Intermittency analysis in heavy ion collisions: a review of the current status and challenges.

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Outline

QCD Phase Diagram and Critical Phenomena

- 2 Intermittency analysis methodology
- Intermittency analysis results
- 4 Critical Monte Carlo Simulations
- 5 Challenges & possible solutions in intermittency analysis
- 6) Statistical significance of the signal
- 7 Conclusions & Outlook

The phase diagram of QCD



- Phase transitions from hadronic matter to quark-gluon plasma:
 - Low μ_B & high $T \rightarrow$ cross-over (lattice QCD)
 - High µ_B & low T → 1st order (effective models)
 - \Rightarrow 1st order transition line ends at Critical Point (CP) \rightarrow 2nd order transition
- At the CP: scale-invariance, universality, collective modes ⇒ good physics signatures



- Detection of the QCD Critical Point (CP): Main goal of many heavy-ion collision experiments (in particular the SPS NA61/SHINE experiment)
- Look for observables tailored for the CP; Scan phase diagram by varying energy and size of collision system.

Critical Observables & the Order Parameter (OP)



Self-similar density fluctuations near the CP



Observing power-law fluctuations through intermittency



[Csorgo, Tamas, PoS CPOD2009 (2009) 035]

Experimental observation of local, power-law distributed fluctuations of net baryon density

Intermittency in transverse momentum space at mid-rapidity (Critical opalescence in ion collisions)

[F.K. Diakonos, N.G. Antoniou and G. Mavromanolakis, PoS (CPOD2006) 010, Florence]

• Net proton density carries the same critical fluctuations as the net baryon density, and can be substituted for it.

[Y. Hatta and M. A. Stephanov, PRL91, 102003 (2003)]

• Furthermore, antiprotons can be ignored (their multiplicity is negligible compared to protons), and we can analyze just the proton density.

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Intermittency analysis review

Observing power-law fluctuations: Factorial moments

- Pioneered by Białas and others, as a method to detect non-trivial dynamical fluctuations in high energy nuclear collisions
- Transverse momentum space is partitioned into *M*² cells
- Calculate second factorial moments F₂(M) as a function of cell size ⇔ number of cells M:

$$F_2(M) \equiv \frac{\left\langle \frac{1}{M^2} \sum_{i=1}^{M^2} n_i(n_i - 1) \right\rangle}{\left\langle \frac{1}{M^2} \sum_{i=1}^{M^2} n_i \right\rangle^2},$$

where $\langle \ldots \rangle$ denotes averaging over events.

[A. Bialas and R. Peschanski, Nucl. Phys. B 273 (1986) 703-718]
 [A. Bialas and R. Peschanski, Nucl. Phys. B 308 (1988) 857-867]
 [J. Wosiek, Acta Phys. Polon. B 19 (1988) 863-869]
 [A. Bialas and R. Hwa, Phys. Lett. B 253 (1991) 436-438]
 [Z. Burda, K. Zalewski, R. Peschanski, J. Wosiek, Phys. Lett. B 314 (1993) 74-78]



 $p_{x,y}$ range in present analysis: -1.5 $\leq p_{x,y} \leq$ 1.5 GeV/c $M^2 \sim$ 10 000

Background subtraction – the correlator $\Delta F_2(M)$

Background of non-critical pairs must be subtracted from experimental data;

Partitioning of pairs into critical/background

$$\langle n(n-1)\rangle = \underbrace{\langle n_c(n_c-1)\rangle}_{\text{critical}} + \underbrace{\langle n_b(n_b-1)\rangle}_{\text{background}} + \underbrace{2\langle n_b n_c\rangle}_{\text{cross term}}$$

$$\underbrace{\Delta F_2(M)}_{\text{correlator}} = \underbrace{F_2^{(d)}(M)}_{\text{data}} - \lambda(M)^2 \cdot \underbrace{F_2^{(b)}(M)}_{\text{background}} - 2 \cdot \underbrace{\lambda(M)}_{\text{ratio}} \cdot (1 - \lambda(M)) f_{bc}$$

• If $\lambda(M) \leq 1$ (dominant background) \Rightarrow cross term negligible & $F_2^{(b)}(M) \sim F_2^{mix}(M)$ (Critical Monte Carlo* simulations) then: φ_2 : intermittency index

$$\Delta F_2(M) \simeq F_2^{\text{data}}(M) - F_2^{\text{mix}}(M)$$

Intermittency restored in $\Delta F_2(M)$:

$$\Delta F_2(M) \sim \left(M^2\right)^{\varphi_2}, \ M \gg 1$$

$$\Rightarrow$$

$$\varphi_{2,cr}^{(p)} = \frac{5}{6} (0.833...)$$

*[Antoniou et al, PRL 97, 032002 (2006)]

NA49 C+C, Si+Si, Pb+Pb @ $\sqrt{s_{NN}} \simeq 17$ GeV – dipions

- 3 sets of NA49 collision systems at 158A GeV/c (√s_{NN} ≃ 17 GeV)
 [T. Anticic *et al*, Phys. Rev. C 81, 064907 (2010); T. Anticic *et al*., Eur. Phys. J. C 75:587 (2015)]
- Intermittent behaviour (φ₂^(σ) ≃ 0.35) of dipion pairs (π⁺, π⁻) in transverse momentum space observed in central Si+Si collisions at 158A GeV.



[T. Anticic et al, Phys. Rev. C 81, 064907 (2010)]

 No such power-law behaviour observed in central C+C and Pb+Pb collisions at the same energy.

NA49 C+C, Si+Si, Pb+Pb @ $\sqrt{s_{NN}} \simeq 17$ GeV – protons

Factorial moments of proton transverse momenta analyzed at mid-rapidity



F₂(M), ΔF₂(M) errors estimated by the bootstrap method
 [W.J. Metzger, "Estimating the Uncertainties of Factorial Moments", HEN-455 (2004).]

• Fit with
$$\Delta F_2^{(e)}(M \ ; \ C, \phi_2) = 10^C \cdot \left(\frac{M^2}{M_0^2}\right)^{\phi_2}$$
, for $M^2 \ge 6000 \ (M_0^2 \equiv 10^4)$

Evidence for intermittency in "Si"+Si – but large statistical errors.

NA61/SHINE intermittency: Be+Be @ $\sqrt{s_{NN}} \simeq 17$ GeV

- Intermittency analysis is pursued within the framework of the NA61/SHINE experiment, inspired by the positive, if ambiguous, NA49 Si+Si result.
 [T. Anticic et al., Eur. Phys. J. C 75:587 (2015), arXiv:1208.5292v5]
- Two NA61/SHINE systems were initially examined:
 ⁷Be + ⁹Be and ⁴⁰Ar + ⁴⁵Sc @ 150A GeV/c (√s_{NN} ~ 17 GeV)



- $F_2(M)$ of data and mixed events overlap \Rightarrow
- Subtracted moments ΔF₂(M) fluctuate around zero ⇒
 No intermittency effect is observed in Be+Be.

NA61/SHINE 40 Ar + 45 Sc @ $\sqrt{s_{NN}} \simeq 17$ GeV

- First indication of intermittency in mid-central Ar+Sc 150A GeV/c collisions presented at CPOD2018; In 2019, an extended event statistics set was analysed;
- A scan in centrality was performed, in the 0-20% range, in 5% and 10% intervals, as centrality may influence the system's freeze-out temperature;
- Event statistics were of the order of ~ 400K events per 10% centrality interval;
- Bootstrap confidence intervals are calculated for $\Delta F_2(M)$ values;
- Due to M-bin correlations, determining confidence intervals for φ₂ is challenging; Various approaches to the problem are being investigated, such as model-weighting;
- Ar+Sc system is still inconclusive.

NA61/SHINE Ar+Sc @150A GeV/c: 5% cent. intervals



NA61/SHINE Ar+Sc 150, cent.5 - 10%, pur > 90%



No signal in c.0-5%, 5-10%; weak signal in c.10-15%, 15-20%.

NA61/SHINE Ar+Sc @150A GeV/c: 10/20% cent. intervals



• Centrality dependence is evident; Some signal indication in c.10-20%.

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NA61/SHINE preliminary

Critical Monte Carlo (CMC) algorithm for baryons

- Simplified version of CMC* code:
 - Only protons produced
 - One cluster per event, produced by random Lévy walk:

 $\tilde{d}_F^{(B,2)} = 1/3 \Rightarrow \phi_2 = 5/6$

- Lower / upper bounds of Lévy walks *p_{min,max}* plugged in.
- Cluster center exponential in p_T, slope adjusted by T_c parameter.
- Poissonian proton multiplicity distribution.



Input parameters

*[Antoniou, Diakonos, Kapoyannis and Kousouris, Phys. Rev. Lett. 97, 032002 (2006).]

Be+Be - Critical Monte Carlo toy model



- Critical Monte Carlo (CMC) + random background in transverse momentum space;
- $\Delta F_2(M)$ retains critical behaviour of pure CMC ($\lambda = 0$), even when their moments differ by orders of magnitude!
- Preliminary analysis with CMC simulation indicates an upper limit of ~ 0.3% critical protons (λ ≈ 0.997) [PoS(CPOD2017) 054]
- CMC results show our approximation (dominant background) is reasonable.

Ar+Sc - Critical Monte Carlo



 Preliminary analysis with CMC simulation indicates an upper limit of ~ 0.5% critical protons

Challenges in proton intermittency analysis

- Particle species, especially protons, cannot be perfectly identified experimentally; candidates will always contain a small percentage of impurities;
- Experimental momentum resolution sets a limit to how small a bin size (large M) we can probe;
- A finite (small) number of usable events is available for analysis; the "infinite statistics" behaviour of ΔF₂(M) must be extracted from these;
- Proton multiplicity for medium-size systems is low (typically ~ 2 3 protons per event, in the window of analysis) and the demand for high proton purity lowers it still more;
- M-bins are correlated the same events are used to calculate all F₂(M)! This biases fits for the intermittency index φ₂, and makes confidence interval estimation hard.

Intermittency analysis tools: the bootstrap

- Random sampling of events, with replacement, from the original set of events;
- k bootstrap samples (k ~ 1000) of the same number of events as the original sample;
- Each statistic (ΔF₂(M), φ₂) calculated for bootstrap samples as for the original; [B. Efron, *The Annals of Statistics* 7,1 (1979)]
- Variance of bootstrap values estimates standard error of statistic.





Intermittency analysis tools: correlated fit

Possible to perform correlated fits for φ₂, with *M*-correlation matrix estimated via bootstrap;

Correlated fit

NA61/SHINE Ar+Sc 150, cent.10 - 20%, pur > 90%

Uncorrelated fit NA61/SHINE Ar+Sc 150, cent.10 - 20%, pur > 90%



- Replication of events means bootstrap sets are not independent of the original: magnitude of variance and covariance estimates can be trusted, but central values will be biased to the original sample;
- Correlated fits for \$\phi_2\$ are known to be unstable;

[B. Wosiek, APP B21, 1021 (1990); C. Michael, PRD49, 2616 (1994)]

Intermittency analysis tools: the AMIAS scheme

Avoid fitting, use model weighting!

• AMIAS algorithm

(A Model-Independent Analysis Scheme) can be used to extract **model-parameter distributions** such as ϕ_2 from sets of data;

[C. N. Papanicolas, E. Stiliaris, arXiv:1205.6505 (2012);

AIP Conf. Proc. 904, 257 (2007)]

• It works by sampling parameter space at random, then weighting selected models by goodness-of-fit function $(e^{-\chi^2/2})$ to the dataset.

Model:
$$\Delta F_2(M; a_0, \phi_2) = 10^{a_0} \left(\frac{M^2}{10^4}\right)^{\phi_2}$$



[N. G. Antoniou et al, Nucl. Phys. A 1003 122018 (2020).]

AMIAS scheme example: Critical Monte Carlo

CMC Ar+Sc 150, cent.10 - 20%, bkg. = 99.3%



- Testing millions of possible solutions on 400 independent F₂(M) samples!
- Also works on **bootstrap** samples (although **not independent**).



[N. G. Antoniou et al, Nucl. Phys. A 1003 122018 (2020).]

Model:
$$\Delta F_2(M; a_0, \phi_2) = 10^{a_0} \left(\frac{M^2}{10^4}\right)^{\phi_2}$$

Ar+Sc 150 $\Delta F_2(M)$ – statistical significance of signal

- Plan: compare ΔF₂(M) bootstrap distribution of Ar+Sc data to uncorrelated proton moments of the same event statistics.
- $\Delta F_2(M)$, NA61/SHINE Ar+Sc @ 150 GeV/c bootstrap distributions

2 $\Delta F_2(M)$, random background sub-sample distributions



• ~ 85 – 95% of $\Delta F_2(M)$ values above zero in Ar+Sc 150

2 ~ 85 – 95% of random background $\Delta F_2(M)$ values below Ar+Sc 150 average

Roughly 5 – 15% chance of random background producing a spurious "effect"

Conclusions & Outlook

- Intermittency analysis of proton density is a promising strategy for detecting the Critical Point;
- However, this analysis is challenging in the context of an actual heavy-ion collision experiment, always constrained in terms of available statistics, particle multiplicity, and proton identification;
- New techniques were developed to better determine $\Delta F_2(M)$ and ϕ_2 uncertainties (bootstrap errors, AMIAS weighting);
- Detailed exploration of refined models with critical & non-critical components is certainly needed, in order to assess experimental data;
- Analysis of different systems and collision energies is ongoing.



Thank You!



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Backup Slides

- NA49 intermittency results
- NA61/SHINE intermittency results
- Oritical Monte Carlo
- Intermittency analysis challenges
- 2 Remedies to intermittency problems

NA49 C+C, Si+Si, Pb+Pb @ $\sqrt{s_{NN}} \simeq 17$ GeV



• No intermittency detected in the "C"+C, Pb+Pb datasets.

NA49 C+C, Si+Si, Pb+Pb @ $\sqrt{s_{NN}} \simeq 17$ GeV

- Evidence for intermittency in "Si"+Si but large statistical errors.
- Distribution of φ_2 values, $P(\varphi_2)$, and confidence intervals for φ_2 obtained by fitting individual bootstrap samples [B. Efron, *The Annals of Statistics* 7,1 (1979)]



- Bootstrap distribution of \$\phi_2\$ values is highly asymmetric (due to closeness of \$F_2^{(d)}(M)\$ to \$F_2^{(m)}(M)\$).
- Uncorrelated fits used, but errors between M are correlated!
- Estimated intermittency index: $\phi_{2,B} = 0.96^{+0.38}_{-0.25}$ (stat.) ± 0.16 (syst.)

[T. Anticic et al., Eur. Phys. J. C 75:587 (2015), arXiv:1208.5292v5]

The NA61/SHINE experiment



- Fixed-target, high-energy collision experiment at CERN SPS;
- Reconstruction & identification of emitted protons in an extended regime of rapidity, with precise evaluation of their momentum vector;
- Centrality of the collision measured by a forward Projectile Spectator Detector (PSD);

- Direct continuation of NA49
- Search for Critical Point signatures



beam momentum [A GeV/c]

Independent bin analysis with cumulative variables

- M-bin correlations complicate uncertainties estimations for ΔF₂(M) & φ₂; one way around this problem is to use independent bins a different subset of events is used to calculate F₂(M) for each M;
- Advantage: correlations are no longer a problem;
 Disadvantage: we break up statistics, and can only calculate F₂(M) for a handful of bins.
- Furthermore, instead of p_x and p_y, one can use cumulative quantities: [Bialas, Gazdzicki, PLB 252 (1990) 483]

$$Q_x(x) = \int_{min}^{x} P(x)dx \Big| \int_{min}^{max} P(x)dx;$$
$$Q_y(x, y) = \int_{ymin}^{y} P(x, y)dy \Big| P(x)$$

- transform any distribution into **uniform** one (0, 1);
- remove the dependence of F₂ on the shape of the single-particle distribution;
- approximately preserves ideal power-law correlation function. [Antoniou, Diakonos, https://indico.cern.ch/event/818624/]



Pb+Pb @ 30 GeV/c analysis ($\sqrt{s_{NN}} \simeq 7.6$ GeV)

- NA61/SHINE presently undergoes the effort of a concentrated analysis of new Pb + Pb data at the energies of $\sqrt{s_{NN}} = 5.1 7.6$ GeV, to verify the latest STAR results (arXiv: 2001.02852) claimed as a possible signal of the Critical Point;
- So far, preliminary analysis shows no indication of intermittency in Pb+Pb.



Exclusion plots: Ar+Sc, Pb+Pb

[T. Czopowicz, Search for critical point via intermittency analysis in NA61/SHINE, C.P.O.D. 2021 Online]

Independent bin analysis – Ar+Sc & Pb+Pb results



[T. Czopowicz, Search for critical point via intermittency analysis in NA61/SHINE, C.P.O.D. 2021 Online]

Simulating fractal sets through random Lévy walks

 In D-dimensional space, we can simulate a fractal set of dimension d_F, D - 1 < d_F < D, through a random walk with step size Δr distribution:



Proton selection

NA61/SHINE preliminary

NA61/SHINE preliminary



- Particle ID through energy loss dE/dx in the Time Projection Chambers (TPCs);
- Employ p_{tot} region where Bethe-Bloch bands do not overlap (3.98 GeV/c ≤ p_{tot} ≤ 126 GeV/c);
- Mid-rapidity region ($|y_{CM}| < 0.75$) selected for present analysis.

Momentum resolution: effect on intermittency





- CMC + background + Gaussian noise (5 MeV radius);
- A 5 MeV Gaussian error in p_x, p_y leads to ~ 10% discrepancy in the value of φ₂.
- For very large backround values (> 99%), momentum resolution matters little to the overall distortion.

AMIAS on NA49 & NA61/SHINE data – ϕ_2 vs $N_{wounded}$

- φ₂ AMIAS confidence intervals calculated for NA49 & NA61/SHINE systems with indications of intermittency
- Corresponding mean number of participating ("wounded") nucleons N_w estimated via geometrical Glauber model simulation



- Peripheral Ar+Sc collisions approach Si + Si criticality ⇒ insight of how the critical region looks as a function of baryon density μ_B.
- Check theoretical predictions* for narrow critical scaling region in *T* & μ_B

*[F. Becattini *et. al.*, arXiv:1405.0710v3 [nucl-th] (2014); N. G. Antoniou, F. K. Diakonos,

arXiv:1802.05857v1 [hep-ph] (2018)]