## Cluster Algebras for Feynman Integrals

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## CLUSTER OF EXCELLENCE QUANTUM UNIVERSE

## CORFU2021 Workshop on the Standard Model and Beyond August 24, 2021

## Scattering amplitudes

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High precision calculations crucial for ${ }^{\text {[Canko Talk] }}$

- Determining Standard Model parameters ${ }^{\text {[De Roeck, Fayard Talks] }}$
- Tell apart New Physics from Standard Model background Especially relevant in light of High-Luminosity LHC 2027-2037. [Zerlauth Talk]

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- Canonical Differential Equations ${ }^{[H e n n]}$

All of them vital for recent state of the art calculation of 2-loop 5-point 1-mass planar master integrals, relevant for W -boson production +2 jets, [Abreu,Ita,Moriello,Page,Tschernow,Zeng] [Canko,Papadopoulos,Syrrakos]


Image Credit: 2005.04195

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Discovered cluster algebras encode singularities of a wealth of physically relevant examples, including QCD corrections to amplitudes for $\mathrm{pp} \rightarrow$ Higgs+jet!

## Outline

Introduction: Cluster Algebras and $\mathcal{N}=4$ SYM

Cluster Algebras for Feynman Integrals
$C_{2} \&$ Higgs amplitudes
Further Examples

Conclusions \& Outlook

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Exchange graph: Clusters=vertices, mutations=edges

## Geometric Interpretation of Cluster Algebras

Finite cluster algebras classified by Dynkin diagrams. For $A_{n}$ :

- Cluster $=$ triangulation of $(n+3)$-gon by noncrossing diagonals
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Example: $A_{3}=$ hexagon

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Example: $A_{3}=$ hexagon exchange graph


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$f_{k}$ is a MPL of weight $k$ if its differential obeys

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Collection of $\phi_{\alpha_{i}}$ : symbol alphabet $\Phi$

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[PoS CORFU2019 Review:Caron-Huot,Dixon,Drummond,Dulat,Foster, Gurdogan,Hippel,McLeod,GP]

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- $n=6: 7$ loops (MHV)
- $n=7$ : 4 loops


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E.g. $A_{3}$ : Crossing diagonals forbidden,
$\ldots \otimes(15) \otimes(26) \otimes \ldots, \ldots \otimes(14) \otimes(26) \otimes \ldots, \ldots \otimes(14) \otimes(36) \otimes \ldots$


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For physical $n=6,7$ functions, equivalent to extended Steinmann relations. Massively reduces size of function space.
[Caron-Huot,Dixon,DulatMcLeod,Hippel, GP]

The Genetic Material of $\mathcal{N}=4$ SYM Amplitudes
Amplitude symbol

The Genetic Material of $\mathcal{N}=4$ SYM Amplitudes
DNA (2)

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## DNA

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Bases A,T,G,C

Amplitude symbol


Letters=cluster variables

The Genetic Material of $\mathcal{N}=4$ SYM Amplitudes

## DNA



Bases A,T,G,C
Base pairs $A-T, G-C$

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Could cluster algebras provide the genetic material of generic quantum field theories?

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Look at scalar Feynman integrals with massless internal propagators in dimensional regularization.

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For any given process, scalar integrals related by integration-by-parts identities. Basis in the vector space they span=master integrals.
[Chetyrkin, Tkachov]

Main Example: Four-point functions with one leg offshell/massive


$$
p_{i}^{2}=0, \quad P^{2} \neq 0
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- Kinematic variables:

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z_{1} \equiv \frac{2 p_{1} \cdot p_{2}}{P^{2}}, \quad z_{2} \equiv \frac{2 p_{2} \cdot p_{3}}{P^{2}}, \quad z_{3} \equiv \frac{2 p_{1} \cdot p_{3}}{P^{2}}
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- Alphabet of all known master integrals: [Gehrmann,Remiddi]
[Di Vita, Mastrolia, Schubert, Yundin]

$$
\Phi_{2 \mathrm{dHPL}}=\left\{z_{1}, z_{2}, z_{3}, 1-z_{1}, 1-z_{2}, 1-z_{3}\right\},
$$

"2-dimensional HPLs" [Gehrmann,Remiddi]

## Identifying Candidate Cluster Algebras

\# independent variables $\left|\begin{array}{c}2 \mathrm{dHPL} \\ 2\end{array}\right|$

## Identifying Candidate Cluster Algebras

\# independent variables | 2 dHPL | $A_{2}$ |
| :---: | :---: | :---: |
| 2 | 2 |

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|  | 2 dHPL | $A_{2}$ |  |
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| \# independent variables | 2 | 2 |  |
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## Identifying Candidate Cluster Algebras

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| 109 | 109 | $\odot$ |

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$$

$$
\Phi_{C_{2}}=\left\{a_{1}, a_{2}, 1+a_{1}, 1+a_{2}^{2}, 1+a_{1}+a_{2}^{2}, 1+2 a_{1}+a_{1}^{2}+a_{2}^{2}\right\} .
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$$
2 \mathrm{dHPLs}=C_{2} \text { polylogarithms! }
$$

## Physical Significance

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Physical Significance: $C_{2}$ Cluster Algebra Underlies Higgs Amplitudes!
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- $p p \rightarrow$ Higgs + jet [Gehrmann, Jaquier, Glover, Koukoutsakis] [Duhr] in heavy top mass limit



## What Do Cluster Algebras Buy Us?

Adjacency properties of Feynman integrals
To analyze this, instructive to recall evaluation of master integrals via differential equations: ${ }^{[K o t i k o v][G e h r m a n n, R e m i d d i] ~}$

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d \mathbf{f}(\vec{z} ; \epsilon)=d \mathbf{M}(\vec{z} ; \epsilon) \cdot \mathbf{f}(\vec{z} ; \epsilon)
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\tilde{\mathbf{A}}=\sum_{i} \mathbf{A}_{i} \log \alpha_{i}(\vec{z})
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Adjacency property significantly reduces size of $C_{2}$ symbol space:

| weight | 1 | 2 | 3 | 4 | 5 | 6 | 7 | 8 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| First entry condition | 3 | 12 | 45 | 165 | 597 | 2143 | 7653 | 27241 |
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Application: In parallel work, used to bootstrap $\mathcal{N}=4$ SYM analogue of Higgs amplitude through 5 loops

More examples of cluster algebras from one-loop Feynman integrals

| family | \# variables | \# letters | cluster algebra |
| :---: | :---: | :---: | :---: |
|  | 3 | 9 | $\subset A_{3}$ |
|  | 3 | 10 | $\subset C_{3}$ |
|  | 4 | 16 | $\subset C_{4}$ |

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Finite 6D integrals | \# variables | \# letters | cluster algebra |  |
| :---: | :---: | :---: | :---: |
|  | 4 | 16 | $D_{4}$ |
|  | 5 | 24 | $\subset D_{5}$ |
|  | 5 | 27 | $\lim \operatorname{Tr}(4,8)$ |

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from relation of $G r(4, n)$ with (dual of) tropical Grassmanian $\operatorname{Tr}(4, n)$.



From appropriately symmetrized limit, obtain all letters appearing in finite (hard) part of two-loop five-gluon amplitude in QCD!

## Conclusions

The beautiful mathematics of cluster algebras underlie the analytic structure of several physically relevant Feynman integrals \& processes!

- Higgs+jet amplitudes to all orders in $\epsilon$
- 5-gluon planar amplitudes in QCD to finite part
- Reveal new, potentially useful properties such as adjacency


## Next Stage

Very recently, predictions for symbol letters of $\mathcal{A}_{9}$ in $\mathcal{N}=4$ SYM. ${ }^{[H e n k e,}{ }^{\text {GP'21] }} \Rightarrow$ Limit to six-gluon amplitude letters in QCD?

- More examples \& first-principle proof? Bootstrap of QCD quantities?


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> Could cluster algebras and generalizations provide an organizing principle ("genetic material") that simplifies future collider physics calculations?

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- In general, variables $a_{m}$ of a Grassmannian $\operatorname{Gr}(4, n)$ cluster algebra! $G r(4,6) \simeq A_{3}, G r(4,7) \simeq E_{6}{ }^{\text {[Golden,Goncharov,Spradlin, Vergu, Volovich] }}$

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Emerge when parametrizing planar $n$-particle massless kinematics in terms of $n$ momentum twistors $Z_{i}$ on $\mathbb{C P}^{3}\left(Z_{i} \sim \lambda Z_{i}\right)$,


$$
\begin{gathered}
p_{i} \equiv x_{i+1}-x_{i}, \quad x_{i} \sim Z_{i-1} \wedge Z_{i} \\
\left(x_{i}-x_{j}\right)^{2} \equiv\left(p_{j}+\ldots+p_{i-1}\right)^{2} \sim \operatorname{det}\left(Z_{i-1} Z_{i} Z_{j-1} Z_{j}\right) \equiv\langle i-1 i j-1 j\rangle
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Cluster $\mathcal{A}$-coordinates $a_{m}$ : Certain homogeneous polynomials of $\langle i j k l\rangle$

Application: The Steinmann Cluster Bootstrap for $\mathcal{N}=4$ SYM Amplitudes Evade Feynman diagrams by exploiting analytic structure

Computation

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| Physical Branch Cuts <br> [Gaiotto,Maldacena, <br> Sever,Vieira] | $\mathcal{A}_{6}^{(L)}, L=3,4$ |

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| Cluster Algebras | $\mathcal{A}_{7, \text { MHV }}^{(3)}$ |
| [Golden, Goncharov, <br> Spradlin,Vergu,Volovich] | [Drummond, GP, <br> Spradlin] |

$$
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\mathcal{A}_{\mathrm{MHV}} & =\mathcal{A}(--+\ldots+) \\
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| [Steinmann] | [Caron-Huot,Dixon,...] <br> [Dixon,..., GP,Spradlin] |
| Cluster Adjacency | $\mathcal{A}_{7, \mathrm{NMHV}}^{(4)}$ |
| [Drummond,Foster, Gurdogan] | [Drummond,Foster, Gurdogan, GP] |
| Extended Steinmann | $\Leftrightarrow \quad \mathcal{A}_{6}^{(6)}, \mathcal{A}_{6, \mathrm{MHV}}^{(7)}$ |
| Coaction Principle | [Caron-Huot,Dixon,Dulat, McLeod,Hippel,GP] |

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See also $S\left(A_{n}^{(2)}\right) \rightarrow A_{n}^{(2)}, S\left(\mathcal{A}_{7}\right) \rightarrow \mathcal{A}_{7}$ work

## Steinmann Relations

Provide physical backing to cluster adjacency, forbidding double discontinuities in overlapping channels. For $A_{3}$, equivalent to

$$
(14) \otimes(36) \otimes \ldots
$$


$v s$.


## Extended Steinmann Relations

## [Caron-Huot,Dixon,DulatMcLeod,Hippel,GP]

Provide physical backing to cluster adjacency, forbidding multiple discontinuities in overlapping channels. For $A_{3}$, equivalent to

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## Extended Steinmann Relations

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- For physical $n=6,7$ functions, imply remaining cluster adjacency restrictions.
- Recently confirmed in planar 5-pt 1-mass master integrals. [Abreu,Ita,Moriello,Page,Tschernow,Zeng]

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& \cdot\left(x_{i+i}-x_{i}\right)^{2}=0 \quad \Rightarrow X_{i}=Z_{i-1} \wedge Z_{i}
\end{aligned}
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## $\operatorname{Conf}_{n}\left(\mathbb{P}^{3}\right)$ and Graßmannians

Can realize $\operatorname{Conf}_{n}\left(\mathbb{P}^{3}\right)$ as $4 \times n$ matrix

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\left(Z_{1}\left|Z_{2}\right| \ldots \mid Z_{n}\right)
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- Under $G L(k)$ transformations, basis vectors change, but still span the same plane.
Comparing the two matrices,

$$
\operatorname{Conf}_{n}\left(\mathbb{P}^{3}\right)=G r(4, n) /\left(C^{*}\right)^{n-1}
$$

## The Rules of the Mutation Game

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Mutation on $a_{k}: \quad a_{k}^{\prime}=a_{k}^{-1}\left(\prod_{i=1}^{d+m} a_{i}^{\left[b_{i k}\right]_{+}}+\prod_{i=1}^{d+m} a_{i}^{\left[-b_{i k}\right]_{+}}\right), k \leq n$,
all other $a_{j}$ unchanged. $b_{i j}$ elements of $B$, which itself mutates as

$$
b_{i j}^{\prime}= \begin{cases}-b_{i j} & \text { for } i=k \text { or } j=k \\ b_{i j}+\left[-b_{i k}\right]_{+} b_{k j}+b_{i k}\left[b_{k j}\right]_{+} & \text {otherwise }\end{cases}
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with $[x]_{+}=\max (0, x)$.

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## Simplest case:



$$
p_{i}^{2}=0,
$$

Single kinematic variable: $z \equiv \frac{p_{2} \cdot p_{3}}{p_{1} \cdot p_{2}}$
Alphabet of known integrals:
[Henn,Smirnov ${ }^{2}$ ][Panzer] [Henn,Mistlberger,Smirnov,Wasser]

$$
\Phi=\{z, 1+z\}
$$

(Nonpositive) Harmonic Polylogarithms [Remiddi,Vermaseren]

Instead, consider next-to simplest case as our main example.

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Can the $\mathcal{N}=4$ world help us understand the observed $C_{2}$ adjacency?

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- Nevertheless, $C_{2}$ is parity (=up-down reflection) invariant surface of $A_{3}$ !


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$C_{2}$ adjacency $=$ extended Steinmann relations for $\mathcal{A}_{6}$ !

Five-particle scattering from $p \operatorname{Tr}(4,8)$ ।

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Specifically, 27 letters of $D=6$ 2-mass hard hexagon $=1$-loop 1-mass pentagon alphabet except $\Delta_{3}=\lambda\left(P^{2}, s_{23}, s_{45}\right), \Delta_{5}=\left.\operatorname{det}\left(2 p_{i} \cdot p_{j}\right)\right|_{i, j \leq 4}$ +8 2-loop letters ( $\sqrt{\Delta_{5}}$ rationalized by mom.twistors).

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In particular, single missing letter observed to drop out from finite (hard) part of five-gluon amplitude in QCD! $\left(\sqrt{\Delta}=\left.\sqrt{\operatorname{det}\left(2 p_{i} \cdot p_{j}\right)}\right|_{i, j \leq 4}\right)$

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- Access to Lorentz-invariant 2-mass pentagon or 1-mass hexagon?

