

Flavourful Z' models for RK(*) Steve King

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Corfu Summer Institute

17th Hellenic School and Warkshops on Elementary Particle Physics and Gravi Corfy, Greece 201



B physics anomalies: R_K, R_{K*}

$$R_K = \frac{\text{BR}\left(B^+ \to K^+ \mu^+ \mu^-\right)}{\text{BR}\left(B^+ \to K^+ e^+ e^-\right)} = 0.745 \pm 0.09_{\text{stat}} \pm 0.036_{\text{syst}}$$

$$R_{K^*} = \frac{\mathrm{BR}(B \to K^* \mu^+ \mu^-)}{\mathrm{BR}(B \to K^* e^+ e^-)} = \begin{cases} 0.660^{+0.110}_{-0.070} \pm 0.024 & (2m_\mu)^2 < q^2 < 1.1 \,\mathrm{GeV}^2 \\ 0.685^{+0.113}_{-0.069} \pm 0.047 & 1.1 \,\mathrm{GeV}^2 < q^2 < 6 \,\mathrm{GeV}^2 \end{cases}$$

New physics in μ 1704.05438



Best fit operator

$$V_{tb}V_{ts}^* \frac{\alpha_{em}}{4\pi v^2} \left(C_{b_L\mu_L}^{\rm SM} + C_{b_L\mu_L}^{\rm BSM} \right) \bar{b}_L \gamma^{\mu} s_L \bar{\mu}_L \gamma_{\mu} \mu_L$$

$$8.64 \quad -1.3$$

$$V_{tb}V_{ts}^* \frac{\alpha_{em}}{4\pi v^2} \approx \frac{1}{(36 \text{ TeV})^2}$$

SM penguins are universal



BSM should be non-universal



Z' models



$$G_{b_L \mu_L}^{\text{BSM}} = g_{b_L}^{Z'} g_{\mu_L}^{Z'} \left(\frac{g'^2}{M_{Z'}^2}\right) \approx -\frac{1}{(33 \text{ TeV})^2}$$

Typically small so Z^\prime mass in LHC range

SFK 1706.06100 Flavourful Z' models

	Representation/charge				
Field	$SU(3)_c$ $SU(2)_L$		$U(1)_Y$	U(1)'	
Q_{Li}	3	2	1/6	q_{Q_i}	
u_{Ri}	3	1	2/3	q_{u_i}	
d_{Ri}	3	1	-1/3	q_{d_i}	
L_{Li}	1	2	-1/2	q_{L_i}	
e_{Ri}	1	1	-1	q_{e_i}	
$ u_{Ri}$	1	1	0	$q_{ u_i}$	

	Representation/charge				
Field	$SU(3)_c$	$SU(2)_L$	$U(1)_Y$	U(1)'	
$Q_{L4}, ilde{Q}_{R4}$	3	2	1/6	q_{Q_4}	
$u_{R4}, ilde{u}_{L4}$	3	1	2/3	q_{u_4}	
$d_{R4},\! ilde{d}_{L4}$	3	1	-1/3	q_{d_4}	
$L_{L4}, ilde{L}_{R4}$	1	2	-1/2	q_{L_4}	
$e_{R4}, \! ilde{e}_{L4}$	1	1	-1	q_{e_4}	
${ u}_{R4}, { ilde u}_{L4}$	1	1	0	$q_{{m u}_4}$	

Three families of quarks and leptons with universal charges

Vector-like fourth family with non-universal U(1)' charges $q_1 = q_2 = q_3 \neq q_4$

When U(1)' is broken, mass mixing of the four families leads to Z' flavour violation This mechanism works for any Z' model



Light quarks contain admixtures of non-universal $\,Q_4\,$

After mass matrix is diagonalised

$$Q'_{\alpha} = V_Q^{\alpha\beta} Q_{\beta}$$
 4x4

Non-universal and flavour changing couplings $V_Q D_Q V_Q^{\dagger}$

non-universal charge matrix

 $D_Q = \operatorname{diag}(q_1, q_2, q_3, q_4)$

 Q'_{β}

Three light quarks

One heavy vector-like quark

$$Q'_{\alpha=1,2,3}$$
 contain some Q_4

 $\tilde{M}_4^Q \overline{Q'}_4 \tilde{Q}_4$

$$Q'_{\alpha=4}$$

QB UB Ua (UB $V_Q D_Q V_Q^{\intercal}$ $\overline{V_u}\overline{D_u}\overline{V_u^\dagger}$ $V_d D_d V_d^{\intercal}$ ea eB ß 4x4 mixing matrices V_Q, V_u, V_d, V_L, V_e $V_L D_L V_L^{\intercal}$ $V_e D_e V_e$ non-universal charge matrices $D_Q, \overline{D_u}, \overline{D_d}, \overline{D_L}, \overline{D_e}$

SFK 1706.06100 SO(10) GUT model $SO(10) \rightarrow SU(5) \times U(1)_X$ Broken at TeV scale Broken at GUT scale

$16_{Fi} \rightarrow (10, 1)_i + (\overline{5}, -3)_i + (1, 5)_i$

Three families of quarks and leptons

 $10_F \rightarrow (5, -2) + (\overline{5}, 2),$ $45_F \rightarrow (10, -4) + (\overline{10}, 4) + (1, 0) + (24, 0)$

> TeV scale vector-like fourth family with non-universal charges

GUT scale masses splitting requires 210_H

SFK 1706.06100 Flavourful Z' from SO(10)

 Q_4

 \tilde{u}_4^c

 \tilde{d}_4^c

 \tilde{L}_4

 \tilde{e}_4^c

	Representation				Representation			
Field	$SU(3)_c$	$SU(2)_L$	$U(1)_Y$	$U(1)_X$	Field	$SU(3)_c$	$SU(2)_L$	U(1)
$egin{array}{c} Q_i \ u_i^c \ d_i^c \ L_i \ o^c \end{array}$	$\frac{3}{\overline{3}}$ $\frac{1}{1}$	2 1 1 2	$1/6 \\ -2/3 \\ 1/3 \\ -1/2 \\ 1$	1 1 -3 -3	$egin{array}{c} Q_4 \ u_4^c \ d_4^c \ L_4 \end{array}$	$\begin{array}{c} 3\\\overline{3}\\\overline{3}\\\overline{3}\\1\end{array}$	2 1 1 2	$1/6 \\ -2/3 \\ 1/3 \\ -1/2$
$e_i u_i^c$	1 1	1	1	1 5	e_4^c	$\frac{1}{2}$	1 	1

Three families of quarks and leptons with universal charges

 $\overline{\mathbf{5}} \to L, d^c, \quad \mathbf{10} \to Q, u^c, e^c$

Vector-like fourth family with non-universal $U(I)_X$ charges

1

 $\overline{2}$

3

3

 $U(1)_Y$

1/6

1/3

-1/2

1 2/3

-1/6

-1/3

1/2

_1

-2/3

 $U(1)_{X}$

-4

-4

2

2

-4

4

-2

-2

4

$$\mathcal{L}_{Z'}^{q} = g' Z'_{\mu} \left(\overline{u}_{L} \quad \overline{c}_{L} \quad \overline{t}_{L} \right) V'_{uL} \tilde{D}'_{Q} V'^{\dagger}_{uL} \gamma^{\mu} \begin{pmatrix} u_{L} \\ c_{L} \\ t_{L} \end{pmatrix} + \text{other quarks and} \\ + g' Z'_{\mu} \left(\overline{d}_{L} \quad \overline{s}_{L} \quad \overline{b}_{L} \right) V'_{dL} \tilde{D}'_{Q} V'^{\dagger}_{dL} \gamma^{\mu} \begin{pmatrix} d_{L} \\ s_{L} \\ b_{L} \end{pmatrix} \\ \mathsf{CKM mixing} \qquad (\tilde{D}'_{Q})_{ij} = (V_{Q_{L}} D_{Q} V^{\dagger}_{Q_{L}})_{ij}$$

Ν

quarks

$$V_{Q_{L}} = V_{34}^{Q_{L}} V_{24}^{Q_{L}} V_{14}^{Q_{L}}, \qquad V_{34}^{Q_{L}} = \begin{pmatrix} 1 & 0 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 & 0 \\ 0 & 0 & c_{34}^{Q_{L}} & s_{34}^{Q_{L}} e^{-i\delta_{34}^{Q_{L}}} \\ 0 & 0 & -s_{34}^{Q_{L}} e^{i\delta_{34}^{Q_{L}}} & c_{34}^{Q_{L}} \end{pmatrix},$$

Vector-like mass mixing matrices

$$P_{24}^{Q_{L}} = \begin{pmatrix} 1 & 0 & 0 & 0 & 0 \\ 0 & c_{24}^{Q_{L}} & 0 & s_{24}^{Q_{L}} e^{-i\delta_{24}^{Q_{L}}} \\ 0 & 0 & 1 & 0 & 0 \\ 0 & -s_{24}^{Q_{L}} e^{i\delta_{24}^{Q_{L}}} & 0 & c_{24}^{Q_{L}} \end{pmatrix}, \quad V_{14}^{Q_{L}} = \begin{pmatrix} c_{14}^{Q_{L}} & 0 & 0 & s_{14}^{Q_{L}} e^{-i\delta_{14}^{Q_{L}}} \\ 0 & 1 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 & 0 \\ -s_{14}^{Q_{L}} e^{i\delta_{14}^{Q_{L}}} & 0 & 0 & c_{14}^{Q_{L}} \end{pmatrix}.$$

 V_{i}

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$$SO(10) \text{ for } R_{K}(*)$$
Assuming only θ_{34}^{Q} and θ_{14}^{L} non-zero
Universal Non-universal

$$\vec{D}_{L}^{\prime} = -\frac{3}{2\sqrt{10}} \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix} + \frac{5}{2\sqrt{10}} \begin{pmatrix} s_{14}^{L} \rangle^{2} & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}$$

$$\vec{D}_{Q}^{\prime} = \frac{1}{2\sqrt{10}} \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix} - \frac{5}{2\sqrt{10}} \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & (s_{34}^{Q})^{2} \end{pmatrix}$$

$$\vec{T}_{bL\mu L}^{\text{BSM}} \ \vec{b}_{L} \gamma^{\lambda} s_{L} \left[\bar{\mu}_{L} \gamma_{\lambda} \mu_{L} + \left(1 - \frac{5}{3} (s_{14}^{L})^{2}\right) \bar{e}_{L} \gamma_{\lambda} e_{L} + \frac{1}{3} \bar{\mu}_{R} \gamma_{\lambda} \mu_{R} + \frac{1}{3} \bar{e}_{R} \gamma_{\lambda} e_{R} + \cdots \right]$$





Antusch, Hohl, SFK, Susic (in prep) SO(I0) GUT model Singlet for inverse/ Continued... linear seesaw $\mu \mathbf{1}_F^2 + m_{10} \mathbf{10}_F^2 + m_{45} \mathbf{45}_F^2 \longleftarrow$ Vector-like fermion masses $+ y_1 \mathbf{16}_F^2 \cdot \mathbf{10}_H \quad \longleftarrow \text{ Quark and lepton Yukawa couplings}$ $+Y_1 \mathbf{16}_F \cdot \mathbf{1}_F \cdot \mathbf{16}_H + Y_2 \mathbf{16}_F \cdot \mathbf{10}_F \cdot \mathbf{16}_H + Y_3 \mathbf{16}_F \cdot \mathbf{45}_F \cdot \mathbf{16}_H$ Quark/lepton mixing with Singlet-neutrino vector-like fermions couplings

Antusch, Hohl, SFK, Susic (in prep) nverse vs linear seesaw
$$\begin{split} m_{\nu}^{inv} &\sim \mu \left(\frac{y_1 v_u}{Y_1 \overline{v}_R}\right)^2 \frac{\left(\nu_{Li}, \nu_{Ri}, S_i, S_i, \nu_{L4}, \tilde{\nu}_{R4}\right)}{\underset{w_{\nu}}{\text{Dirac}} \underset{V_1 v_u}{\overset{\text{Linear}}{\text{V}_1 \overline{v}_L}} \\ m_{\nu}^{lin} &\sim \frac{y_1 v_u \overline{v}_L}{\overline{v}_R} \\ M_{\nu} &= \begin{pmatrix} 0 & y_1 v_u & \Gamma_{1} \overline{v}_L \\ y_1 v_u & 0 & Y_1 \overline{v}_L \\ Y_1 \overline{v}_L & Y_1 \overline{v}_R \\ \mu \end{pmatrix} \overset{\text{Nverse}}{\overset{\text{Dirac}}{\text{Dirac}}} \begin{pmatrix} 0 & Y_2 v_R \\ 0 & Y_2 v_L \\ 0 & 0 \end{pmatrix} \end{split}$$
 $\overline{16}_H \to (\overline{10}, -1)_H + (5, 3)_H + (1, -5)_H$

Antusch, Hohl, SFK, Susic (in prep) Lepton flavour violation



$$\begin{aligned} &\operatorname{Br}(\mu \to 3e) \approx \frac{1}{8^2} \cdot (\theta_{12}^e)^2 (s_{14}^L)^4 \cdot \frac{g'^4 M_W^4}{g_2^4 M_{Z'}^4} \\ &\operatorname{Near \ current \ bound \ | \ 0^{-12}}_{&\approx 1.64 \cdot 10^{-12}} \left(\frac{\theta_{12}^e}{3^\circ}\right)^2 \left(\frac{3.2 \cdot 10^3 \ \operatorname{GeV}}{M_{Z'}}\right)^4 \end{aligned}$$

Crispim Romao, SFK, Leontaris (in prep) F-theory Z' models $SU(5)_{GUT}$ **E**8 Higgs 10_{M} SU(5)GUT X SU(5)Perp Higgs SU(5)_{GUT} x U(1)⁴_{Perp} Heckman, Tavanfar, Vafa Flux | Z₂ monodromy $SU(3)xSU(2)xU(1)_{Y} \times U(1)^{3}_{Perp}$ Singlets 4 Non-universal $SU(3)xSU(2)xU(1)_{Y}xU(1)_{X}$

Crispim Romao, SFK, Leontaris (in prep) F-theory Z'models Three chiral families with non-universal Z'

Model 1: $(H_u)_{2\frac{1}{\sqrt{15}}} + (H_d)_{-2\frac{1}{\sqrt{15}}} + 3 \times \overline{5}_{\frac{1}{2}\frac{1}{\sqrt{15}}} + 2 \times 10_{-\frac{1}{\sqrt{15}}} + 10_{\frac{3}{2}\frac{1}{\sqrt{15}}}$ Model 2: $(H_u)_{\frac{1}{2}\frac{1}{\sqrt{15}}} + (H_d)_{-\frac{1}{2}\frac{1}{\sqrt{15}}} + 2 \times \overline{5}_{2\frac{1}{\sqrt{15}}} + \overline{5}_{-\frac{7}{4}\frac{1}{\sqrt{15}}} + 3 \times 10_{-\frac{1}{4}\frac{1}{\sqrt{15}}}$ Model 3: $(H_u)_{\frac{3}{2}\frac{1}{\sqrt{10}}} + \overline{5}_{-\frac{1}{4}\frac{1}{\sqrt{10}}} + \overline{5}_{\frac{1}{\sqrt{10}}} + (2L + d^c)_{-\frac{3}{2}\frac{1}{\sqrt{10}}} + 2 \times 10_{-\frac{3}{4}} + 10_{\frac{7}{4}\frac{1}{\sqrt{10}}}$ Model 4: $(H_u)_{\frac{3}{2}\frac{1}{\sqrt{10}}} + (H_d)_{-\frac{3}{2}\frac{1}{\sqrt{10}}} + \overline{5}_{-\frac{1}{4}\frac{1}{\sqrt{10}}} + \overline{5}_{\frac{1}{\sqrt{10}}} + \overline{5}_{\frac{9}{4}\frac{1}{\sqrt{10}}} + 2 \times 10_{-\frac{3}{4}\frac{1}{\sqrt{10}}} + 10_{\frac{1}{2}\frac{1}{\sqrt{10}}}$

Three chiral plus one vector-like family w/ non-universal Z'

Example 1: $(H_u)_{4\frac{1}{\sqrt{85}}} + (H_d)_{-4\frac{1}{\sqrt{85}}} + 3 \times \overline{5}_{\frac{7}{2}\frac{1}{\sqrt{85}}} + \overline{5}_{\frac{3}{2}\frac{1}{\sqrt{85}}} + 5_{6\frac{1}{\sqrt{85}}} + 3 \times 10_{2\frac{1}{\sqrt{85}}} + 10_{-\frac{11}{2}\frac{1}{\sqrt{85}}} + \overline{10}_{\frac{1}{2}\frac{1}{\sqrt{85}}}$ Example 2: $(H_u)_{\frac{3}{2}\frac{1}{\sqrt{10}}} + (H_d)_{-\frac{3}{2}\frac{1}{\sqrt{10}}} + \overline{5}_{\frac{1}{\sqrt{10}}} + \overline{5}_{-\frac{1}{4}\frac{1}{\sqrt{10}}} + \overline{5}_{-\frac{9}{4}\frac{1}{\sqrt{10}}} + L_{\frac{1}{\sqrt{10}}} + d_{-\frac{1}{4}\frac{1}{\sqrt{10}}}^c + 5_{-\frac{1}{\sqrt{10}}} + 3 \times 10_{-\frac{3}{4}\frac{1}{\sqrt{10}}} + 10_{-\frac{3}{4}\frac{1}{\sqrt{10}}} + \overline{10}_{-\frac{1}{2}\frac{1}{\sqrt{10}}}$ Example 3: $(H_u)_{-\frac{1}{2}} + (H_d)_{\frac{1}{2}} + 3 \times \overline{5}_{-\frac{1}{4}} + \overline{5}_{\frac{3}{4}} + 5_0 + 3 \times 10_{\frac{1}{4}} + 10_{-\frac{1}{2}} + \overline{10}_{\frac{1}{4}}$

Conclusion

- R_{K(*)} motivates non-universal Z' models
- Any Z' model can be made non-universal by adding fourth vector-like family with non-universal charges which mixes
- We studied an SO(10) GUT example with U(1)_X at TeV and a vector-like family coming from 45_F and 10_F
- SO(10) GUT for $R_{K(*)}$ requires Z' near 3 TeV LHC bound with mu->eee near current bound 10⁻¹²
- SO(10) with $U(1)_X$ at TeV requires inverse/linear seesaw
- F-theory for non-universal Z' with optional vector-like family

