Recent results from proton intermittency analysis in nucleus-nucleus collisions from NA61/SHINE at CERN SPS

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1 QCD Phase Diagram and Critical Phenomena

- 2 Method of intermittency analysis
- Previously released results at 150/158A GeV/c
- 4 New results on Ar+Sc at 150A GeV/c
- 5 Summary and outlook

Phase diagram of QCD

• Objective: Detection / existence of the QCD Critical Point (CP)



• Look for observables tailored for the CP; Scan phase diagram by varying energy and size of collision system.

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Critical Observables; the Order Parameter (OP)



*[Y. Hatta and M. A. Stephanov, PRL91, 102003 (2003)]

Self-similar density fluctuations near the CP



Experimental observation of local, power-law distributed fluctuations Intermittency in transverse momentum space (net protons at mid-rapidity) (Critical opalescence in ion collisions*)

Net proton density carries the same critical fluctuations as the net baryon density, and can be substituted for it.
 [Y. Hatta and M. A. Stephanov, PRL91, 102003 (2003)]

• Furthermore, antiprotons can be dropped to the extent that their multiplicity is much lower than of protons, and proton density analyzed.

[J. Wosiek, Acta Phys. Polon. B 19 (1988) 863-869]
 [A. Bialas and R. Hwa, Phys. Lett. B 253 (1991) 436-438]
 *[F.K. Diakonos, N.G. Antoniou and G. Mavromanolakis, PoS (CPOD2006) 010, Florence]

Observing power-law fluctuations: Factorial moments

- Transverse momentum space is partitioned into M^2 cells
- Calculate second factorial moments F₂(M) as a function of cell size ⇔ number of cells M:

$$F_2(M) \equiv \frac{\left\langle \frac{1}{M^2} \sum_{i=1}^{M^2} n_i(n_i - 1) \right\rangle}{\left\langle \frac{1}{M^2} \sum_{i=1}^{M^2} n_i \right\rangle^2},$$

where $\langle \ldots \rangle$ denotes averaging over events.



Subtracting the background from factorial moments

- Experimental data is noisy ⇒ a background of non-critical pairs must be subtracted at the level of factorial moments.
- Intermittency will be revealed at the level of subtracted moments $\Delta F_2(M)$.

Partitioning of pairs into critical/background

$$\langle n(n-1) \rangle = \underbrace{\langle n_c(n_c-1) \rangle}_{\text{critical}} + \underbrace{\langle n_b(n_b-1) \rangle}_{\text{background}} + \underbrace{2\langle n_b n_c \rangle}_{\text{cross term}}$$

$$\underbrace{\Delta F_2(M)}_{\text{correlator}} = \underbrace{F_2^{(d)}(M)}_{\text{data}} - \lambda(M)^2 \cdot \underbrace{F_2^{(b)}(M)}_{\text{background}} - 2 \cdot \underbrace{\lambda(M)}_{\text{ratio} \frac{\langle n \rangle_b}{\langle n \rangle_d}} \cdot (1 - \lambda(M)) f_{bc}$$

- The cross term can be neglected under certain conditions (non-trivial! Justified by Critical Monte Carlo* simulations)
- * [Antoniou, Diakonos, Kapoyannis and Kousouris, Phys. Rev. Lett. 97, 032002 (2006).]

Scaling of factorial moments - Subtracting mixed events

- For $\lambda \lesssim 1$ (background domination), two approximations can be applied:
- Cross term can be neglected
- On-critical background moments can be approximated by (uncorrelated) mixed event moments; then,

$$\Delta F_2(M) \simeq \Delta F_2^{(e)}(M) \equiv F_2^{\mathsf{data}}(M) - F_2^{\mathsf{mix}}(M)$$

For a critical system, ΔF_2 scales with cell size (number of cells, M) as:

$$\Delta F_2(M) \sim \left(M^2\right)^{\varphi_2}$$

where φ_2 is the intermittency index.

Theoretical prediction for φ_2

$$\begin{cases} \sup_{\substack{v \in V \\ v \in V}} \sup_{\substack{v \in V \\ v \in V}} \begin{cases} \varphi_{2,cr}^{(p)} = \frac{5}{6} (0.833...) \\ \text{net baryons (protons)} \\ [N. G. Antoniou, F. K. Diakonos, A. S. Kapoyannis, K. S. Kousouris, Phys. Rev. Lett. 97, 032002 (2006)] \end{cases}$$

Statistical uncertainties & systematic effect estimation

- Bootstrap method used to calculate statistical uncertainties
- Bootstrap samples of events created by sampling of events with replacement
- $\Delta F_2(M)$ calculated for each bootstrap sample; variance of sample values provides statistical error of $\Delta F_2(M)$

[W.J. Metzger, "Estimating the Uncertainties of Factorial Moments", HEN-455 (2004).]

• Distribution of φ_2 values, $P(\varphi_2)$, and confidence intervals for φ_2 obtained by fitting individual bootstrap samples

[B. Efron, The Annals of Statistics 7,1 (1979)]

- Systematic uncertainties arise from:
 - Misidentification of protons & detector effects (e.g. acceptance)
 - The fact that $F_2(M)$ are correlated for different bin sizes M
 - Selection of *M*-range to fit for power-law
- Bin correlations are partially handled by the bootstrap φ_2 distribution
- Other systematic uncertainties are estimated by varying proton and *M*-range selection

NA49: "C"+C, "Si"+Si, Pb+Pb at 158A GeV/c

- 3 sets of NA49 collision systems were analysed, at 158A GeV/c [T. Anticic et al., Eur. Phys. J. C 75:587 (2015), arXiv:1208.5292v5]
- Factorial moments of proton transverse momenta analyzed at mid-rapidity
- Fragmentation beams used for C and Si ("C"=C,N; "Si"=Si,AI,P) components were merged to enhance statistics



• Fit with $\Delta F_2^{(e)}(M \ ; \ \mathcal{C}, \phi_2) = e^{\mathcal{C}} \cdot \left(M^2\right)^{\phi_2}$, for $M^2 \ge 6000$

• No intermittency detected in the "C"+C, Pb+Pb datasets.

NA49: "C"+C, "Si"+Si, Pb+Pb at 158A GeV/c

• Evidence for intermittency in "Si" + Si – but large statistical errors.



- Bootstrap distribution of ϕ_2 values is highly asymmetric due to closeness of $F_2^{(d)}(M)$ to $F_2^{(m)}(M)$.
- \bullet Based on CMC simulation, we estimate a fraction of $\sim 1\%$ critical protons are present in the sample.
- Estimated intermittency index: $\phi_{2,B} = 0.96^{+0.38}_{-0.25}(\text{stat.}) \pm 0.16(\text{syst.})$ [T. Anticic *et al.*, Eur. Phys. J. C 75:587 (2015), arXiv:1208.5292v5]

NA61/SHINE: Be+Be at 150A GeV/c



- *F*₂(*M*) of data and mixed events overlap ⇒
- Subtracted moments $\Delta F_2(M)$ fluctuate around zero \Rightarrow
- No intermittency effect is observed.
- Preliminary analysis with CMC simulation indicates an upper limit of $\sim 0.3\%$ critical protons [PoS(CPOD2017) 054]

NA61/SHINE: ${}^{40}Ar + {}^{45}Sc$ at 150A GeV/c

- First released results of preliminary analysis in Ar+Sc at 150A GeV/c CPOD 2018.
- Intermittency analysis process:
 - Proton selection via particle energy loss dE/dx
 - Removal of split tracks q_{inv} distribution & cut of proton pairs
 - Probe Δp_T distribution of proton pairs for power-law like behaviour in the limit of small p_T differences
 - Calculate factorial moments $F_2(M)$, $\Delta F_2(M)$ for selected protons
 - Calculate intermittency index ϕ_2 (when possible) & estimate its statistical uncertainty
- Results were obtained for:
 - $\bullet~0\mathchar`-5\%,~5\mathchar`-10\%$ and 10-15% centrality bins
 - 80%, 85% and 90% minimum proton purity selections

Proton selection



- Employ p_{tot} region where Bethe-Bloch bands do not overlap (3.98 GeV/c ≤ p_{tot} ≤ 126 GeV/c)
- Fit dE/dx distribution with 4-gaussian sum for $\alpha = \pi$, K, p, e Bins: p_{tot} , p_T
- 30 Bins in $Log_{10}(p_{tot}): 10^{0.6} \rightarrow 10^{2.1} \text{ GeV/c}$
- 20 Bins in p_T : 0.0 \rightarrow 2.0 GeV/c
- Proton purity: probability for a track to be a proton, $\mathcal{P}_p = p/(\pi + K + p + e)$
- Additional cut along Bethe-Blochs (avoid low-reliability region between p and K curves)

Split tracks & the q_{inv} cut

- Events may contain split tracks: sections of the same track erroneously identified as a pair of tracks that are close in momentum space.
- Three cuts to root them out:
 - Ratio of points / potential points in a track (removes most)
 - Minimum track distance in the detector (pair cut)
 - q_{inv} cut (pair cut, physics-significant)
- q_{inv} distribution of track pairs probed in order to root the rest out: $q_{inv}(p_i, p_j) \equiv \frac{1}{2} \sqrt{-(p_i - p_j)^2}$, p_i : 4-momentum of i^{th} track.
- We calculate the ratio of $q_{inv}^{data}/q_{inv}^{mixed}$.



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Split tracks & the q_{inv} cut

- A peak at low q_{inv} (below 20 MeV/c) indicates a possible split track contamination that must be removed.
- Anti-correlations due to F-D effects and Coulomb repulsion must be removed before intermittency analysis ⇒ "dip" in low q_{inv}, peak predicted around 20 MeV/c [Koonin, PLB 70, 43-47 (1977)]
- Universal cutoff of $q_{inv} > 7 \text{ MeV/c}$ applied to all sets before analysis.



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Δp_T distributions: NA61 data vs EPOS

• Ar+Sc at 150A GeV/c: $\Delta p_T = 1/2 \sqrt{(p_{X_1} - p_{X_2})^2 + (p_{Y_1} - p_{Y_2})^2}$

distributions of protons selected for intermittency analysis



• In NA61 data, we see strong correlations in $\Delta p_T \to 0 \Rightarrow$ indication of intermittent behaviour

Δp_T distributions & $F_2(M)$: NA61 data vs EPOS



NA61/SHINE: Ar+Sc at 150A GeV/c: $F_2(M)$



NA61/SHINE: Ar+Sc at 150A GeV/c: $\Delta F_2(M)$



NA61/SHINE: Ar+Sc at 150A GeV/c: ϕ_2 bootstrap dist.



NA61/SHINE: Ar+Sc at 150A GeV/c: Summary



Ar+Sc EPOS: $F_2(M)$, $\Delta F_2(M)$, ϕ_2 bootstrap distribution



Intermittency analysis at 150/158A GeV/c: Summary



- Indication of intermittency effect in middle-central NA61/SHINE Ar+Sc collisions
- First possible evidence of CP signal in NA61/SHINE
- Effect quality increases with increased proton purity selection, up to 90% proton purity; EPOS does not reproduce observed effect.

Outlook



• Expanding the analysis to other NA61/SHINE systems (Xe+La, Pb+Pb) and SPS energies (Ar+Sc) will hopefully lead to a more reliable interpretation of the observed intermittency signal in terms of the critical point.

Thank you!

Acknowledgements

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Scaling of factorial moments - Subtracting mixed events

For $\lambda \lesssim 1$ (background domination), $\Delta F_2(M)$ can be approximated by:

$$\Delta F_2^{(e)}(M) = F_2^{\mathsf{data}}(M) - F_2^{\mathsf{mix}}(M)$$

For a critical system, ΔF_2 scales with cell size (number of cells, M) as:

$$\Delta F_2(M) \sim \left(M^2\right)^{\varphi_2}$$

where φ_2 is the intermittency index.

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Critical Monte Carlo (CMC) algorithm for baryons

• Simplified version of CMC* code:

- Only protons produced
- One cluster per event, produced by random Lévy walk:

$$\tilde{d}_F^{(B,2)} = 1/3 \Rightarrow \phi_2 = 5/6$$

- Lower / upper bounds of Lévy walks *p_{min,max}* plugged in.
- Cluster center exponential in *p*_T, slope adjusted by *T*_c parameter.
- Poissonian proton multiplicity distribution.



Input parameters

Parameter	$p_{\min}\left(MeV\right)$	p_{\max} (MeV)	$\lambda_{Poisson}$	T_c (MeV)
Value	0.1 ightarrow 1	$800 \rightarrow 1200$	$\langle p angle_{non-empty}$	163

* [Antoniou, Diakonos, Kapoyannis and Kousouris, Phys. Rev. Lett. 97, 032002 (2006).]

NA61/SHINE data analysis – ${}^{40}Ar + {}^{45}Sc$ at 150A GeV/c

- NA49 analysis encourages us to look for intermittency in medium-sized nuclei, in the NA61 experiment.
- Intermittency analysis requires:
 - Large event statistics $\Rightarrow \sim 100 K$ events min., ideally $\sim 1 M$ events.
 - Reliable particle ID \Rightarrow proton purity should be $\sim 80\% 90\%$.
 - Central collisions.
 - Adequate mean proton multiplicity in midrapidity (≥ 2)
- A preliminary analysis for Be+Be data at 150A GeV/c was previously performed [PoS(CPOD2017) 054]; no intermittency signal was observed.
- We now expand on it with our preliminary analysis in Ar+Sc at 150A GeV/c.
- Simulation through EPOS* (detector effects included) would suggest:

$$\left.\frac{dN_p}{dy}\right|_{|y_{CM}| \le 0.75, \, p_T \le 1.5} \sim 1.6 - 2$$

for $\sim 0-15\%$ centrality; adequate for an intermittency analysis

• We perform a 2D scan in proton purity (80-90%) and centrality of collisions

*[K. Werner, F. Liu, and T. Pierog, Phys. Rev. C 74, 044902 (2006)]

- Production used: Ar_Sc_150_15/026_17c_v1r8p0_pA_slc6_phys_PP (miniSHOE, unofficial)
- Runs: 20328 20345 , 20368 20380
- Bad runs rejected almost 2/3rds of total!
- miniSHOE files with Potential Point information provided by B. Maksiak not an official production yet
- SHINE code to select events (primary vertex charged particles)
- Event & Track cuts based on Maciej Lewicki's and Michal Naskret's h^- analysis.
- Non-bias event cuts: used Andrey Seryakov's NonBiasEventCutsArSc class.
- 0%-20% most central events in 5% bin intervals selected via cut in energy sum of PSD selected modules (based on Andrey Seryakov's Moscow meeting presentation on centrality determination).

Ar+Sc at 150A GeV/c: NA61 data vs EPOS

EPOS – proton p_T statistics

Centrality	#events	$\left_{ p_{\mathcal{T}} \leq 1.5} c$ Non-empty	GeV, _{YCM} ≤0.75 With empty	$\Delta p_{x,y}$
0- 5%	293,412	3.06 ± 1.60	2.89 ± 1.70	0.35 - 0.43
5-10%	252,362	2.72 ± 1.45	2.49 ± 1.58	0.35 - 0.43
10-15%	274,072	$\textbf{2.45} \pm \textbf{1.33}$	2.16 ± 1.48	0.35 - 0.43

$^{40}Ar + ^{45}Sc$ NA61 data – proton p_T statistics

Centrality	#events	$\left_{ p_T \le 1.5} c$ Non-empty	GeV, _{YCM} ≤0.75 With empty	$\Delta p_{x,y}$
0- 5%	144,362	$\textbf{3.44} \pm \textbf{1.79}$	$\textbf{3.30} \pm \textbf{1.89}$	0.46 - 0.58
5-10%	148,199	3.00 ± 1.61	2.79 ± 1.73	0.46 - 0.58
10-15%	142,900	2.81 ± 1.53	2.58 ± 1.66	0.45 - 0.57

$p_{X,Y}$ spectra comparison – NA61 vs EPOS (0 - 15%)



Event & Track cuts

Event cuts

- Target IN/OUT,
- BPD status,
- WFA particles (4.5 µs),
- WFA interaction (25 µs),
- BPD3X(Y) charge,
- S5 (0
 ightarrow 170),
- T2 trigger (eAll),
- Vertex track fitted to the main vertex,
- Vertex fit quality = ePerfect,
- $\bullet\,$ Fitted vertex position -580 ± 10 cm,
- PSD Module Energy Sum cut (inner/outer),
- Centrality 0-20% (based on PSD)
- nTracksFit/nTracksAll > 0.25 if nTracksFit ≤ 50 (Andrey)

Track cuts

- Track status,
- Charge ± 1 ,
- Impact point [±4cm; ±2cm],
- Total number of clusters \geq 30,
- VTPCs clusters \geq 15,
- NO GTPC clusters,
- dE/dx clusters \geq 30,
- $0.5 \le \frac{\#Points}{\#Potential Points} \le 1.0$
- TTD cut $> 2 \ cm$
- dE/dx \leq 1.8 (dE/dx fit issue)
- proton selection (scan)
- 3.98 GeV/c $\leq p_{tot} \leq$ 126 GeV/c (for dE/dx proton ID - scan)

$^{40}Ar + {}^{45}Sc - EPOS MC production overview$

- Production used: Simulation/Ar_Sc_150_15/ 15_011_v14e_v1r2p0_pA_slc6_phys/EPOS_with_potential_points/
- \bullet An estimated ${\sim}300K$ simulated events per 5% centrality bin.
- Potential Point information included for limited events subset.
- SHINE code to select events (primary vertex charged particles)
- Event & Track cuts (hastily) adapted to match Ar+Sc @150 data analysis (where applicable).
- No PSD simulation centrality selection based on # of forward spectators, nFSpec = 40 simEvent.GetPrimaryInteraction().GetProjectileParticipants()

(see Andrey Seryakov's centrality determination information on twiki).

Event & Track cuts – EPOS

Event cuts

- Target IN/OUT,
- BPD status,
- Vertex track fitted to the main vertex,
- Vertex fit quality = ePerfect,
- \bullet Fitted vertex position -580 ± 10 cm,
- Centrality 10% (based on nFSpec)

Track cuts

- Track status,
- Charge ± 1 ,
- Impact point [±4cm; ±2cm],
- Total number of clusters \geq 30,
- VTPCs clusters \geq 15,
- NO GTPC clusters,
- TTD cut > 2 cm,
- proton selection matching closest simTrack,
- 3.98 GeV/c $\leq p_{tot} \leq$ 126 GeV/c (to match effect of dE/dx p_{tot} cut),
- acceptance cut

Centrality selection via # forward spectators

• A probabilistic selection based on nFSpec percentiles used to select centrality bin.



- A discrepancy observed in multiplicity distribution between data (above) & EPOS (below).
- Acceptance cut fixes the problem.



All Event cuts – statistics



Centrality – statistics



Track cuts – statistics



Cuts (plots)



Cuts (plots)

PSD energy, mods: 1-16,21,22,27,28



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nP/nPP ratio

dE/dx vs p_{tot} (proton ID)



- Avoid p_{tot} region where Bethe-Bloch curves overlap (3.98 GeV/c $\leq p_{tot} \leq 126$ GeV/c)
- Using Hans Dembinski/Raul R Prado's dE/dx fitting software Bins: ptot, pT
- Presented in Moscow meeting by Prado, Herve & Unger
- 30 Bins in $Log_{10}(p_{tot}): 10^{0.6} \rightarrow 10^{2.1} \text{ GeV/c}$
- 20 Bins in p_T : 0.0 \rightarrow 2.0 GeV/c
- Preliminary p selection: 90% purity removing deuterons from the model
- Cut along Bethe-Blochs: $BB_p + 0.15(BB_K BB_p)$

dE/dx simulation & proton purity assignment in EPOS



- Used *dE/dx* spectra from Ar+Sc @150 data in the 6% 18% centrality interval
- For each track, assign a dE/dx value based on particle species and phase space bin
- Apply dE/dx & purity cuts identical to NA61/SHINE data

Improving calculation of $F_2(M)$ via lattice averaging

- Problem: With low statistics/multiplicity, lattice boundaries may split pairs of neighboring points, affecting $F_2(M)$ values (see example below).
- Solution: Calculate moments several times on different, slightly displaced lattices (see example)
- Average corresponding $F_2(M)$ over all lattices. Errors can be estimated by variance over lattice positions.
- Lattice displacement is larger than experimental resolution, yet maximum displacement must be of the order of the finer binnings, so as to stay in the correct p_T range.



Improved confidence intervals for ϕ_2 via resampling

- In order to estimate the statistical errors of $\Delta F_2(M)$, we need to produce variations of the original event sample. This, we can achieve by using the statistical method of resampling (bootstrapping) \Rightarrow
 - Sample original events with replacement, producing new sets of the same statistics (# of events)
 - Calculate $\Delta F_2(M)$ for each bootstrap sample in the same manner as for the original.
 - The variance of sample values provides the statistical error of $\Delta F_2(M)$.

[W.J. Metzger, "Estimating the Uncertainties of Factorial Moments", HEN-455 (2004).]

• Furthermore, we can obtain a distribution $P(\varphi_2)$ of φ_2 values. Each bootstrap sample of $\Delta F_2(M)$ is fit with a power-law:

$$\Delta F_2(M; \mathcal{C}, \varphi_2) = e^{\mathcal{C}} \cdot (M^2)^{\varphi_2}$$

and we can extract a confidence interval for φ_2 from the distribution of values. [B. Efron, *The Annals of Statistics* 7,1 (1979)]

Split tracks; the q_{inv} cut in analysed datasets

- Split tracks can create false positive for intermittency ⇒ must be reduced or removed.
- q_{inv} -test distribution of track pairs: $q_{inv}(p_i, p_j) \equiv \frac{1}{2}\sqrt{-(p_i p_j)^2}$, p_i : 4-momentum of i^{th} track.
- Calculate ratio $q_{inv}^{data}/q_{inv}^{mixed} \Rightarrow \text{peak at low } q_{inv}$ (below 20 MeV/c): possible split track contamination.



• Anti-correlations due to F-D effects and Coulomb repulsion must be removed before intermittency analysis \Rightarrow "dip" in low q_{inv} , peak predicted around 20 MeV/c [Koonin, PLB 70, 43-47 (1977)]

• Universal cutoff of $q_{inv} > 25$ MeV/c applied to all sets before analysis.

NA49 analysis – Δp_T distributions

• We measure correlations in relative p_T of protons via

$$\Delta p_T = 1/2 \sqrt{(p_{X_1} - p_{X_2})^2 + (p_{Y_1} - p_{Y_2})^2}$$



- Strong correlations for $\Delta p_T \rightarrow 0$ indicate power-law scaling of the density-density correlation function \Rightarrow intermittency presence
- We find a strong peak in the "Si" +Si dataset
- A similar peak is seen in the Δp_T profile of simulated CMC protons with the characteristics of "Si" +Si.

- Events may contain split tracks: sections of the same track erroneously identified as a pair of tracks that are close in momentum space.
- Intermittency analysis is based on pairs distribution \Rightarrow split tracks can create a false positive, and so must be reduced or removed.
- Standard cuts remove part of split tracks. In order to estimate the residual contamination, we check the *q_{inv}* distribution of track pairs:

$$q_{inv}(p_i,p_j) \equiv \frac{1}{2}\sqrt{-(p_i-p_j)^2},$$

 p_i : 4-momentum of i^{th} track.

We calculate the ratio of q^{data}_{inv} / q^{mixed}_{inv}. A peak at low q_{inv} (below 20 MeV/c) indicates a possible split track contamination that must be removed.

Noisy CMC (baryons) - estimating the level of background

- $F_2(M)$ of noisy CMC approximates "Si" +Si for $\lambda \approx 0.99$
- ΔF₂^(e)(M) reproduces critical behaviour of pure CMC, even though their moments differ by orders of magnitude!



• Noisy CMC results show our approximation is reasonable for dominant background.

q_{inv} proton distributions – NA61/SHINE





Δp_T proton distributions – NA61/SHINE







$q_{inv} \& \Delta p_T$ distributions – EPOS

