Cluster Algebras for Feynman Integrals

Georgios Papathanasiou



#### **CLUSTER OF EXCELLENCE**

QUANTUM UNIVERSE

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The arena where perturbative quantum field theory confronts experiment. High precision calculations crucial for <sup>[Canko Talk]</sup>





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- Tell apart New Physics from Standard Model background





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Especially relevant in light of High-Luminosity LHC 2027-2037. [Zerlauth Talk]

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All of them vital for recent state of the art calculation of 2-loop 5-point 1-mass planar master integrals, relevant for W-boson production + 2 jets, [Abreu,Ita,Moriello,Page,Tschernow,Zeng][Canko,Papadopoulos,Syrrakos]



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Discovered cluster algebras encode singularities of a wealth of physically relevant examples, including QCD corrections to amplitudes for  $pp \rightarrow Higgs+jet!$ 

# Outline

Introduction: Cluster Algebras and  $\mathcal{N}$  = 4 SYM

### Cluster Algebras for Feynman Integrals

 $C_2$  & Higgs amplitudes Further Examples

Conclusions & Outlook

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Exchange graph: Clusters=vertices, mutations=edges

Finite cluster algebras classified by Dynkin diagrams. For  $A_n$ :

• Cluster = triangulation of (n + 3)-gon by noncrossing diagonals

 $\langle n \rangle$ 

Cluster coordinates = diagonals of this triangulation

Example:  $A_3 = hexagon$ 

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 $f_k$  is a MPL of weight k if its differential obeys

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Collection of  $\phi_{\alpha_i}$ : symbol alphabet  $\Phi$ 

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E.g.  $A_3$ : Crossing diagonals forbidden, ...  $\otimes$  (15)  $\otimes$  (26)  $\otimes$  ..., ...  $\otimes$  (14)  $\otimes$  (26)  $\otimes$  ..., ...  $\otimes$  (14)  $\otimes$ 

6

2

3



5

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For physical n = 6, 7 functions, equivalent to extended Steinmann relations. Massively reduces size of function space. [Caron-Huot,Dixon,DulatMcLeod,Hippel,GP]

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Bases A, T, G, C

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Could cluster algebras provide the genetic material of generic quantum field theories?

Look at scalar Feynman integrals with massless internal propagators in dimensional regularization.

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$$p_{2} = \int \frac{d^{D}k}{k^{2}(k+p_{1})^{2}(k+p_{1}+p_{2})^{2}(k-P)^{2}}$$

For any given process, scalar integrals related by integration-by-parts identities. Basis in the vector space they span=master integrals. [Chetyrkin,Tkachov]
Main Example: Four-point functions with one leg offshell/massive



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Kinematic variables:

$$z_1 \equiv \frac{2p_1 \cdot p_2}{P^2}$$
,  $z_2 \equiv \frac{2p_2 \cdot p_3}{P^2}$ ,  $z_3 \equiv \frac{2p_1 \cdot p_3}{P^2}$ ,

with  $z_1 + z_2 + z_3 = 1$ .

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Alphabet of all known master integrals: <sup>[Gehrmann,Remiddi]</sup>
[Di Vita, Mastrolia, Schubert, Yundin]

$$\Phi_{2\mathsf{dHPL}} = \left\{ z_1, z_2, z_3, 1 - z_1, 1 - z_2, 1 - z_3 \right\},\$$

"2-dimensional HPLs" [Gehrmann, Remiddi]

# independent variables 2

$$\begin{array}{c|c} & 2 \text{dHPL} & A_2 \\ \# \text{ independent variables} & 2 & 2 \end{array}$$

	2dHPL	$A_2$	
# independent variables	2	2	
# letters	6	5	٢

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$$\Phi_{2\mathsf{dHPL}} = \{z_1, z_2, 1 - z_1 - z_2, 1 - z_1, 1 - z_2, z_1 + z_2\},\$$

$$\Phi_{C_2} = \left\{a_1, a_2, 1 + a_1, 1 + a_2^2, 1 + a_1 + a_2^2, 1 + 2a_1 + a_1^2 + a_2^2\right\}.$$

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 $2dHPLs = C_2$  polylogarithms!

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## Physical Significance: C<sub>2</sub> Cluster Algebra Underlies Higgs Amplitudes!

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- ▶  $pp \rightarrow Z$ -boson + jet [Gehrmann,Tancredi,Weihs]
- *pp* → Higgs + jet [Gehrmann, Jaquier, Glover, Koukoutsakis] [Duhr]

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to all orders in  $\epsilon$ . Important structural information for manipulating f, e.g. analytic continuation.



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Adjacency property significantly reduces size of  $C_2$  symbol space:

weight	1	2	3	4	5	6	7	8
First entry condition	3	12	45	165	597	2143	7653	27241
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Application: In parallel work, used to bootstrap  $\mathcal{N} = 4$  SYM analogue of Higgs amplitude through 5 loops <sup>[Dixon,Mcleod,Wilhelm]</sup>

#### More examples of cluster algebras from one-loop Feynman integrals

family	# variables	# letters	cluster algebra
	3	9	$\subset A_3$
	3	10	$\subset C_3$
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Finite 6D integrals	# variables	# letters	cluster algebra
	4	16	$D_4$
$\rightarrow$	5	24	$\subset D_5$
	5	27	$\lim \operatorname{Tr}(4,8)$

From  $n\text{-}\mathsf{particle}$  amplitudes in  $\mathcal{N}$  = 4 SYM know they cannot be the end of story:

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From appropriately symmetrized limit, obtain all letters appearing in finite (hard) part of two-loop five-gluon amplitude in QCD!

## Conclusions

The beautiful mathematics of cluster algebras underlie the analytic structure of several physically relevant Feynman integrals & processes!

- $\blacktriangleright$  Higgs+jet amplitudes to all orders in  $\epsilon$
- 5-gluon planar amplitudes in QCD to finite part
- Reveal new, potentially useful properties such as adjacency

# Next Stage

Very recently, predictions for symbol letters of  $A_9$  in  $\mathcal{N} = 4$  SYM. <sup>[Henke,GP'21]</sup>  $\Rightarrow$  Limit to six-gluon amplitude letters in QCD?

More examples & first-principle proof? Bootstrap of QCD quantities?

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Could cluster algebras and generalizations provide an organizing principle ("genetic material") that simplifies future collider physics calculations?
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Emerge when parametrizing planar *n*-particle massless kinematics in terms of *n* momentum twistors  $Z_i$  on  $\mathbb{CP}^3$  ( $Z_i \sim \lambda Z_i$ ), <sup>[Hodges]</sup>



$$p_{i} \equiv x_{i+1} - x_{i}, \quad x_{i} \sim Z_{i-1} \wedge Z_{i}$$
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Cluster A-coordinates  $a_m$ : Certain homogeneous polynomials of  $\langle ijkl \rangle$ 

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QFT Property Computation

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$\mathcal{A}$	

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Cluster Adjacency	${\cal A}_{7,{\sf NMHV}}^{(4)}$
[Drummond,Foster, Gurdogan]	[Drummond,Foster, Gurdogan, GP]
Extended Steinmann	$\Leftrightarrow  \mathcal{A}_{6}^{(6)}, \mathcal{A}_{6,MHV}^{(7)}$
Coaction Principle	[Caron-Huot,Dixon,Dulat, McLeod,Hippel,GP]

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See also  $S(A_n^{(2)}) \to A_n^{(2)}$ ,  $S(\mathcal{A}_7) \to \mathcal{A}_7$  work [Golden, Paulos, Spradlin, Volovich)] [Dixon, Liu] [Golden, McLeod]

#### Steinmann Relations

Provide physical backing to cluster adjacency, forbidding double discontinuities in overlapping channels. For  $A_3$ , equivalent to



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- Recently confirmed in planar 5-pt 1-mass master integrals. [Abreu,Ita,Moriello,Page,Tschernow,Zeng]

# Momentum Twistors $Z^{I \ [\mathrm{Hodges}]}$

▶ Represent dual space variables  $x^{\mu} \in \mathbb{R}^{1,3}$  as projective null vectors  $X^{M} \in \mathbb{R}^{2,4}$ ,  $X^{2} = 0$ ,  $X \sim \lambda X$ .

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Can realize  $\operatorname{Conf}_n(\mathbb{P}^3)$  as  $4 \times n$  matrix

 $(Z_1|Z_2|\ldots|Z_n)$ 

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Comparing the two matrices,

$$\operatorname{Conf}_n(\mathbb{P}^3) = Gr(4,n)/(C^*)^{n-1}$$

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Mutation on 
$$a_k$$
:  $a'_k = a_k^{-1} \left( \prod_{i=1}^{d+m} a_i^{[b_{ik}]_+} + \prod_{i=1}^{d+m} a_i^{[-b_{ik}]_+} \right), \ k \le n,$ 

all other  $a_j$  unchanged.  $b_{ij}$  elements of B, which itself mutates as

$$b'_{ij} = \begin{cases} -b_{ij} & \text{for } i = k \text{ or } j = k \\ b_{ij} + \left[ -b_{ik} \right]_+ b_{kj} + b_{ik} \left[ b_{kj} \right]_+ & \text{otherwise} \end{cases},$$

with  $[x]_{+} = \max(0, x)$ .

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 $p_i^2 = 0,$ Single kinematic variable:  $z \equiv \frac{p_2 \cdot p_3}{p_1 \cdot p_2}$ Alphabet of known integrals: [Henn,Smirnov<sup>2</sup>][Panzer][Henn,Mistlberger,Smirnov,Wasser]  $\Phi = \{z, 1 + z\}$ 

(Nonpositive) Harmonic Polylogarithms [Remiddi,Vermaseren]

Instead, consider next-to simplest case as our main example.

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• Nevertheless,  $C_2$  is parity (=up-down reflection) invariant surface of  $A_3!$ 



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P(*ij*) = P(*i*+3, *j*+3), e.g. {(35), (36), (26)} → {(36), √(26)(35)}

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 $C_2$  adjacency = extended Steinmann relations for  $\mathcal{A}_6$ !

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Specifically, 27 letters of D = 6 2-mass hard hexagon = 1-loop 1-mass pentagon alphabet except  $\Delta_3 = \lambda(P^2, s_{23}, s_{45}), \Delta_5 = \det(2p_i \cdot p_j)|_{i,j \le 4}$ + 8 2-loop letters ( $\sqrt{\Delta_5}$  rationalized by mom.twistors).

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- Access to Lorentz-invariant 2-mass pentagon or 1-mass hexagon?