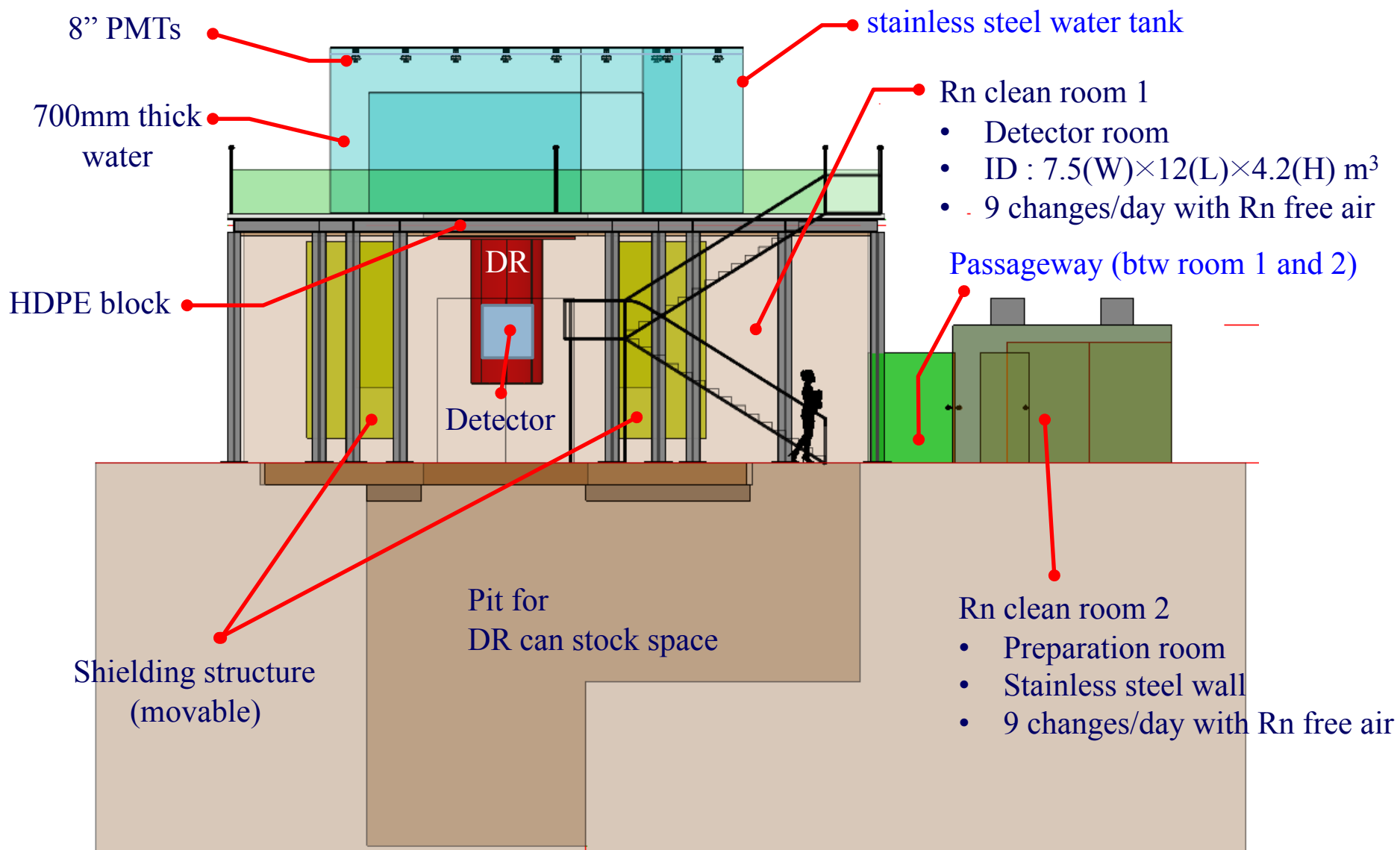


H.J. Kim et al., Crystal Research & Technology, Nov. 2019

- LMO has light output smaller than CMO by a factor larger than 10.
- But, DP is similar between these two crystals.
- LMO crystal is chosen for AMoRE-II
- Crystal growing is easier

Crystals	Scintillation				Mechanical		Thermal		Pro Con
	$\lambda_{em}$ (nm)	$E_g$ (eV)	$\tau$ ( $\mu$ s) @10K	$E_{scin}$ (Rel.)	Dens. (g/cc)	Mo Fraction	$T_D$ (K)	$T_M$ (C)	
CMO (CARAT)	540	3.78[1]	240	100	4.32	0.49	446	1445	High light out High melt T, difficult growing, high bkg, 48Ca
NMO-I (NIIC)	663	3.50	750	9	3.62	0.558			Good light out Cleavage plane
LMO (CUP)	535	4.26.[2]	23	5	3.03	0.562	765	705	Low melt. T, easy growing, low bkg, high $T_D$ Low light, hygroscopic
PbMoO <sub>4</sub>	592	3.20[4]	20	105	6.95	0.269			High light out Low Mo fraction, higher bkg

# Transparent front view of dry & radon clean room



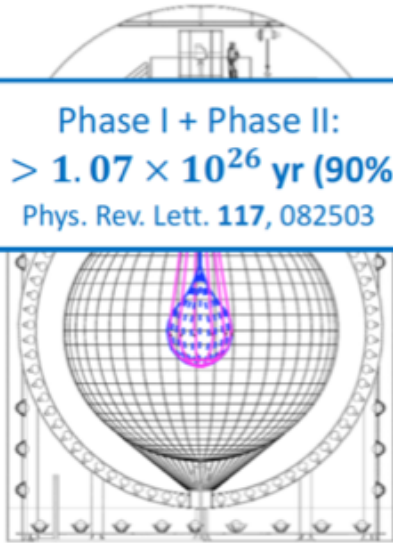
# Near future experiments

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## Evolution of KamLAND-Zen

Past

Phase I + Phase II:  
 $T_{1/2} > 1.07 \times 10^{26}$  yr (90% C.L.)  
Phys. Rev. Lett. 117, 082503



KamLAND-Zen 400

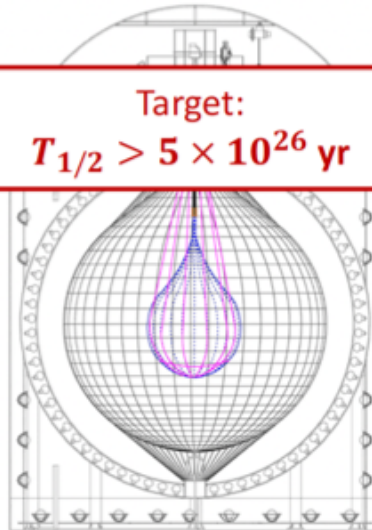
Mini-balloon Radius = 1.54 m

Xenon mass = 320 ~ 380 kg

2011 ~ 2015

Current

Target:  
 $T_{1/2} > 5 \times 10^{26}$  yr



KamLAND-Zen 800

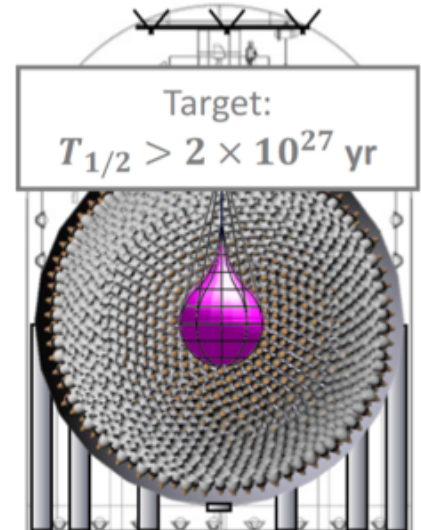
Mini-balloon Radius = 1.90 m

Xenon mass = 745 kg

Started January 2019

Future

Target:  
 $T_{1/2} > 2 \times 10^{27}$  yr



KamLAND2-Zen

Xenon mass ~ 1 ton

× 5 increase in light collection

Scintillation balloon film

# What is the origin of neutrino mass ?

Balantekin, *Ann.Rev.Nucl.Part.Sci.* 68  
(2018) 313

## Dirac Masses

- Standard model (SM) has only left-handed neutrinos and right-handed antineutrinos.
- Lepton #:  $L(\nu) = 1, L(\bar{\nu}) = -1 \rightarrow L$  is conserved.
- A Dirac neutrino mass can be generated with the same Higgs mechanism giving masses to quarks and charged leptons.
- Need to extend SM to add right-handed neutrinos,  $\nu_R$ .

$$L_D = -m_D(\bar{\nu}_R \nu_L + \bar{\nu}_L \nu_R)$$

$$\text{If } \nu \equiv \nu_L + \nu_R, \quad L_D = -m_D \bar{\nu} \nu$$

$$m_D = y^\nu \langle H^0 \rangle_0$$

$\langle H^0 \rangle_0$  : vacuum expectation value of  $H^0$ , 174 GeV

$y^\nu$ : Yukawa Coupling  $\sim 10^{-12}$

No explanation why Higgs-neutrino Yukawa coupling is so small.

## Majorana Masses

- Once a right-handed neutrino,  $\nu_R$ , is introduced, “right-handed Majorana mass term can be ;

$$L_R = -m_R/2 \left[ (\nu_R)^c \nu_R + \bar{\nu}_R (\nu_R)^c \right]$$

If  $\nu \equiv \nu_R + (\nu_R)^c$ ,  $L_R = -(m_R/2) \bar{\nu} \nu + h.c.$

Here,  $(\nu)^c = \nu \rightarrow$  Majorana particle

$$\mathcal{L}_R = -\frac{m_R}{2} \overline{(\nu_R)^c} \nu_R + h.c. ,$$

### See-Saw Mechanis...

- If we add both Dirac and Majorana mass terms, two mass eigenstates are obtained. Both are Majorana particles.



$$m_1 \simeq \frac{m_D^2}{m_R}$$

$$m_2 \simeq m_R$$

For,  $m_D = m_\mu = 10^2 \text{ MeV}$

$$m_1 = 0.1 \text{ eV}$$

$$\rightarrow m_R = 10^8 \text{ GeV}$$